



# CMS $H^+$ searches, part 2:

## $H^+ \rightarrow \tau^+ \nu_\tau$ in fully hadronic final state\*



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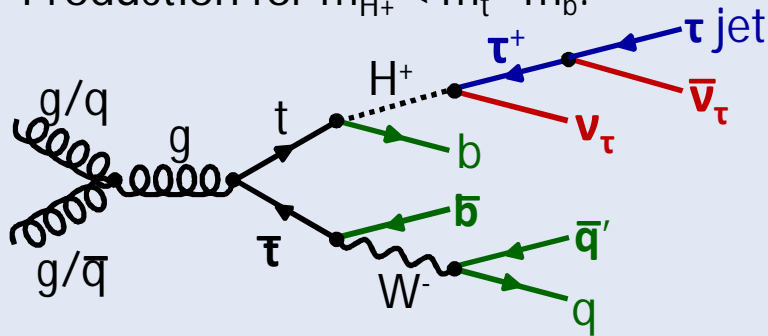
Overview of the latest CMS  $H^+$  analyses:

- $H^+ \rightarrow \tau^+ \nu_\tau$  and  $H^+ \rightarrow t\bar{b}$  in leptonic final states,  $m_{H^+} > m_t - m_b$  (see talk by P. Vischia)
- $H^+ \rightarrow \tau^+ \nu_\tau$  in fully hadronic final state,  $m_{H^+} < m_t - m_b$  and  $m_{H^+} > m_t - m_b$  (this talk)
- $H^+ \rightarrow c\bar{s}$ ,  $m_{H^+} < m_t - m_b$  (see talk by G. Kole)
- Data-driven background measurements in the  $H^+$  searches (see talk by M. Kortelainen)

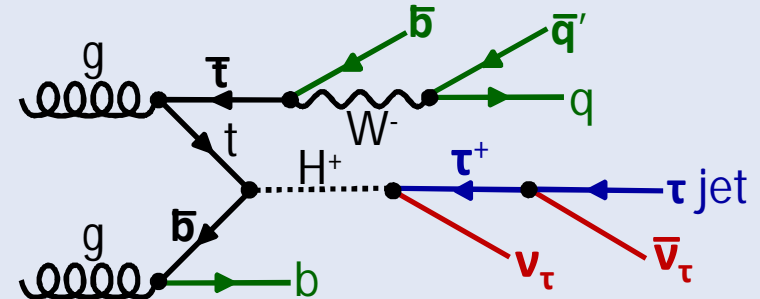
\* Charge conjugation is implied throughout this talk

# Overview of the $H^+ \rightarrow \tau^+ \nu_\tau$ (fully hadr.) analysis

Production for  $m_{H^+} < m_t - m_b$ :



Production for  $m_{H^+} > m_t - m_b$  (5FS):

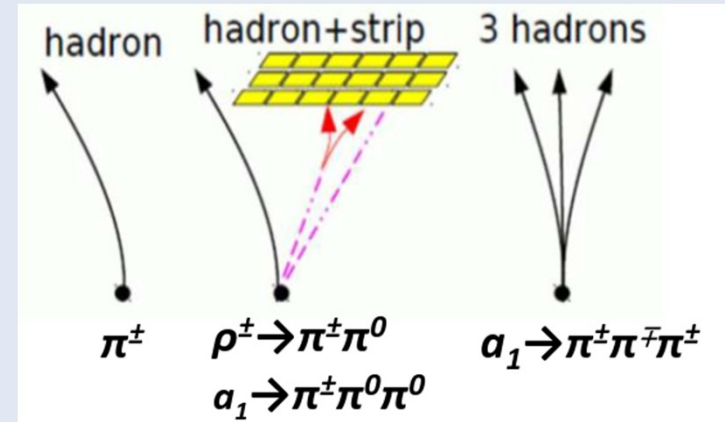


- For  $H^+ \rightarrow \tau^+ \nu_\tau$  decay mode, the final state is the same for light and heavy  $H^+$ 
  - Same methods can be used and only the thresholds need to be optimized
- Main backgrounds: QCD multijet,  $t\bar{t}$ , and single top
  - Multijet background (reducible): shape and normalization measured from data
  - EWK+ $t\bar{t}$  with  $\tau_h$ : shape and normalization measured from data
  - EWK+ $t\bar{t}$  no  $\tau_h$ : small background, mainly jet  $\rightarrow$   $\tau$  misid., estimated from simulation
- Speciality of the fully hadronic final state:
  - All neutrinos in the event come from  $H^+$  decay
    - I.e. can separate EWK+ $t\bar{t}$  from signal with transverse mass,  $m_T(\tau_h, E_t^{\text{miss}})$
    - This feature makes the fully hadr. final state for  $H^+ \rightarrow \tau^+ \nu_\tau$  the most sensitive final state!
  - In this analysis, we are able to separate the  $H^+ \rightarrow \tau^+ \nu_\tau$  signal from other possible  $H^+$  signals since in the other  $H^+$  signals the tau always comes from W decay and they are measured as part of EWK+ $t\bar{t}$  with  $\tau_h$
- In this talk, results are shown for 2012 data for both light and heavy  $H^+$

} More in talk by M. Kortelainen



- Single tau +  $E_T^{\text{miss}}$  trigger
  - $\tau p_T > 35 \text{ GeV}$ , loose isolation on  $\tau$ , and calorimetric  $E_T^{\text{miss}} > 70 \text{ GeV}$
- Particle-flow (PF) technique used
  - Combine the CMS subdetector information in an optimal way to identify electrons, muons, photons, charged hadrons and neutral hadrons
  - Construct composite objects such as taus, jets, and  $E_T^{\text{miss}}$  from these particles
- Primary vertex: choose vertex with highest  $\sum_{\text{tracks}} (p_T)^2$
- Taus with hadron-plus-strips (HPS) algorithm
  - $\tau$  candidates are required to match to the trigger  $\tau$  object
  - $p_T > 41 \text{ GeV}$ ,  $|\eta| < 2.1$ , and  $p_T(\text{leading ch. particle}) > 20 \text{ GeV}$
  - Reject electrons and muons
  - Require isolation
  - 1 charged particle associated to the  $\tau$  object
  - Helicity cut:  $R_\tau (p_{\text{leading ch. particle}} / p_{\tau \text{ jet}}) > 0.7$
- Electrons and muons for vetoing:
  - Require  $p_T > 15 \text{ GeV}$  (10 GeV for muons) and  $|\eta| < 2.5$
  - Require isolation and identification



- Jets:
  - Ignore jets overlapping with selected with  $\tau$ , i.e.  $\Delta R(\text{jet}, \tau_h) < 0.5$
  - Require jet  $p_T > 30$  GeV and  $|\eta| < 2.4$
  - Jet energy corrections are applied
  - Require at least 3 jets passing above criteria
- b tag:
  - Discriminator based on secondary vertices and track impact parameter information
  - 0.1 % rate for  $udsg \rightarrow b$  (30-60 % efficiency for identifying b jets)
  - Require at least 1 of the selected jets passes b tagging
- $E_T^{\text{miss}}$ :
  - PF  $E_T^{\text{miss}}$  used with jet energy corrections propagated to it
  - Require  $E_T^{\text{miss}} > 60$  GeV
- Pile-up mitigation:
  - Correct for the presence of charged and neutral particles from pile-up vertices in electron/muon/tau isolation and jet energy and identification

- In QCD multijet events:
  - The  $E_T^{\text{miss}}$  comes typically from overestimation of the  $\tau_h$  energy
  - The  $p_T$  vector of the misidentified  $\tau_h$  is typically almost back-to-back to  $E_T^{\text{miss}}$  vector in  $\phi$ 
    - This configuration maximizes the  $m_T$  and contaminates therefore the signal region
  - The  $p_T$  vector of the misidentified  $\tau_h$  is typically almost back-to-back to a jet (recoiling jet) in  $\phi$ 
    - I.e.  $\Delta\phi(E_T^{\text{miss}}, \tau_h) \sim 180^\circ$  and  $\min(\Delta\phi(E_T^{\text{miss}}, \text{jet}_N)) \sim 0^\circ$

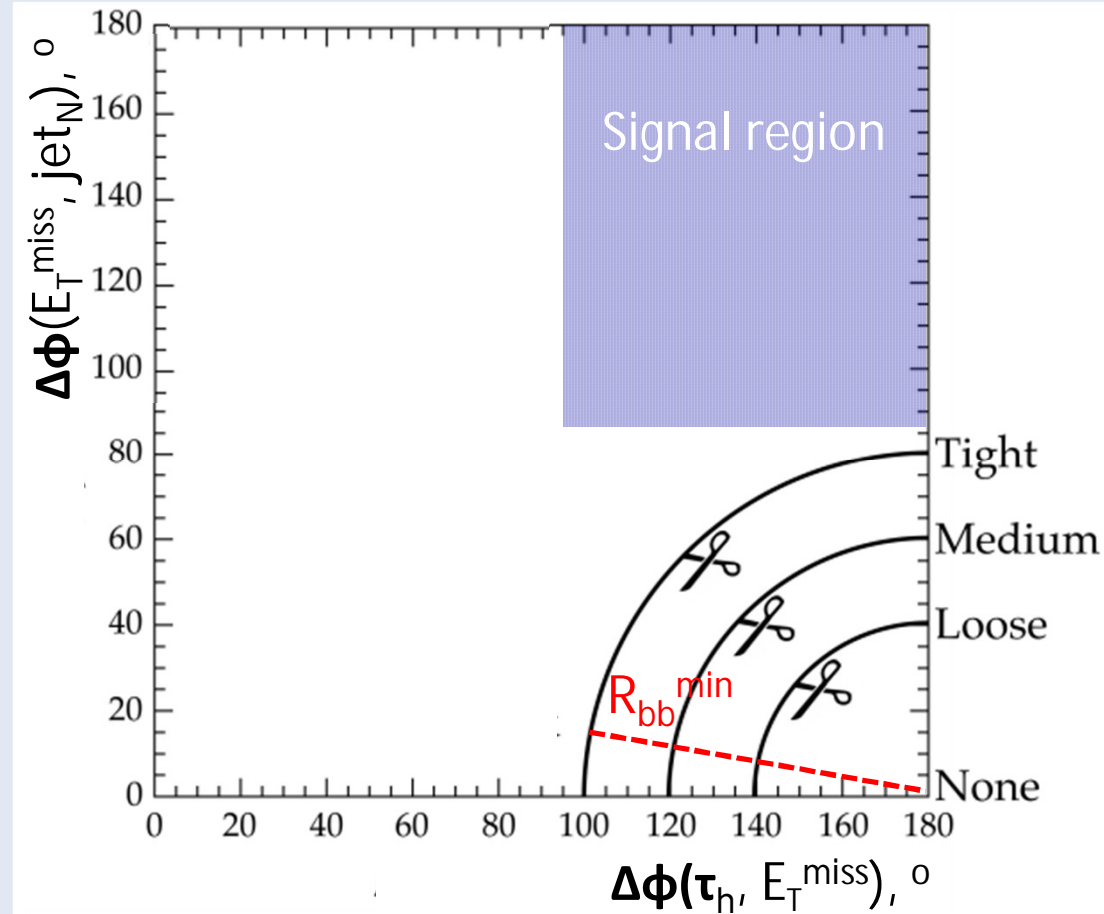
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  - I.e.  $\Delta\phi(E_T^{\text{miss}}, \tau_h) \sim 180^\circ$  and  $\min(\Delta\phi(E_T^{\text{miss}}, \text{jet}_N)) \sim 0^\circ$

- This feature can be used to suppress multijet events by cut on:

$$R_{bb}^{\text{min}} = \min\{\sqrt{\Delta\phi_{\text{MET}, \text{jet}_N}^2 + (180^\circ - \Delta\phi_{\tau, \text{MET}})^2}\}$$

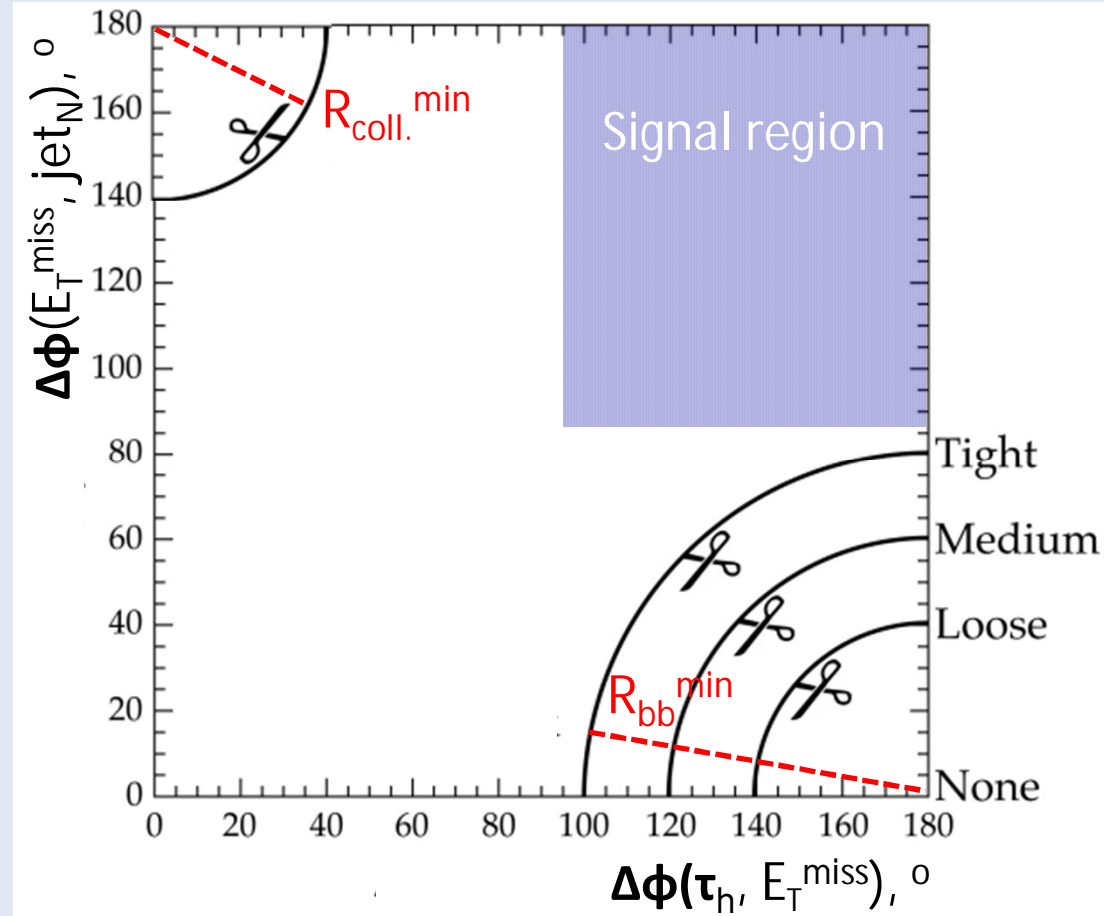
where  $N = 1..3$  highest- $p_T$  jet in the event



$$m_T = \sqrt{2 E_T^{\text{miss}} p_T^\tau (1 - \cos \Delta\phi(E_T^{\text{miss}}, \tau_h))}$$

• Additionally:

- The  $E_T^{\text{miss}}$  can also come from the from underestimation of the  $\tau_h$  energy
- The  $p_T$  vector of the misidentified  $\tau_h$  is almost collinear to  $E_T^{\text{miss}}$  vector in  $\phi$ 
  - This configuration minimizes the  $m_T$ , i.e. It does not contaminate the signal region, but causes a small correlation in the multijet measurement if a cut on  $R_{bb}^{\text{min}}$  is applied
- The  $p_T$  vector of the misidentified  $\tau_h$  is typically almost back-to-back to a jet (recoiling jet) in  $\phi$ 
  - I.e.  $\Delta\phi(E_T^{\text{miss}}, \tau_h) \sim 0^\circ$  and  $\min(\Delta\phi(E_T^{\text{miss}}, \text{jet}_N)) \sim 180^\circ$



• This feature can be used to make a cut on:

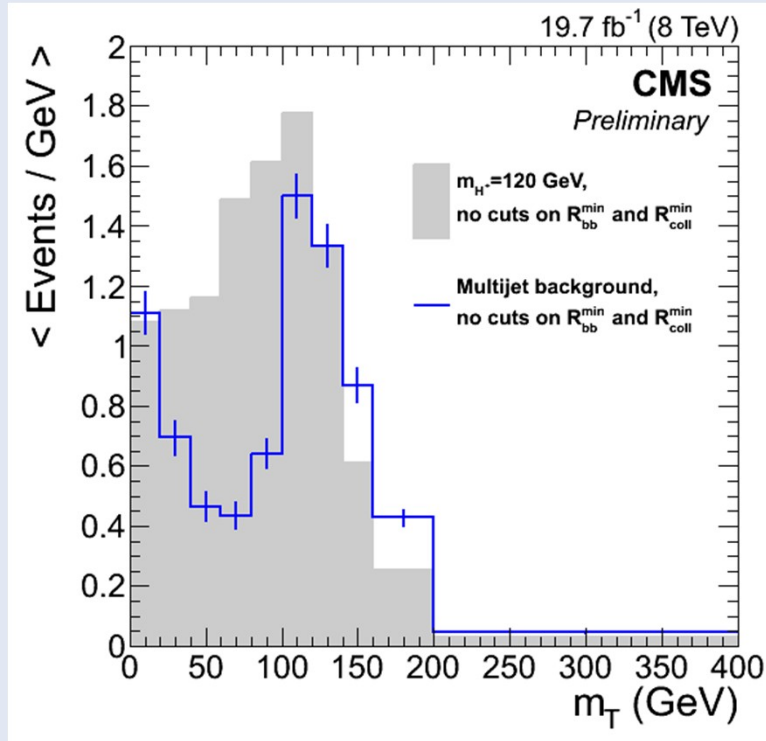
$$R_{\text{coll.}}^{\text{min}} = \min\{\sqrt{(180^\circ - \Delta\phi_{\text{MET}, \text{jet}_N})^2 + \Delta\phi_{\tau, \text{MET}}^2}\}$$

where  $n = 1..3$  highest- $p_T$  jet in the event

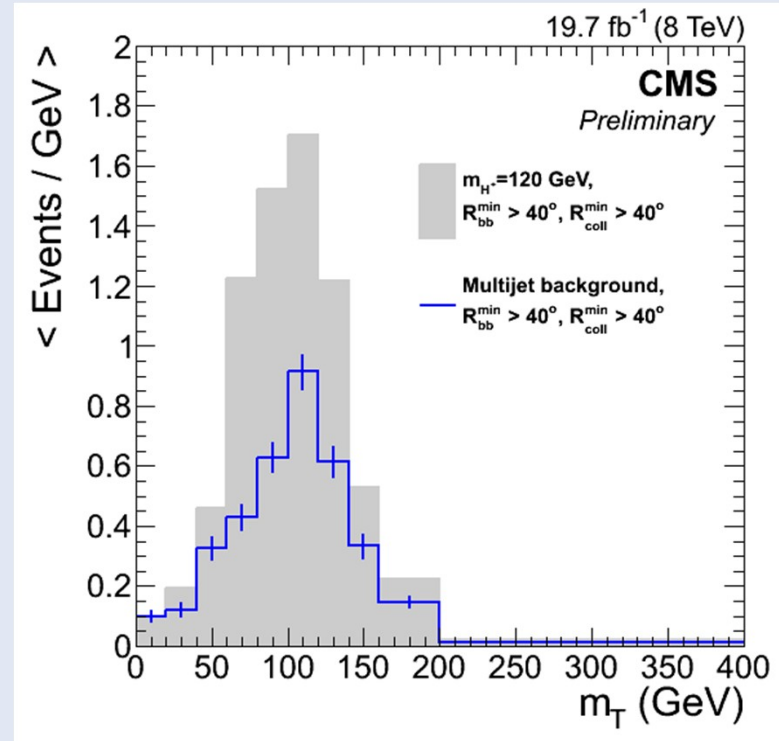
$$m_T = \sqrt{2 E_T^{\text{miss}} p_T^\tau (1 - \cos \Delta\phi(E_T^{\text{miss}}, \tau_h))}$$



Number of selected events from signal and multijet background with no angular cuts applied



Number of selected events from signal and multijet background with  $R_{\text{coll}}^{\text{min}} > 40^\circ$  and  $R_{bb}^{\text{min}} > 40^\circ$  applied



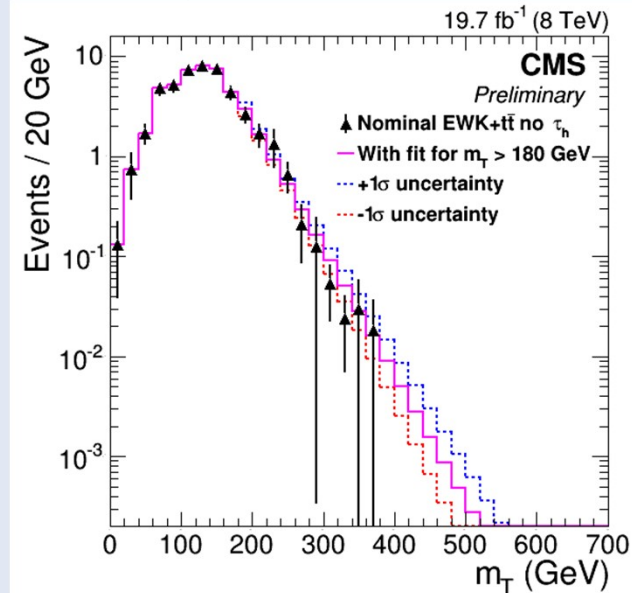
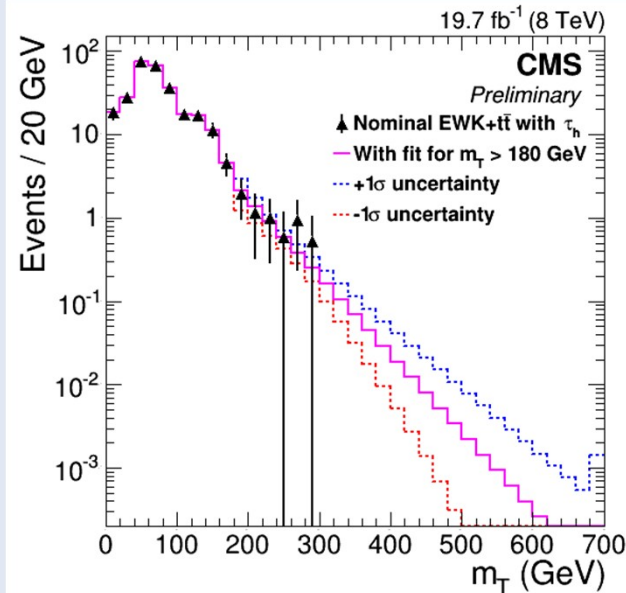
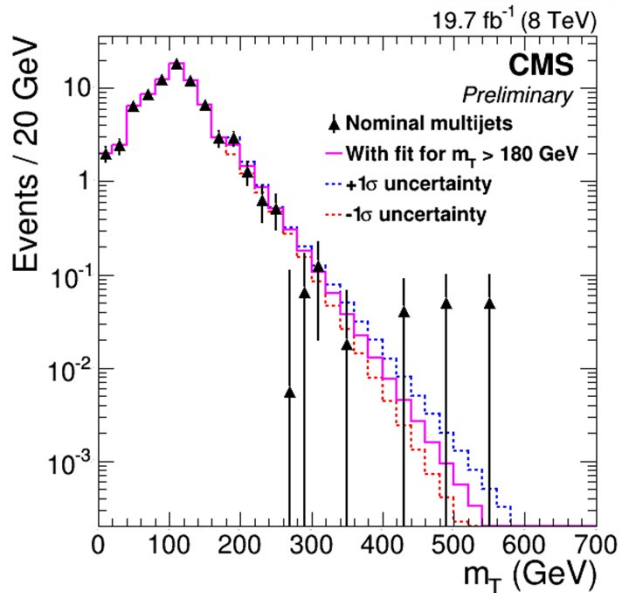
- $R_{\text{coll}}^{\text{min}} > 40^\circ$  and  $R_{bb}^{\text{min}} > 40^\circ$  used in the analysis

	Signal $H^+H^-$	Signal $H^+W^-$	Signal $H^+$	Multi- jets	EWK+ $t\bar{t}$ with $\tau_h$	EWK+ $t\bar{t}$ no $\tau_h$
Trigger $\tau$ part, data (S)	1.5–1.8	1.3–1.5	1.8–3.0	0.5	1.2	1.4
Trigger $\tau$ part, sim. (S)	0.7–0.8	0.6–0.7	0.8–1.1	0.2		0.8
Trigger $E_T^{\text{miss}}$ part, data (S)	2.6–3.3	2.5–2.8	2.9–4.2	1.2	2.5	2.8
Trigger $E_T^{\text{miss}}$ part, sim. (S)	0.1	0.1	0.1	0.1		0.4
Trigger, approximation in $E_T^{\text{miss}}$					12	
Trigger, single $\mu$ (S)					0.1	
$\tau_h$ identification (S)	6.0	6.0	5.9–6.0	0.8	6.0	
$e \rightarrow \tau_h$ misident. (S)	< 0.1	< 0.1	< 0.1	0.1		3.3
$\mu \rightarrow \tau_h$ misident. (S)	< 0.1	< 0.1	< 0.1	< 0.1		1.1
Jet $\rightarrow \tau_h$ misident. (S)	0.1	0.1–0.3	0.1	6.9		17
Veto of events with $e/\mu$	0.4–0.5	0.6–0.7	0.5–0.6	< 0.1		1.6
$\tau$ energy scale (S)	0.3–2.6	2.7–5.2	0.3–2.7	1.8	5.8	2.0
Jet energy scale (S)	2.6–5.2	2.0–3.0	1.6–2.1	1.4		3.2
Jet energy resolution (S)	1.1–1.8	0.5–1.3	0.7–1.5	1.4		3.2
Unclustered $E_T^{\text{miss}}$ ES (S)	0.1–0.4	0.1–0.9	0.1–0.4	0.5		1.5
b tagging (S)	5.9–20	4.7–5.3	4.6–5.4	3.5		5.0
Probabilistic $m_T$ in single top						6.8
Top $p_T$ reweighting (S)				+5.6 –6.8		+11 –6.6
Multijet $m_T$ distribution shape (S)				4.6		
Multijet template fit				3.0		
$\mu$ identification (S)					< 0.1	
Multijet contamination					2.0	
$W \rightarrow \tau\nu_\tau \rightarrow \mu\nu_\mu\nu_\tau$ frac. (S)					1.2	
Non-emb./emb. difference					+14 –12	
Pileup modelling (S)	0.1–0.9	0.1–0.8	0.1–0.6	0.1		2.9
$t\bar{t}$ cross section	+5.2 –6.0	+5.2 –6.0		+1.8 –1.5		+4.5 –5.2
Single top cross section						1.0
W+jets, DY, VV cross section						0.1
Luminosity	2.6	2.6	2.6	0.8		2.6

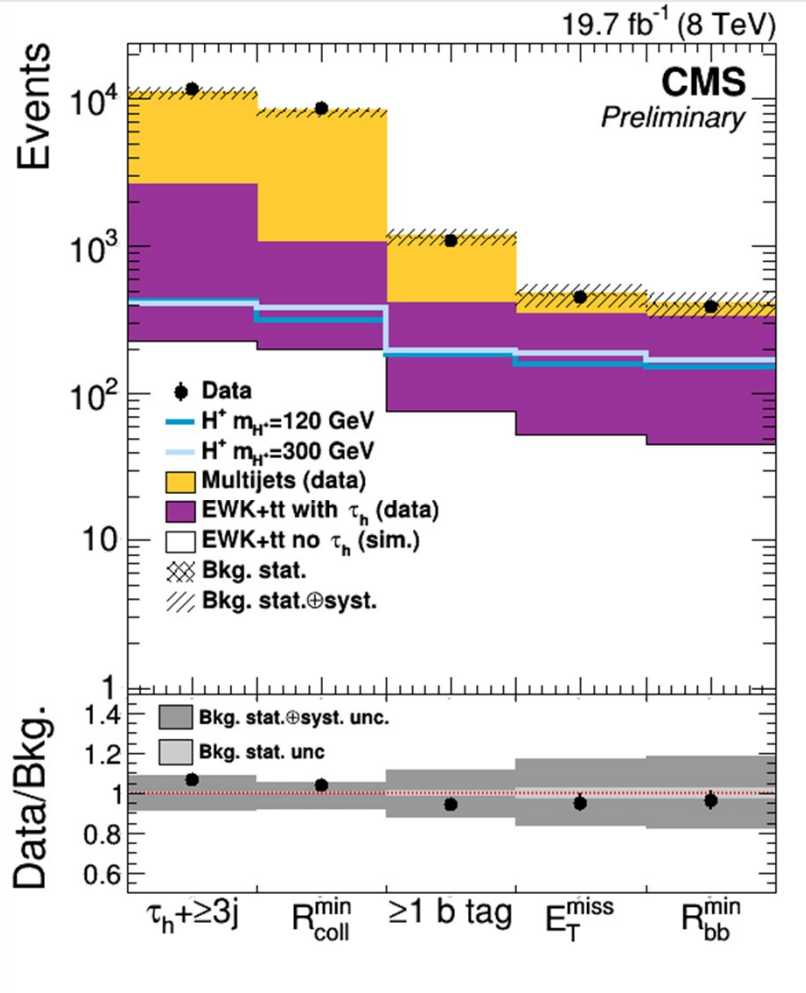
- Note that some of the backgrounds include simulation related uncertainties
  - Multijets: through subtraction of impurity from simulated EWK+tt (subtraction causes these uncertainties to be anti-correlated with other samples)
  - EWK+ttbar with taus: simulation of tau decay
- Most uncertainties change the  $m_T$  shape and are therefore treated as shape nuisances denoted by (S)
- Correlation is implied on the same line
- Dominating uncertainties:
  - Signal ( $t\bar{t} \rightarrow b\bar{b}H^+H^-$ ,  $t\bar{t} \rightarrow b\bar{b}H^+W^-$ ,  $pp \rightarrow t(b)H^+$ )
    - $\tau$  ID and b tag
  - QCD multijets:
    - Jet  $\rightarrow \tau$  mis-ID
    - top  $p_T$  reweighting
    - shape difference between non-isol./isol. sample
  - EWK+tt genuine taus:
    - Approximation in  $E_T^{\text{miss}}$  part of trigger
    - shape difference between non-emb./emb. sample
  - EWK+tt mis-ID taus:
    - jet  $\rightarrow \tau$  mis-ID
    - top  $p_T$  reweighting

- In the signal or background distributions, an uncorrelated nuisance parameter is assigned for each bin of the  $m_T$  templates
- In addition, for  $m_{H^+}=180-600$  GeV, the falling part of the  $m_T$  distribution in the backgrounds is fitted to  $p_0 * \exp\{-p_1 * (m_T - 180 \text{ GeV})\}$ 
  - The relative syst. uncertainties are kept constant
  - The stat. uncertainties in the fit region are given by the uncertainties on the fit parameters

Fit on the falling part of the  $m_T$  distributions for the different backgrounds



Number of events after each selection step



Number of selected events after all selections

	$N_{\text{events}} \pm \text{stat.} \pm \text{syst.}$
Signal, $m_{H^+} = 120 \text{ GeV}$	$151 \pm 4 \begin{smallmatrix} +17 \\ -18 \end{smallmatrix}$
Signal, $m_{H^+} = 300 \text{ GeV}$	$168 \pm 2 \pm 16$
Multijet background (data)	$78 \pm 3 \pm 17$
EWK+ $t\bar{t}$ with $\tau_h$ (data)	$283 \pm 12 \begin{smallmatrix} +55 \\ -54 \end{smallmatrix}$
EWK+ $t\bar{t}$ no $\tau_h$ (sim.)	$47 \pm 2 \begin{smallmatrix} +11 \\ -10 \end{smallmatrix}$
Total expected from the SM	$407 \pm 12 \begin{smallmatrix} +59 \\ -58 \end{smallmatrix}$
Observed:	392

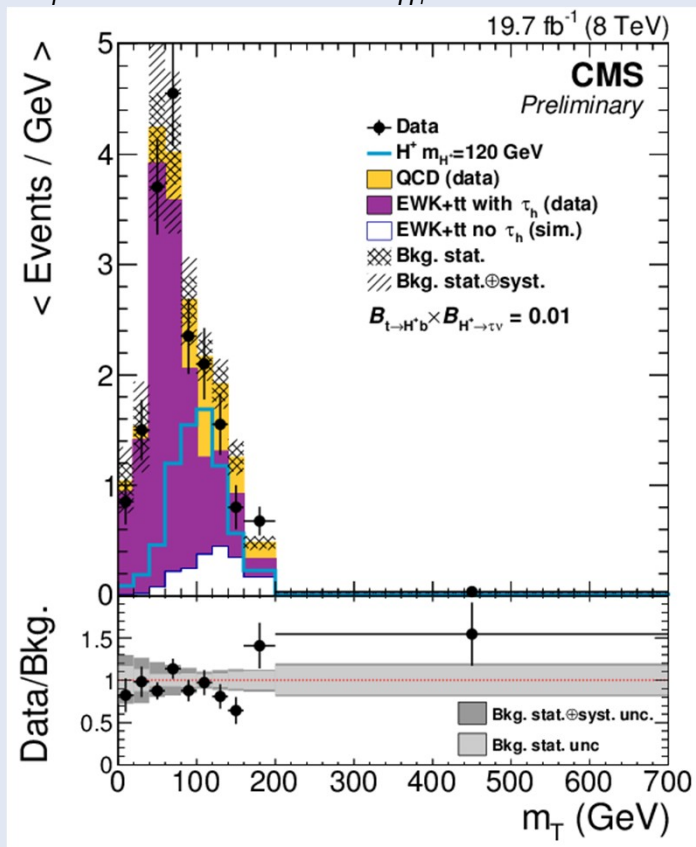
- Signal is normalized here for illustration purposes to:

$$B(t \rightarrow b H^+) * B(H^+ \rightarrow \tau^+ \nu_\tau) = 0.01$$

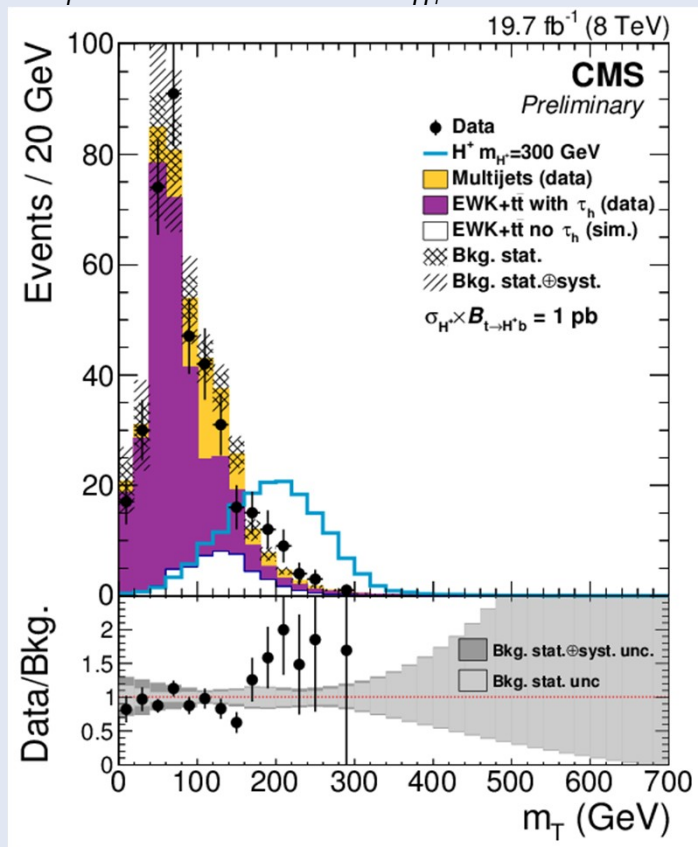
or

$$\sigma(pp \rightarrow t(b) H^+) * B(H^+ \rightarrow \tau^+ \nu_\tau) = 1 \text{ pb}$$

$m_T$  distributions for  $m_{H^\pm}=80-160$  GeV

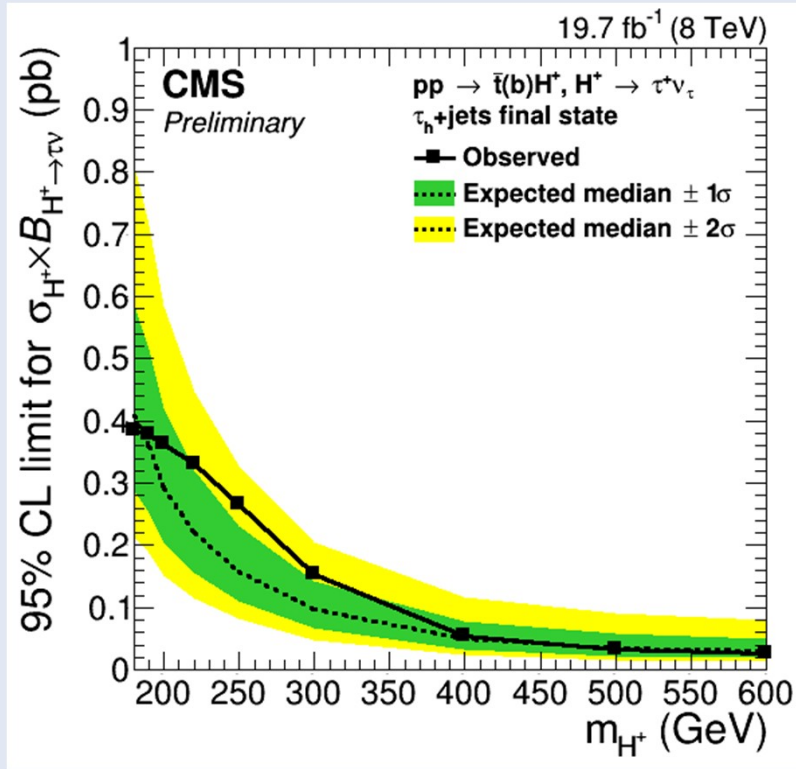
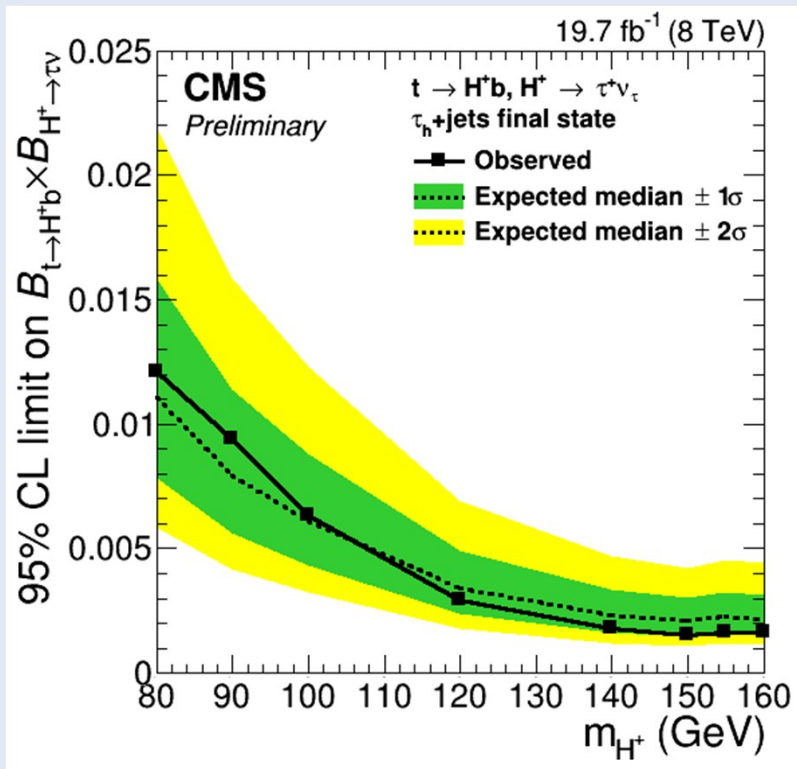


$m_T$  distributions for  $m_{H^\pm}=180-600$  GeV



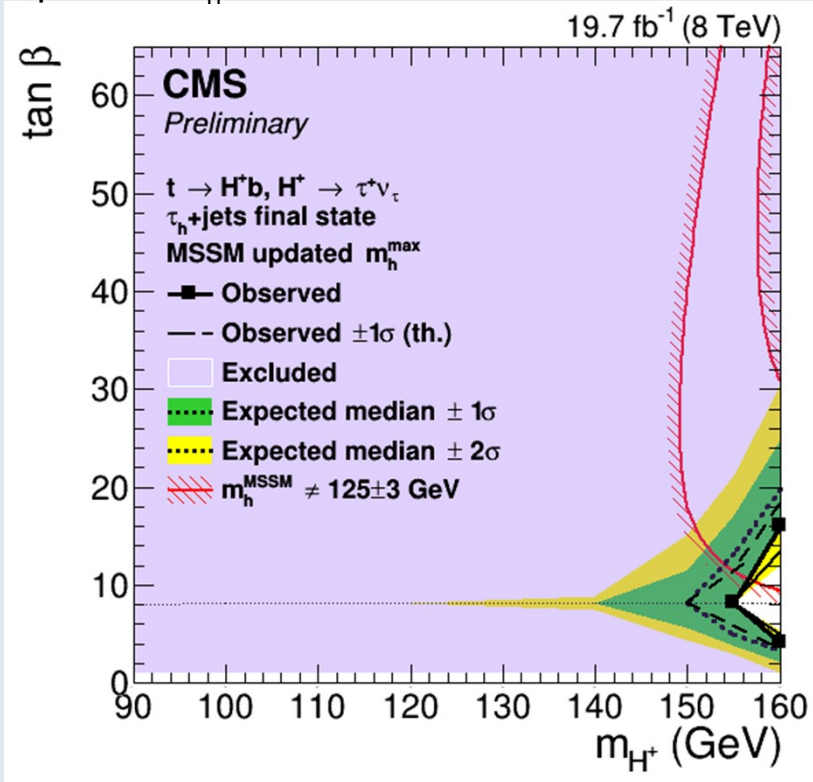
- Signal is normalized here for illustration purposes to  $B(t \rightarrow bH^\pm) \times B(H^\pm \rightarrow \tau \nu_\tau) = 0.01$  or  $\sigma(pp \rightarrow t(b)H^\pm) \times B(H^\pm \rightarrow \tau \nu_\tau) = 1$  pb
- For  $m_{H^\pm}=180-600$  GeV, the fit on the  $m_T$  tail is used
- Signal separation from backgrounds improves as a function of  $m_{H^\pm}$

- The statistical analysis is done with a binned maximum-likelihood fit
- The 95 % CL limits are derived with modified frequentist  $CL_s$  criterion
- Model-independent limits:
  - $B(t \rightarrow bH^+) * B(H^+ \rightarrow \tau^+ \nu_\tau)$  for  $m_{H^+} = 80-160$  GeV
  - $\sigma(pp \rightarrow t(b)H^+) * B(H^+ \rightarrow \tau^+ \nu_\tau)$  for  $m_{H^+} = 180-600$  GeV
- Model-dependent limits:
  - Calculated for MSSM benchmark scenarios [LHCHXSWG YR3 CERN-2013-004, Carena et al., EPJ C 73 (2013) 2552]
    - Updated  $m_h^{\max}$ ,  $m_h^{\text{mod}+}$ ,  $m_h^{\text{mod}-}$ , light stop, light stau,  $\tau$ -phobic, and low  $m_H$
  - Theoretical uncertainties of 21 % used for  $m_{H^+} = 80-160$  GeV and 32 % for  $m_{H^+} = 180-600$  GeV
    - Includes uncertainty from higher order  $t\bar{t}$  corrections for  $t\bar{t}$  production and 30 % uncertainty for  $pp \rightarrow t(b)H^+$  production
    - Includes uncertainty on one-loop EWK diagrams, missing two-loop QCD diagrams, and  $\Delta_b$  corrections for  $B(t \rightarrow bH^+)$  and  $B(H^+ \rightarrow \tau^+ \nu_\tau)$  [LHCHXSWG YR2 CERN-2012-002]

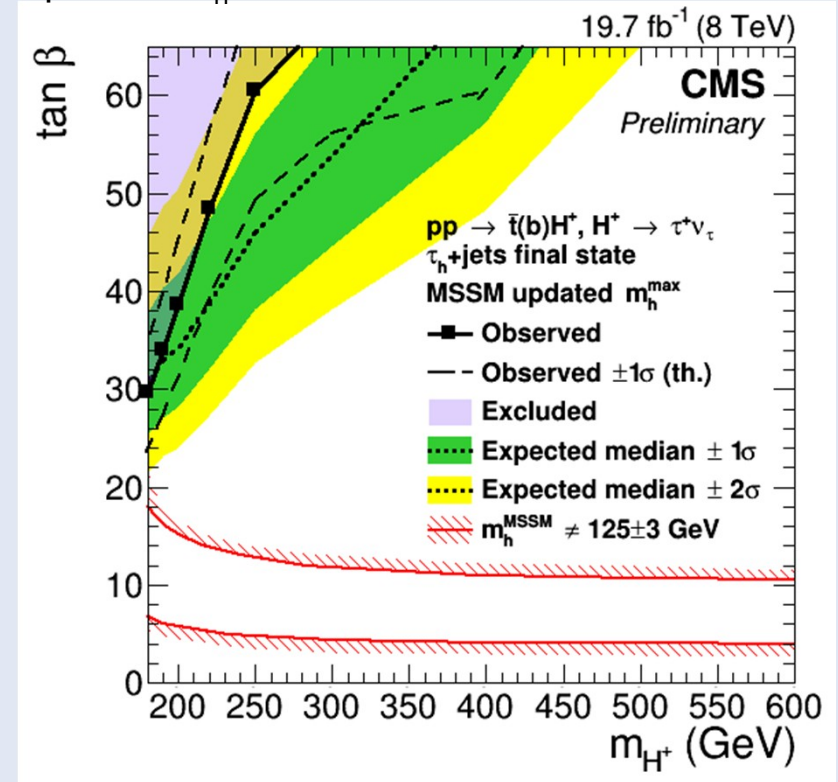


- 95 % CL upper limit on  $B(t \rightarrow b H^+) \times B(H^+ \rightarrow \tau^+ \nu_\tau)$ : 1.2–0.16 % for  $m_{H^+} = 80\text{--}160$  GeV
- 95 % CL upper limit on  $\sigma(pp \rightarrow t b H^+) \times B(H^+ \rightarrow \tau^+ \nu_\tau)$ : 0.38–0.026 pb for  $m_{H^+} = 180\text{--}600$  GeV

Updated  $m_h^{\max}$

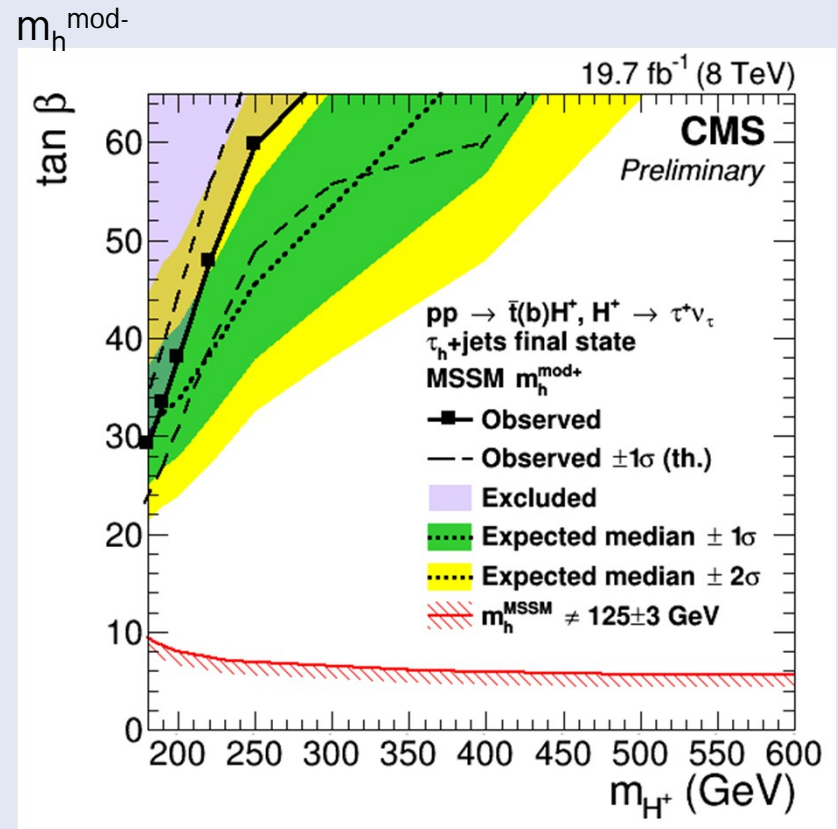
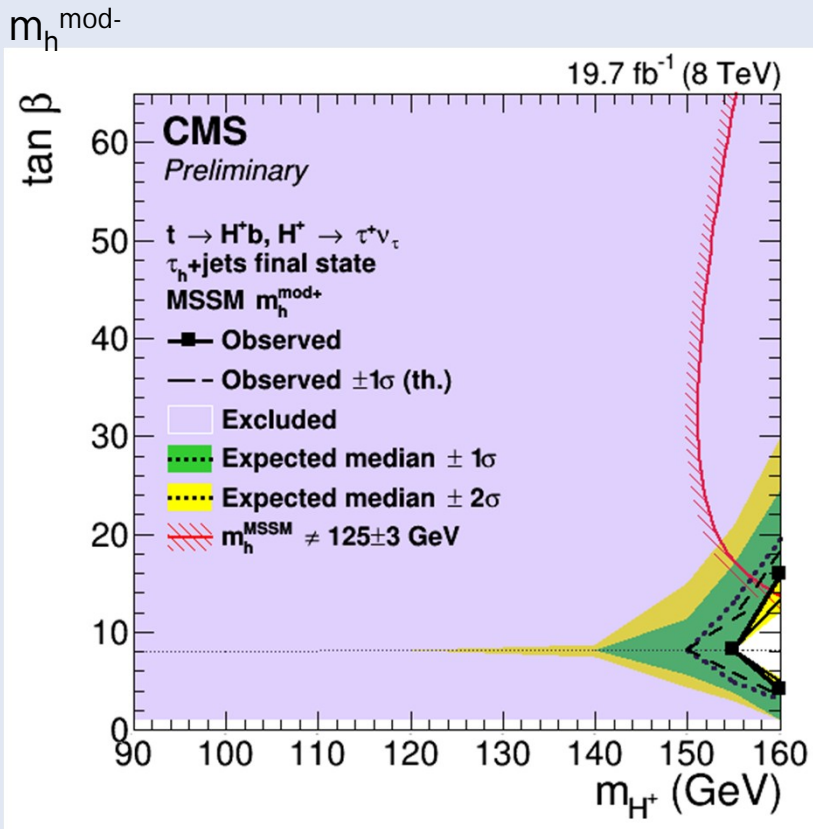


Updated  $m_h^{\max}$

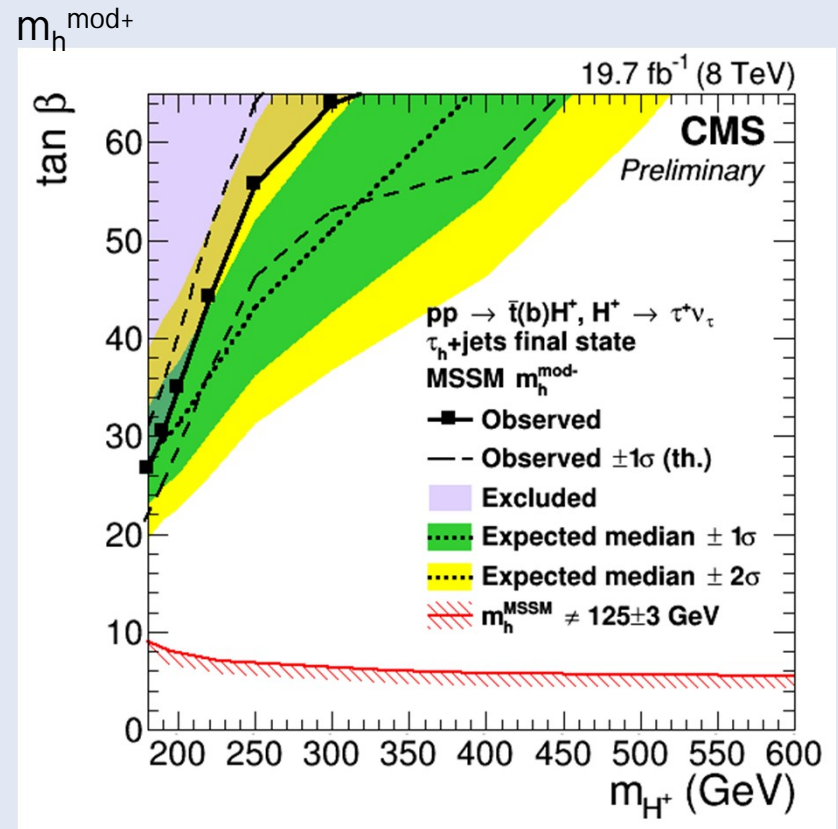
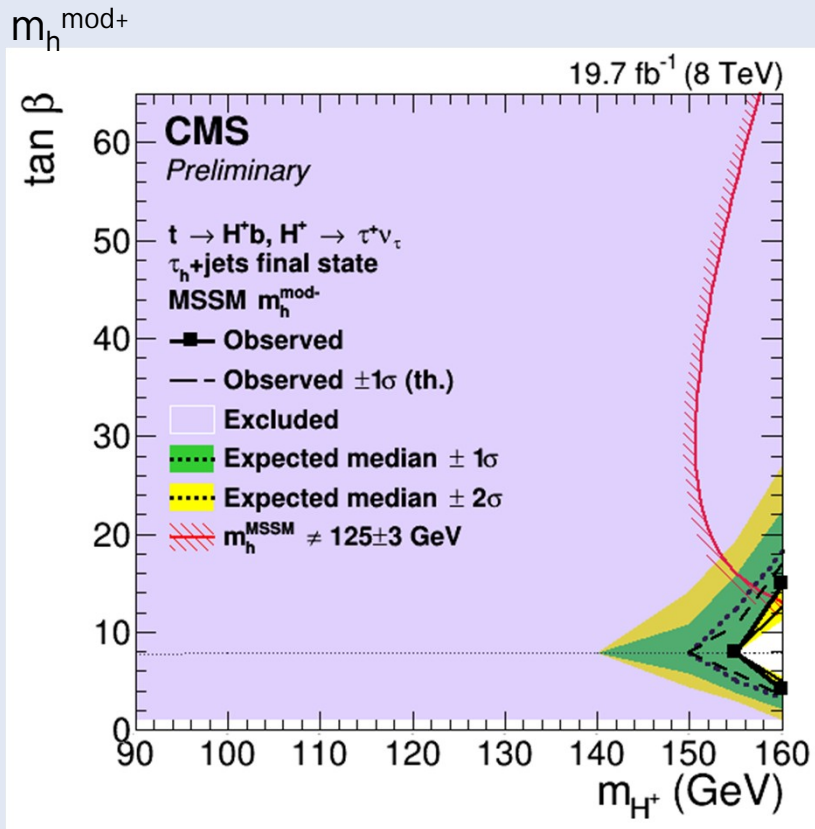


- In this scenario, the  $X_t$  value is chosen to maximize  $m_h$  at large values of  $m_A$
- The area between the red lines is the  $(m_{H^+}, \tan \beta)$  area, where  $m_h = 125.0 \pm 3.0$  GeV
- $m_{H^+}$  excluded for  $m_{H^+} < 155$  GeV



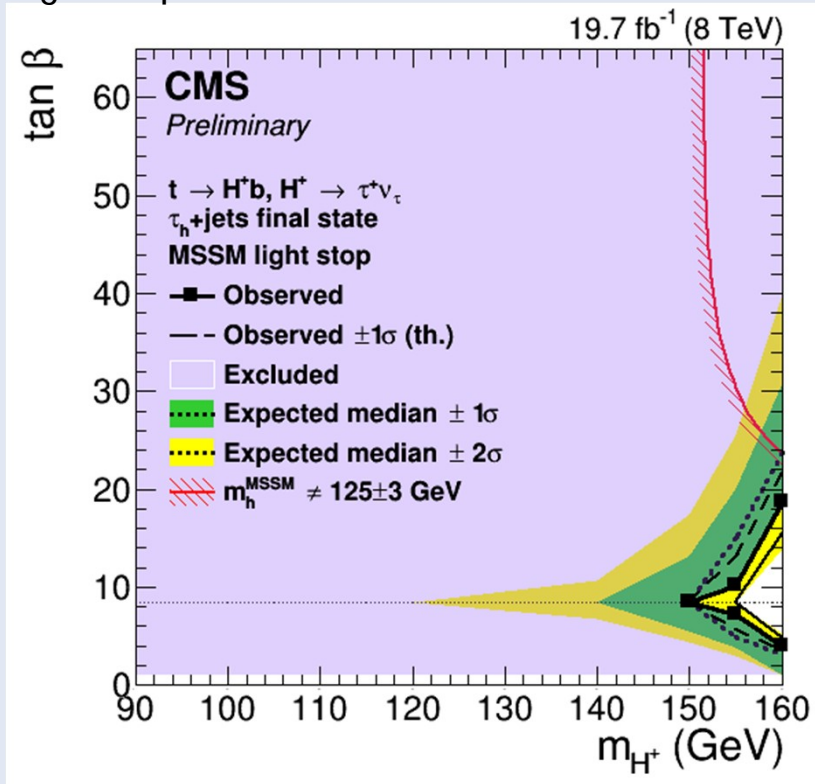


- In this scenario,  $|X_t/M_{\text{SUSY}}|$  is reduced compared to the  $m_h^{\text{max}}$  scenario and  $X_t > 0$
- The area above the red line is the  $(m_{H^+}, \tan \beta)$  area, where  $m_h = 125.0 \pm 3.0$  GeV
- $m_{H^+}$  almost completely excluded for  $m_{H^+} < 160$  GeV

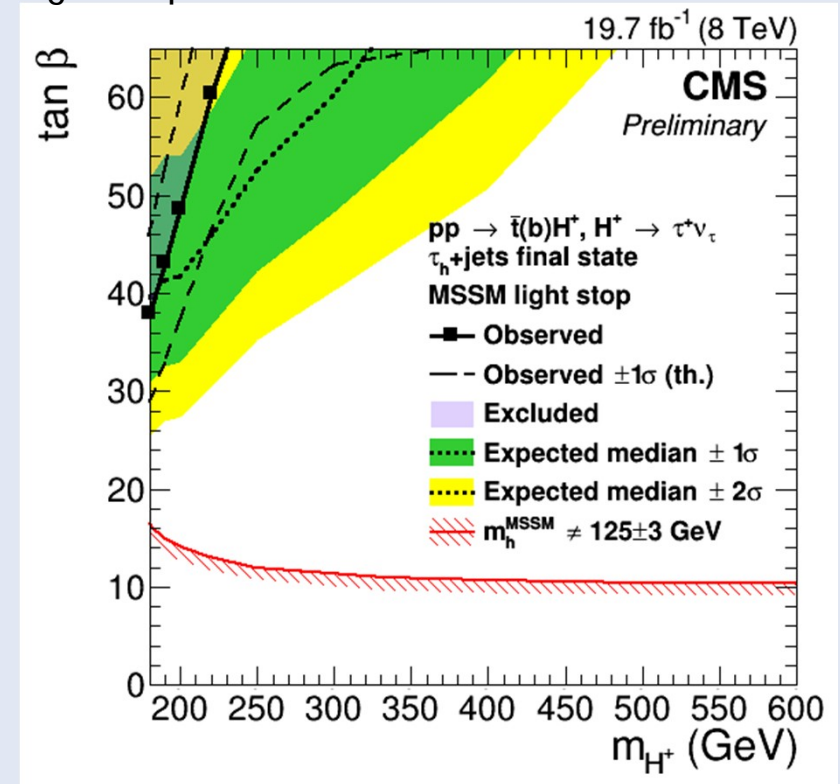


- In this scenario,  $|X_t/M_{\text{SUSY}}|$  is reduced compared to the  $m_h^{\text{max}}$  scenario and  $X_t < 0$
- The area above the red line is the  $(m_{H^+}, \tan \beta)$  area, where  $m_h = 125.0 \pm 3.0$  GeV
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Light stop

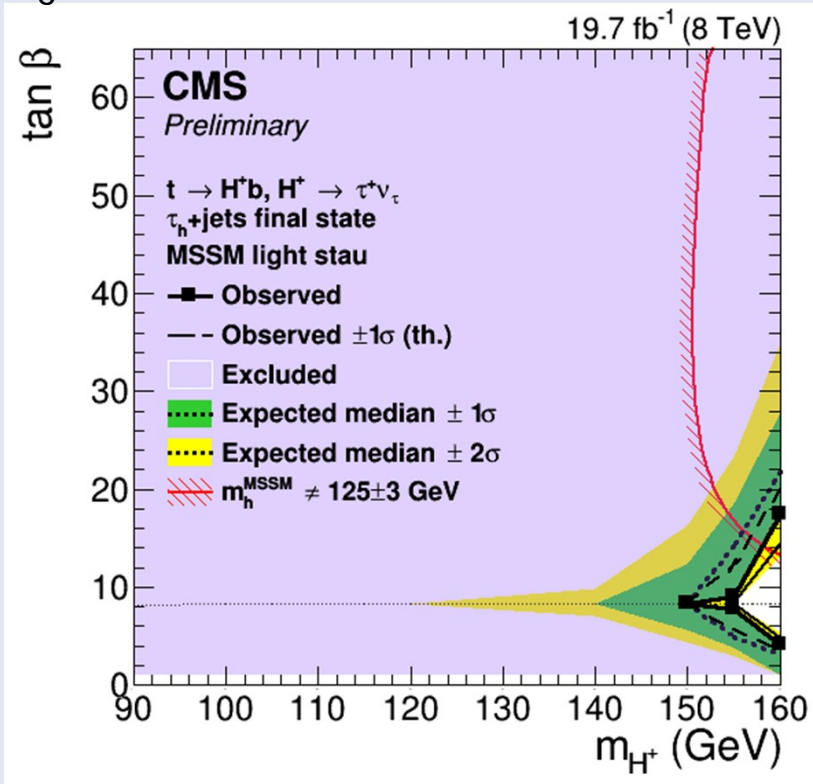


Light stop

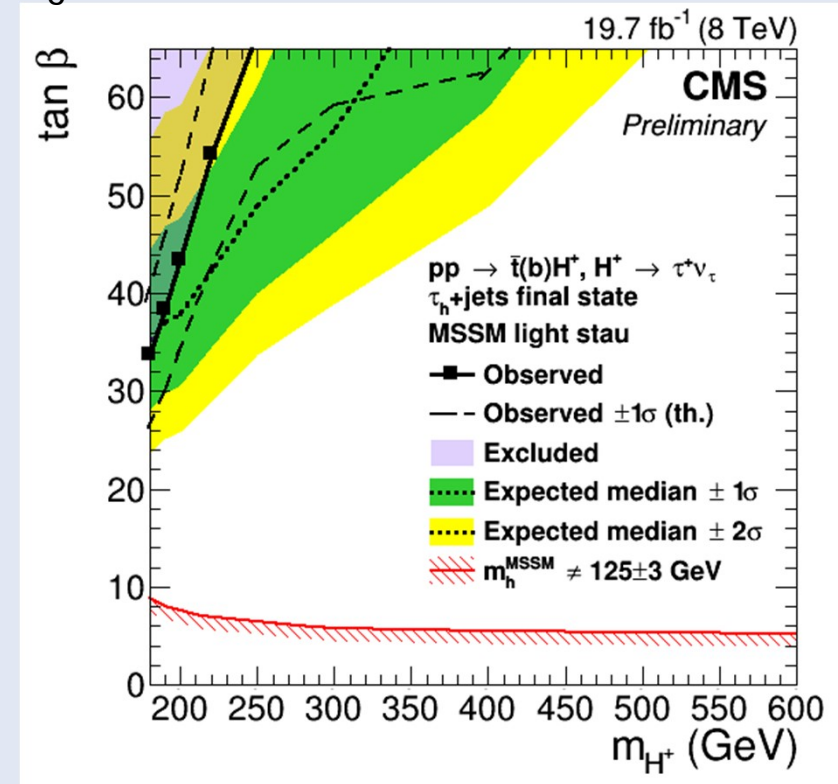


- In this scenario, the gluon-gluon fusion Higgs boson production is suppressed
- The area above the red line is the  $(m_{H^+}, \tan \beta)$  area, where  $m_h = 125.0 \pm 3.0$  GeV
- $m_{H^+}$  excluded for  $m_{H^+} < 160$  GeV

Light stau

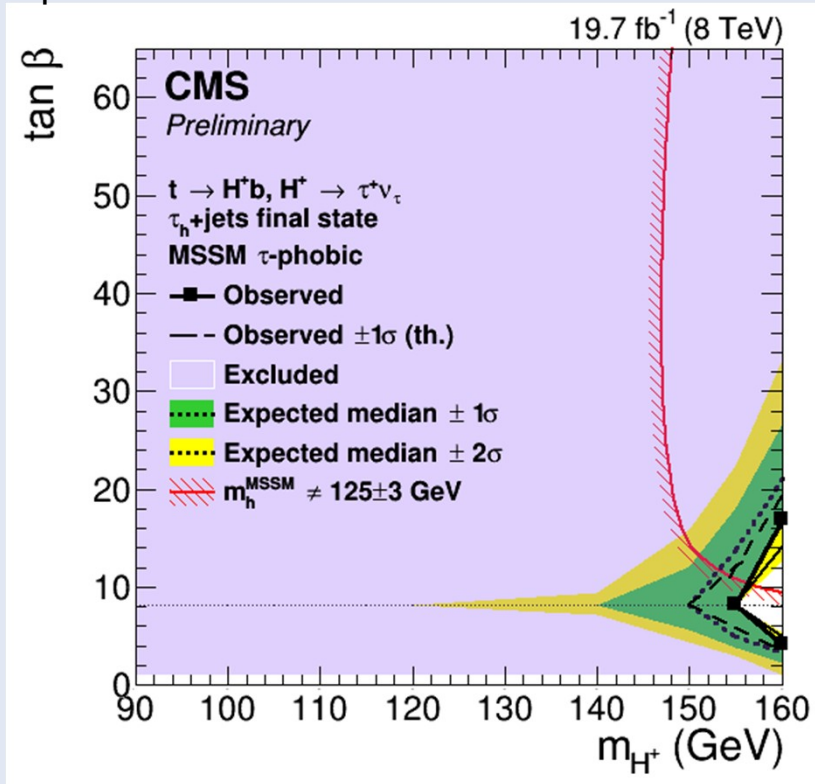


Light stau

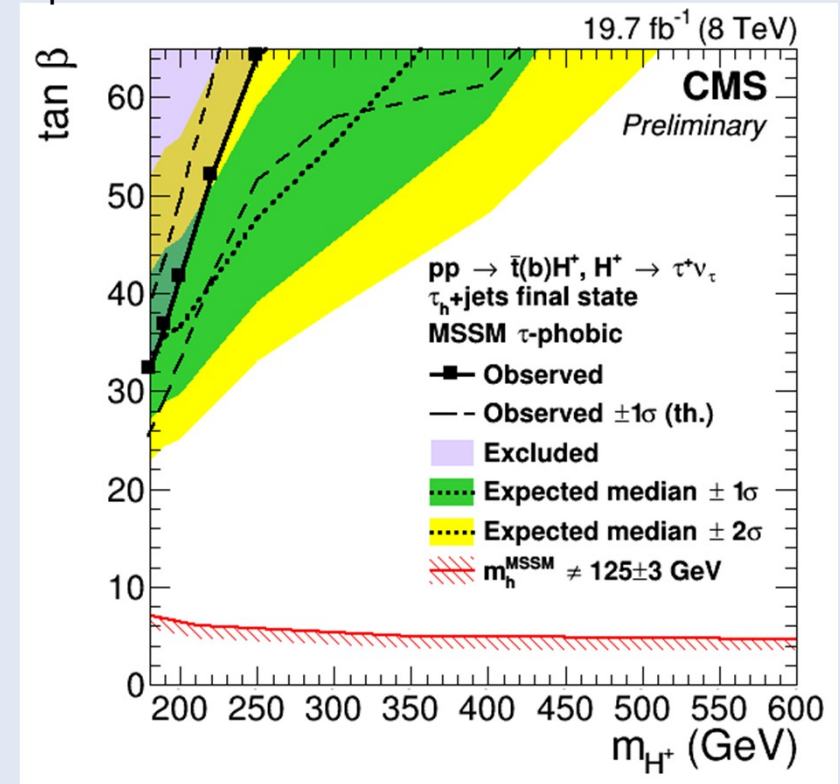


- In this scenario, the rate for  $h \rightarrow \gamma\gamma$  decays is enhanced
- The area above the red line is the  $(m_{H^+}, \tan \beta)$  area, where  $m_h = 125.0 \pm 3.0$  GeV
- $m_{H^+}$  almost completely excluded for  $m_{H^+} < 160$  GeV

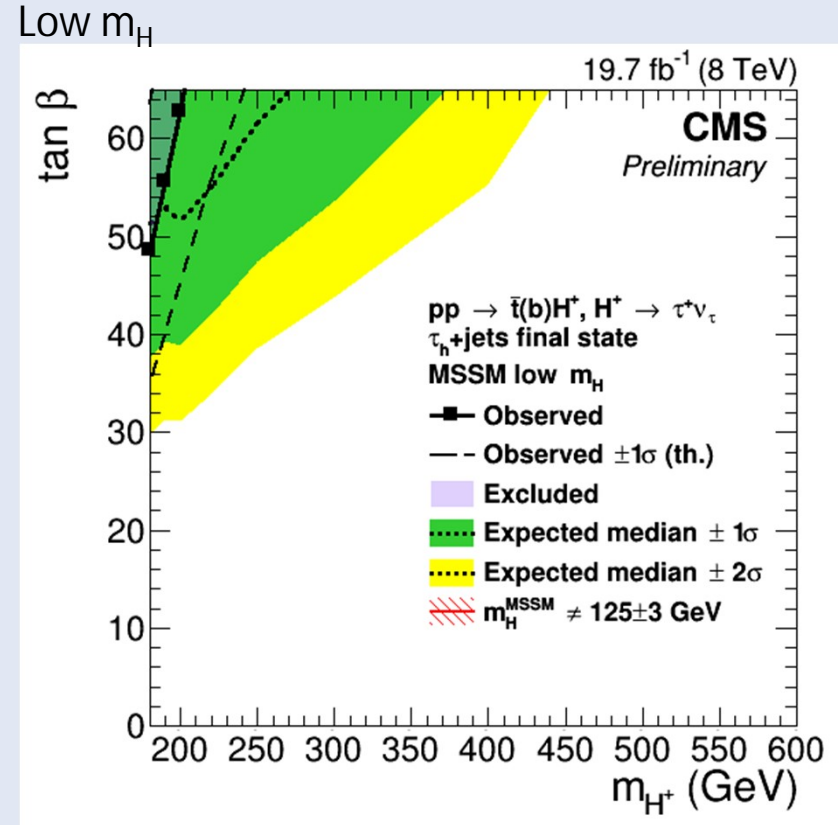
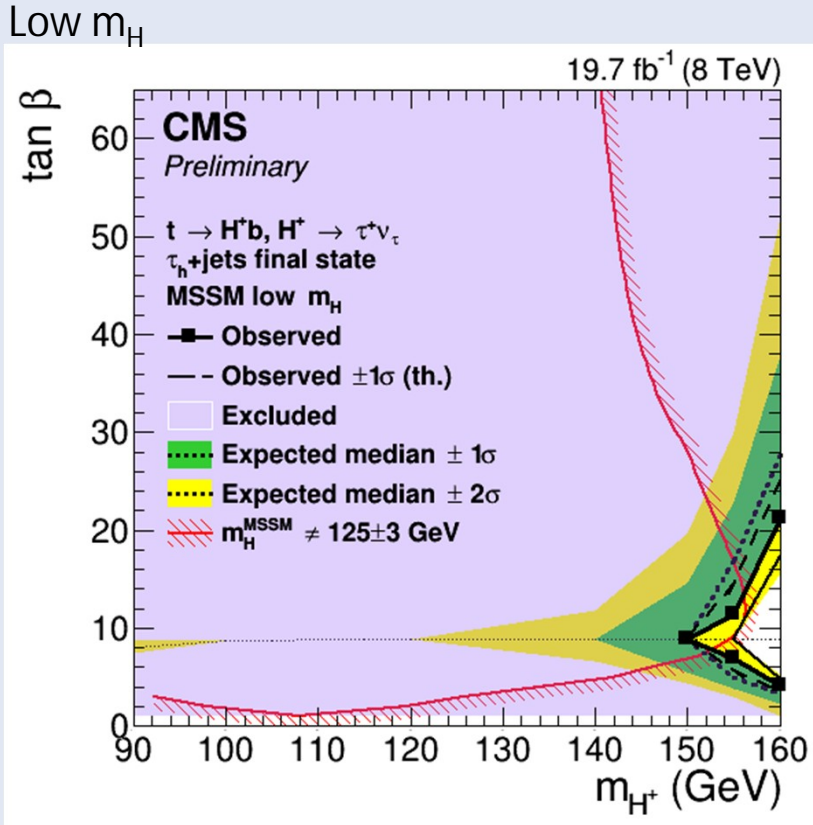
$\tau$ -phobic



$\tau$ -phobic



- In this scenario, the couplings to down-type fermions are suppressed
- The area above the red line is the  $(m_{H^+}, \tan \beta)$  area, where  $m_h = 125.0 \pm 3.0$  GeV
- $m_{H^+}$  excluded for  $m_{H^+} < 155$  GeV

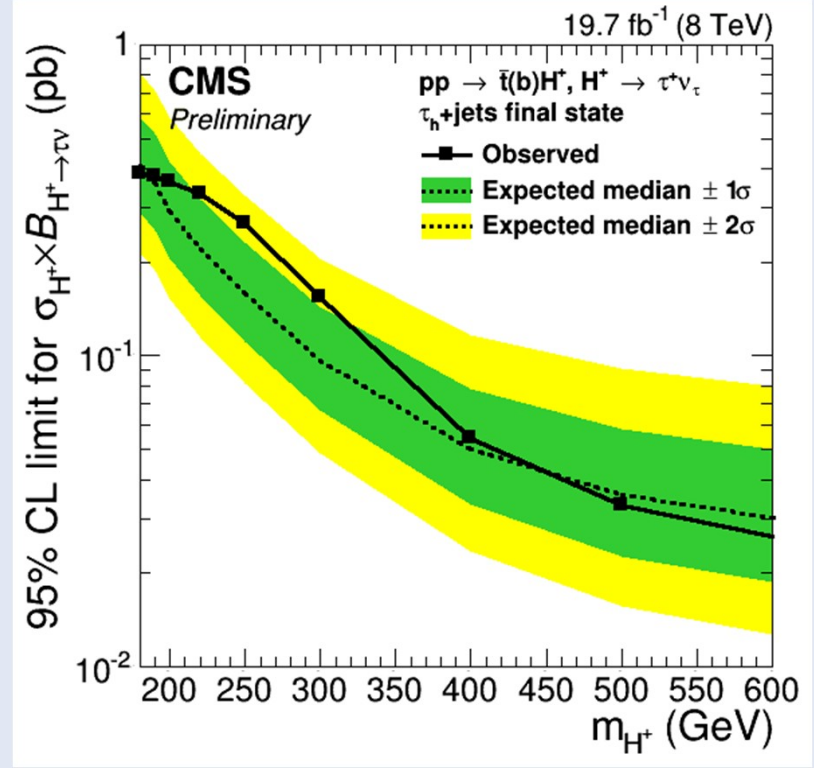
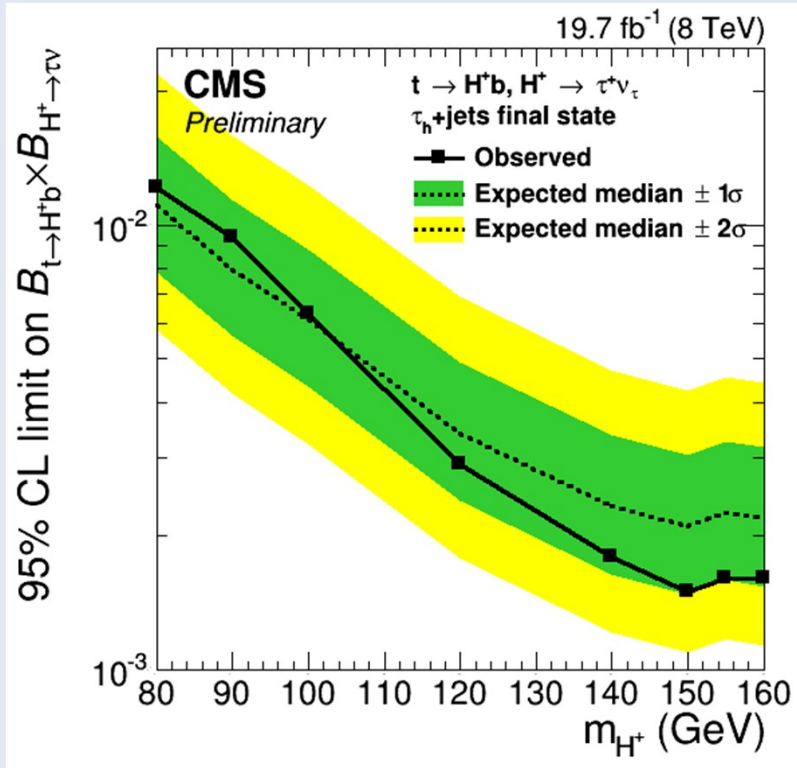


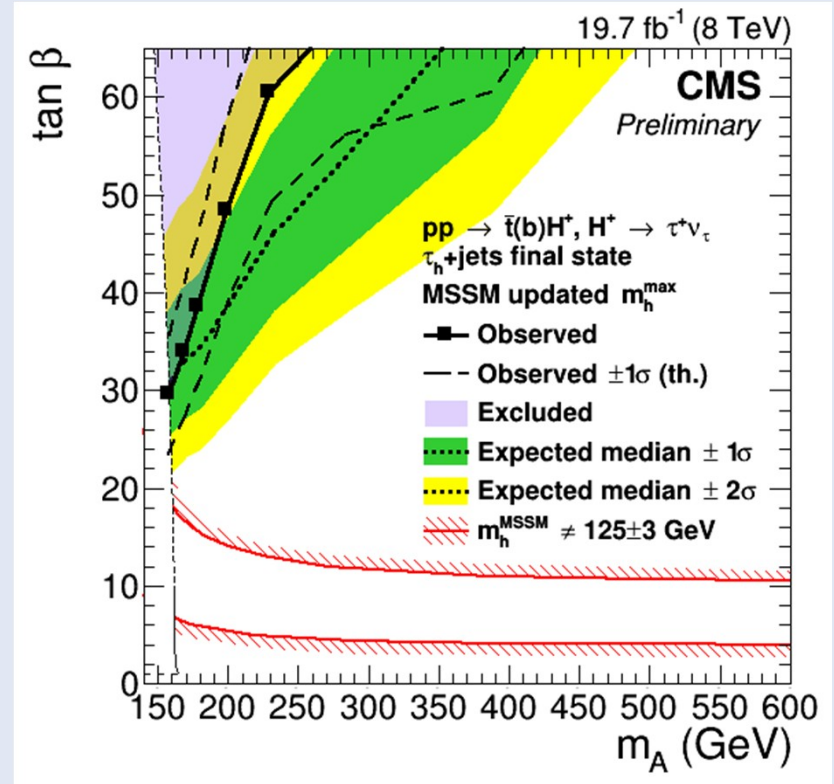
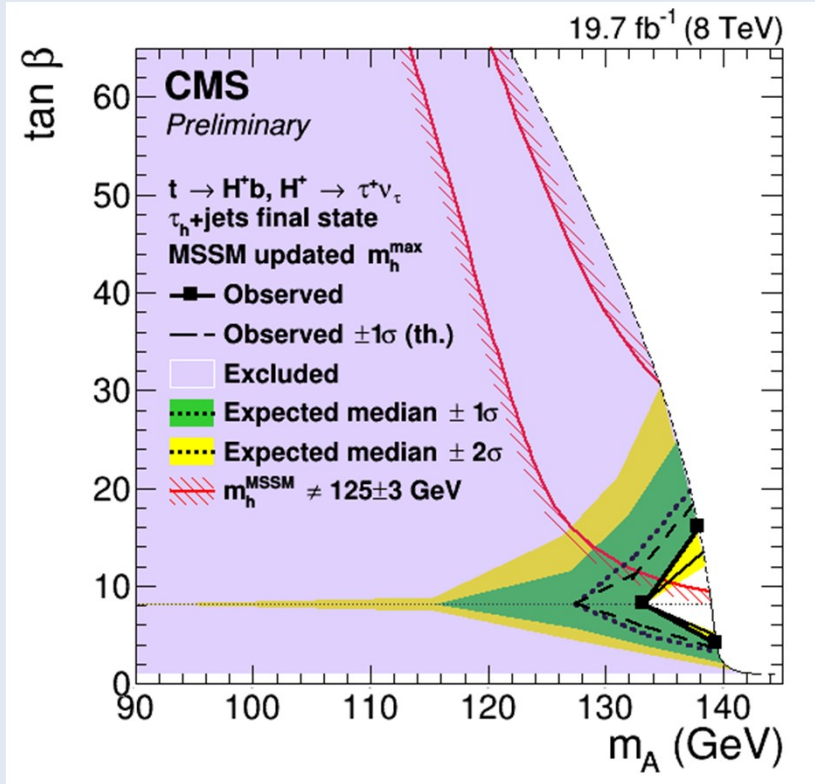
- In this scenario, the discovered Higgs boson is the heavy CP-even Higgs boson
- The area between the red lines is the  $(m_{H^+}, \tan \beta)$  area, where  $m_h = 125.0 \pm 3.0$  GeV
- $m_{H^+}$  excluded except for a tiny area between  $m_{H^+} = 150$ - $155$  GeV at  $\tan \beta \sim 10$

- Presented the search for charged Higgs bosons with  $H^+ \rightarrow \tau^+ \nu_\tau$  in fully hadronic final state and  $20 \text{ fb}^{-1}$  of  $\sqrt{s}=8 \text{ TeV}$  data
  - Majority of backgrounds are measured in data-driven way
  - Enhanced selection on angular cuts w.r.t. previous public results
  - First CMS public result on this final state for  $m_{H^+} > m_t - m_b$
  - Falling part of  $m_\tau$  distribution fitted to account better for low number of events
  - Set new model-independent limits:
    - 95 % CL upper limit on  $B(t \rightarrow bH^+) * B(H^+ \rightarrow \tau^+ \nu_\tau)$ : 1.2 - 0.16 % for  $m_{H^+} = 80..160 \text{ GeV}$
    - 95 % CL upper limit on  $\sigma(pp \rightarrow t b H^+) * B(H^+ \rightarrow \tau^+ \nu_\tau)$ : 0.38 - 0.026 pb for  $m_{H^+} = 180..600 \text{ GeV}$
  - Limits also calculated in the MSSM benchmark scenarios
    - In all scenarios, light  $H^+$  almost excluded ( $m_{H^+} > 150 \text{ GeV}$  or  $m_{H^+} > 155 \text{ GeV}$ )
    - Low  $m_H$  scenario almost completely excluded
- Documentation for the preliminary results is available in CMS PAS HIG-14-020

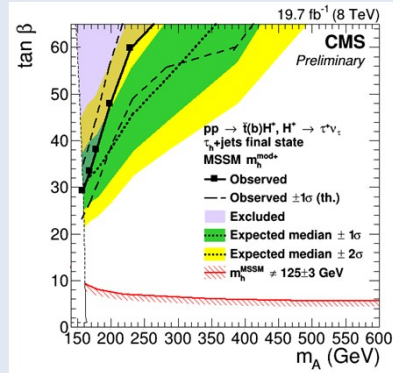
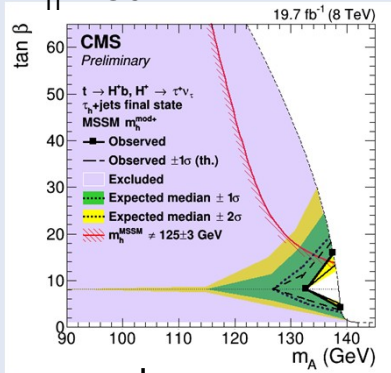
Additional slides



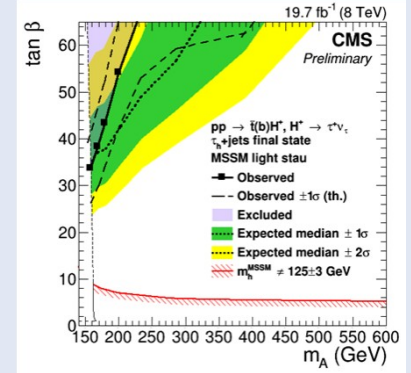
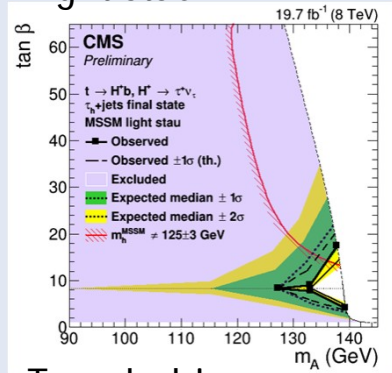




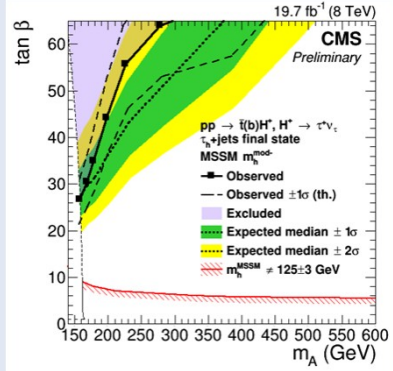
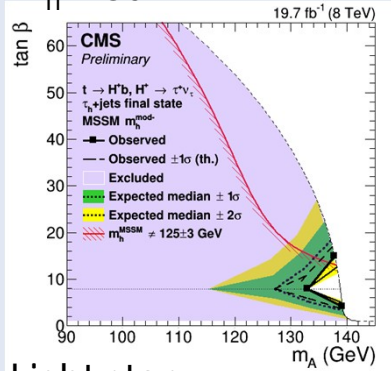
## $m_h$ mod+



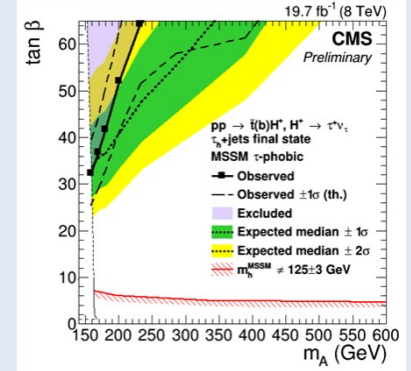
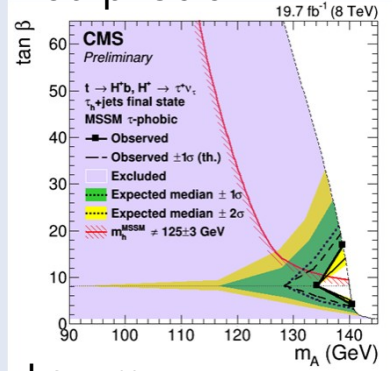
## Light stau



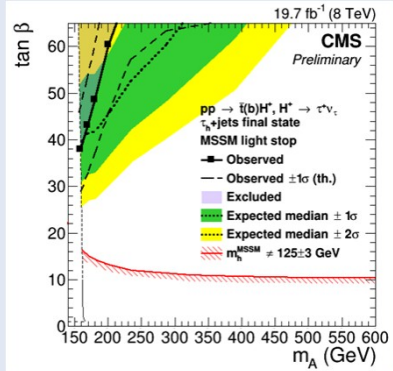
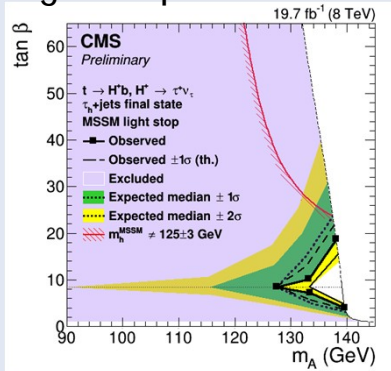
## $m_h$ mod-



## Tau-phobic



## Light stop



## Low $m_H$

