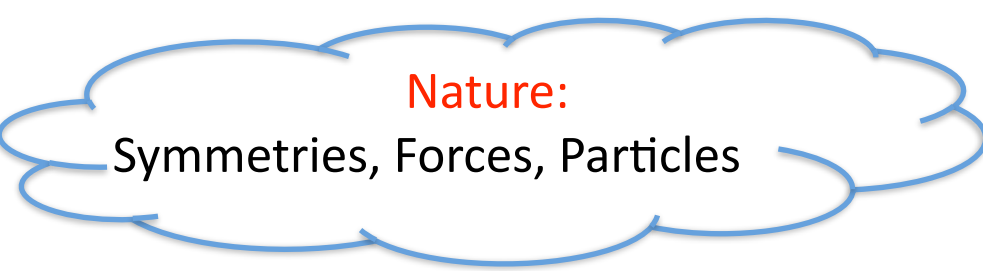


# Matrix Element Method in Higgs Phenomenology

Michael Spannowsky

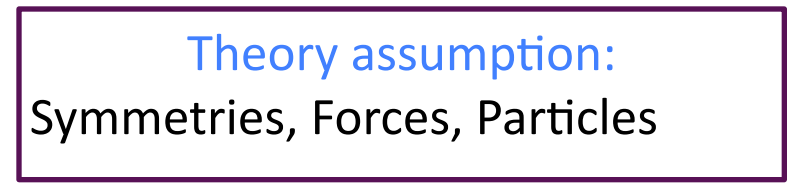
University of Durham



Result in measurable objects, e.g.  
Jets, stable leptons, photons



Experiments measure radiation

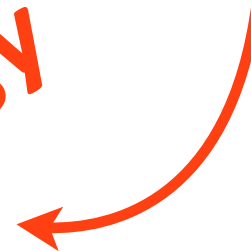
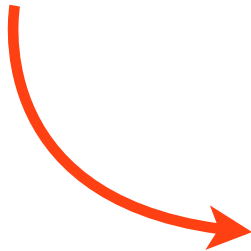


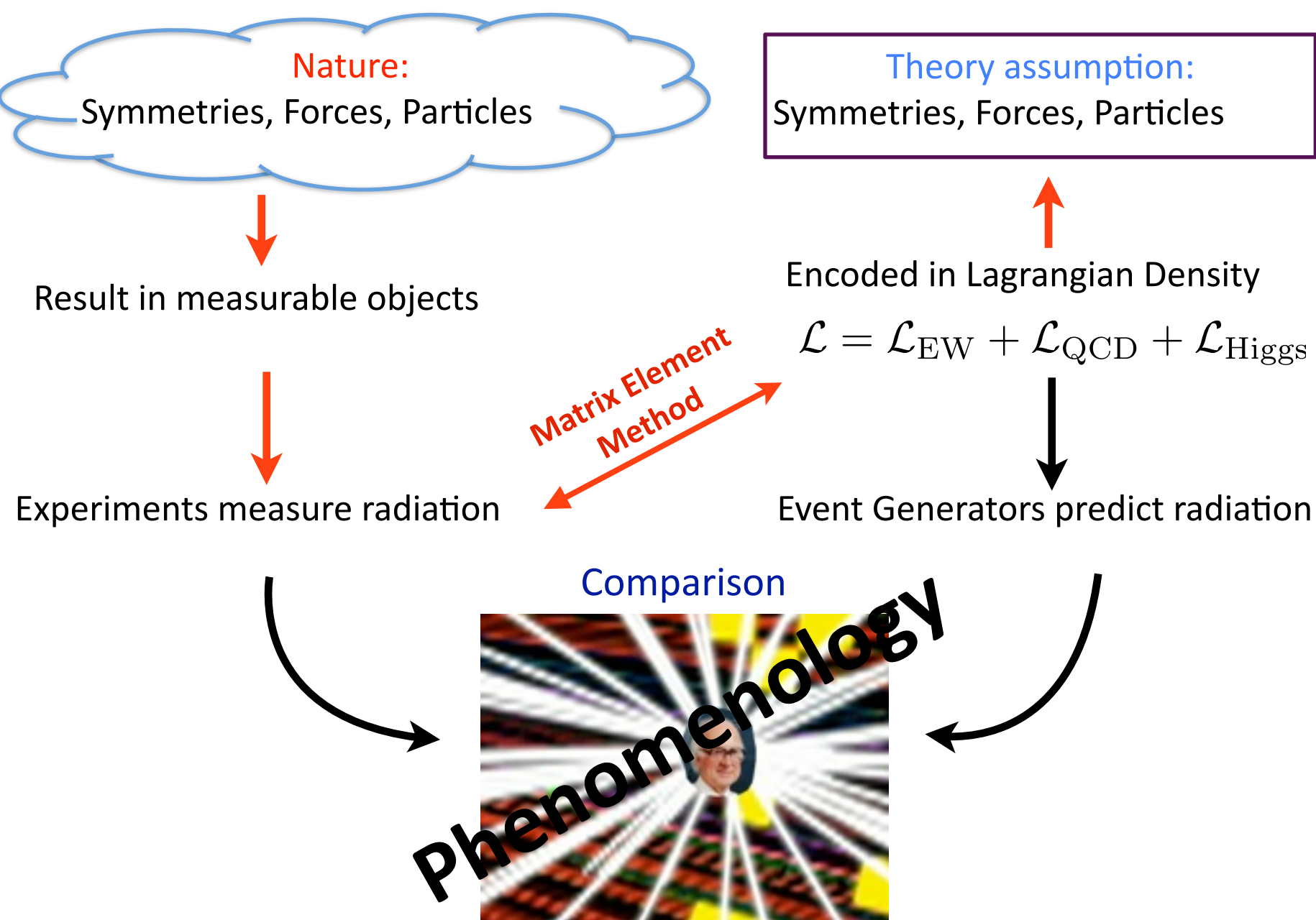
Encoded in Lagrangian Density  
 $\mathcal{L} = \mathcal{L}_{EW} + \mathcal{L}_{QCD} + \mathcal{L}_{Higgs}$



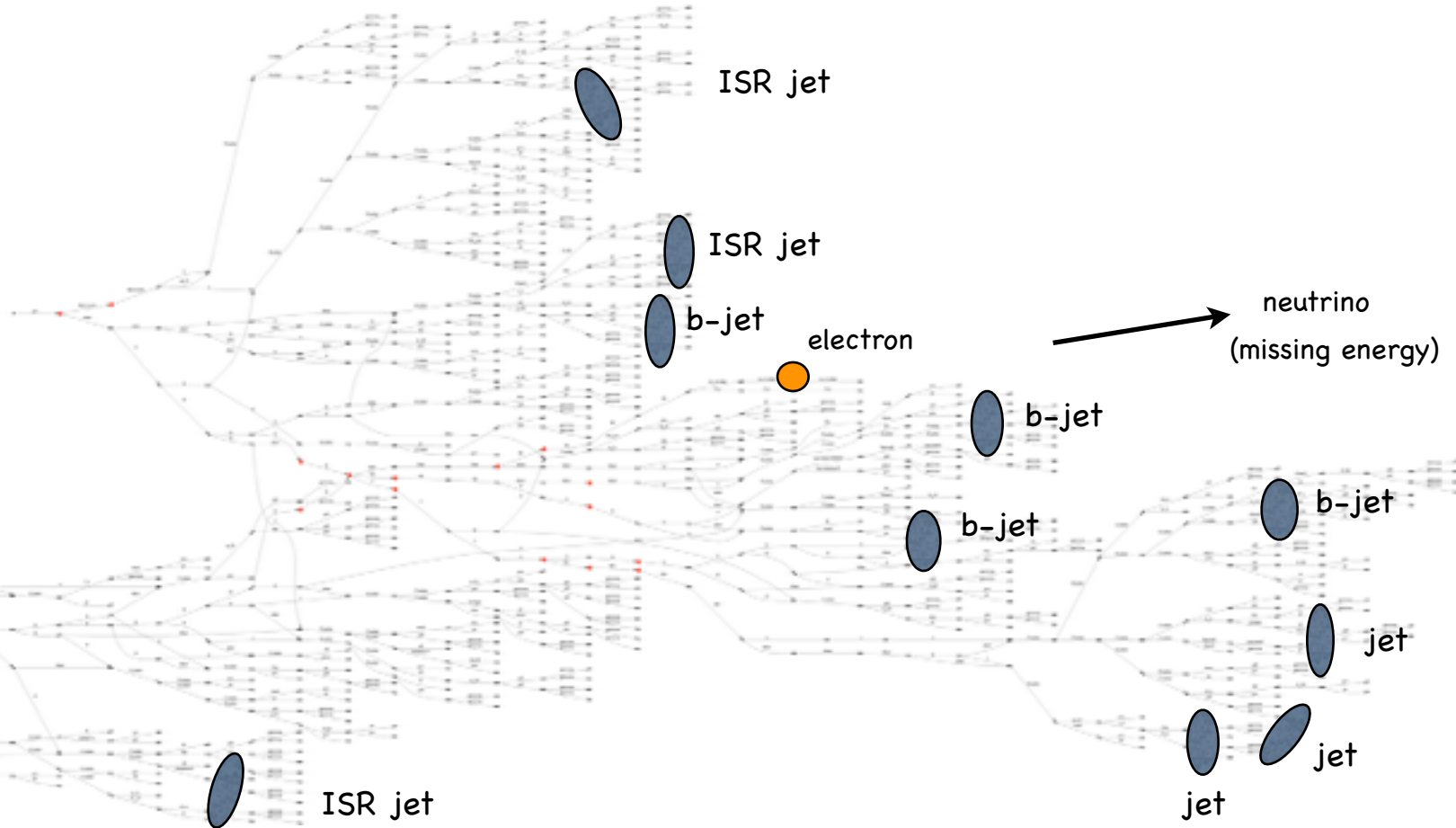
Event Generators predict radiation

Comparison





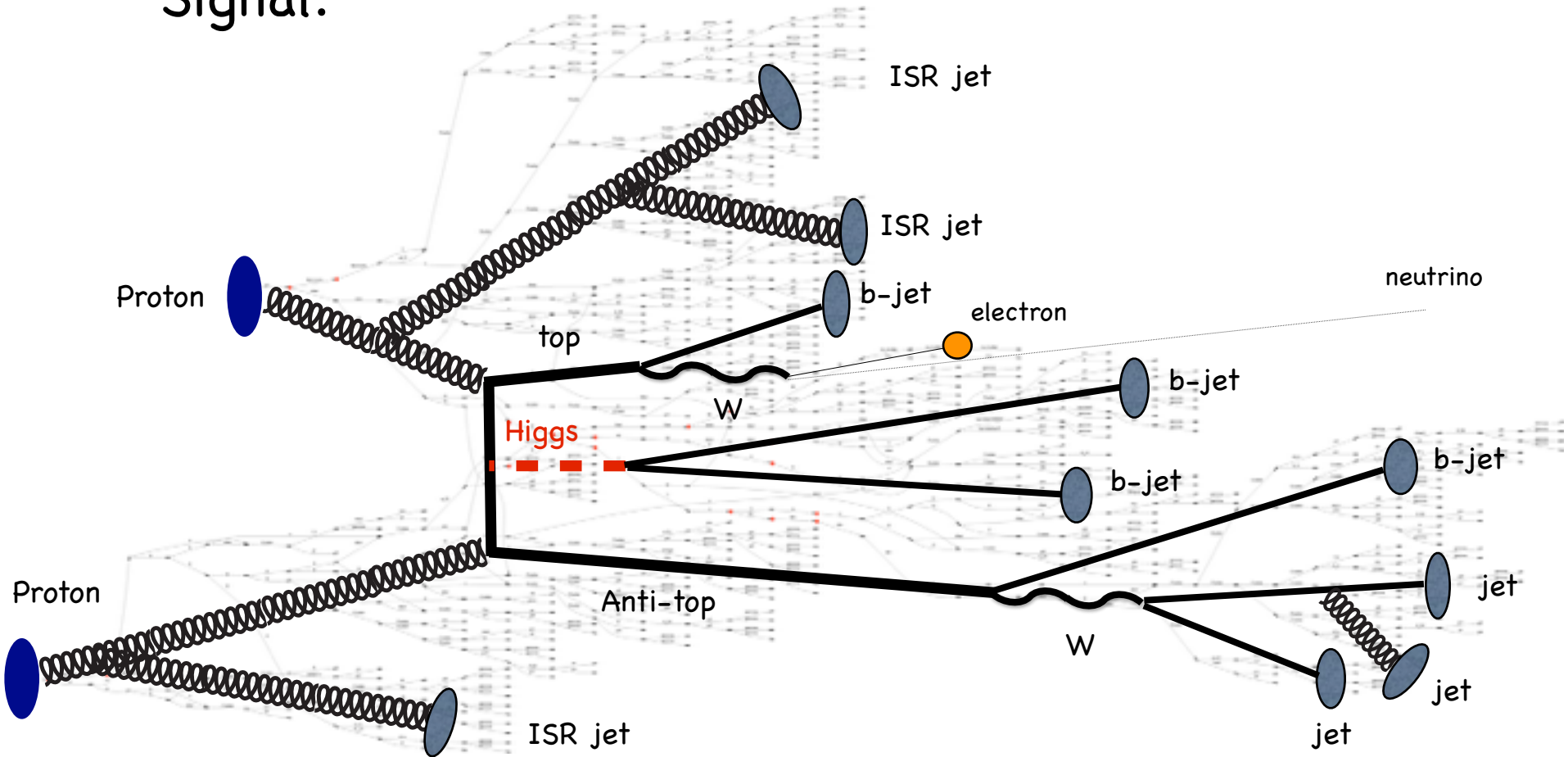
# Inverse Problem: Final state measured (‘phase space point chosen’)





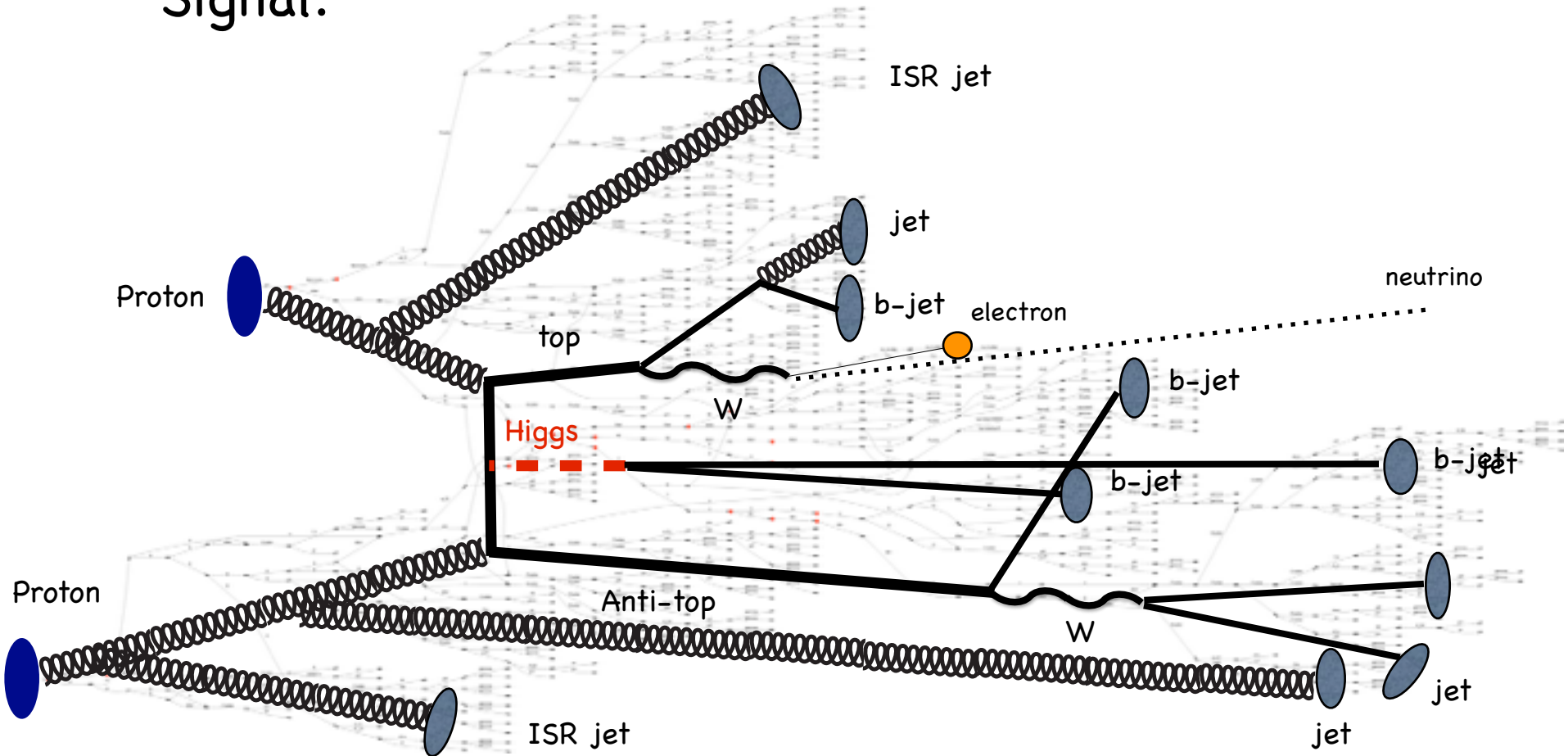
Give final state radiation distinctive meaning in terms of hypothesis

Signal:



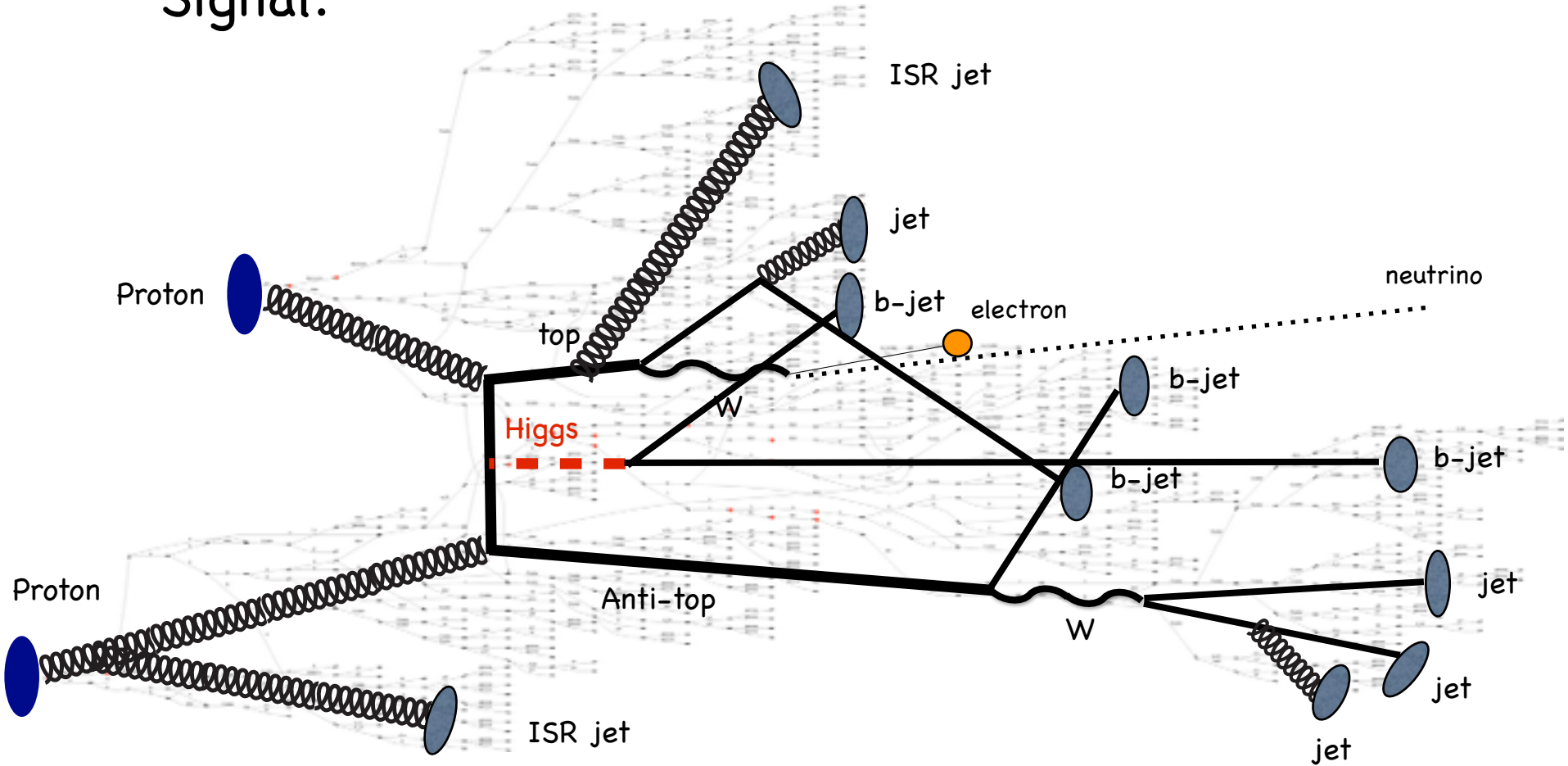
Give final state radiation distinctive meaning in terms of hypothesis

Signal:



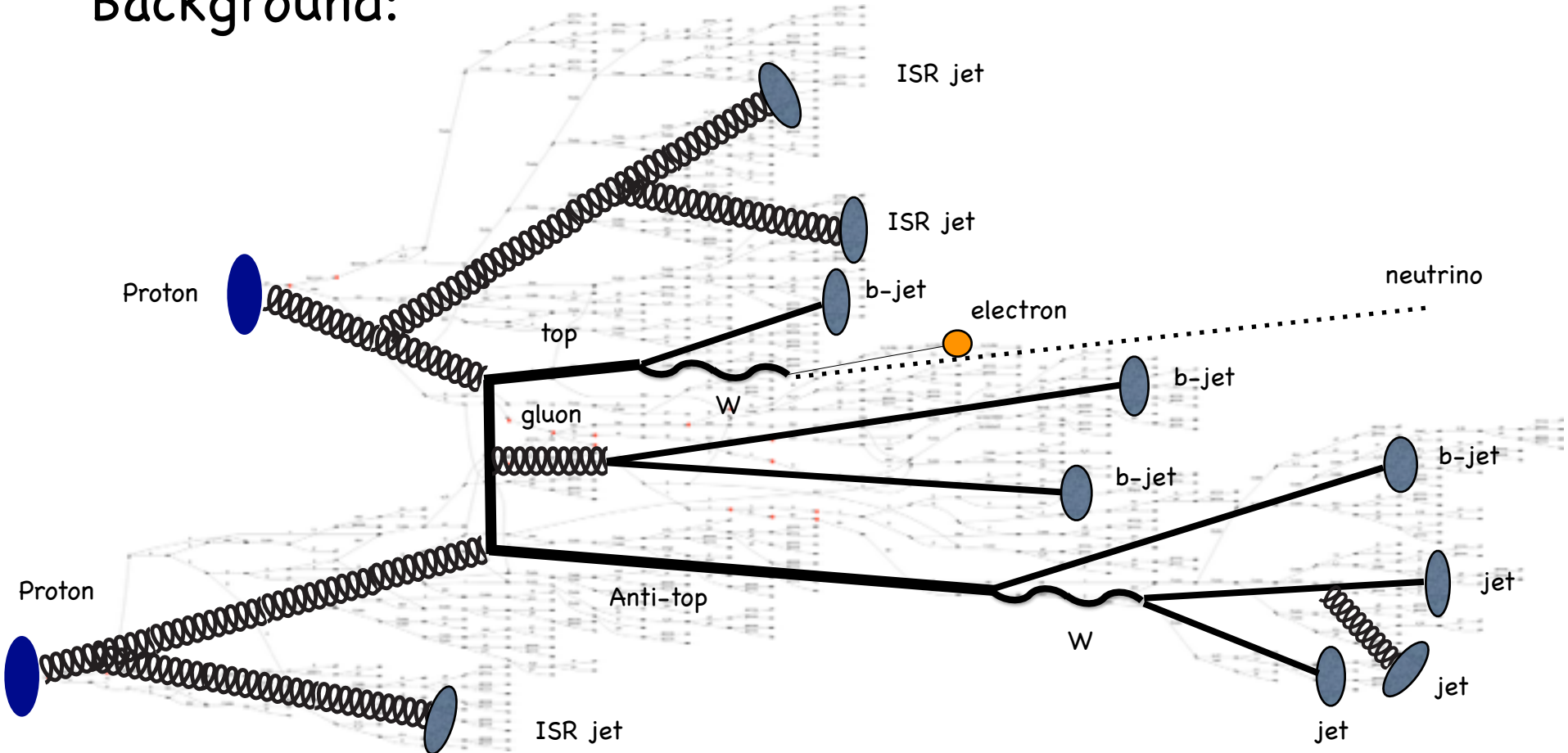
# Give final state radiation distinctive meaning in terms of hypothesis

Signal:

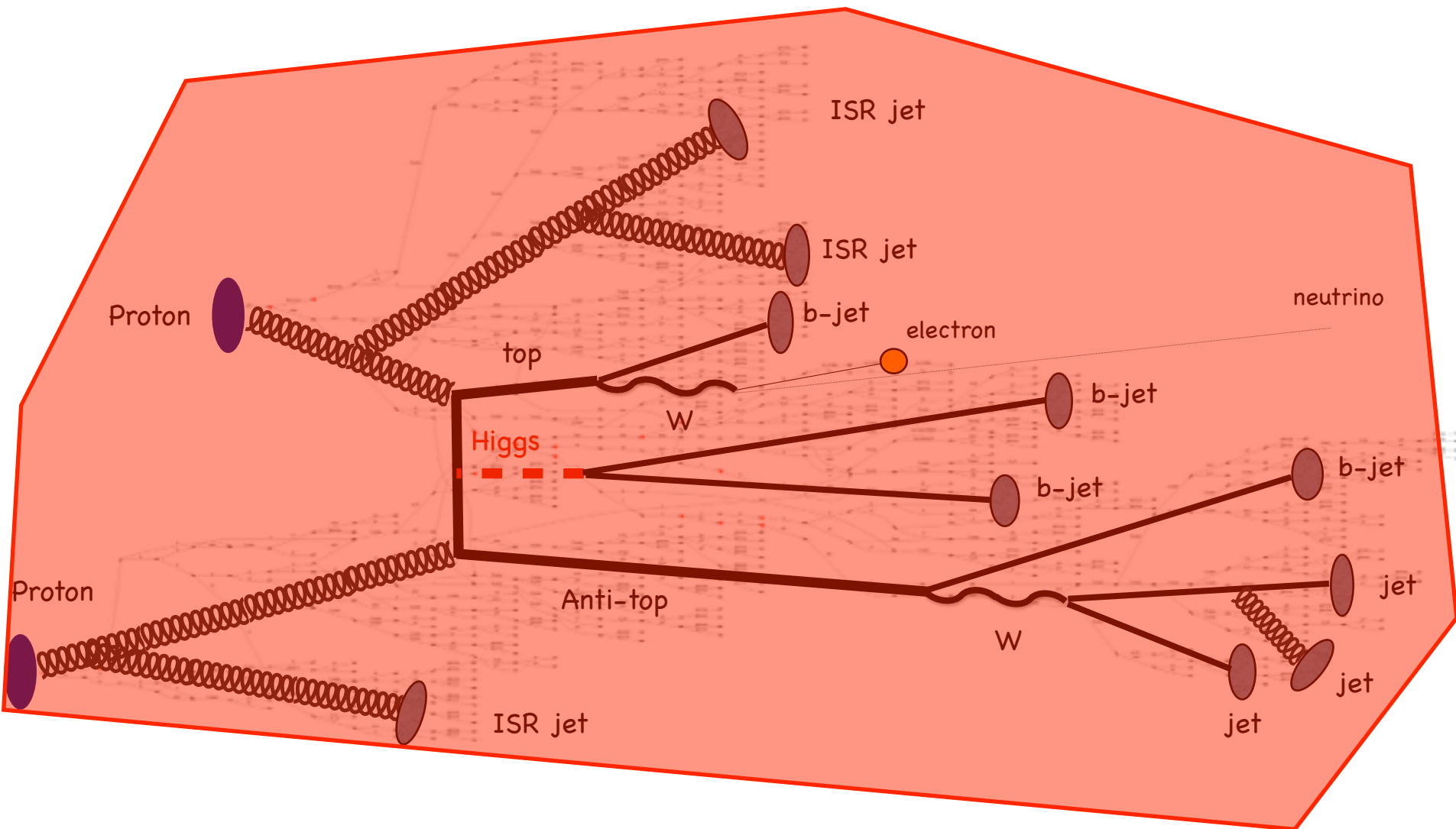


Give final state radiation distinctive meaning in terms of hypothesis

Background:



Ideally one would like to use all radiation related to hard process to discriminate signal from background



# Applications of Matrix Element Method:

- 1988 Rec. of events with MET [Kondo, J.Phys.Soc.Jap. (1988)]
- 1998 Anomalous gauge couplings [Diehl, Nachtmann Eur. Phys. J. C1 (1998)]
- 2005 top quark physics [Abazov et al., Nature (2004), D0 Collab.]  
[Abulencia et al., PRD 73 (2005), CDF Collab.]  
[Abazov et al., PLB 617 (2005), D0 Collab.]
- 2010 Automated implementation in MadWeight  
[Artoisenet et al, JHEP 1012 (2010)]

Plenty of recent applications in Higgs physics:

- $H \rightarrow \mu^+ \mu^-$  [Cranmer, Plehn EPJC 51 (2007)]
- $H \rightarrow b\bar{b}$  [Soper, MS PRD 84 (2011)]
- $H \rightarrow \gamma\gamma$  [Andersen, Englert, MS PRD 84 (2013)]
- $pp \rightarrow t\bar{t}H$  [Artoisenet et al. PRL 111 (2013)]
- $H \rightarrow ZZ^*/WW^*/Z\gamma$  [Campbell et al JHEP 1211 (2012)]  
[Freitas et al PRD 88 (2013)] [Campbell et al PRD 87 (2013)]
- Spin/Parity [Avery, et al. PRD 87 (2013)] [Gao et al. PRD 81 (2010)]

## The matrix element method in a nutshell:

Given a theoretical assumption  $\alpha$ , attach a weight  $P(\mathbf{x}, \alpha)$  to each experimental event  $\mathbf{x}$  quantifying the validity of the theoretical assumption  $\alpha$  for this event.

$$P(\mathbf{x}, \alpha) = \frac{1}{\sigma} \int d\phi(\mathbf{y}) |M_\alpha|^2(\mathbf{y}) W(\mathbf{x}, \mathbf{y})$$

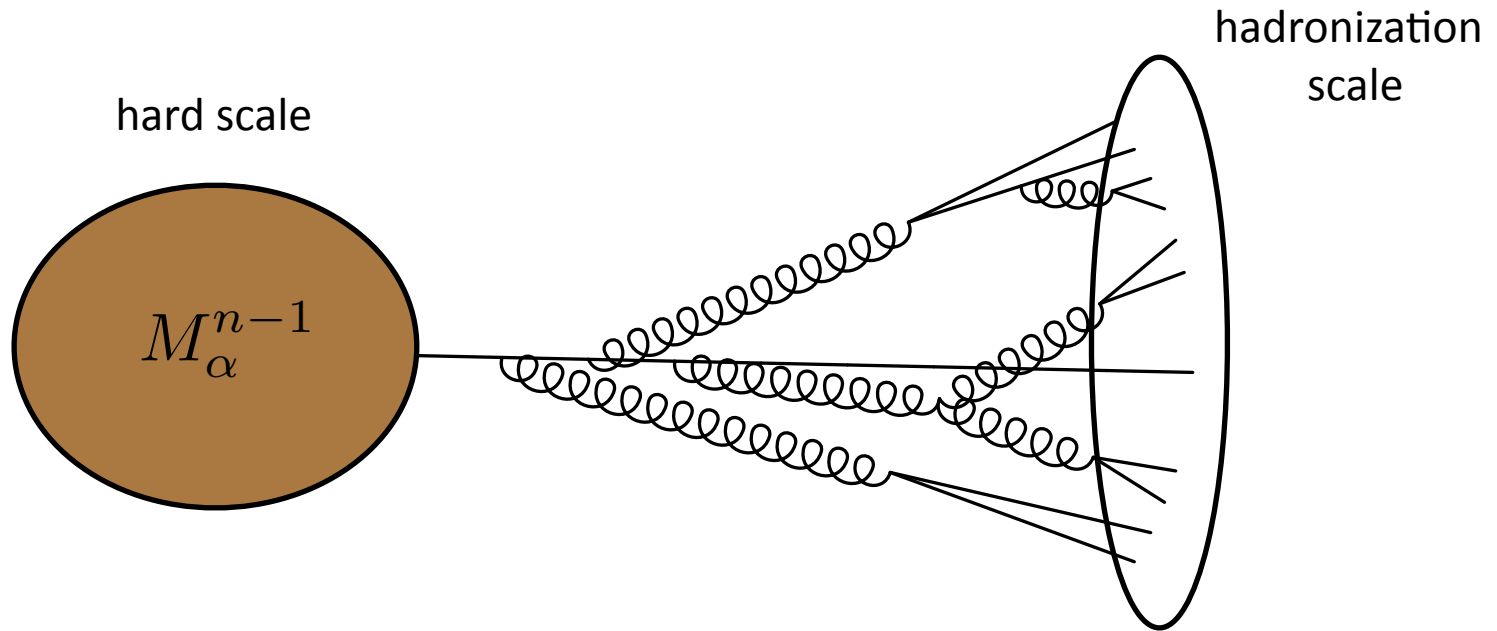
$|M_\alpha|^2$  is squared matrix element

$W(\mathbf{x}, \mathbf{y})$  is the resolution or transfer function

$d\phi(\mathbf{y})$  is the parton-level phase-space measure

The value of the weight  $P(\mathbf{x}, \alpha)$  is the probability to observe the experimental event  $\mathbf{x}$  in the theoretical frame  $\alpha$

Purpose of the transfer function is to match jets to partons



Probability density function:  $\int d\mathbf{y} W(\mathbf{x}, \mathbf{y}) = 1$



The form of the transfer function:

$$W(\mathbf{x}, \mathbf{y}) \approx \prod_i \frac{1}{\sqrt{2\pi}\sigma_{E,i}} e^{-\frac{(E_i^{rec} - E_i^{gen})^2}{2\sigma_{E,i}^2}}$$

resolution in  
Energy

$$\times \frac{1}{\sqrt{2\pi}\sigma_{\phi,i}} e^{-\frac{(\phi_i^{rec} - \phi_i^{gen})^2}{2\sigma_{\phi,i}^2}}$$

azimuthal angle

$$\times \frac{1}{\sqrt{2\pi}\sigma_{y,i}} e^{-\frac{(y_i^{rec} - y_i^{gen})^2}{2\sigma_{y,i}^2}}$$

rapidity

Complex, high-dimensional gaussian distribution!

Transfer function introduces new peaks on top of propagators

## Subtleties of the convolution $|M(y)|^2 \times W(y, x)$

### 1) $|M(y)|^2$

- Can be calculated at different order in pert. series (LO, NLO)
- Final state multiplicity fixed (exclusive process)
- Some kinematic configurations induce large logs (need resummation)

### 2) $W(y, x)$

- Number of final state objects limited to exclusive process
- Integration very time consuming  $\rightarrow$  limits final state multiplicity
- Transfer function fit dependent (input from experiment)

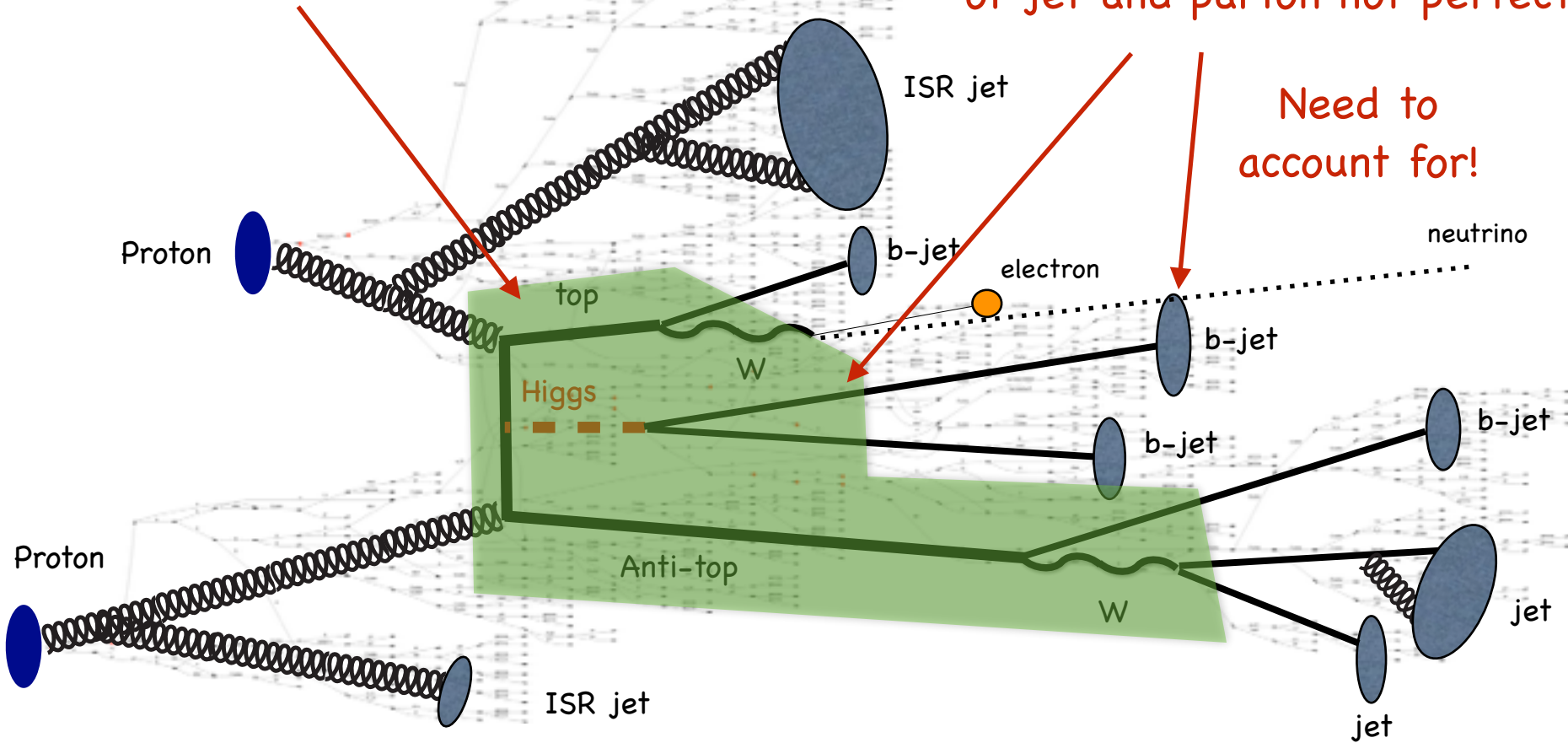
# tth: di-lepton vs semileptonic channel

[Artoisenet et al. PRL 111 (2013)]

Hard Interaction

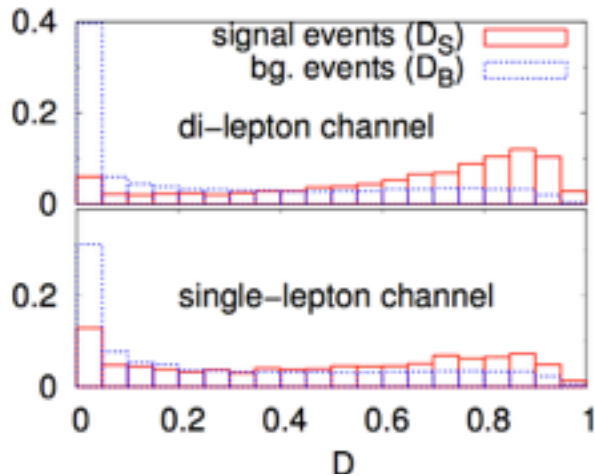
Matching between 4-momentum of jet and parton not perfect

Need to account for!



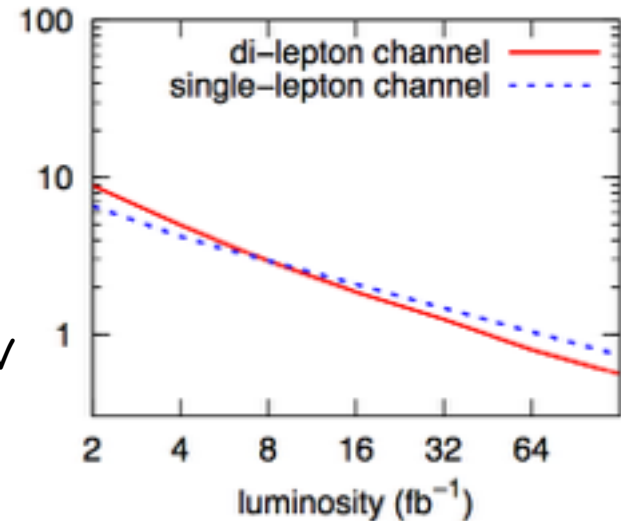
- Analysis with 4 b-jets and std reconstruction as input to MEM
- Full integration over invisible particles

process	incl. $\sigma$	efficiency	$\sigma^{\text{rec}}$
$t\bar{t}h$ , single-lepton	111 fb	0.0485	5.37 fb
$t\bar{t}h$ , di-lepton	17.7 fb	0.0359	0.634 fb
$t\bar{t}$ +jets, single-lepton	256 pb	$0.463 \times 10^{-3}$	119 fb
$t\bar{t}$ +jets, di-lepton	40.9 pb	$0.168 \times 10^{-3}$	6.89 fb



Projection at 14 TeV

$$D_i = \frac{P(x_i|S)}{P(x_i|S) + P(x_i|B)}$$



- Using Matrix Element Method di-lepton channel as sensitive than single-lepton channel

- However, single-lepton channel uses standard input, **boosted region not captured** [Plehn, Salam, MS PRL 104 (2009)]

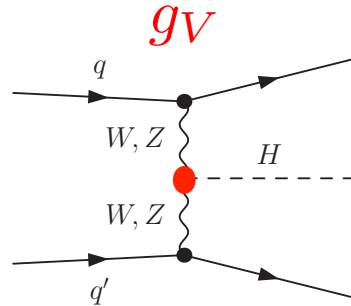
We want to study more objects in final state ->

Transfer function limits us. Do we always need it?

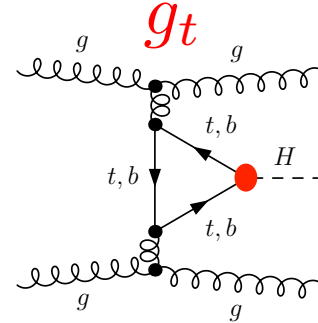
Transfer functions only important if matrix element varies quickly:

Example

[Andersen, Englert,  
MS PRD 84 (2013)]

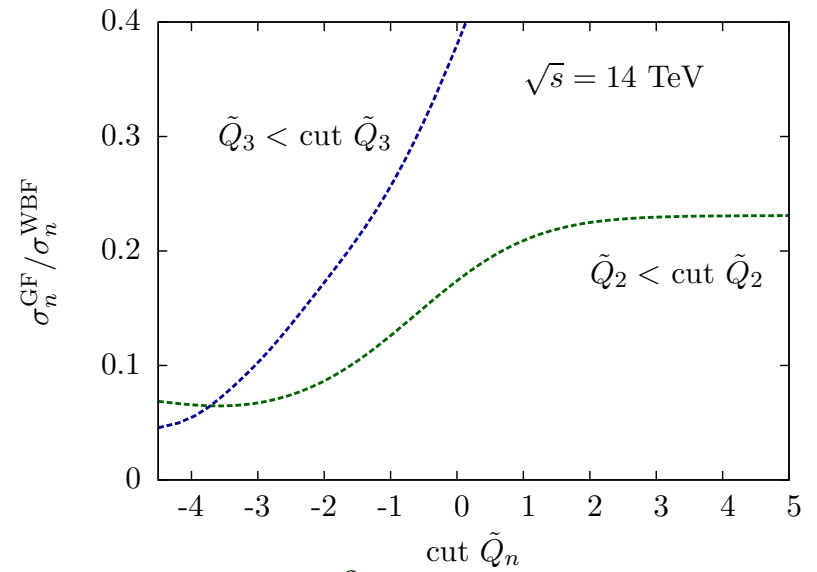
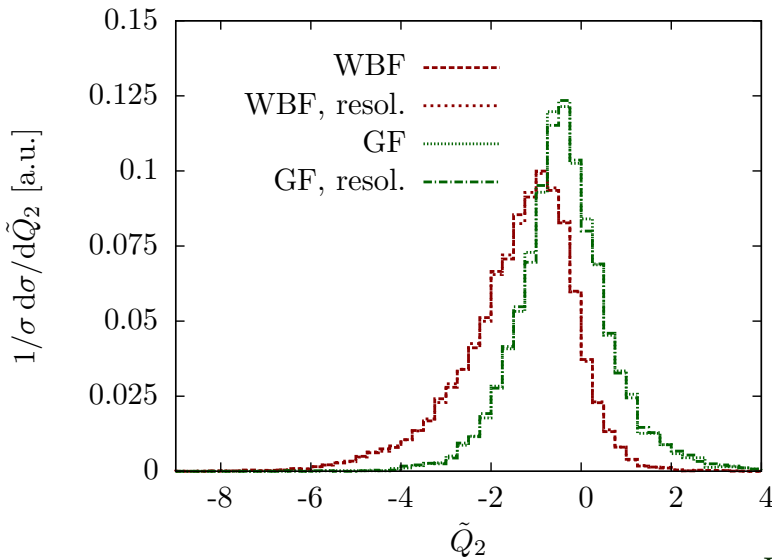


VS



in  $H \rightarrow \gamma\gamma$

Higgs reconstructed, but no transfer function for jets:

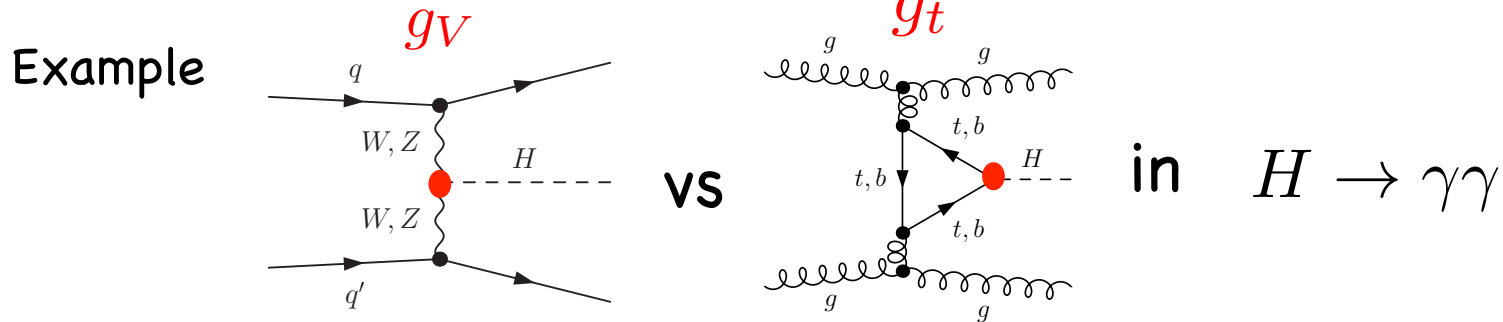


$$\tilde{Q}_n = -\log \left[ \frac{|\mathcal{M}^{\text{WBF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2}{|\mathcal{M}^{\text{GF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2} \right]$$

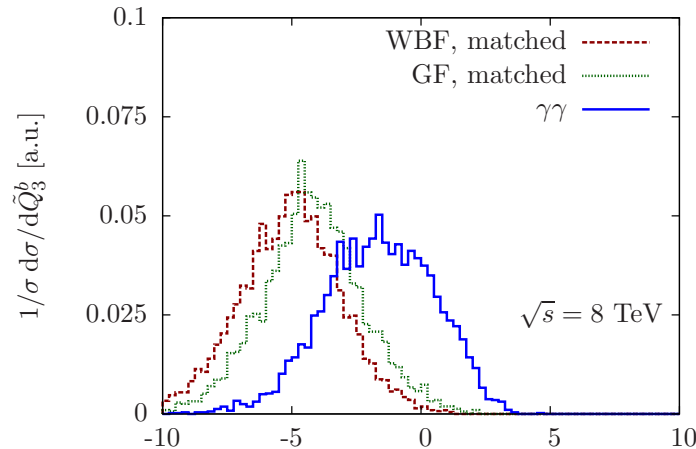
We want to study more objects in final state ->

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Transfer functions only important if matrix element varies quickly:



Higgs reconstructed, but no transfer function for jets:



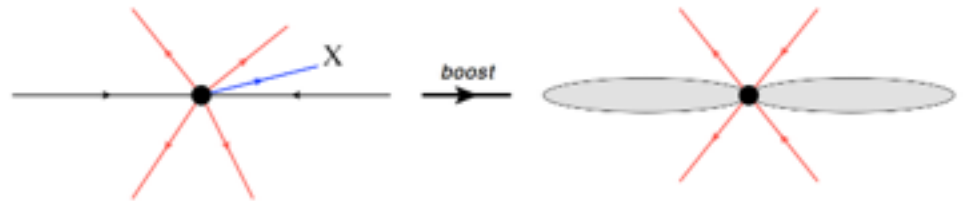
$S/B \nearrow 100\%$

$$\tilde{Q}_n^b(p_1^\gamma, p_2^\gamma, \{p_n^j\}) = -\log \left[ \frac{\{|\mathcal{M}^{\text{WBF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2 + |\mathcal{M}^{\text{GF}}(pp \rightarrow (h \rightarrow \gamma\gamma)j^n)|^2\}}{|\mathcal{M}^{2\gamma}(pp \rightarrow \gamma\gamma j^n)|^2} \right]$$

# After removing transfer function we can improve on precision of matrix element $|M(y)|^2$

Matrix element method at NLO: [Campbell, Giele, Williams JHEP 1211 (2012)]

Boost along transverse and longitudinal direction such that LO final state multiplicity momenta balance



↳ Born phase space, but long. boost not unique, need longitud. integration

$$\mathcal{P}_{NLO}^{MEM}(\{Q_n\}) = \frac{1}{\sigma_{NLO}} \int_{x_{min}}^{x_{max}} dx_1 \mathcal{P}_{NLO}(\Phi_B)$$

↳ Calculate virtual for born topology  
real for jet function

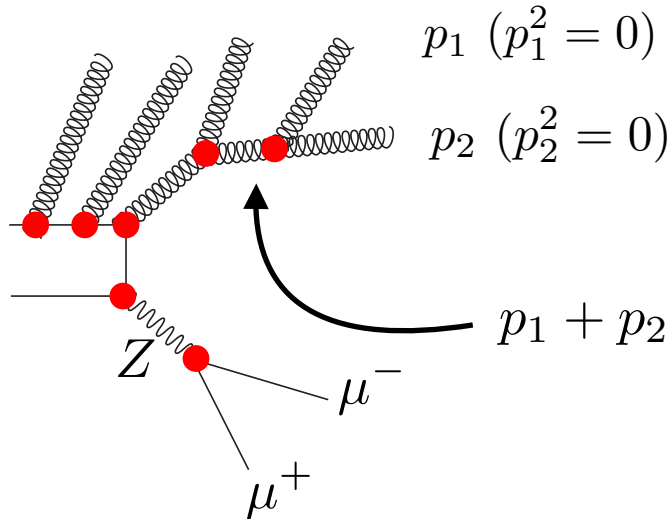
$$\eta^{lab,i} = \frac{1}{2} \log \left( \frac{x_a^2 s_{ib}}{s_{ab} s_{ai}} \right)$$

↳ Application to  $H \rightarrow 4l$   
(boost easier to identify)

sensitivity LO vs NLO improvement  $\sim 10\%$

# Remove limitation of final state objects on $|M(y)|^2$

Shower approximation for matrix element, i.e. **shower deconstruction**:



partons from the hard interaction emit other partons (gluons and quarks)

These emissions are enhanced if they are collinear and/or soft with respect to the emitting parton

Probability enhanced in soft and collinear region due to  $\sim 1/(p_1 + p_2)^2$

- If  $p_1 \rightarrow 0$ , then  $1/(p_1 + p_2)^2 \rightarrow \infty$
- If  $p_2 \rightarrow 0$ , then  $1/(p_1 + p_2)^2 \rightarrow \infty$
- If  $p_2 \rightarrow \lambda p_1$ , then  $1/(p_1 + p_2)^2 \rightarrow \infty$



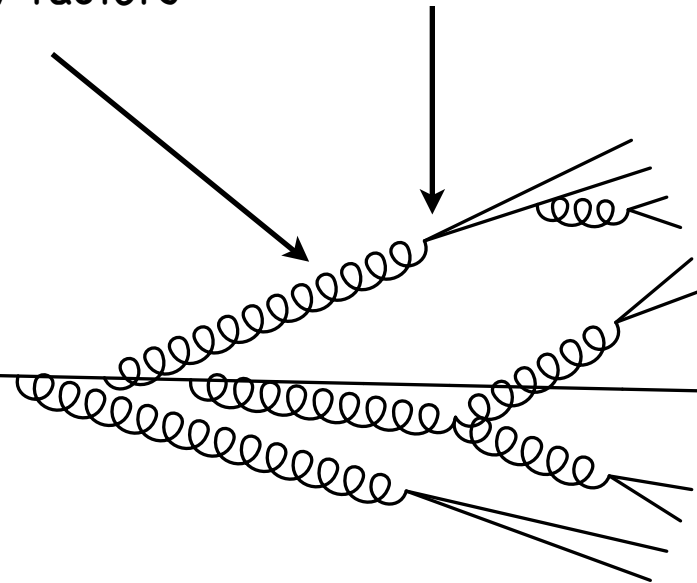
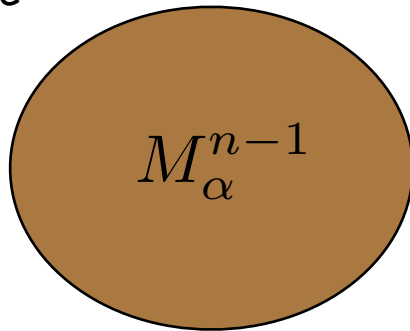
# Factorization of emissions in soft/collinear limit

and Sudakov factors allow semiclassical approximation of quantum process:

propagator-lines = Sudakov factors

vertices = Splitting functions

hard scale



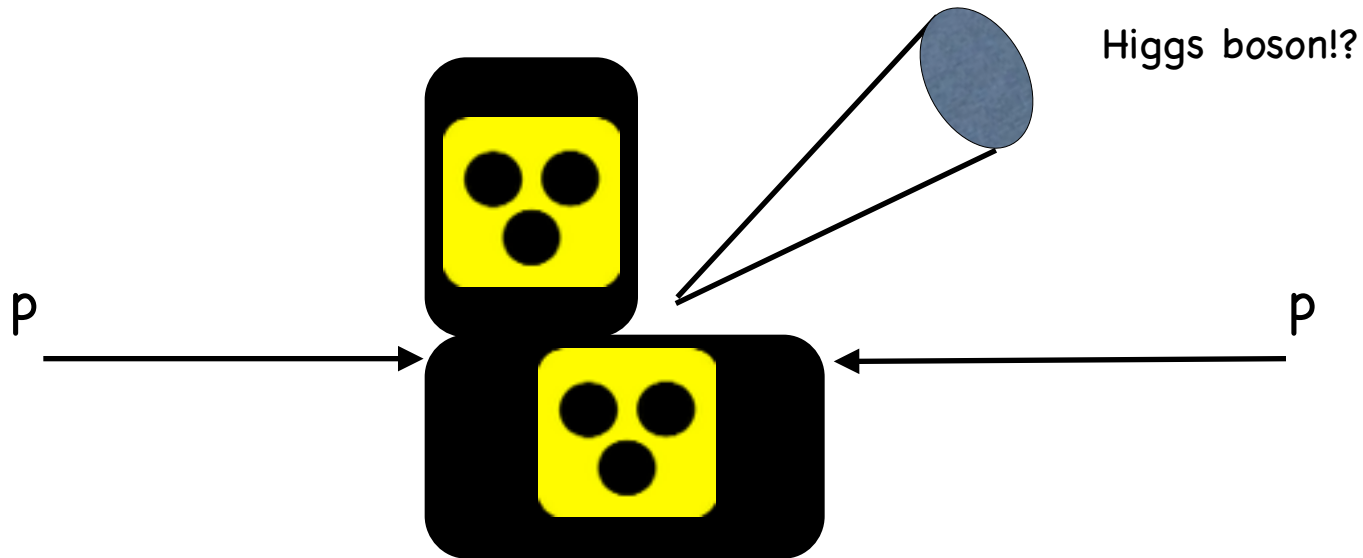
hadronization scale



Can calculate weight for shower history iteratively

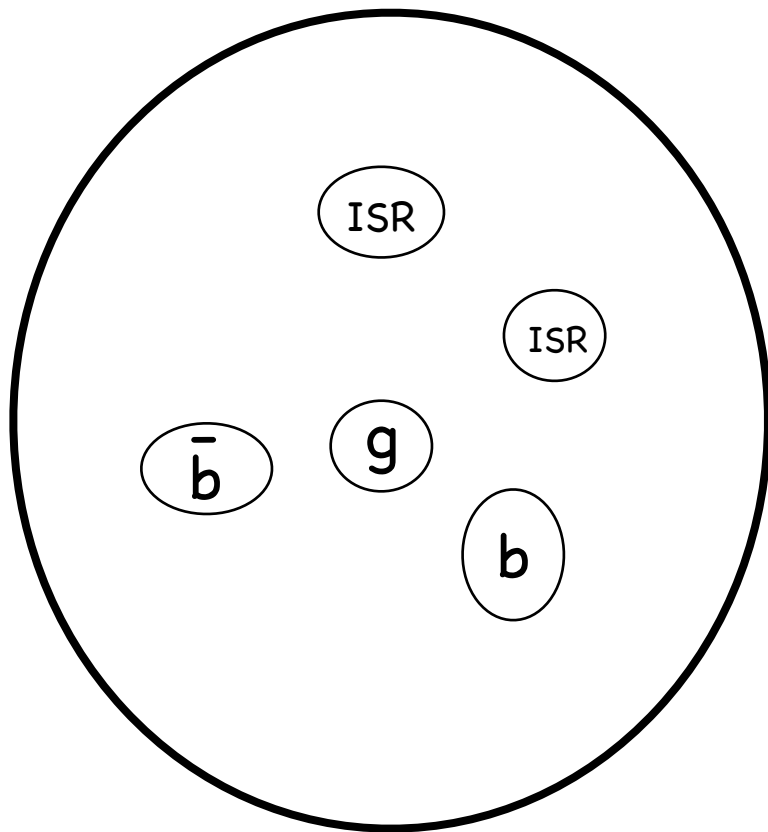
Can use smaller objects and more objects (more information)

This approach can be used as a tagger for  
Higgs bosons in  $H \rightarrow bb$



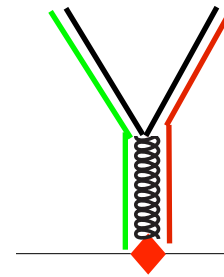
In boosted resonances radiation collimated,  
need Sudakov factors for valid description

# Fat jet: $R=1.2$ , anti- $k_T$

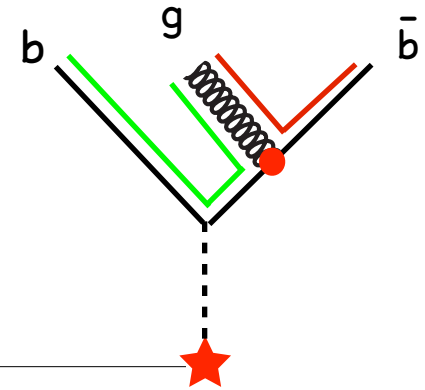


microjets

ISR/UE



hard interaction

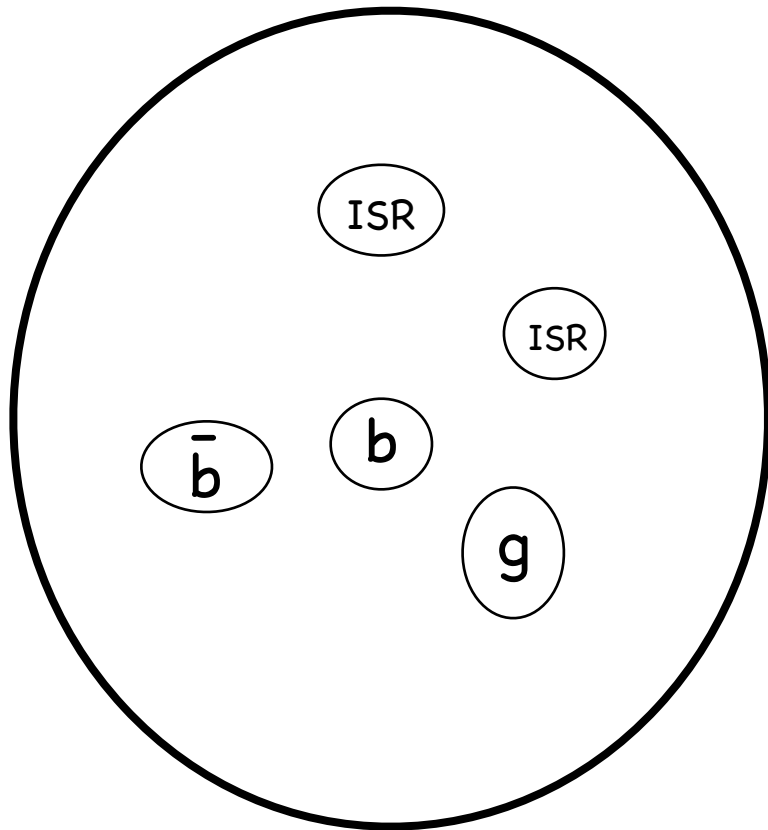


Build all possible shower histories

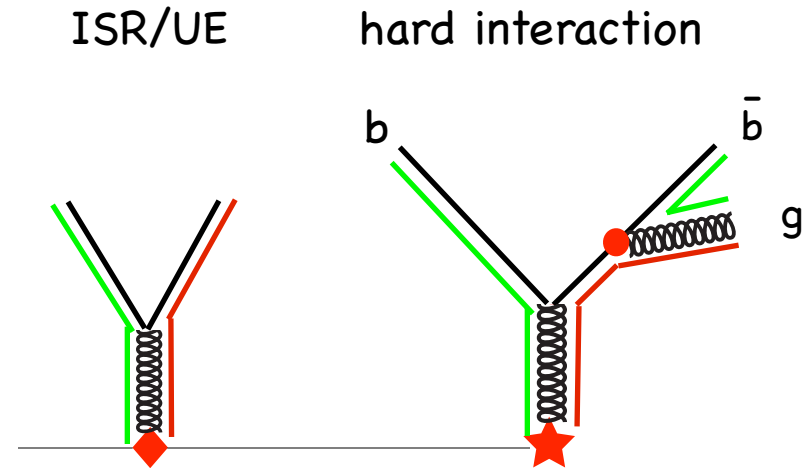
signal vs background hypothesis based on:

- ▶ Emission probabilities
- ▶ Color connection
- ▶ Kinematic requirements
- ▶  $b$ -tag information

Fat jet:  $R=1.2$ , anti- $k_T$



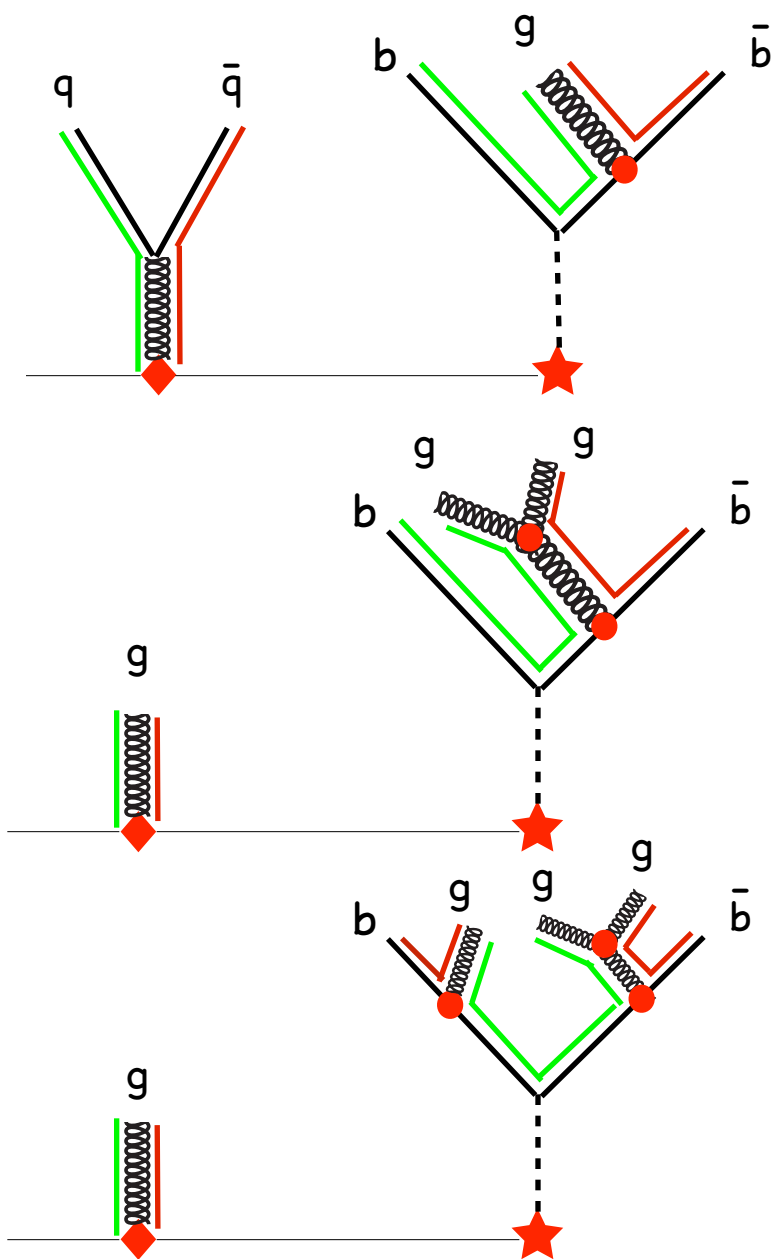
microjets



Build all possible shower histories

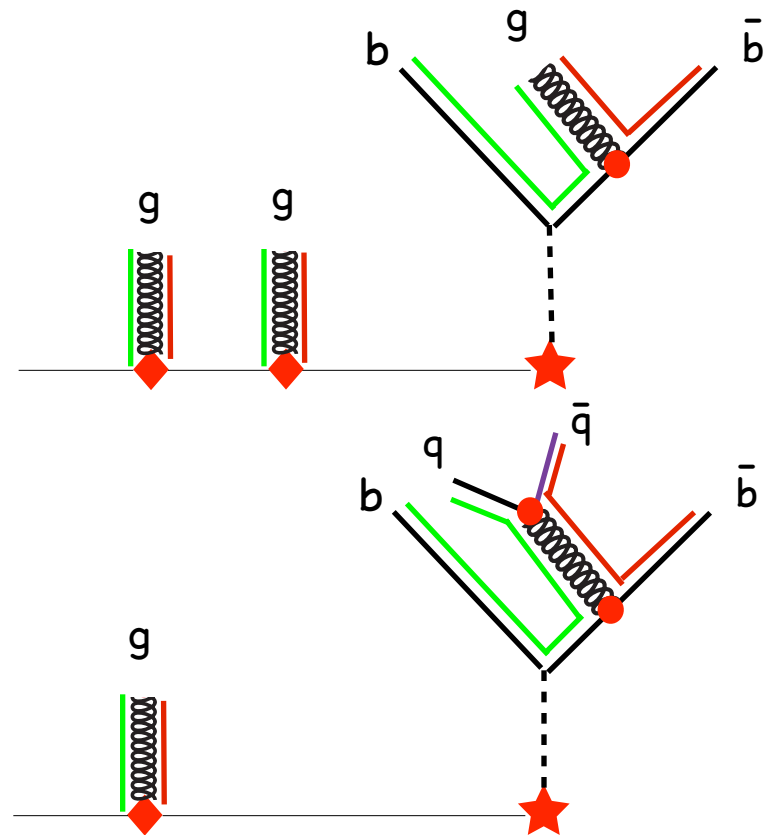
signal vs background hypothesis based on:

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Higgs couplings

Turin

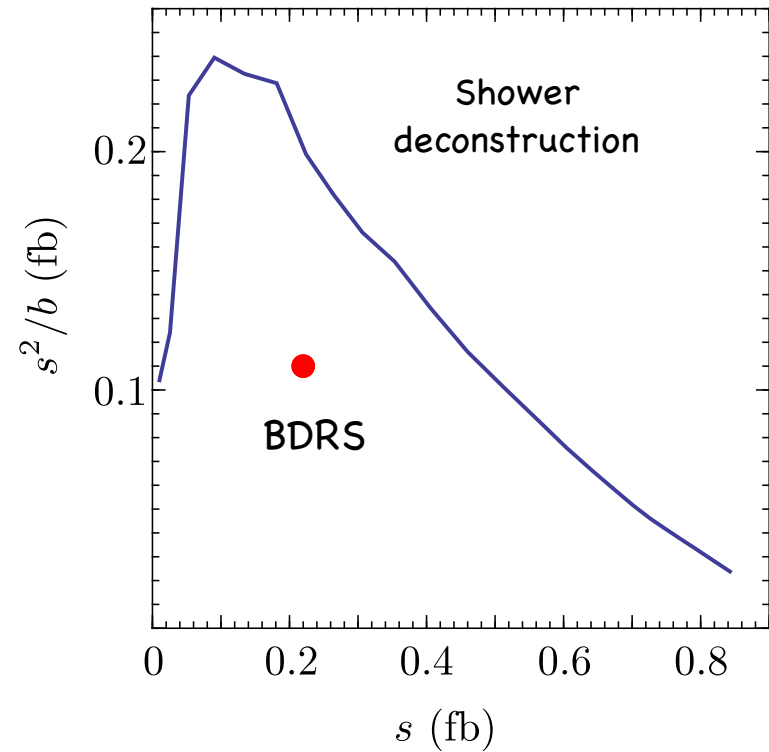
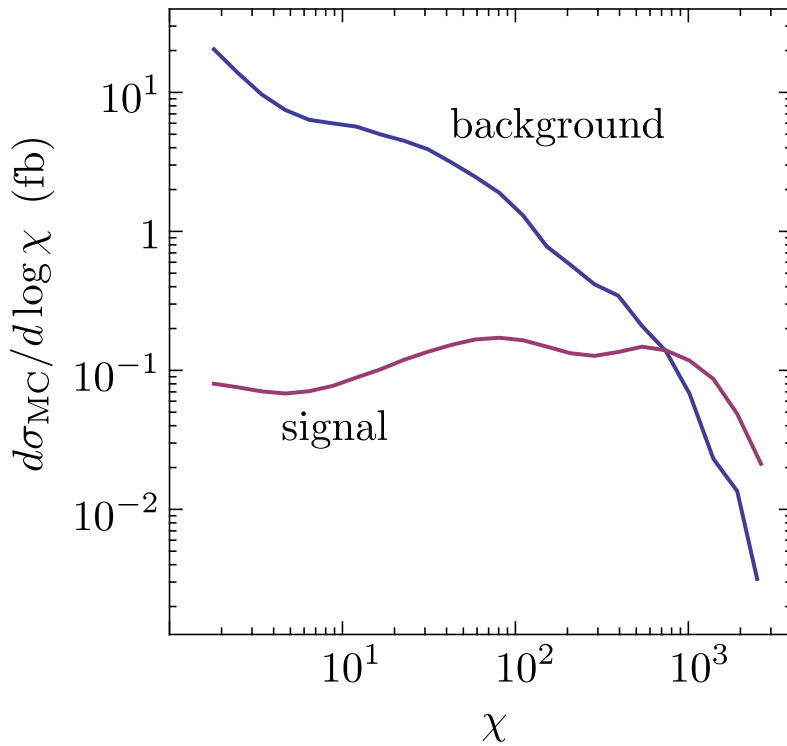


- And many more...
- And for all backgrounds...

# Results for Higgs boson:

$$\chi(\{p, t\}_N) = \frac{P(\{p, t\}_N|S)}{P(\{p, t\}_N|B)}$$

[Soper, MS PRD 84 (2011)]

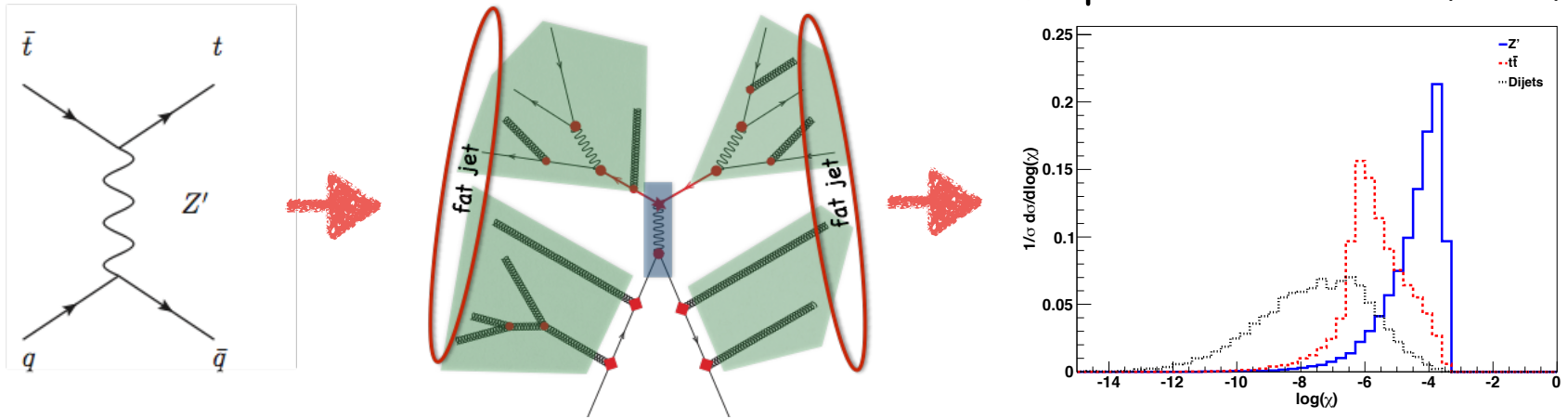


imperfect b-tagging (60%,2%) no b-tag required

# Next step, merge hard matrix element with shower approx.:

First attempt of Event Deconstruction in  $pp \rightarrow Z' \rightarrow tt$

[Soper, MS PRD 89 (2014)]

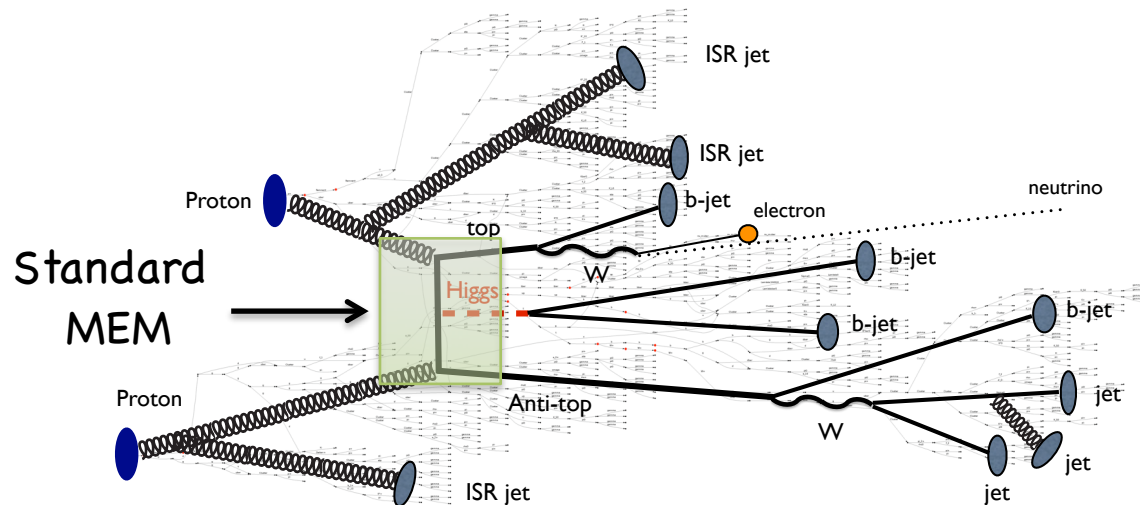


For full Event Deconstruction many steps missing:

- Include full model of Initial State Radiation
- Merge high-mult. matrix elements (CKKW)
- Add model for soft/non-pert radiation
- Improved sampling over histories

# Summary

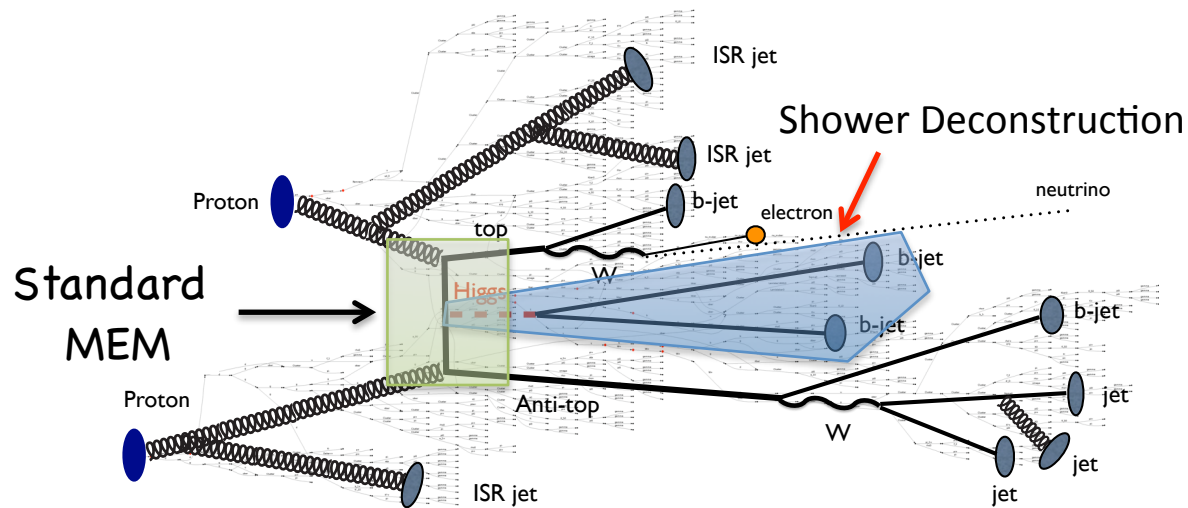
- Matrix Element Method is active field of research  
[see also MEM Workshops in Louvain (2013) and Zurich (2014)]
- Measurement of Higgs properties relies on reconstruction  
MEM can be important tool in many processes
- My personal view:





# Summary

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# Summary

- Matrix Element Method is active field of research  
[see also MEM Workshops in Louvain (2013) and Zurich (2014)]
- Measurement of Higgs properties relies on reconstruction  
MEM can be important tool in many processes
- My personal view:  
Event Deconstruction, i.e. Pattern Recognition for full event

Future MEM

