

New W -prime Signals at the LHC

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Based on:

ArXiv: 1404.5558

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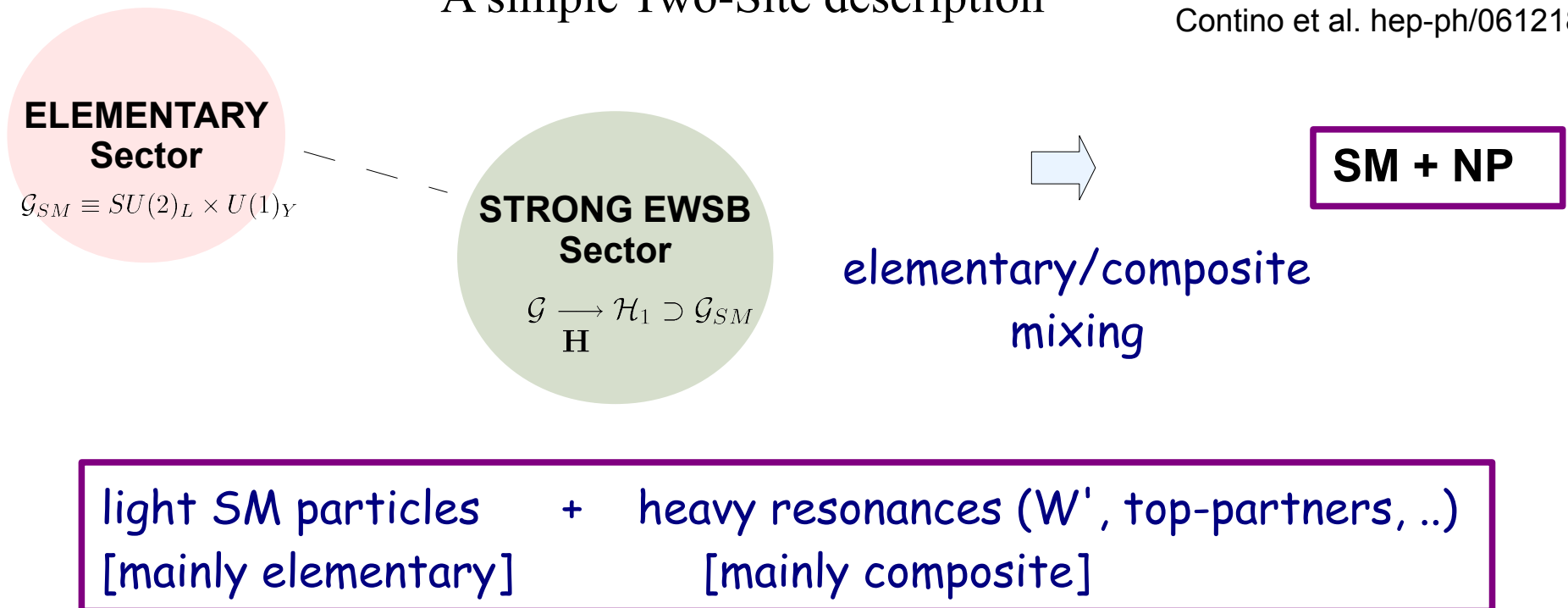
W-prime in Composite Higgs / RS Models

- EWSB triggered by a new Strong Dynamics, composite at the TeV scale
- Higgs: composite + pGB of global invariance of the strong sector

Georgi, Kaplan, 1984

A simple Two-Site description

Contino et al. hep-ph/0612180



W-prime in Composite Higgs / RS Models

Custodial symmetry

$$SU(2)_L \times U(1)_Y$$

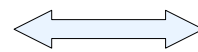
$$g_2^{el} \simeq g_2$$

$$SO(5) \rightarrow SO(4) \sim SU(2)_L \times SU(2)_R$$

$$g_2^* > 1$$

~ rho-photon
mixing

$$W_{el}^a \in (\mathbf{3}, \mathbf{1})$$



$$W_{com}^{\prime a} \in (\mathbf{3}, \mathbf{1})$$

$$\cot \theta_2 = \frac{g_2^*}{g_2^{el}}$$

$$g_2 = g_2^{el} \cos \theta_2 = g_2^* \sin \theta_2$$

$$W' = \cos \theta_2 W_{com}' + \sin \theta_2 W_{el}$$

Controls the strength of W' interactions

More strongly-coupled EW sectors correspond to larger $\cot \theta_2$

top-partners in Composite Higgs / RS Models

Linear mixing terms

$$\begin{array}{l}
 t_R^{el} \quad \longleftrightarrow \quad \tilde{T} = (\mathbf{1}, \mathbf{1})_{2/3} \\
 \left(\begin{array}{c} t_L^{el} \\ b_L^{el} \end{array} \right) \quad \longleftrightarrow \quad \left[\begin{array}{c} T \\ B \end{array} \right] \begin{array}{c} T_{5/3} \\ T_{2/3} \end{array} \quad = \quad (\mathbf{2}, \mathbf{2})_{2/3}
 \end{array}
 \quad \left. \vphantom{\begin{array}{l} t_R^{el} \\ \left(\begin{array}{c} t_L^{el} \\ b_L^{el} \end{array} \right) \end{array}} \right\} \mathbf{5}_{2/3}$$

Rotation
(Mixing angles: s_L, s_R)

$$\left\{ \begin{array}{l} t_L = c_L t_L^{el} + s_L T_L \\ t_R = c_R t_R^{el} + s_R \tilde{T}_R \end{array} \right. \quad s_L / s_R \text{ are the } t_L / t_R \text{ degree of compositeness}$$

$$m_t = Y_* s_L s_R v$$

Heavier particles have larger degrees of compositeness

top-partners in Composite Higgs / RS Models

Linear mixing terms

$$t_R^{el} \iff \tilde{T} = (\mathbf{1}, \mathbf{1})_{2/3}$$

$$\begin{pmatrix} t_L^{el} \\ b_L^{el} \end{pmatrix} \iff \begin{bmatrix} T & T_{5/3} \\ B & T_{2/3} \end{bmatrix} = (\mathbf{2}, \mathbf{2})_{2/3}$$

$\mathbf{5}_{2/3}$

Custodians

$$m_C = c_L m_T$$

Rotation
(Mixing angles: s_L, s_R)

$$\begin{cases} t_L = c_L t_L^{el} + s_L T_L \\ t_R = c_R t_R^{el} + s_R \tilde{T}_R \end{cases}$$

s_L / s_R are the t_L / t_R degree of compositeness

$$m_t = Y_* s_L s_R v$$

Heavier particles have larger degrees of compositeness

W-primes and top-partners in MCHMs

naturalness

Top-partners cut-off Higgs mass divergencies

A ~ 125 GeV Composite Higgs implies Top-partners below 1 TeV

G. Panico, M. Redi, A. Tesi and A. Wulzer, [arXiv:1210.7114 [hep-ph]]

.....

EWPT

Constraints on S parameter imply W-prime above ~ 2 TeV

M. Ciuchini, E. Franco, S. Mishima and L. Silvestrini, [arXiv:1306.4644 [hep-ph]].

C. Grojean, O. Matsedonskyi and G. Panico, [arXiv:1306.4655 [hep-ph]].

We will consider $m_c = 0.9$ TeV { CMS bound [arXiv:1312.2391]: $m_c > 0.8$ TeV }

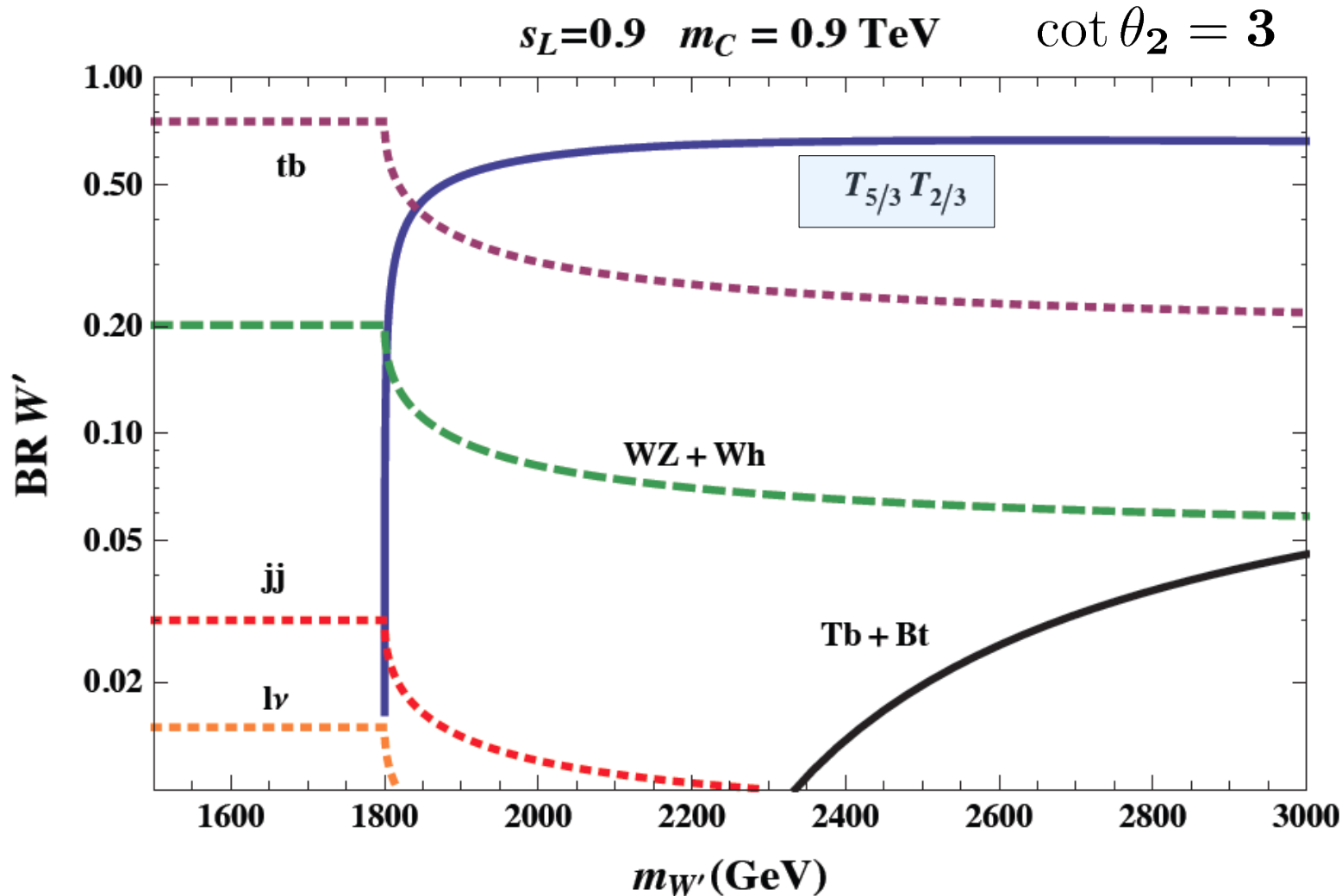
And we will focus our analysis on the W' decays to vector-like top partners

W' couplings to elementary modes
(lv, jj)

$$\propto g_2 \tan \theta_2$$

W' couplings to composite modes
($W_L h, W_L Z, \text{top-partners}$)

$$\propto g_2 \cot \theta_2$$

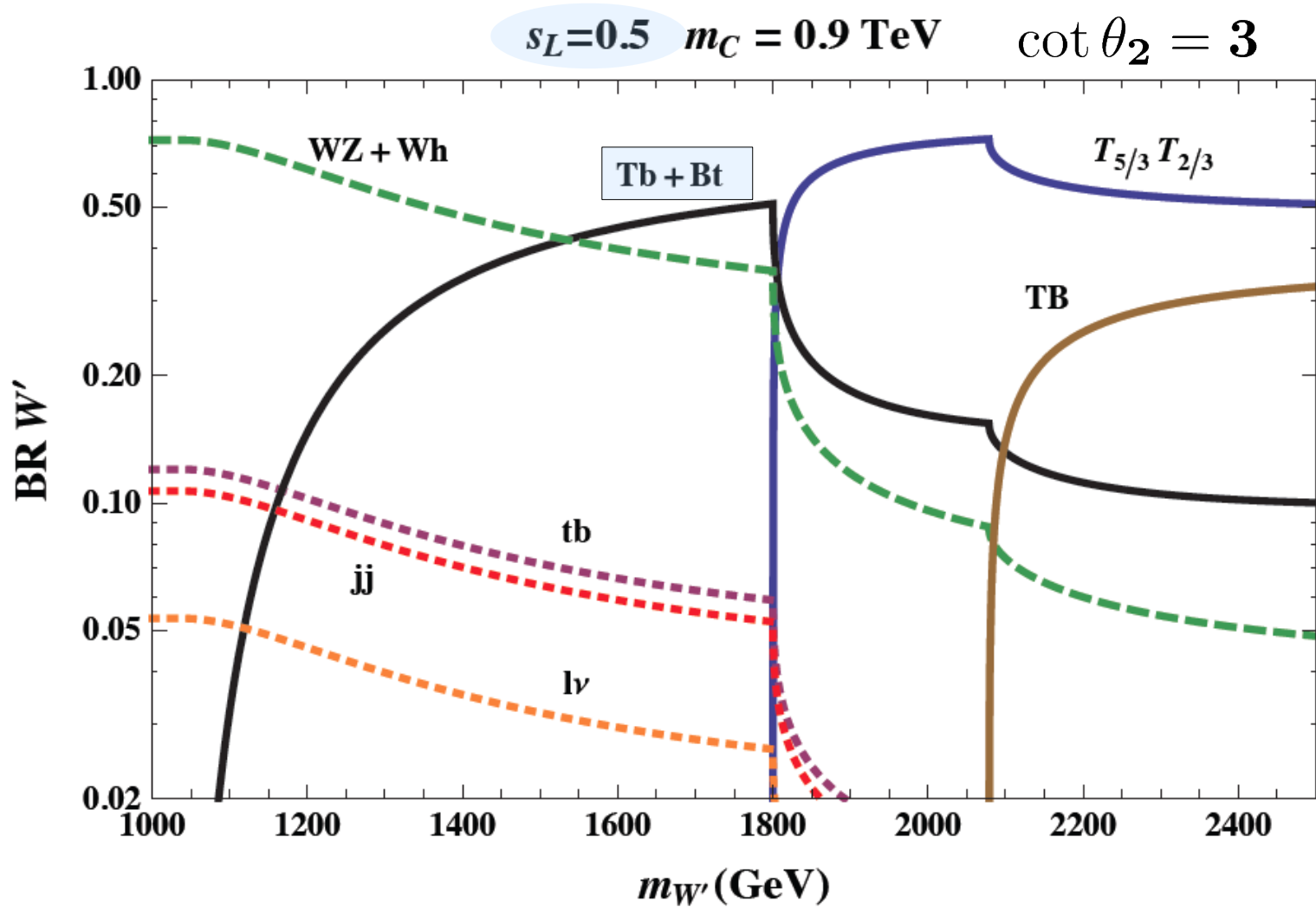


W' couplings to elementary modes
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$$\propto g_2 \tan \theta_2$$

W' couplings to composite modes
(W_Lh, W_LZ, top-partners)

$$\propto g_2 \cot \theta_2$$



LHC-8 LIMITS

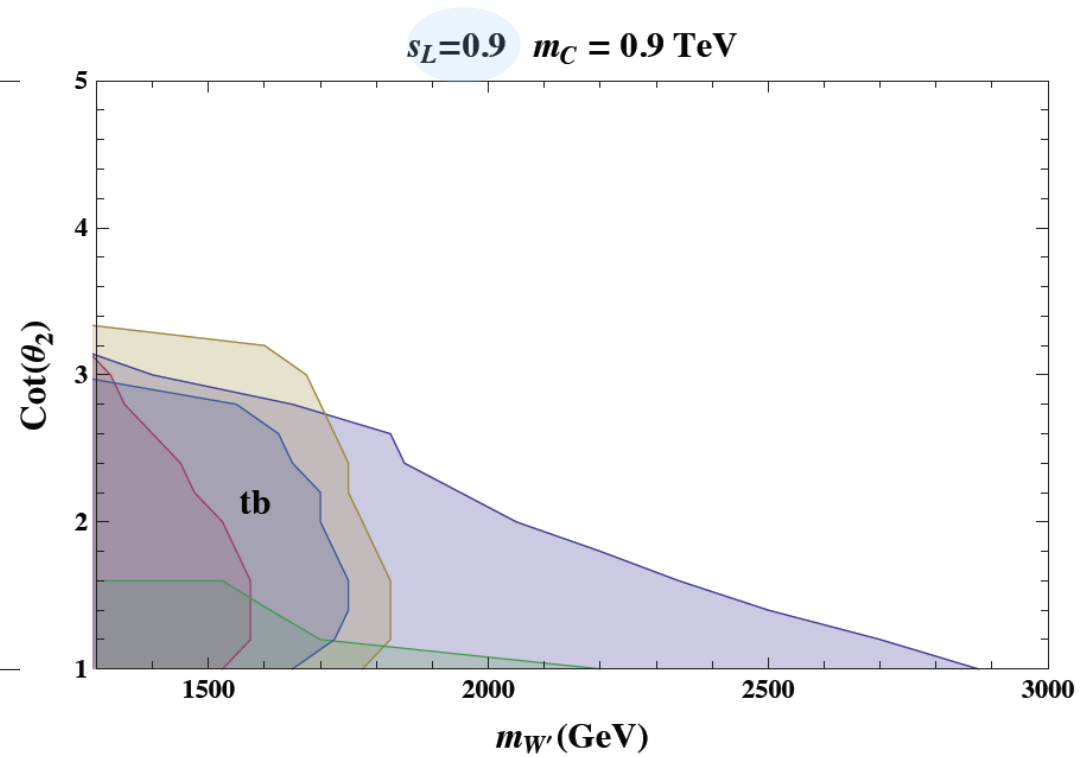
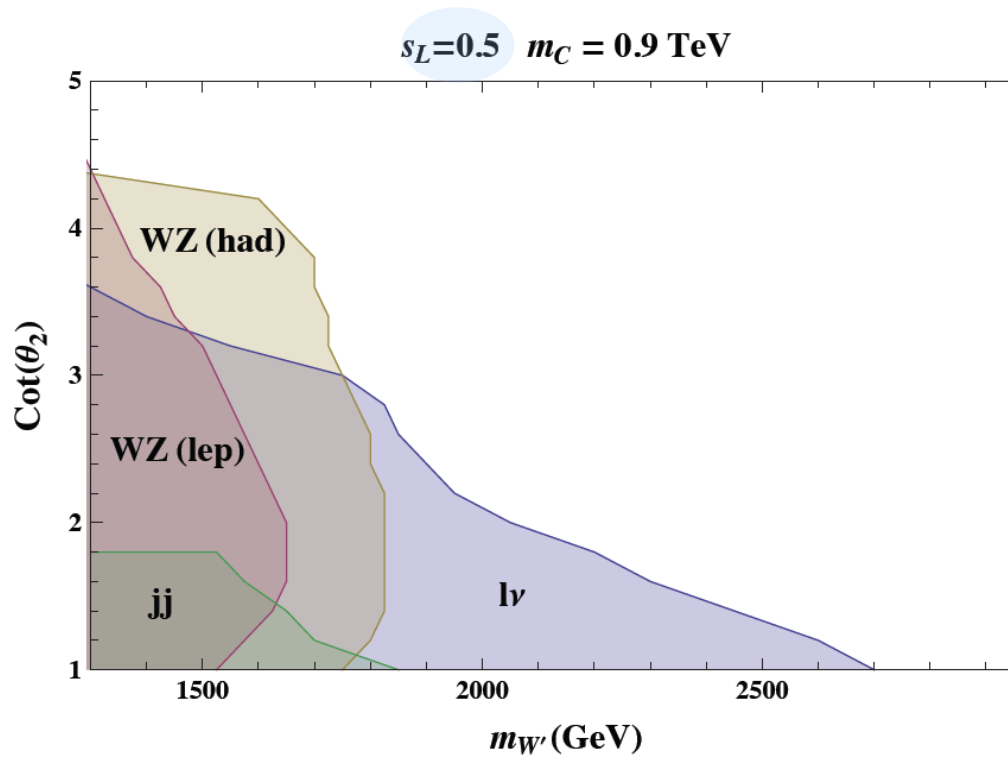
lv CMS-PAS-EXO-12-060; ATLAS-CONF-2014-017

jj CMS-PAS-EXO-12-059; ATLAS-CONF-2012-148

WZ (had) CMS-PAS-EXO-12-024

tb CMS-PAS-B2G-12-010; ATLAS-CONF-2013-050

WZ(lep) CMS-PAS-EXO-12-025; ATLAS-CONF-2014-015

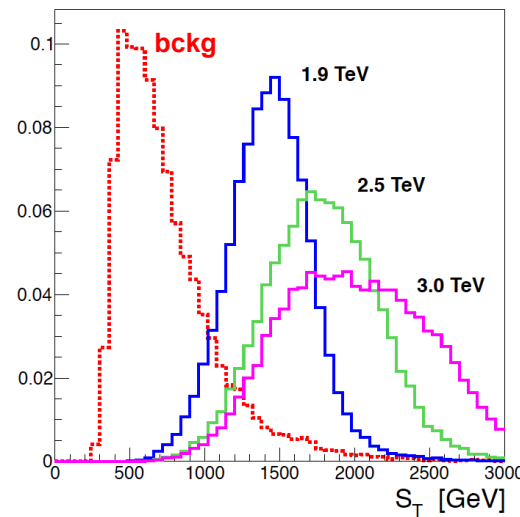
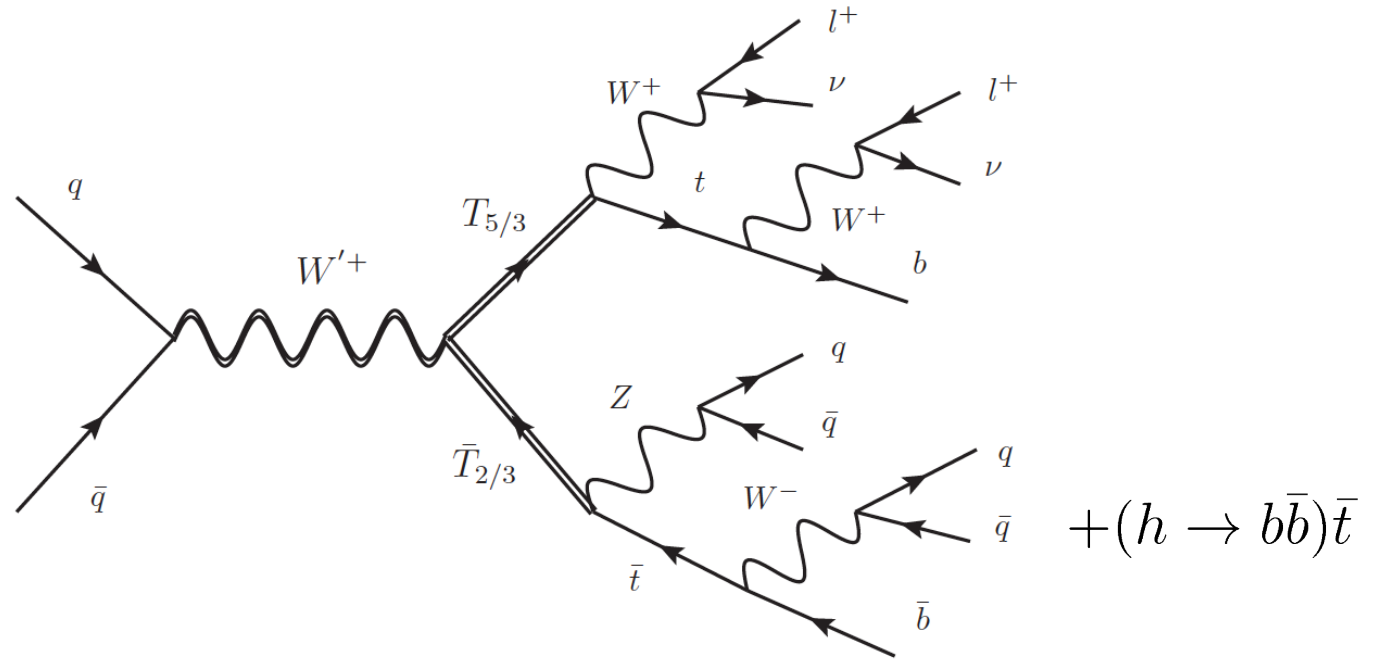


The custodian channel in the same-sign dilepton final state

LHC-14

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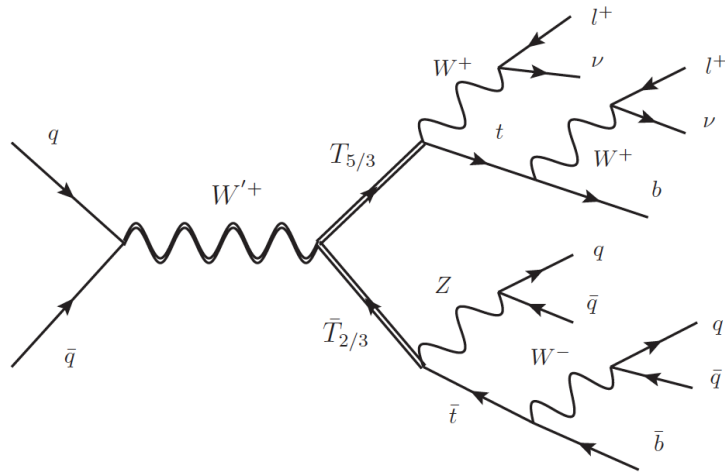
Same-Sign Dilep	acceptance
$m_{W'} = 1.9 \text{ TeV}$	0.82
$m_{W'} = 2.2 \text{ TeV}$	0.52
$m_{W'} = 2.5 \text{ TeV}$	0.29
$m_{W'} = 3.0 \text{ TeV}$	0.11
$m_{W'} = 3.5 \text{ TeV}$	0.041
$W^+ t \bar{t}$	4.1
$W^+ W^+$	1.5
$W^+ W^+ W^-$	0.6
Total background	6.2



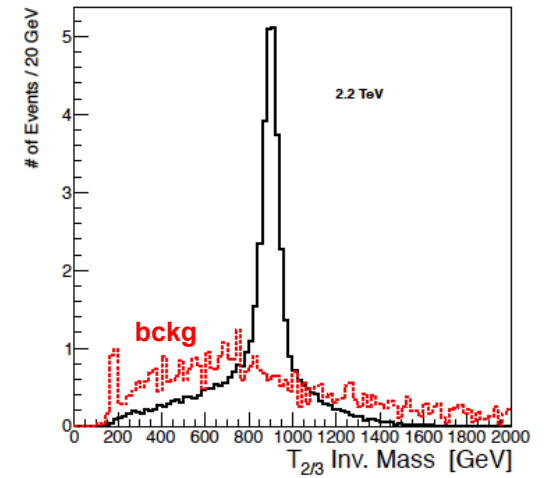
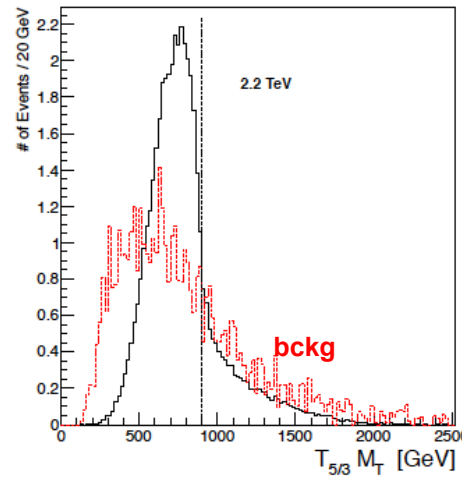
$$S_T \equiv H_T + \cancel{p}_T$$

	CUT-1	CUT-2
$p_T l(1)$	90	
$p_T l(2)$	30	
$p_T j(1)$	160	
$p_T j(2)$	100	
H_T	550	700
S_T	1100	1400

The custodian channel in the same-sign dilepton final state

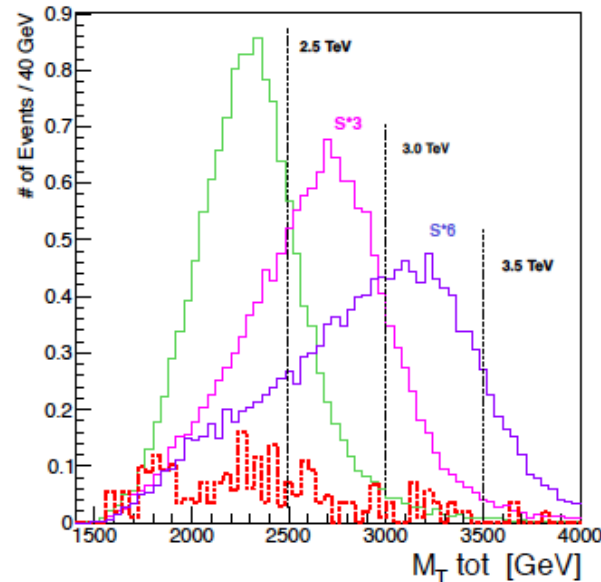
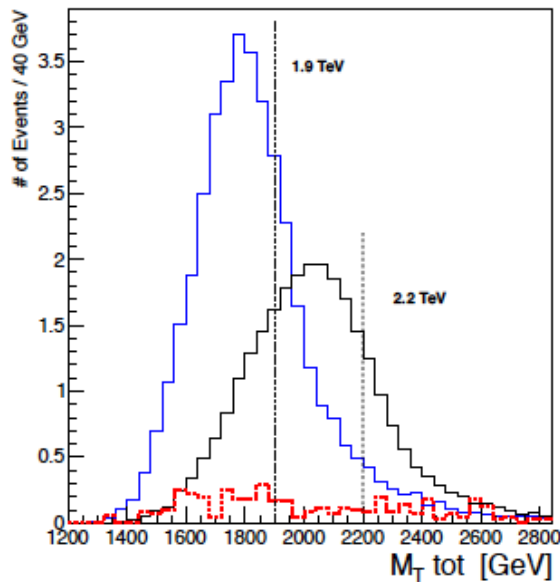


Custodians



Total transverse mass

(smearing in jet momentum included)

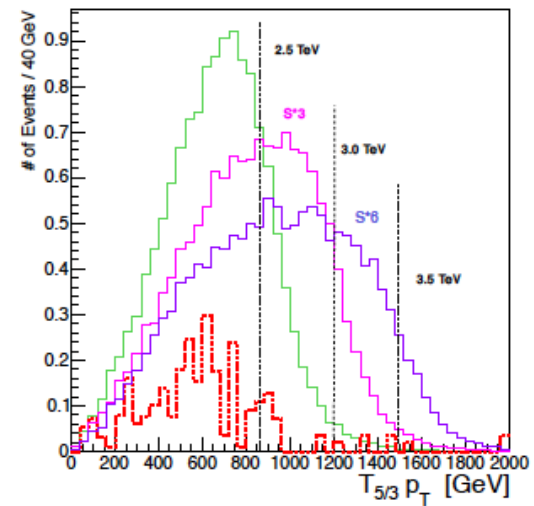
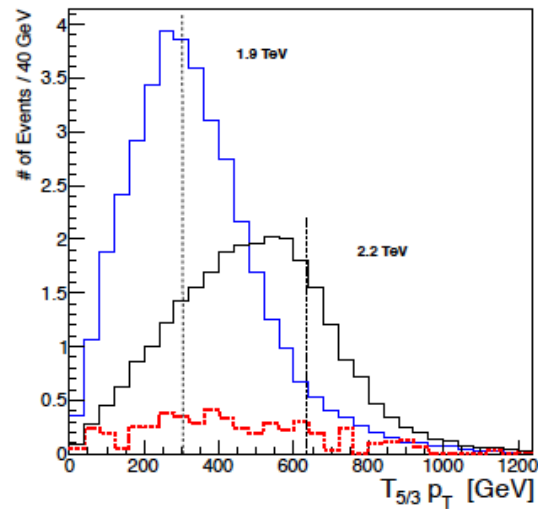
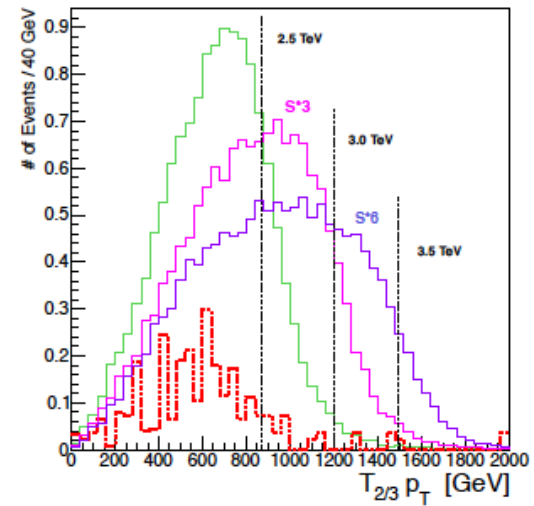
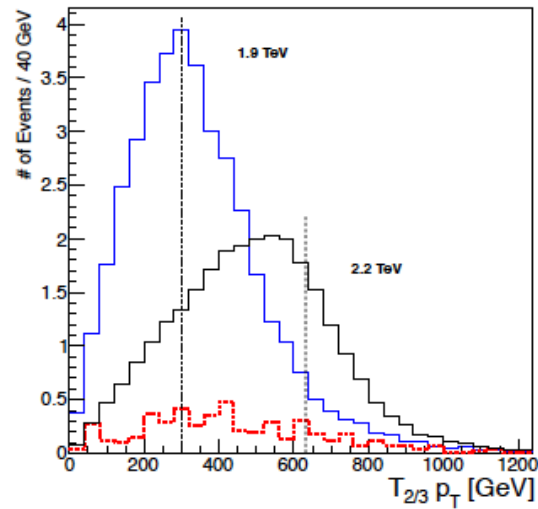


$$M_T(T_{5/3}) > 400 \text{ GeV}$$

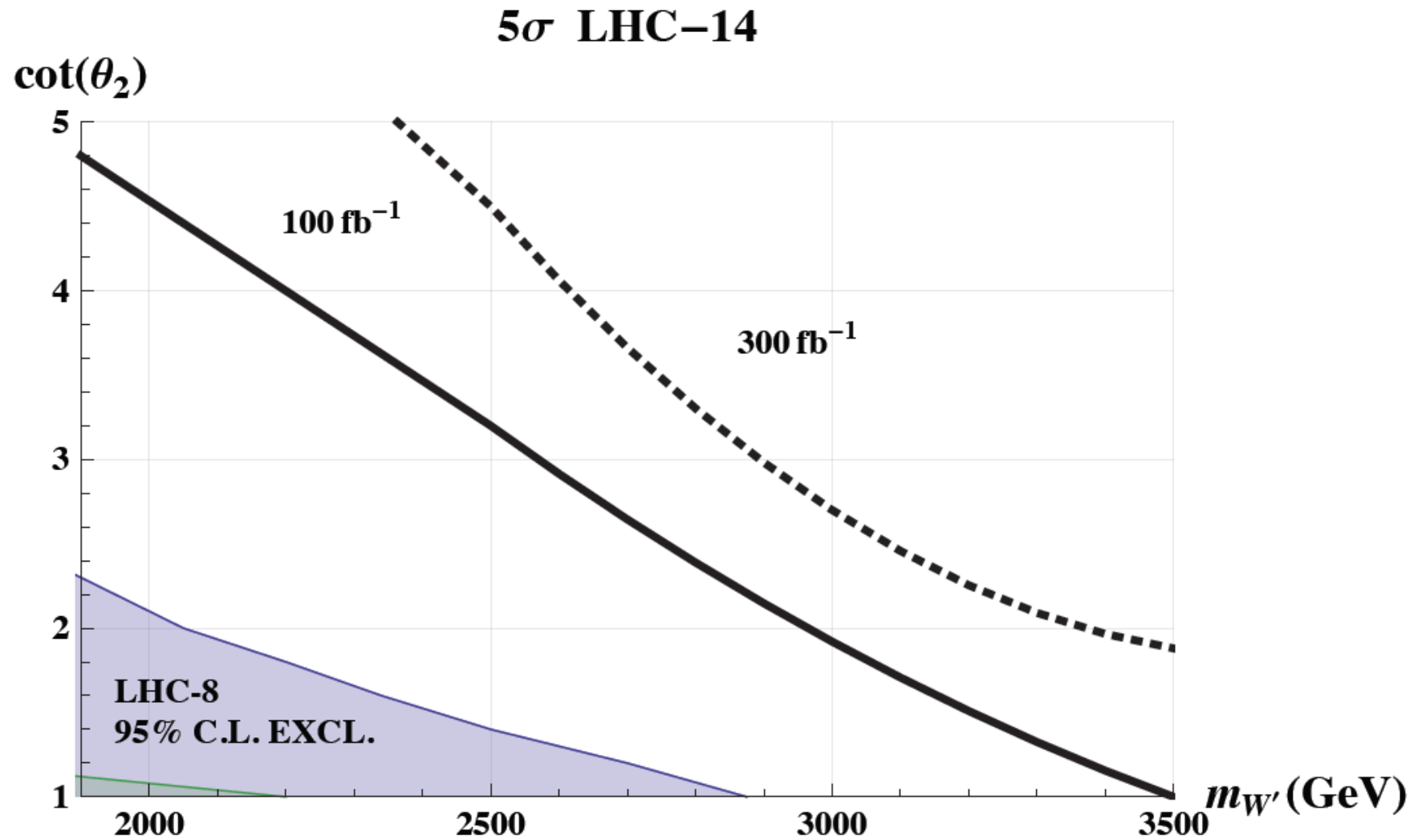
$$M_{T_{2/3}} \in [0.8, 1.0] \text{ TeV}$$

The custodian channel in the same-sign dilepton final state

Custodian
 p_T distributions



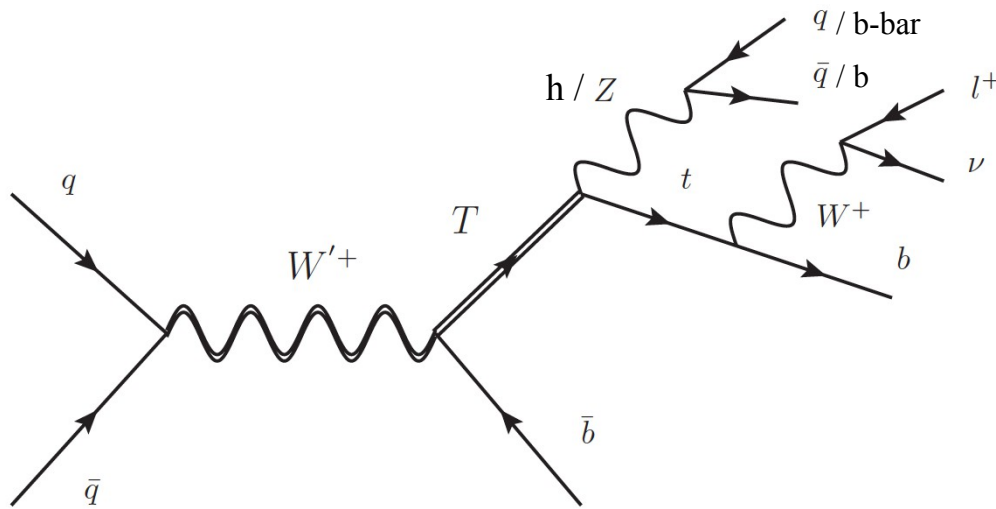
The custodian channel in the same-sign dilepton final state



The heavy-light Tb channel, for $s_L=0.5$

$$pp \rightarrow l^+ + n \text{ jets} + \cancel{E}_T$$

$$n \geq 3, \quad \text{at least 2 } b\text{-tag}$$

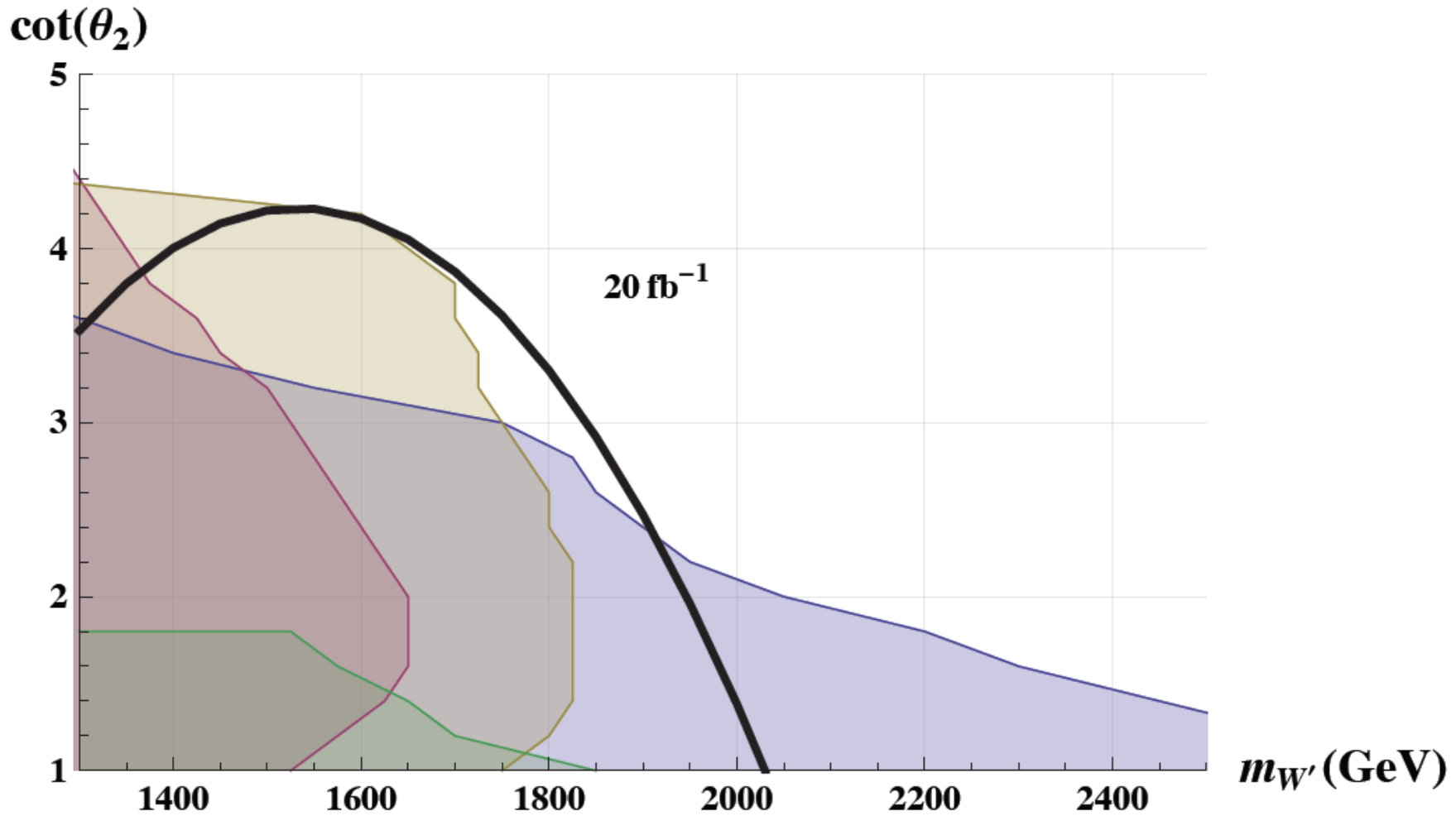


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LHC-14	acceptance
$m_{W'} = 1.3 \text{ TeV}$	9.1
$m_{W'} = 1.5 \text{ TeV}$	7.5
$m_{W'} = 1.7 \text{ TeV}$	5.0
$m_{W'} = 2.0 \text{ TeV}$	0.70
$m_{W'} = 2.5 \text{ TeV}$	0.11
$WWbb$	19000
$Wbb + jets$	1600
$W + jets$	560
Total background	21000

- Neutrino, top and T reconstruction
- $S_T > 1.1 \text{ TeV}$
- $(t, b, Z/h, T) p_T > 150 \text{ GeV}$
- T inv mass cut

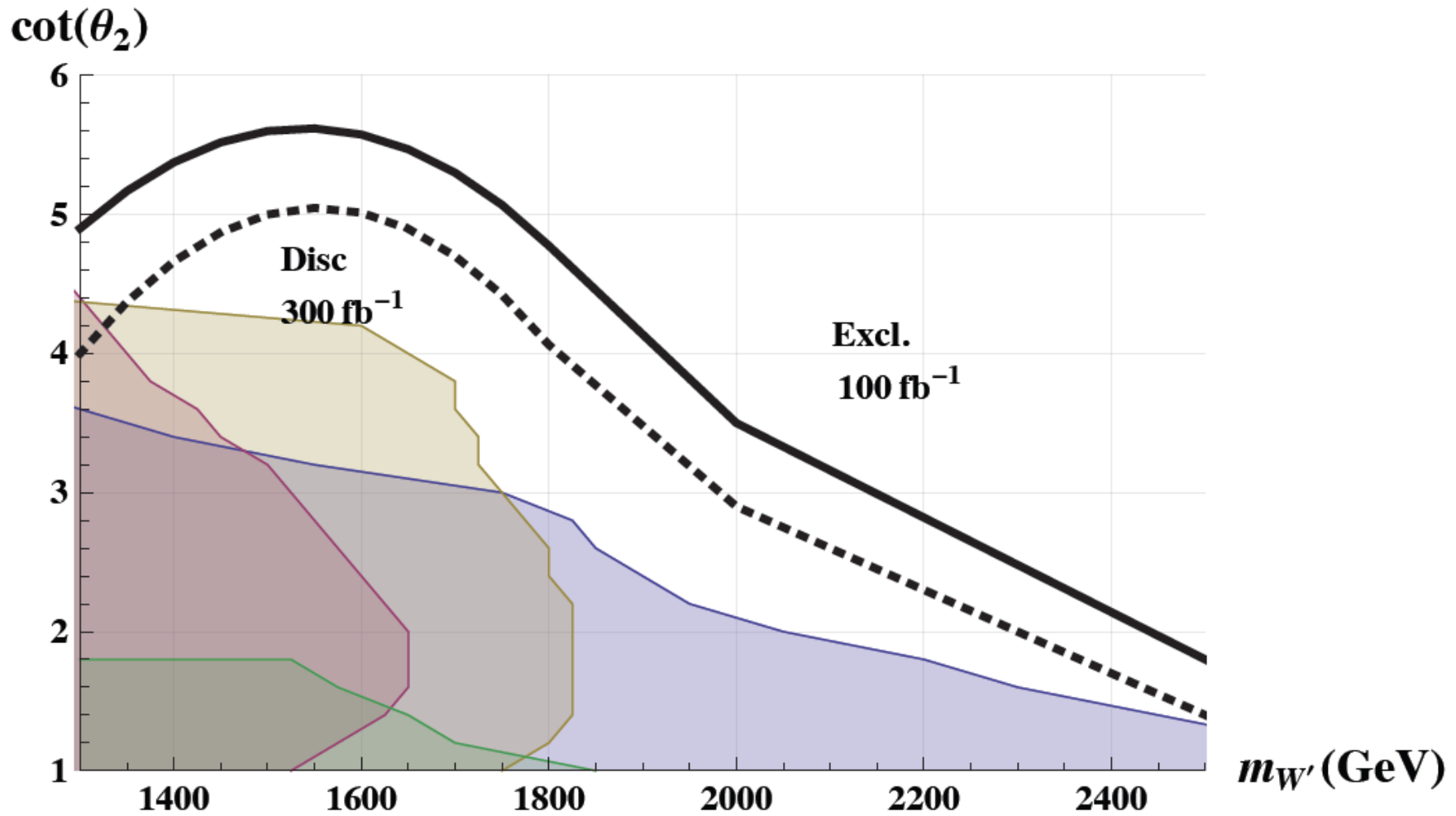
The heavy-light Tb channel, for $s_L=0.5$

95% C.L. EXCL LHC-8



The heavy-light Tb channel, for $s_L=0.5$

LHC-14



Conclusions

Naturalness+EWPT favor the scenario where the W' is above the threshold for decays into vector-like top-partners

Current LHC-8 analyses in the standard channels exclude a consistent portion of the W' mass-coupling parameter space at lower W' masses but have a low sensitivity to the higher mass region $m_{W'} > 2 \text{ TeV}$

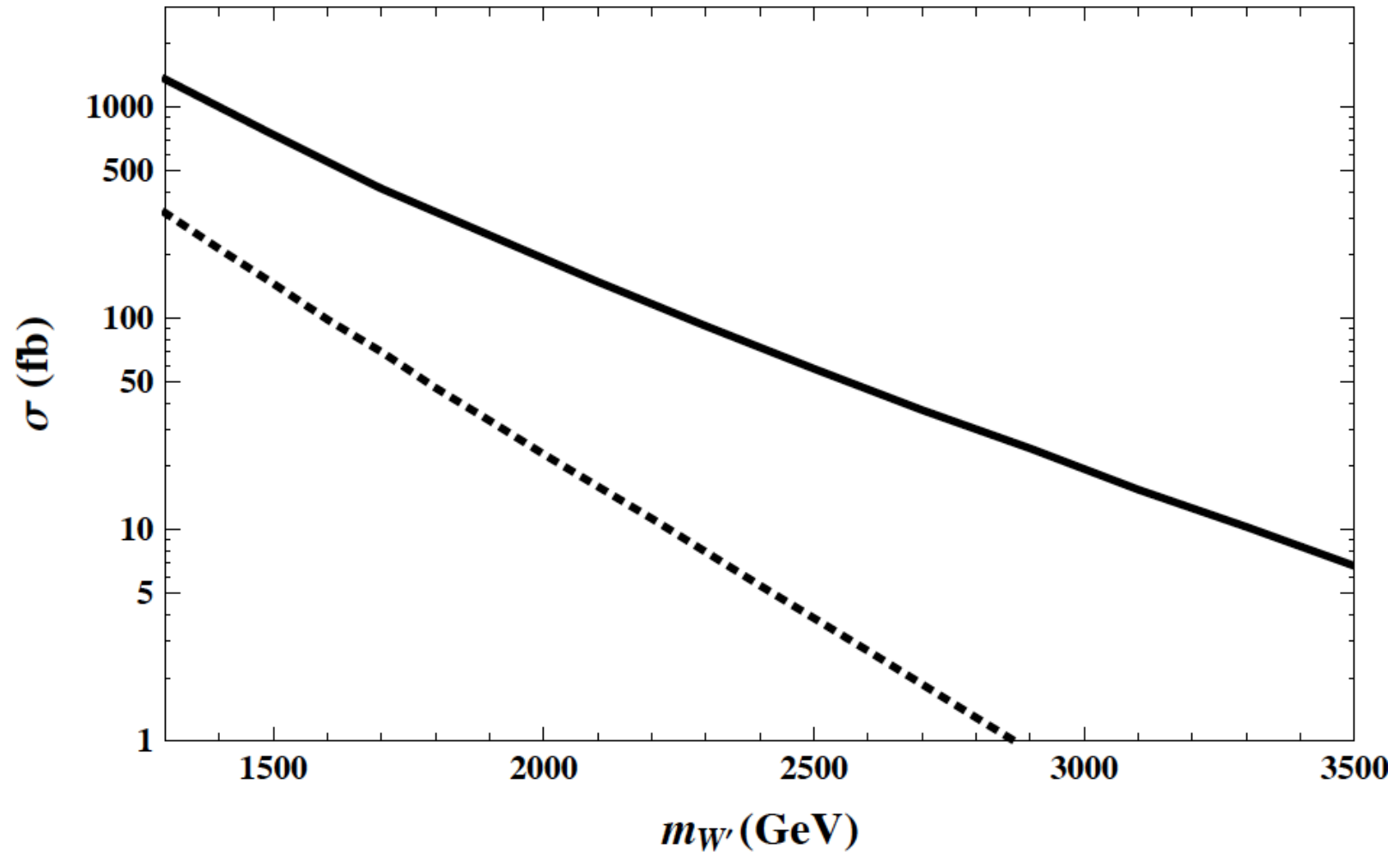
The custodian channel $W' \rightarrow T_{5/3} T_{2/3}$ can extensively probe this interesting region at the 14 TeV LHC

The heavy-light channel $W' \rightarrow T b$ is also a promising channel to test the intermediate W' mass region. It could extend the bounds from LHC-8 data.

Extra Slides

W' Drell-Yan production

$$\text{Cot}(\theta_2) = 3$$



W' decay width

