

# Obtaining a nonzero $\theta_{13}$ in lepton models based on $SO(3) \rightarrow A_4$

Yuval Grossman and Wee Hao Ng,  
Cornell University

[arXiv:1404.1413 \[hep-ph\]](https://arxiv.org/abs/1404.1413)

Phenomenology 2014

# Outline

- Tri-bimaximal mixing found and lost; how do  $A_4$  models cope?
- Basics of  $SO(3) \rightarrow A_4$  model
- Nonzero  $\theta_{13}$  in  $SO(3) \rightarrow A_4$  model
- Summary

# Tri-bimaximal mixing and $A_4$ models

- For long time, PMNS matrix thought to be consistent with TBM

$$U_{\text{TBM}} = \begin{pmatrix} \frac{\sqrt{2}}{\sqrt{3}} & \frac{1}{\sqrt{3}} & 0 \\ -\frac{1}{\sqrt{6}} & \frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{6}} & -\frac{1}{\sqrt{3}} & \frac{1}{\sqrt{2}} \end{pmatrix} \leftarrow \sin(\theta_{13}) e^{i\delta_{13}}$$

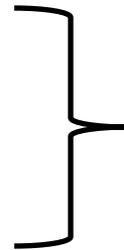
- Highly specific pattern... discrete lepton flavor symmetries?

# Tri-bimaximal mixing and $A_4$ models

- Simplest  $A_4$  lepton models

- Components

- SM Higgs and leptons
    - Right handed neutrinos
    - Scalar flavons  $\phi, \phi'$



Representations of  $A_4$

- Features

- Additional  $Z_2$  symmetry so flavons sectorized
    - Flavons gain VEV  $\langle \phi \rangle = (v, v, v), \langle \phi' \rangle = (v', 0, 0)$

**$\Rightarrow$  TBM!**

# Tri-bimaximal mixing... lost

- Recently,  $\theta_{13}$  found to be much larger!

- Daya Bay (2012):

$$\sin^2(2\theta_{13}) = 0.092 \pm 0.016(\text{stat}) \pm 0.005(\text{syst})$$

- RENO (2012):

$$\sin^2(2\theta_{13}) = 0.113 \pm 0.013(\text{stat}) \pm 0.019(\text{syst})$$

- TBM ruled out!

# How do $A_4$ models cope?

- Actually many  $A_4$  models already allow this!
- Approaches:
  - Higher dimension operators, e.g. [Altarelli and Meloni (2009)]
  - More flavons, e.g. [Chen et al (2013)]
  - Perturb flavon alignments, e.g. [King (2011)]
  - Radiative corrections, e.g. [Antusch et al (2003)]

# Continuous symmetry $\rightarrow A_4$

- Class of models based on continuous symmetry  $\rightarrow A_4$ .
- Originally motivated to explain origin of  $A_4$ .
- Specific example:  $SO(3) \rightarrow A_4$  [Berger and Grossman (2009)]
- Turns out  $\theta_{13}$  already nonzero at tree level!
- Even better:  $\theta_{13}$  related to  $\frac{\cancel{A_4} \text{ scale}}{\cancel{SO(3)} \text{ scale}}$

# Basics of $SO(3) \rightarrow A_4$ model

| Field    | $SU(2)_L$ | $U(1)_Y$       | $SO(3)_F$ | $Z_2$ |                        |
|----------|-----------|----------------|-----------|-------|------------------------|
| $\psi_l$ | <b>2</b>  | $-\frac{1}{2}$ | <b>3</b>  | -     | } Left-handed leptons  |
| $\psi_f$ | <b>1</b>  | -1             | <b>3</b>  | -     |                        |
| $\psi_e$ | <b>1</b>  | -1             | <b>1</b>  | +     | } Right-handed leptons |
| $\psi_m$ | <b>1</b>  | -1             | <b>5</b>  | +     |                        |
| $\psi_n$ | <b>1</b>  | 0              | <b>3</b>  | -     |                        |
| $H$      | <b>2</b>  | $\frac{1}{2}$  | <b>1</b>  | +     | } Scalars              |
| $\phi$   | <b>1</b>  | 0              | <b>3</b>  | -     |                        |
| $\phi'$  | <b>1</b>  | 0              | <b>3</b>  | +     |                        |
| $\phi_5$ | <b>1</b>  | 0              | <b>5</b>  | -     |                        |
| $T$      | <b>1</b>  | 0              | <b>7</b>  | -     |                        |

# Basics of $SO(3) \rightarrow A_4$ model

- Most general Lagrangian for leptons

$$\begin{aligned}
 L_l = & -y_e \bar{\Psi}_l^a \frac{H}{\Lambda} \phi^a \psi_e - y_m \bar{\Psi}_l^a \frac{H}{\Lambda} \phi^b \psi_m^{ab} - y_m^T \bar{\Psi}_l^a \frac{H}{\Lambda} T^{abc} \psi_m^{bc} - y_m^5 \bar{\Psi}_l^a \frac{H}{\Lambda} \phi_5^{bd} \psi_m^{cd} \\
 & -y_e' \bar{\Psi}_f^a \phi^a \psi_e - y_m' \bar{\Psi}_f^a \phi^b \psi_m^{ab} - y_m^{T'} \bar{\Psi}_f^a T^{abc} \psi_m^{bc} - y_m^{5'} \bar{\Psi}_f^a \phi_5^{bd} \psi_m^{cd} \\
 L_\nu = & -M \bar{\Psi}_n^{ca} \Psi_n^a - \frac{X_\nu}{\Lambda} \bar{\Psi}_n^{ca} \Psi_n^b \phi'^c T^{abc} - y_\nu \bar{\Psi}_l^a H \Psi_n^a
 \end{aligned}$$

- Flavons gain VEV

3D coordinate basis vectors 

$$\langle T \rangle = v_T x^a y^b z^c, \quad \langle \phi \rangle = \begin{pmatrix} v \\ v \\ v \end{pmatrix}, \quad \langle \phi_5 \rangle = \begin{pmatrix} 0 & v_5 & v_5 \\ v_5 & 0 & v_5 \\ v_5 & v_5 & 0 \end{pmatrix}, \quad \langle \phi' \rangle = \begin{pmatrix} v' \\ 0 \\ 0 \end{pmatrix}$$

$\Rightarrow$  **6 × 6 mass matrices**

# Obtaining $U_{\text{PMNS}}$

- $U_{\text{PMNS}}$  characterize charged-current interactions between light eigenstates.
- $6 \times 6$  mass matrices:
  - Block-diagonalise  $6 \times 6$  mass matrices.
  - Then diagonalise  $3 \times 3$  mass matrices.
- Neutrinos:

$$M_{\nu}^{3 \times 3} = \begin{pmatrix} -\frac{y_{\nu}^2 v_H^2}{M} & 0 & 0 \\ 0 & -\frac{y_{\nu}^2 M v_H^2 \Lambda^2}{M^2 \Lambda^2 - x_{\nu}^2 v'^2 v_T^2} & \frac{y_{\nu}^2 x_{\nu} v_H^2 v' v_T \Lambda}{M^2 \Lambda^2 - x_{\nu}^2 v'^2 v_T^2} \\ 0 & \frac{y_{\nu}^2 x_{\nu} v_H^2 v' v_T \Lambda}{M^2 \Lambda^2 - x_{\nu}^2 v'^2 v_T^2} & -\frac{y_{\nu}^2 M v_H^2 \Lambda^2}{M^2 \Lambda^2 - x_{\nu}^2 v'^2 v_T^2} \end{pmatrix}, \text{ of form } \begin{pmatrix} a & 0 & 0 \\ 0 & b & c \\ 0 & c & b \end{pmatrix}$$

# Obtaining $U_{\text{PMNS}}$

- Charged leptons:  $M_l^{6 \times 6} \equiv \begin{pmatrix} \frac{v_H}{\Lambda} A & \frac{v_H}{\Lambda} B \\ C & D \end{pmatrix}$

$$\omega = e^{i \frac{2\pi}{3}}$$

$$A = \begin{pmatrix} y_e v & [y_m v + y_m^5 v_5 (\omega^2 - \omega)] & [y_m v + y_m^5 v_5 (\omega - \omega^2)] \\ y_e v & [y_m v + y_m^5 v_5 (\omega^2 - \omega)] \omega & [y_m v + y_m^5 v_5 (\omega - \omega^2)] \omega^2 \\ y_e v & [y_m v + y_m^5 v_5 (\omega^2 - \omega)] \omega^2 & [y_m v + y_m^5 v_5 (\omega - \omega^2)] \omega \end{pmatrix}, \text{ of form } \begin{pmatrix} a & b & c \\ a & b\omega & c\omega^2 \\ a & b\omega^2 & c\omega \end{pmatrix}$$

$$B = \begin{pmatrix} y_m v + 2y_m^T v_T & y_m v + y_m^5 v_5 & -y_m^5 v_5 \\ y_m v & 2y_m^T v_T & y_m v \\ y_m v + y_m^5 v_5 + y_m^T v_T & y_m v - y_m^5 v_5 & y_m^T v_T \end{pmatrix}$$

$$C = \begin{pmatrix} y'_e v & [y'_m v + y_m^{5'} v_5 (\omega^2 - \omega)] & [y'_m v + y_m^{5'} v_5 (\omega - \omega^2)] \\ y'_e v & [y'_m v + y_m^{5'} v_5 (\omega^2 - \omega)] \omega & [y'_m v + y_m^{5'} v_5 (\omega - \omega^2)] \omega^2 \\ y'_e v & [y'_m v + y_m^{5'} v_5 (\omega^2 - \omega)] \omega^2 & [y'_m v + y_m^{5'} v_5 (\omega - \omega^2)] \omega \end{pmatrix}, \text{ of form } \begin{pmatrix} a & b & c \\ a & b\omega & c\omega^2 \\ a & b\omega^2 & c\omega \end{pmatrix}$$

$$D = \begin{pmatrix} y'_m v + 2y_m^{T'} v_{T'} & y'_m v + y_m^{5'} v_5 & -y_m^{5'} v_5 \\ y'_m v & 2y_m^{T'} v_{T'} & y'_m v \\ y'_m v + y_m^{5'} v_5 + y_m^{T'} v_{T'} & y'_m v - y_m^{5'} v_5 & y_m^{T'} v_{T'} \end{pmatrix}$$

# Nonzero $\theta_{13}$

- Block diagonalize  $M_l^{6 \times 6} (M_l^{6 \times 6})^\dagger$

$$M_l^{3 \times 3} (M_l^{3 \times 3})^\dagger = \frac{v_H^2}{\Lambda^2} \left[ AA^\dagger + BB^\dagger - (AC^\dagger + BD^\dagger)(CC^\dagger + DD^\dagger)^{-1}(CA^\dagger + DB^\dagger) \right]$$

- Useful to define  $E \equiv B - (y_m^T / y_m^{T'})D$
- Then  $A, C, E \sim O(v, v_5)$ ,  $B, D \sim O(v_T)$
- Expanding in small parameter  $\delta = O(v / v_T) \sim O(v' / v_T)$

– LO:  $M_l^{3 \times 3} = \frac{v_H}{\Lambda} \left( A - \frac{y_m^T}{y_m^{T'}} C \right)$

$\theta_{13}$  still 0!

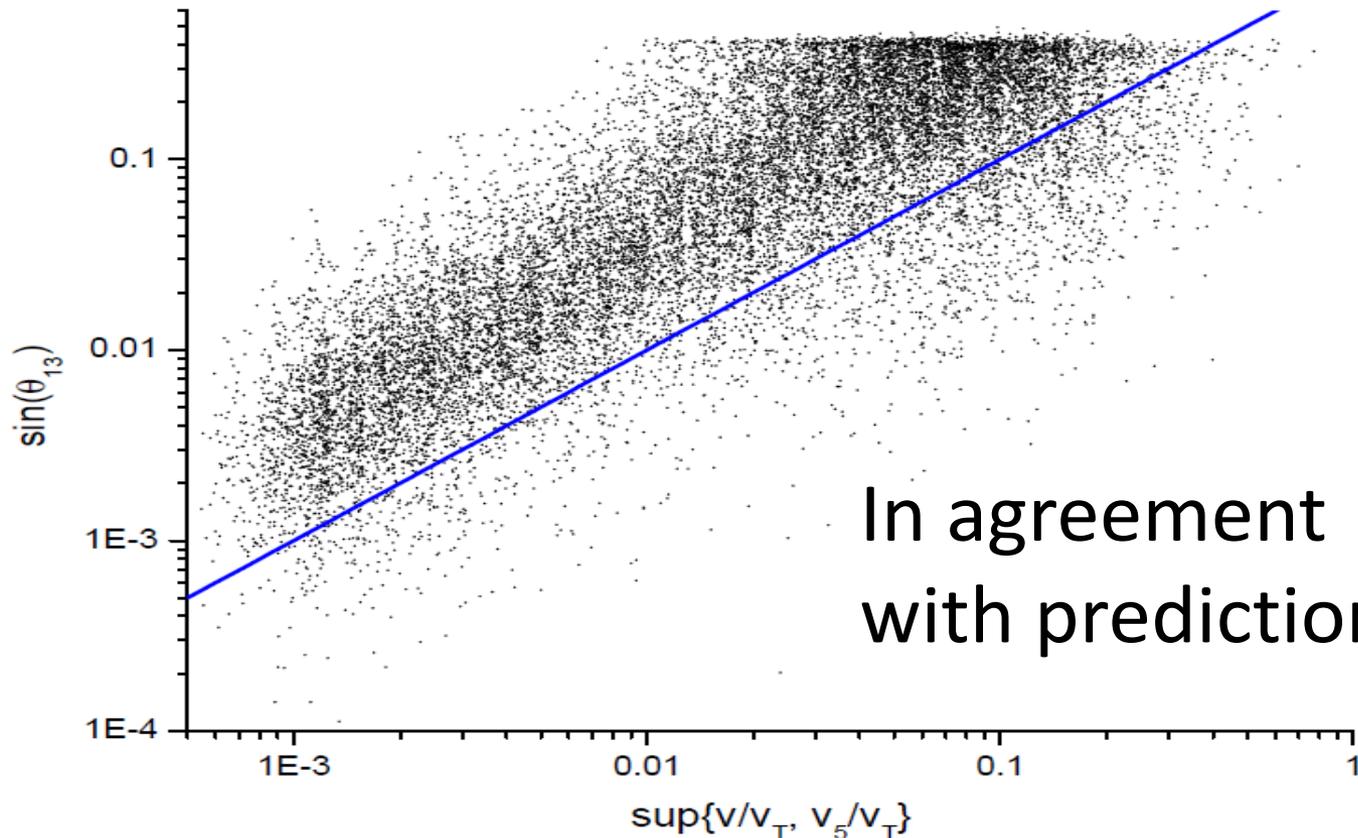
– NLO:  $M_l^{3 \times 3} = \frac{v_H}{\Lambda} \left( A - \frac{y_m^T}{y_m^{T'}} C - ED^{-1}C \right)$

Suggests  
 $\theta_{13} \sim \epsilon!$

~~A4~~ scale  
~~SO(3)~~ scale

# Nonzero $\theta_{13}$

- Actual contribution to  $\theta_{13}$  from deviation turns out to be  $\sim (m_\tau/m_\mu) \varepsilon \sim 10\varepsilon$
- Random simulation results



# Summary and further work

- TBM ruled out, but many  $A_4$  models allow for this scenario
- In particular,  $SO(3) \rightarrow A_4$  model allow nonzero tree-level  $\theta_{13}$  of size  $\sim \frac{m_\tau}{m_\mu} \frac{A_4 \text{ scale}}{SO(3) \text{ scale}}$
- % level separation of  $A_4$  and  $SO(3)$  breaking scales

**Thank you!**