## Exclusive distributions in the Regge limit

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- Basic Concepts
- Phenomenology
- Obstributions generated by Monte Carlo integration

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- At very high energies new degrees of freedom arise
- Present in QCD, SUSY and gravity
- They allow to calculate many processes to all orders in the coupling

#### Hadron-hadron total cross-section rises (also at LHC):



Consistent with Regge theory (soft Pomeron exchange):

 $\sigma_{\rm tot} \sim s^{\alpha(0)-1} = s^{0.1}$  (Donnachie-Landshoff)<sup>2013</sup>

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Regge theory preludes QCD.

Microscopic description of Pomeron in terms of quarks & gluons?

We need a large scale  $Q > \Lambda_{\rm QCD}$  to use perturbation theory in  $\alpha_s(Q) \ll 1$ .

In the limit  $s \gg t$ ,  $Q^2$  we have  $\alpha_s(Q) \log(\frac{s}{t}) \sim \mathcal{O}(1)$ . These dominate the amplitudes and must be resummed to all orders. Kinematic origin:



where  $y_A - y_B$  is the difference in rapidity of particles A and B.

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When  $s \to \infty$  we should resum  $\alpha_s^n \log^n(s) \sim \alpha_s^n (y_A - y_B)^n$ .

$$\sigma_{\text{tot}}^{\text{LL}} = \sum_{n=0}^{\infty} \mathcal{C}_n^{\text{LL}}(\mathbf{k}_i) \alpha_s^n \int_{y_B}^{y_A} dy_1 \int_{y_B}^{y_1} dy_2 \dots \int_{y_B}^{y_{n-1}} dy_n$$
$$= \sum_{n=0}^{\infty} \frac{\mathcal{C}_n^{\text{LL}}(\mathbf{k}_i)}{n!} \underbrace{\alpha_s^n (y_A - y_B)^n}_{\text{LL}}$$

LL BFKL formalism allows us to calculate the coefficients  $C_n^{\text{LL}}(\mathbf{k}_i)$ . NLL is more complicated, sensitive to the running & choice of energy scale:

$$\sigma_{\text{tot}} = \sum_{n=1}^{\infty} \frac{\mathcal{C}_n^{\text{LL}}(\mathbf{k}_i)}{n!} \left(\alpha_s - \mathcal{A}\alpha_s^2\right)^n (y_A - y_B - \mathcal{B})^n$$
$$= \sigma_{\text{tot}}^{\text{LL}} - \sum_{n=1}^{\infty} \frac{\left(\mathcal{B} \, \mathcal{C}_n^{\text{LL}}(\mathbf{k}_i) + (n-1) \, \mathcal{A} \, \mathcal{C}_{n-1}^{\text{LL}}(\mathbf{k}_i)\right)}{(n-1)!} \underbrace{\alpha_s^n \left(y_A - y_B\right)^{n-1}}_{\text{NLL}}$$

besides, quarks enter the game ...

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Linked to elastic amplitudes due to optical theorem:



- Phenomenology
- Obstributions generated by Monte Carlo integration

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Associated phenomenology is very rich and present in different colliders: lepton-lepton (LEP,  $\sigma_{\gamma^*\gamma^*}$ ), lepton-hadron (DIS at low x), hadron-hadron (Tevatron, LHC), heavy ions (RHIC, LHC)

Diagram with a cut: events with high multiplicity :





Two key elements in the calculations:

• Coupling of gluon ladder to external states (Process dependent) High energy effective action at NLL.

April 15, 2014

11 / 26

Control of universal exchange (gluon ladder): Monte Carlo event generator at NLL.

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LHC phenomenology : examples with a cut



Production of Jets, W, Z y Drell-Yan in different topologies:



Production of quark pairs & Higgs in different topologies:



LHC phenomenology : examples without a cut



Diffractive production in different topologies:



Gap = region in the detector without hadronic activity. Clear signal.

Observable proposed in [ASV][NPB746 2006] [ASV,Schwennsen][NPB 2007,PRD77 2008]

#### as perfect to isolate BFKL



Confirmation of this idea in 2013 [Caporale, Murdaca, ASV, Salas][NPB\_875 2013],

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NLL vertices already calculated :

Central Jet production [Bartels, ASV, Schwennsen][JHEP0611 2006]:

Forward Jet production [Hentschinski,ASV][PRD85 2012] + mini Jets

[Chachamis,Hentschinski,Madrigal,ASV] [NPB861 2012,PRD87,NPB876 2013]:

Forward Jet production + rapidity Gap

[Hentschinski, Madrigal, Murdaca, ASV] [1404.2937]:



Working on numerical implementation in Monte Carlo generator...

Together with a model for the coupling of Pomeron to the proton ...

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DIS data:  $F_2(x, Q^2) \simeq x^{-\lambda}$ Hadron-Pomeron coupling based on NLL BFKL



#### Transition from Hard to Soft Pomeron well described.

[Hentschinski,ASV,Salas][PRL110, PRD87 2013]

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#### Monte Carlo

Effective Feynman rules: Simplest case, minijet production at LL.

Gluon Regge trajectory:  $\omega\left(\vec{q}\right) = -\frac{\alpha_s N_c}{\pi} \log \frac{q^2}{\lambda^2}$ 

Modified propagators in the *t*-channel:

$$\left(\frac{s_i}{s_0}\right)^{\omega(t_i)} = e^{\omega(t_i)(y_i - y_{i+1})}$$



$$\left(\frac{\alpha_{s}N_{c}}{\pi}\right)^{2} \int d^{2}\vec{k}_{1} \frac{\theta\left(k_{1}^{2}-\lambda^{2}\right)}{\pi k_{1}^{2}} \int d^{2}\vec{k}_{2} \frac{\theta\left(k_{2}^{2}-\lambda^{2}\right)}{\pi k_{2}^{2}} \delta^{(2)}\left(\vec{k}_{A}+\vec{k}_{1}+\vec{k}_{2}-\vec{k}_{B}\right) \\ \times \int_{0}^{Y} dy_{1} \int_{0}^{y_{1}} dy_{2} e^{\omega\left(\vec{k}_{A}\right)\left(Y-y_{1}\right)} e^{\omega\left(\vec{k}_{A}+\vec{k}_{1}\right)\left(y_{1}-y_{2}\right)} e^{\omega\left(\vec{k}_{A}+\vec{k}_{1}+\vec{k}_{2}\right)y_{2}}$$

$$\sigma(Q_{1}, Q_{2}, Y) = \int d^{2}\vec{k}_{A}d^{2}\vec{k}_{B} \underbrace{\phi_{A}(Q_{1}, \vec{k}_{A})\phi_{B}(Q_{2}, \vec{k}_{B})}_{\text{PROCESS-DEPENDENT}} \underbrace{f(\vec{k}_{A}, \vec{k}_{B}, Y)}_{\text{UNIVERSAL}}$$

$$f(\vec{k}_{A}, \vec{k}_{B}, Y) = \sum_{n} \left| \int_{\gamma_{e}=0, k_{B}}^{\gamma_{e}=\gamma, k_{A}} \int_{\gamma_{e}=0, k_{B}}^{\gamma_{e}=\gamma, k_{A}} \right|^{2}$$

$$= e^{\omega(\vec{k}_{A})Y} \left\{ \delta^{(2)}(\vec{k}_{A} - \vec{k}_{B}) + \sum_{n=1}^{\infty} \prod_{i=1}^{n} \frac{\alpha_{s}N_{c}}{\pi} \int d^{2}\vec{k}_{i} \frac{\theta(k_{i}^{2} - \lambda^{2})}{\pi k_{i}^{2}} \right\}$$

$$\times \int_{0}^{\gamma_{i-1}} dy_{i} e^{(\omega(\vec{k}_{A} + \sum_{l=1}^{i} \vec{k}_{l}) - \omega(\vec{k}_{A} + \sum_{l=1}^{i-1} \vec{k}_{l}))y_{i}} \delta^{(2)}(\vec{k}_{A} + \sum_{l=1}^{n} \vec{k}_{l} - \vec{k}_{B})$$

3



<sup>[</sup>Chachamis,ASV] [PLB709 2012]

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April 15, 2014 20 / 26

Gluon  $|p_T|$  in a given rapidity bin?



[Chachamis, ASV][PLB717 2012] [...+Salas][PRD87 2013][...+Murdaca, Caporale, Madrigal][PLB724 2013]

April 15, 2014 21 / 26

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Need to extend the formalism to include collinear regions [Salam] [ASV][NPB722 2005]

$$f = e^{\omega(\vec{k}_{A})Y} \left\{ \delta^{(2)}\left(\vec{k}_{A} - \vec{k}_{B}\right) + \sum_{n=1}^{\infty} \prod_{i=1}^{n} \frac{\alpha_{s} N_{c}}{\pi} \int d^{2}\vec{k}_{i} \frac{\theta\left(k_{i}^{2} - \lambda^{2}\right)}{\pi k_{i}^{2}} \\ \times \int_{0}^{y_{i-1}} dy_{i} e^{\left(\omega\left(\vec{k}_{A} + \sum_{l=1}^{i} \vec{k}_{l}\right) - \omega\left(\vec{k}_{A} + \sum_{l=1}^{i-1} \vec{k}_{l}\right)\right)y_{i}} \delta^{(2)}\left(\vec{k}_{A} + \sum_{l=1}^{n} \vec{k}_{l} - \vec{k}_{B}\right) \right\}$$
  
Key at NLL:  $\theta\left(k_{i}^{2} - \lambda^{2}\right) \rightarrow \theta\left(k_{i}^{2} - \lambda^{2}\right) - \underbrace{\frac{\bar{\alpha}_{s}}{4} \ln^{2}\left(\frac{\vec{k}_{A}^{2}}{\left(\vec{k}_{A} + \vec{k}_{i}\right)^{2}}\right)}_{\text{NLL}}$   
Drigin: We must resum collinear emissions to have convergent observables

in a more general kinematics (less constrained impact factors):

$$\theta\left(k_{i}^{2}-\lambda^{2}\right) \rightarrow \theta\left(k_{i}^{2}-\lambda^{2}\right) + \sum_{n=1}^{\infty} \frac{\left(-\bar{\alpha}_{s}\right)^{n}}{2^{n}n!(n+1)!} \ln^{2n} \left(\frac{\vec{k}_{A}}{\left(\vec{k}_{A}+\vec{k}_{i}\right)^{2}}\right) = 2^{n} \sqrt{2^{n}n!(n+1)!}$$

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April 15, 2014

24 / 26





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April 15, 2014 25 / 26

### Exclusive distributions in the Regge limit

- Rich phenomenology, all ingredients are ready. New observables?
- Workshops Series in Madrid on MRK (October 2012, February 2014)



