

# Review of SM Higgs properties measured by ATLAS and CMS: couplings, spin, mass

Pavel Jež

Centre for Cosmology, Particle Physics and Phenomenology - CP3 Université catholique de Louvain



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# Higgs boson in the SM

#### cross-section and BR



- ٠ Result of spontaneous symmetry breaking
- Mass to gauge bosons + unitarity at high energy
- ٠ Mass to fermions through Yukawa



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September 16, 2014 2 / 33

# The collider and detectors





- proton-proton collisions
   @ 7 TeV in 2010 and 2011
- pp collisions @ 8 TeV in 2012
- Up to 15/30 interactions per beam collision in 2011/2012
- 5.08 +21.3 fb<sup>-1</sup> (ATLAS)
- 5.55 + 21.79 fb<sup>-1</sup> (CMS)
- O(10<sup>5</sup>) decays of H(125) $ightarrow b\bar{b}$
- O(10<sup>3</sup>) decays of H(125) $\rightarrow \gamma\gamma$







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Higgs properties (ATLAS+CMS)

September 16, 2014

3 / 33

#### Mass measurement

#### Motivation

- Mass of the boson is the last free parameter in the Standard Model
- Fundamental position in the SM:
  - Calculation of the H production and decay rates
  - Precise knowledge necessary to test the coupling structure







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#### Overview

- Exploit  ${\sf H} 
  ightarrow {\sf ZZ} 
  ightarrow {\sf 4I}$  and  ${\sf H} 
  ightarrow \gamma\gamma$  channels
- full mass reconstruction, clean signal, excellent resolution
- Model independent measurement
  - Fit of a peak over smooth background
  - No assumption on production or decay yields
- Improved analyses of Run I data  $\Rightarrow$  twice better precision
- [Phys.Rev.D.90,052004(2014)](ATLAS), [CMS-PAS-HIG-14-009](CMS)



# Extracting mass for individual channels Phys.Rev.D.90,052004(2014), CMS-PAS-HIG-14-009 $H \rightarrow \gamma\gamma$

- split into several (10 20) categories
- fit  $m_{\gamma\gamma}$  distributions simultaneously
- signal shape model determined from MC
- background modeled by smooth function
- bias from choice of background model studied with MC







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#### $H \rightarrow ZZ \rightarrow 4I$

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- using matrix element approach to build discriminants against ZZ background
- mass is obtained via parameter estimation with multi-dimensional unbinned likelihoods

Higgs properties (ATLAS+CMS)







5 / 33

# Results of 2D fits in ( $\mu \times m_H$ ) plane Phys.Rev.D.90,052004(2014), CMS-PAS-HIG-14-009



- fix the relative signal yields between  $\gamma\gamma$  and  $Z\!Z$
- let the overall signal strength ( $\mu = \sigma/\sigma_{
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- Fixed  $\sigma_{SM}$  at ATLAS, mass dependent in CMS





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6 / 33

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#### Mass combination

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- Assume single state with mass m<sub>H</sub>
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Compatibility of measurements in 2 channels Phys.Rev.D.90,052004(2014), CMS-PAS-HIG-14-009

- Using test statistics  $q(m_H^{\gamma\gamma}-m_H^{4\prime})$
- 2 degrees of freedom  $(m_{H}^{\gamma\gamma}$  and  $\Delta m), m_{H}^{\gamma\gamma}$  is profiled



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# Spin and CP of the H boson

- SM: H boson is scalar  $(J^P = 0^+)$
- testing compatibility of data with  $J^P = 0^+, 0^-, 1^{+/-}, 2^{+/-}$
- various models for non-SM spin and parity
- exploiting kinematical observables in  $\gamma\gamma$ ,  $ZZ \rightarrow 4I$  and  $WW \rightarrow I\nu I\nu$ channels
- first measurements of anomalous couplings



• Test statistics:  $q = -2 \ln \frac{\mathcal{L}(J_{alt}^{P})}{\mathcal{L}(J^{P}=0^{+})}$ • exclusion level  $1 - \alpha$ :  $CL_s = \frac{P(q > q_{obs}|J_{alt}^P + bkg)}{P(q > q_{obs}|0^+ + bkg)} < \alpha$ CMS 9 / 33

$$\begin{aligned} \mathcal{A}(X_{J=0} \to V_1 V_2) &\sim \quad v^{-1} \left( \left[ a_1 - e^{i\phi\Lambda_1} \frac{q_{Z_1}^2 + q_{Z_2}^2}{(\Lambda_1)^2} \right] m_Z^2 \epsilon_{Z_1}^* \epsilon_{Z_2}^* \right. \\ &+ \quad \sum_{V_1 V_2} \left( a_2^{V_1 V_2} f_{\mu\nu}^{*(V_1)} f^{*(V_2),\mu\nu} + a_3^{V_1 V_2} f_{\mu\nu}^{*(V_1)} \tilde{f}^{*(V_2),\mu\nu} \right) \right) \end{aligned}$$





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•  $v$ : vev,  $f^{(i)\mu\nu}$ ,  $(\tilde{f}^{(i)\mu,\nu}) \sim$  (dual) field strength tensor





CMS

10 / 33

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CMS



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 $a_i$ 's can in general be complex (loop contributions from light particles)



# Spin 0 hypotheses testing

- Using  $ZZ \rightarrow 4I$  channel
- [Phys.Lett.B726(2013),120-144](ATLAS), [Phys.Rev.D89(2014)092007](CMS)







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CMS

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CMS also excluded scalar decoupled from EWSB( $0_{h}^{+}$ ) with CL<sub>s</sub>=95.5%

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11 33



### Spin 0 measurements in $ZZ \rightarrow 4I$

- Constraints on mixture of scalar and other spin 0 states CMS-PAS-HIG-14-014
- Defining effective fractions:  $f_{a_i} = \frac{|a_i|^2 \sigma_i}{|a_1|^2 \sigma_1 + |a_2|^2 \sigma_2 + |a_3|^2 \sigma_3 + \tilde{\sigma}_{\Lambda_1}/(\Lambda_1)^4}$
- Best fit assumes real phase  $\phi_{a_i}=0$  or  $\phi_{a_i}=\pi$
- Similar results with different methods



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- Landau-Yang theorem forbids spin 1 if  $H \rightarrow \gamma \gamma$  decay exists BUT:
  - Different bosons for different final states
  - Multiple narrow states with different  $J^P$

#### Parametrization:

 $A(X_{J=1} \rightarrow V_1 V_2) \sim \frac{b_1}{\left[(\epsilon_{V_1}^* q)(\epsilon_{V_2}^* \epsilon_X) + (\epsilon_{V_2}^* q)(\epsilon_{V_1}^* \epsilon_X)\right]} + \frac{b_2}{b_2} \epsilon_{\alpha \mu \nu \beta} \epsilon_X^{\alpha} \epsilon_{V_1}^{*\mu} \epsilon_{V_2}^{*\nu} \tilde{q}^{\beta}$ 





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$$f_{b_2} = \frac{|b_2|^2 \sigma_2}{|b_1|^2 \sigma_1 + |b_2|^2 \sigma_2}$$
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#### Spin 1 results



- Using  $ZZ \rightarrow 4I$  and  $WW \rightarrow I\nu I\nu$  channels
- [Phys.Lett.B726(2013),120-144](ATLAS), [CMS-PAS-HIG-14-014](CMS)



#### Non-interfering spin 1 states

- Composite particles can have multiple narrow states with different J<sup>P</sup> and nearly degenerate masses
  - ortho/para-positronium,  $\chi_b, \chi_c$
- $\bullet$  CMS analyzed for presence of second resonance with non-SM  $J^P$  close to dominant  $0^+_{\rm m}$ 
  - $\Gamma_{J^P}$  and  $\Gamma_{0^+_{\mathrm{m}}} << |m_{J^P} m_{0^+_{\mathrm{m}}}| << 1 \text{ GeV}$
  - fractional cross-section  $f(J^P) = \frac{\sigma_{J^P}}{\sigma_{a^+} + \sigma_{J^P}}$



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#### Interaction of general spin-2 resonance with a ZZ or WW pair:

$$\begin{split} \mathcal{A}(X_{J=2} \to V_{1}V_{2}) &\sim \Lambda^{-1} \left[ 2\mathbf{c_{1}}t_{\mu\nu}f^{*1,\mu\alpha}f^{*2,\nu\alpha} + 2\mathbf{c_{2}}t_{\mu\nu}\frac{q_{\alpha}q_{\beta}}{\Lambda^{2}}f^{*1,\mu\alpha}f^{*2,\nu\beta} \right. \\ &+ \mathbf{c_{3}}\frac{\tilde{q}^{\beta}\tilde{q}^{\alpha}}{\Lambda^{2}}t_{\beta\nu}\left(f^{*1,\mu\nu}f^{*2}_{\mu\alpha} + f^{*2,\mu\nu}f^{*1}_{\mu\alpha}\right) + \mathbf{c_{4}}\frac{\tilde{q}^{\nu}\tilde{q}^{\mu}}{\Lambda^{2}}t_{\mu\nu}f^{*1,\alpha\beta}f^{*(2)}_{\alpha\beta} \right. \\ &+ m_{V}^{2}\left(2\mathbf{c_{5}}t_{\mu\nu}\epsilon^{*\mu}_{V_{1}}\epsilon^{*\nu}_{V_{2}} + 2\mathbf{c_{6}}\frac{\tilde{q}^{\mu}q_{\alpha}}{\Lambda^{2}}t_{\mu\nu}(\epsilon^{*\nu}_{V_{1}}\epsilon^{*\alpha}_{V_{2}} - \epsilon^{*\alpha}_{V_{1}}\epsilon^{*\nu}_{V_{2}}) + \mathbf{c_{7}}\frac{\tilde{q}^{\mu}\tilde{q}^{\nu}}{\Lambda^{2}}t_{\mu\nu}\epsilon^{*}_{V_{1}}\epsilon^{*}_{V_{2}}\right) \\ &+ \mathbf{c_{8}}\frac{\tilde{q}^{\mu}\tilde{q}^{\nu}}{\Lambda^{2}}t_{\mu\nu}f^{*1,\alpha\beta}\tilde{f}^{*(2)}_{\alpha\beta} + \mathbf{c_{9}}t^{\mu\alpha}\tilde{q}_{\alpha}\epsilon_{\mu\nu\rho\sigma}\epsilon^{*\nu}_{V_{1}}\epsilon^{*\rho}_{V_{2}}q^{\sigma} \\ &+ \frac{\mathbf{c_{10}}t^{\mu\alpha}\tilde{q}_{\alpha}}{\Lambda^{2}}\epsilon_{\mu\nu\rho\sigma}q^{\rho}\tilde{q}^{\sigma}\left(\epsilon^{*\nu}_{V_{1}}(q\epsilon^{*}_{V_{2}}) + \epsilon^{*\nu}_{V_{2}}(q\epsilon^{*}_{V_{1}})\right) \right] \end{split}$$



September 16, 2014
# Spin 2 hypotheses



Interaction of general spin-2 resonance with a ZZ or WW pair:

$$\begin{split} A(X_{J=2} \to V_{1}V_{2}) &\sim \Lambda^{-1} \left[ 2c_{1}t_{\mu\nu}f^{*1,\mu\alpha}f^{*2,\nu\alpha} + 2c_{2}t_{\mu\nu}\frac{q_{\alpha}q_{\beta}}{\Lambda^{2}}f^{*1,\mu\alpha}f^{*2,\nu\beta} \right. \\ &+ c_{3}\frac{\tilde{q}^{\beta}\tilde{q}^{\alpha}}{\Lambda^{2}}t_{\beta\nu} \left( f^{*1,\mu\nu}f^{*2}_{\mu\alpha} + f^{*2,\mu\nu}f^{*1}_{\mu\alpha} \right) + c_{4}\frac{\tilde{q}^{\nu}\tilde{q}^{\mu}}{\Lambda^{2}}t_{\mu\nu}f^{*1,\alpha\beta}f^{*(2)}_{\alpha\beta} \right. \\ &+ m_{V}^{2} \left( 2c_{5}t_{\mu\nu}\epsilon^{*\mu}_{V_{1}}\epsilon^{*\nu}_{V_{2}} + 2c_{6}\frac{\tilde{q}^{\mu}q_{\alpha}}{\Lambda^{2}}t_{\mu\nu}(\epsilon^{*\nu}_{V_{1}}\epsilon^{*\alpha}_{V_{2}} - \epsilon^{*\alpha}_{V_{1}}\epsilon^{*\nu}_{V_{2}}) + c_{7}\frac{\tilde{q}^{\mu}\tilde{q}^{\nu}}{\Lambda^{2}}t_{\mu\nu}\epsilon^{*}_{V_{1}}\epsilon^{*}_{V_{2}} \right) \\ &+ c_{8}\frac{\tilde{q}^{\mu}\tilde{q}^{\nu}}{\Lambda^{2}}t_{\mu\nu}f^{*1,\alpha\beta}\tilde{f}^{*(2)}_{\alpha\beta} + c_{9}t^{\mu\alpha}\tilde{q}_{\alpha}\epsilon_{\mu\nu\rho\sigma}\epsilon^{*\nu}_{V_{1}}\epsilon^{*\rho}_{V_{2}}q^{\sigma} \\ &+ \frac{c_{10}t^{\mu\alpha}\tilde{q}_{\alpha}}{\Lambda^{2}}\epsilon_{\mu\nu\rho\sigma}q^{\rho}\tilde{q}^{\sigma} \left( \epsilon^{*\nu}_{V_{1}}(q\epsilon^{*}_{V_{2}}) + \epsilon^{*\nu}_{V_{2}}(q\epsilon^{*}_{V_{1}}) \right) \bigg] \end{split}$$

•  $V_1V_2 \in \{ZZ, WW\}$ 

- t<sub>μν</sub>: wave function of X
- $c_1 = c_5 \neq 0$ : graviton with minimal couplings  $(2_m^+)$
- $c_1 \ll c_5$ : graviton + SM fields can propagate to extra dimensions  $(2_b^+)$
- $c_i \neq 0$ : models with higher-dimension operators

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Higgs properties (ATLAS+CMS)

# Spin 2 results: $2_m^+$ model

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- CMS: only in  $\gamma\gamma$  channel [CERN-PH-EP/2014-117]
- ATLAS: combination of  $\gamma\gamma$ , ZZ and WW [Phys.Lett.B726(2013),120-144]
- Angular distribution in diphoton rest frame  $(|\cos \theta^*|)$  sensitive to spin:
  - ATLAS: simultaneous fit to  $m_{\gamma\gamma}$  and  $m_{\gamma\gamma} \times |\cos \theta^*|$  distributions
  - CMS: divide events in bins of  $|cos\theta^*|$  and fit  $m_{\gamma\gamma}$  in each of them



# Other spin-2 results and combinations

- C
- CMS tested also the spin 2 model with ED and with higher-dimension operators
- results combined from ZZ and WW channels



- Combination includes also  $1^{\pm}$
- All models excluded with  $\mathsf{CL}_s \geq 99.9\%$

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Higgs properties (ATLAS+CMS)

## Non-interfering spin 2 states

C

- As for spin 1, CMS studied the presence of narrow spin-2 states close to main resonance
  - ► fractional cross-section  $f(J^{P}) = \frac{\sigma_{J^{P}}}{\sigma_{0_{m}^{+}} + \sigma_{J^{P}}}$



# Conclusion from $J^{\mathrm{P}}$ measurements



- Spin and parity tested in  $\gamma\gamma,$  ZZ and WW channels
- Scalar hypothesis favoured by data
  - All alternatives rejected by > 99.9%
- CMS was studying also mixtures:
  - scalar-pseudoscalar
  - non-interfering spin 1 or 2 states
- Some limits on mixed states set but still large space for BSM physics



# Combination of all channels



- Combination of many channels is used to check Higgs properties (signal strength, couplings)
- Most analyses use full Run I statistics (5+20 fb<sup>-1</sup>)
  - CMS uses only 8 TeV data from  $t\bar{t}H \rightarrow$ leptons
  - ATLAS uses only 8 TeV data from H 
    ightarrow au au
  - final ATLAS results for most channels coming out this autumn

• [ATLAS-CONF-2014-009](ATLAS), [CMS-PAS-HIG-14-009](CMS)

Decay/ production tag	untagged	VBF	VH	tτH
$H \rightarrow \gamma \gamma$	both	both	both	CMS
$H  ightarrow b ar{b}$			both	CMS
$H \rightarrow \tau^+ \tau^-$	both	both	both	CMS
$H \rightarrow W^+ W^-$	both	both	CMS	CMS
$H \rightarrow ZZ$	both	both	both	CMS

. / 33

### SM compatibility: signal strength ATLAS



#### Method

- Use all channels
- Test statistics  $q_{\mu}$ ,  $\hat{\mu} = \sigma / \sigma_{\rm SM}$

### CMS



13

# SM compatibility: 2D signal strength



- Test statistics  $q(\mu_{\rm ggH+ttH}, \mu_{\rm qqH+VH})$ , 2 + 2 production modes grouped together
- Decays as in SM, cross-channel contamination evaluated from MC



### Compatibility of couplings Scaling factors



$$N(xx \to H \to yy) \sim \sigma(xx \to H) \cdot \mathcal{B}(H \to yy) \sim \frac{\Gamma_{xx}\Gamma_{yy}}{\Gamma_{\text{tot}}}$$

- 8 independent parameters relevant for current searches
- $\Gamma_{ZZ}$ ,  $\Gamma_{WW}$ ,  $\Gamma_{\tau\tau}$ ,  $\Gamma_{bb}$ ,  $\Gamma_{\gamma\gamma}$ ,  $\Gamma_{gg}$ ,  $\Gamma_{tt}$ ,  $\Gamma_{tot}$
- Not possible to extract those parameters at the moment
- Scaling factors for couplings:  $\mathbf{g}_{i} = \kappa_{i} \cdot \mathbf{g}_{i}^{SM}$
- Introducing  $\Gamma_{\rm BSM}$
- Following slides are **compatibility tests**, not measurements
- Significant deviation of  $\kappa$ 's from 1 would mean BSM physics

Re-fit of event yields in particular BSM framework would be also needed



# Custodial symmetry



#### Testing $\lambda_{WZ} = \kappa_W / \kappa_Z$ , $\kappa_Z$ and $\kappa_f$





ATLAS

• fitting  $\lambda_{WZ}$ ,  $\kappa_Z$  and  $\kappa_F Z$ 

• 
$$\lambda_{WZ} = 0.94^{+0.14}_{-0.29}$$

- $\kappa_{ZZ} = \kappa_Z^2 / \kappa_H = 1.41^{+0.49}_{-0.34}$
- $\lambda_{FZ} = \kappa_F / \kappa_Z \in [-0.91, -0.63] \cup [0.65, 1.00]$

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# Custodial symmetry



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•  $\lambda_{FZ} = \kappa_F / \kappa_Z \in [-0.91, -0.63] \cup [0.65, 1.00]$ 

Data are consistent with custodial symmetry

Further tests assume

$$\kappa_W = \kappa_Z = \kappa_V$$

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# Couplings to fermions and W/Z: 2D contours

- Assume common scaling factors for fermion and W/Z couplings:  $\kappa_f,\,\kappa_V$
- $\Gamma_{\rm BSM}=0$
- $\Gamma_{gg} \sim \kappa_f^2$
- $\Gamma_{\gamma\gamma} \sim |\alpha\kappa_V + \beta\kappa_f|^2$  (W and t loop)  $\Rightarrow \gamma\gamma$  sensitive to relative sign of  $\kappa_V$  and  $\kappa_f$



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# New physics in the loops: $\kappa_g$ and $\kappa_\gamma$

- C
- Loop diagrams sensitive to new particles,  $\kappa_g$  and  $\kappa_\gamma$  allow contributions from new particles
- $\Gamma_{\rm BSM} = 0$ , all other  $\kappa_i = 1$



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# Non SM Higgs decays

C

- Assume tree-level couplings are SM
- fit for  $\Gamma_{\rm BSM}$ ,  $\kappa_{\gamma}$  and  $\kappa_{g}$







# Non SM Higgs decays

C

- Assume tree-level couplings are SM
- fit for  $\Gamma_{\rm BSM}$ ,  $\kappa_{\gamma}$  and  $\kappa_{g}$



28 / 33

MS

# Fermion coupling asymmetries



### **ATLAS**



#### CMS



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# C6 model @ CMS



#### • Assume 6 independent parameters: $\kappa_V$ , $\kappa_t$ , $\kappa_b$ , $\kappa_\tau$ , $\kappa_\gamma$ , $\kappa_g$ ; $\Gamma_{BSM} = 0$



### 5 parameter model @ ATLAS

- Assume 5 independent parameters:  $\kappa_Z$ ,  $\kappa_W$ ,  $\kappa_t$ ,  $\kappa_b$ ,  $\kappa_\tau$ ; Only SM particles in the loops
- Results on k<sub>t</sub> from gg production





# Coupling compatibility: summary

Tested also generic 7 parameter model (free total width)



**ATLAS** 







# Coupling compatibility: summary







• Coupling to W,Z and 3rd generation fermions tested with 10-15% precision

In all tests no statistically significant deviation from the SM observed

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# Summary



- Final Run 1 mass measurements available
  - measured with precision 2-3 ‰
  - uncertainty dominated by statistics
  - $\gamma\gamma$  and ZZ 
    ightarrow 4/ results compatible within  $2\sigma$
- No statistically significant deviations from the SM couplings observed in any decay channels at both experiments
  - Constraints on 10-50% level
  - room for BSM
- Pure non-scalar hypotheses excluded
- Mixtures of states still allowed



# Summary



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Tightening the constraints on Higgs boson couplings and  $J^P$  mixtures will be main goal of Run II, LHC upgrades and future machines



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### **Additional material**



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Higgs properties (ATLAS+CMS

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34 / 33

CMS

# References and further reading



- ATLAS preliminary combination: ATLAS-CONF-2014-009
- ATLAS mass measurement: CERN-PH-EP-2014-122
- ATLAS spin testing: Phys.Lett.B726(2013),120-144
- ATLAS results in H $\rightarrow \gamma\gamma$  channel: CERN-PH-EP-2014-198
- ATLAS results in H $\rightarrow$  ZZ  $\rightarrow$  4/ channel: CERN-PH-EP-2014-170
- CMS preliminary combinations: CMS PAS HIG-14-009
- CMS anomalous spin 0 in HZZ: CMS PAS HIG-14-014
- CMS anomalous spin 0 in HWW: CMS PAS HIG-14-012
- CMS results in H $\rightarrow$  ZZ  $\rightarrow$  4/ channel: Phys.Rev.D89(2014)092007
- CMS results in H  $\rightarrow \gamma \gamma$  channel: CERN-PH-EP/2014-117
- Procedure for the LHC Higgs boson search combination: ATL-PHYS-PUB 2011-11, CMS NOTE 2011/005



Higgs cross-sections and BR's: CERN Yellow Report



### Statistical combination methodology Based on the approach agreed by ATLAS and CMS in http://cdsweb.cern.ch/record/1379837

Likelihood

$$\mathcal{L}(\text{data}|\mu \cdot s + b, \theta) = \mathcal{P}(\text{data}|\mu \cdot s + b, \theta) \cdot p(\tilde{\theta}|\theta)$$

*P*... Product of probabilities over all channels and all bins (or all events)
 *p*(θ̃|θ)... Probability of observing measured value θ̃ of nuisance parameter θ

#### Excess

LS

Test statistics: 
$$q_0 = -2 \ln \frac{\mathcal{L}(\mathrm{obs}|b,\hat{\theta}_0)}{\mathcal{L}(\mathrm{obs}|\hat{\mu}\cdot s+b,\hat{\theta})}, \ \hat{\mu} > 0$$

- $\mathcal{L}(\mathrm{obs}|b,\hat{ heta}_0)\dots$  maximal likelihood for background only hypothesis
- $\mathcal{L}(\operatorname{obs}|\hat{\mu} \cdot s + b, \hat{\theta}) \dots$  global maximal likelihood
- local p-value:  $p_0 = \mathsf{P}(q_0 \geq q_0^{\mathrm{obs}} | b)$
- significance Z:  $p_0 = \int_Z^\infty \frac{1}{\sqrt{2\pi}} e^{-\frac{x^2}{2}} dx$

Statistical combination methodology

C

Based on the approach agreed by ATLAS and CMS in http://cdsweb.cern.ch/record/1379837

Signal model parameters limits

Test statistics: 
$$q(a) = -2 \ln rac{\mathcal{L}(\mathrm{obs}|s(a)+b,\hat{ heta}_a)}{\mathcal{L}(\mathrm{obs}|s(\hat{a})+b,\hat{ heta})}$$

- The 68% (95%) CL on a given parameter of interest  $a_i$ :  $q(a_i) = 1$  (3.84)
- For 2D contours, The 68% (95%) CL on a given parameter of interest a<sub>i</sub>: q(a<sub>i</sub>, a<sub>j</sub>) = 2.3 (6.99)



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## Higgs cross-section and BR







38 / 33

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## Higgs cross-section and width









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### Mass in the ZZ subchannels







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### Limits on width







# Constraints on H $\rightarrow$ Z $\gamma^*/\gamma^*\gamma^*$



- $\bullet$  Constraints on the fractional presence of  $Z\gamma^*$  and  $\gamma^*\gamma^*$
- Similar approach as before (assuming real phases)



## Pair of spin-0 couplings



- Simultaneous presence of 2 anomalous ZZ couplings
- Profiling the phases of studied pair
- Other parameters have SM values



# Combination with $WW \rightarrow I \nu I \nu$

C

- Also in CMS-PAS-HIG-14-012
- $\bullet$  Assume same ratio of couplings in ZZ and WW channels
- Improving constraints wrt ZZ alone







# Signal strength by production tag









CMS




