

Maximum Mass for Particle Dark Matter

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August 14, 2014

Dark Matter Candidates

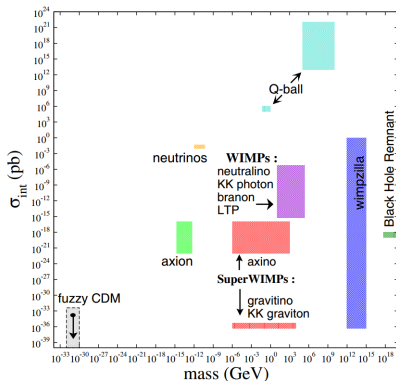


Figure 1 : Some dark matter candidates, taken from a July 2007 report from the Dark Matter Scientific Assessment Group

Typical Mass Scales

- Supersymmetry
 - Particle zoo: neutralinos, sneutrinos, axinos, etc ...
 - Dark matter: in phenomenological SUSY, $\sim 1 - 10$ TeV
- Little Higgs models
 - Motivation: stabilization of electroweak physics up to energies $\sim 1 - 10$ TeV
 - Dark matter: $\lesssim 1 - 10$ TeV
- Light scalar dark matter
 - Motivation: experimental hints, e.g., the 511 keV γ -ray signal; theoretical position somewhat *ad hoc*
 - Dark matter: mass scales $\sim 1 - 100$ MeV
- Kaluza-Klein and extra-dimensional theories
 - Motivation: unification, the hierarchy problem(s)
 - Dark Matter: mass scales ~ 100 GeV – 10 TeV

Typical Mass Scales

General Features - Expt.

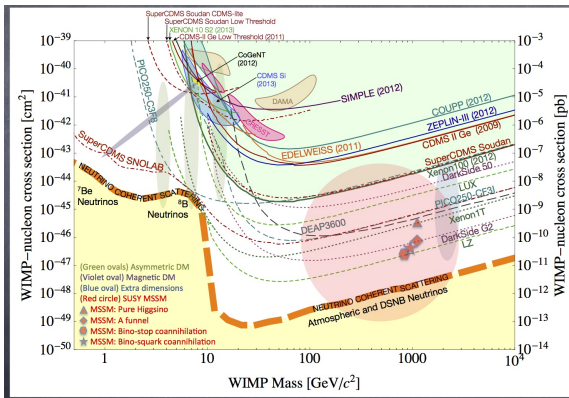


Figure 2 : Typical experimental mass ranges - CF1 Snowmass report

Unitarity (1/3)

- *Unitarity arguments* give the strongest model-independent theoretical bounds on thermal dark matter.
- Unitarity \iff QM “probability ≤ 1 ”
Unitarity \iff Scattering probability bounded from above
- Recall from non-relativistic scattering theory the partial wave expansion of the scattering amplitude $f(\theta)$:

$$f(\theta) = \sum_l^{\infty} (2l+1) \left(\frac{e^{2i\delta_l} - 1}{2ik} \right) P_l(\cos \theta)$$

- “Morally speaking” the Jacob-Wick partial wave expansion plays an analogous role for the helicity amplitude \mathcal{M} :

$$\mathcal{M}(\lambda_3 \lambda_4; \lambda_1 \lambda_2) = \frac{8\pi\sqrt{s}}{(p_i p_f)^{1/2}} e^{i(\lambda_i - \lambda_f)\phi} \sum_J (2J+1) \mathcal{M}_{\lambda}^J(s) d_{\lambda_i \lambda_f}^J(\theta)$$

- Partial wave unitarity $\Rightarrow |\mathcal{M}^J|^2 \leq |\text{Im} \mathcal{M}^J| \leq 1$

Unitarity (2/3)

- Fact from cosmology:

$$\Omega_X h^2 \sim \frac{3 \times 10^{-27} \text{cm}^3/\text{sec}}{\langle \sigma v_{\text{rel}} \rangle_f}$$

- Partial wave unitarity \implies *Maximal* cross section $\langle \sigma v_{\text{rel}} \rangle_f$
- Maximal cross section $\implies \sigma_J \leq \pi(2J + 1)/p_i^2$
- Observational data constrains Ωh^2
- Early universe kinematics: $p_i^2 = E^2 - m_X^2 \approx m_X^2 v_{\text{rel}}^2/4$

Unitarity (3/3)

- $(\sigma_J)_{max} v_{rel} \approx \frac{3 \times 10^{-22} (2J+1) \text{cm}^3/\text{sec}}{m_X / (1\text{TeV})}$
- Detailed analysis: $J = 0$ term dominates + $\mathcal{O}(v_{rel}^2/4)$
(Recall that $v_{rel} \approx 1/4 \ll 1$ at freezeout)

Solve for m_X and insert numbers:

$$m_X \lesssim 30 \text{ TeV}$$

- Details: K. Griest and M. Kamionkowski, Phys. Rev. Lett. 64 (1990) 615.
- Note that 30 TeV bound includes more recent PLANCK constraints on $\Omega_X h^2$

Evading Unitarity

- Non-thermal relics - evade unitarity bounds most easily
 - Axions
 - Superheavy Dark matter - WIMPzillas etc.
 - Q-Balls
- The precise mass bound from unitarity depends a bit on particle nature ...
- Some subtleties about the applicability of partial wave expansion - issues of non-convergence - very unlikely.

Nota bene:

Unitarity itself still should hold in all cases; the previous argument simply does not always apply.

Axions

- Well motivated theoretically - QCD/Stringy origins
- Non-thermal relics - mass constrained to be very small

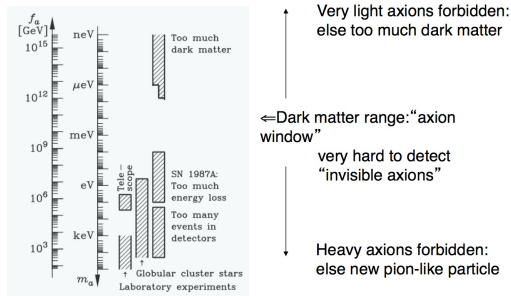


Figure 3 : Axion Constraints

Super Heavy Dark Matter

- Gravitationally produced at the end of inflation in cosmologically interesting amounts
- Stability guaranteed by Global/Discrete/(Gauged)-R symmetry
- Observation of Ultra high energy cosmic rays above the GZK cutoff is another motivation to look at SHDMs.
- Above this $\sim 5 \times 10^{19}$ GeV cutoff, protons interact at resonance with CMB photons with large cross sections, making the universe opaque to ultra high energy protons over cosmological distances.
- Exotic SHDM candidates offer plausible explanation

WIMPzillas

- Non-thermal in nature - relic abundance set by production cross section
- Produced naturally at the end of inflation by the expansion of the background spacetime acting on vacuum quantum fluctuations of the dark matter field.
- Natural mass scale for these particles at inflaton-GUT scale $\sim 10^{10} - 10^{16}$ GeV
- Stability guaranteed by discrete gauge symmetry - weakly broken
- D-matter (non-perturbative particle like states from compactified D-branes) provides a somewhat larger mass
- Thermal production with $m_X \geq 10^{13}$ GeV requires a very high T_{reh}

Experimental Constraints on SHDM

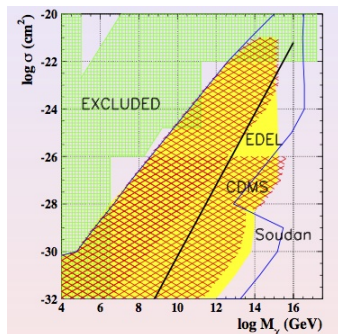


Figure 4 : Direct detection constraints

Experimental Constraints on SHDM

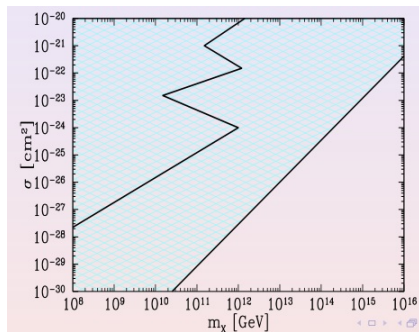


Figure 5 : Indirect detection constraints

Q-Balls

- Non-topological solitonic configurations - stability guaranteed by a global $U(1)$ Noether charge.
- Distinct from a topological soliton which is stable because there is no continuous transformation that homotopically maps it to the vacuum.
- Arise generically in supersymmetric field theories.
- Can be produced in the early universe via phase transitions, fusion or pair production at high temperatures.
- Can decay via massless fermions that carry the same global charge.
- Mass $\sim 10^{10}$ GeV

Primordial Black Holes (PBHs)

- PBHs are predicted to form in the early Universe from large density fluctuations.
- There is a small window for PBHs as CDM.
- The upper window of PBH masses allowed can be probed with Kepler and WFIRST
- Ideally, the entire range for CDM can be probed with gravitational wave detectors. ($f_{GW} \sim 10^{-5} - 10^{-2}\text{Hz}$)

Expt. Constraints on Primordial Black Holes

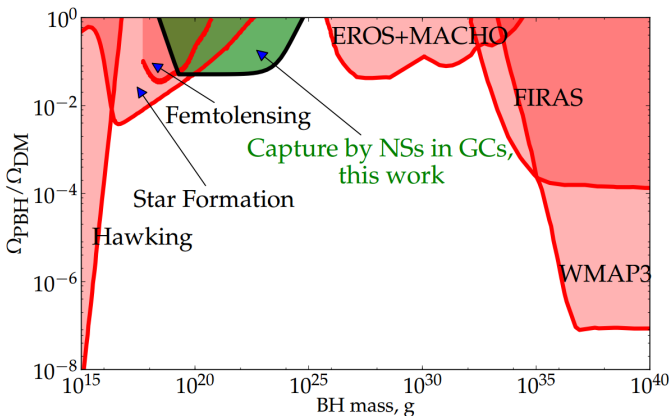


Figure 6 : Limits on Ω_{DM} from Capela et al., Phys. Rev. D 87, 123524 (2013)

Summary

- Unitarity restricts thermal dark matter - and thus a large class of explicit theories - to masses less than ~ 30 TeV.
- Non-thermal dark matter evades the unitarity argument and is often much heavier.
- Some well-motivated models, like axions, often involve significantly lighter masses.

Question: Is there a maximum mass for a dark matter particle?

Answer: Yes, with caveats.

Thank You