

# Probing anomalous quartic gauge couplings using proton tagging at the Large Hadron Collider

LHC Working Group on Forward Physics and Diffraction

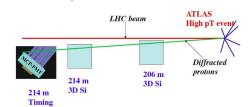
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E. Chapon, S. Fichet, G. von Gersdorff,
O. Kepka, B. Lenzi, C. Royon

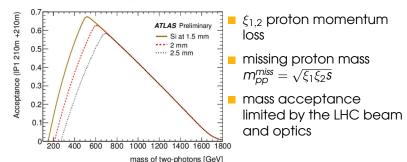
<sup>1</sup>CEA Saclay - Irfu/SPP

16/04/2014

### Forward proton detectors at the LHC

The ATLAS Forward Physics Project (AFP) and the Precision Proton Spectrometer (PPS, CMS/TOTEM)





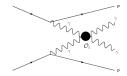






# Exclusive production via photon induced processes



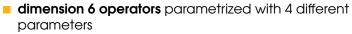


- All particles at the final state are detected: two protons in the forward detectors and two high energy particles in the central detector → strong kinematics constraints
- Requirement of two intact protons + kinematics constraints → strong background reduction
- Final state can be  $\gamma\gamma$ , WW, ZZ  $\rightarrow$  ideal to study anomalous quartic gauge couplings (aQGC)
- aQGC important for various topics: electroweak symmetry breaking, extra-dimension models, ...
- Drawback: reduced cross-sections (intact protons)

# $WW\gamma\gamma$ and $ZZ\gamma\gamma$ anomalous couplings

C. Royon, O. Kepka, Phys. Rev. D 78 (2008)

E. Chapon, C. Royon, O. Kepka, Phys. Rev. D 81 (2010)



$$\mathcal{L}_{6}^{0} \sim \frac{-e^{2}}{8} \frac{a_{0}^{W}}{\Lambda^{2}} F_{\mu\nu} F^{\mu\nu} W^{+\alpha} W_{\alpha}^{-} - \frac{e^{2}}{16 \cos^{2}(\theta_{W})} \frac{a_{0}^{Z}}{\Lambda^{2}} F_{\mu\nu} F^{\mu\nu} Z^{\alpha} Z_{\alpha}$$

$$\mathcal{L}_{6}^{C} \sim \frac{-e^{2}}{16} \frac{a_{C}^{W}}{\Lambda^{2}} F_{\mu\alpha} F^{\mu\beta} (W^{+\alpha} W_{\beta}^{-} + W^{-\alpha} W_{\beta}^{+})$$

$$- \frac{e^{2}}{16 \cos^{2}(\theta_{W})} \frac{a_{C}^{Z}}{\Lambda^{2}} F_{\mu\alpha} F^{\mu\beta} Z^{\alpha} Z_{\beta}$$

- Only the leptonic decays of the heavy bosons are considered as final states (clean signal)
- **Background considered:** ND WW/ZZ production, dilepton photoproduction, DPE dilepton, DPE WW/ZZ
- Generation and simulation performed with the Forward Physics MC generator (FPMC) interfaced with the fast simulation of the ATLAS detector (ATLFast++ package) ATLAS full simulation also performed and gave similar results



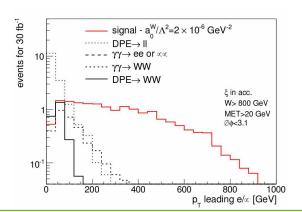


# $WW\gamma\gamma$ and $ZZ\gamma\gamma$ anomalous couplings

E. Chapon, C. Royon, O. Kepka, Phys. Rev. D 81 (2010)



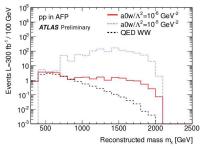
■ WW $\gamma\gamma$  signal and background composition passing all the selection but the p<sub>T</sub> cut on the leading lepton

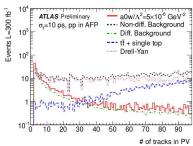


# Dealing with pile-up at the LHC

#### ATLAS full simulation study

- The LHC is operated at very high luminosity → high event multiplicites in a single bunch-crossing (pile-up)
  - Cut on the number of tracks fitted to the primary vertex: very efficient to remove remaining pile up after requesting a high mass object to be produced
  - Full simulation results





### $WW\gamma\gamma$ and $ZZ\gamma\gamma$ sensitivities

E. Chapon, C. Royon, O. Kepka, Phys. Rev. D 81 (2010)



- Best limits from LEP, OPAL (Phys. Rev. D **70** (2004) 032005) of the order of 0.02-0.04 GeV $^{-2}$  for each coupling
- Recent papers from D0 and CMS for  $WW\gamma\gamma$  with reach of the order of  $10^{-4}$  GeV<sup>-2</sup> (CMS-PAS-FSQ-12-010)
- Sensitivites predicted at the LHC (P. J. Bell, ArXiV:0907.5299) of the order of a few 10<sup>-4</sup> GeV<sup>-2</sup>
- Sensitivities predictions with AFP (30 and 200 fb<sup>-1</sup>) improvement by a factor  $\simeq 100$

		limits $[10^{-6}  \mathrm{GeV}^{-2}]$				limits $[10^{-6}  \text{GeV}^{-2}]$			
	form factor	$\left a_0^W/\Lambda^2\right $	$\left a_C^W/\Lambda^2\right $	$\left a_0^Z/\Lambda^2\right $	$\left a_C^Z/\Lambda^2\right $	$\left a_0^W/\Lambda^2\right $	$\left a_C^W/\Lambda^2\right $	$\left a_0^Z/\Lambda^2\right $	$\left a_C^Z/\Lambda^2\right $
95% c.1 {	$\Lambda_{cut} = \infty$	1.2	4.2	2.8	10	0.7	2.4	1.1	4.1
	$\Lambda_{cut} = \infty$ $\Lambda_{cut} = 2  \text{TeV}$	2.6	9.4	6.4	24	1.4	5.2	2.5	9.2
$3\sigma$ evidence $\Big\{$	$\Lambda_{cut} = \infty$ $\Lambda_{cut} = 2 \text{ TeV}$	1.6	5.8	4.0	14	0.85	3.0	1.6	5.7
	$\Lambda_{cut} = 2\text{TeV}$	3.6	13	9.0	34	1.8	6.7	3.5	13
$5\sigma$ discovery $\Big\{$	$\Lambda_{cut} = \infty$ $\Lambda_{cut} = 2 \text{ TeV}$	2.3 5.4	9.7	6.2	23	1.2	4.3	4.1	8.9
	$\Lambda_{cut} = 2  \text{TeV}$	5.4	20	14	52	2.7	9.6	5.5	20

# $\gamma\gamma\gamma\gamma$ anomalous couplings (2013-2014)

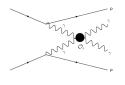




- Direct coupling absent from the SM\*
  \*loop induced production measurable at the LHC with heavy ions (D. D'Enterria, G. Da Silveira, PRL 111 (2013)
- New couplings predicted by strongly-interacting composite states, extra-dimensions, ...
- No constraints from collider experiments

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- New couplings predicted by strongly-interacting composite states, extra-dimensions, ...
- No constraints from collider experiments
- Small couplings, high luminosity required,  $\mu > 50$ . (300 fb<sup>-1</sup> of data expected at the LHC at  $\sqrt{s} = 14$  TeV)
- **Huge background** if only 2 high energy  $\gamma$  required (SM  $\gamma\gamma$  production + fakes from electrons and jets)
- Additional requirement of **two intact protons** with forward detectors highly suppresses the background



# Operators of the $\gamma\gamma\gamma\gamma$ couplings

R.S. Gupta, Phys. Rev. D 85 (2012) 014006

S. Fichet and G. von Gersdorff, arXiv:1311.6815



$$L_{4\gamma} = \frac{\zeta_1^{\gamma}}{F_{\mu\nu}} F^{\mu\nu} F_{\rho\sigma} F^{\rho\sigma} + \frac{\zeta_2^{\gamma}}{F_{\mu\nu}} F^{\nu\rho} F_{\rho\sigma} F^{\sigma\mu} \text{ (dimension 8)}$$

Unitary breaks down at large scale  $\Lambda' \to \mathbf{use}$  of a form factor (f.f.) at the amplitude level

We use 
$$f.f = \frac{1}{1+(\frac{\hat{s}\gamma\gamma}{2})^2}$$
 with  $\Lambda' = 1$  TeV



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$$f.f = \frac{1}{1 + (\frac{\hat{s}_{\gamma\gamma}}{\lambda t^2})^2}$$
 with  $\Lambda' = 1$  TeV

f.f. not necessary when the new particles have a large enough mass



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■ Unitary breaks down at large scale  $\Lambda' \rightarrow$  use of a form factor (f.f.) at the amplitude level

We use 
$$f.f = \frac{1}{1 + (\frac{\hat{s}_{\gamma\gamma}}{d^2})^2}$$
 with  $\Lambda' = 1$  TeV

- f.f. not necessary when the new particles have a large enough mass
- The signal showed in the plots of this presentation are for a signal with  $\zeta_1 > 0$  and  $\zeta_2 = 0$  and with f.f.  $\zeta_1$  and  $\zeta_2$  have the same angular behaviour
- A table with sensitivities for both  $\zeta_1$  and  $\zeta_2$  are given at the end of the presentation



# New physics contributions to $4\gamma$ couplings

#### New charged particles via loops

- Effective coupling only depends on the mass, charge and spin :  $\zeta_i^{\gamma} \propto c_i^s Q^4 m^{-4}$
- Example: top partners

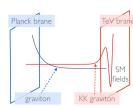
#### New neutral particles at tree level

- Effective coupling depends on mass, spin and the non-renormalizable  $\gamma\gamma X$  coupling  $\zeta_i^{\gamma} \propto b_i^s f^{-2} m^{-2}$
- Example: KK gravitons, dilaton/radion (warped extra-dim)

if coupling  $\simeq$  TeV and  $m_{KK} \simeq$  few TeV  $\rightarrow$  predictions of  $\zeta_i^{\gamma} \simeq 10^{-14} \cdot 10^{-13}$  GeV<sup>-4</sup>, which we are sensitive





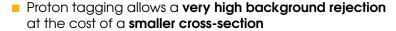


# Where does proton tagging do better?



- Proton tagging allows a very high background rejection at the cost of a smaller cross-section
  - A single observation has a high significance
  - Ideal to probe small deviations from the Standard Model like aQGC
  - ex: new charged particles via loops, ADD gravity effects, ...
  - Should be less competitive on resonances searches ex: new neutral particles at tree level
  - Very difficult to quantify precisely the improvements compared to the central detector alone (in progress)

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- A single observation has a high significance
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- Should be less competitive on resonances searches ex: new neutral particles at tree level
- Very difficult to quantify precisely the improvements compared to the central detector alone (in progress)
- For the first time, we reach sensitivities allowing to probe directly a large class of new models
  - Warped extra-dimensions: KK gravitons, radion/dilaton, high  $\kappa$  untested domain (Randall-Sundrum model)
  - Strongly-interacting composite states, monopoles: generic searches of new heavy charged particles





- Evaluate the LHC potential to probe  $4\gamma$  couplings using proton tagging
  - **4** $\gamma$  aQGC operators implemented in the **FPMC generator** in summer 2013
  - Pile-up simulation with Pythia8 minimum bias events
  - Background estimation (very small)
  - The backgrounds from pile-up should dominate
  - 2 scenarios were considered
    - LHC full stat (ATLAS or CMS) : 300 fb<sup>-1</sup>,  $\mu = 50$
    - HL-LHC (ATLAS) : 3000 fb $^{-1}$ ,  $\mu = 200$

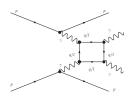
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- New: update of the exclusive  $\gamma\gamma$  SM production, adding the W loop contribution and the fermion masses



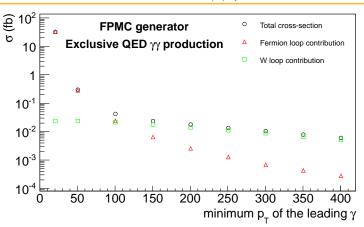
# SM QED exclusive $\gamma\gamma$ production





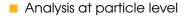
- Different loop contributions: fermions (quarks, leptons), vectors (W)
- W loop contribution and massive fermions added to the process in FPMC rev.913 (negligible at low mass but not at high mass)
- Interferences SM/Exotics added for the full amplitude calculation

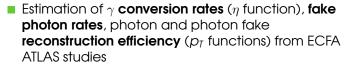
# SM QED exclusive $\gamma\gamma$ production



- W loop non negligeable for  $p_{T,\gamma} > 50$  GeV
- QCD and DPE contributions to be added
- $\blacksquare$  Same plot for the  $\gamma\gamma$  mass is being made
- Mass of the fermions taken into account







- **Smearing** of 1% in  $\gamma$  energies, 0.001 in  $\eta$  and  $\phi$  (absolute), 2% for  $\xi$  to mimic detector resolution
- Requirement of at least one converted photon  $\rightarrow$  constraint on the  $\gamma$  vertex, possibility to combine with forward proton timing measurement
- Selection on high  $p_T^{\gamma}$ , high diphoton mass,  $\Delta \Phi^{\gamma\gamma}$ , match proton missing/ $\gamma\gamma$  mass (see S12-S17)

#### Final outputs

- $5\sigma$  and 95% C.L sensitivities on the  $\gamma\gamma\gamma\gamma$  vertex (effective field theory)
- M-Q exclusion plane for generic exotic fermions/vectors (full amplitude)



# Conversion, fake and efficiency reconstruction rates



- Inputs from the ECFA ATLAS studies
- Photon conversion factors: 15% in the barrel, 30% in the end-caps
- Photon and electron reconstruction efficiency:  $Eff(p_T) = 0.76 1.98 \ exp^{\frac{-p_T}{16.1(GeV)}}$
- **Photon fake factors:** 1% for electron European Strategy studies
- **Fake photon p**<sub>T</sub> **for jets:** gaussian draw (Mean=75%, $\sigma$ =13%) on the jet p<sub>T</sub> and use of

$$Eff_{fake}(p_T) = 0.0093 \text{ exp}^{\frac{-min(p_T, 200 \text{GeV})}{17.5(\text{GeV})}}$$

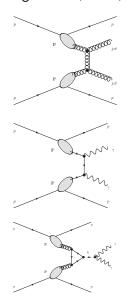
almost no fake  $\gamma$  from jets at very high  $p_T$ 

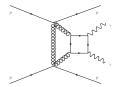
# Backgrounds (FPMC, ExHuME)

#### 

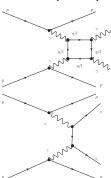








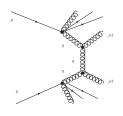
#### Exclusive QED (FPMC)



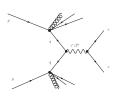
# Pile-up backgrounds (HERWIG 6.5)



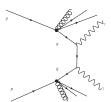
#### Dijet



#### Drell-Yan



#### Diphoton



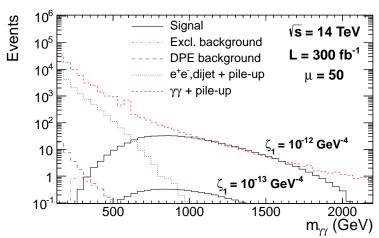
+ protons generated from **minimum bias events** (Pythia 8)

transported to the forward detectors through the LHC magnets with FPTracker/MADX

# Mass distribution of signal and backgrounds

 $\blacksquare$  0.015 <  $\xi$  < 0.15,  $|\eta|$  < 2.37,  $p_{71,2}^{\gamma}$  > 50 GeV ONLY

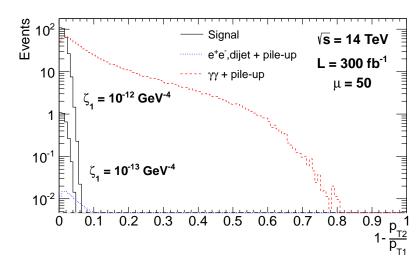




 $\blacksquare$  By requesting  $m_{\gamma\gamma} >$  600 GeV, Only pile-up backgrounds remain

# Exclusive signal: $p_T$ ratio

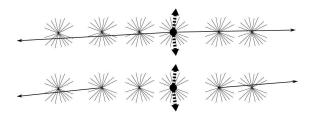




 $\blacksquare$  p<sub>T</sub> ratio distribution after p<sub>T</sub> and  $m_{\gamma\gamma}$  cuts

#### Forward detectors measurements

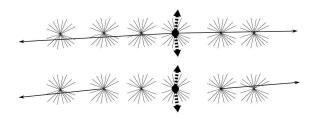




■ **Proton missing mass** measurement with 3% resolution in case of double tag

#### Forward detectors measurements

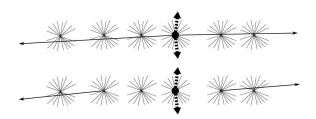




- Proton missing mass measurement with 3% resolution in case of double tag
- It matches the central  $\gamma\gamma$  mass for signal. Can match as well for pile-up backgrounds as a statistical fluctuation

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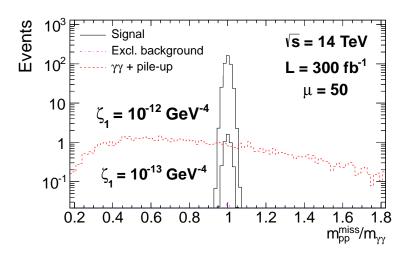


- **Proton missing mass** measurement with 3% resolution in case of double tag
- It matches the central  $\gamma\gamma$  mass for signal. Can match as well for pile-up backgrounds as a statistical fluctuation
- Double tag probability from pile-up protons on forward detectors (no mass requirement):

32% (
$$\mu$$
 = 50) 66% ( $\mu$  = 100) 93% ( $\mu$  = 200)

# Mass matching and pile-up

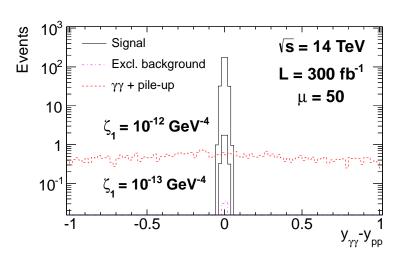




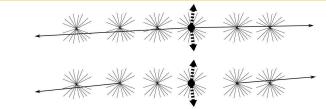
■ A mass window of 3% (= resolution) is required in the event selection

# Rapidity cut





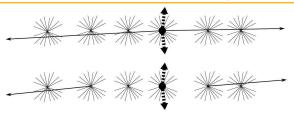
- $|y_{\gamma\gamma} y_{pp}| < 0.03$  with  $y_{pp} = (0.5 * ln(\frac{\xi_1}{\xi_2}))$
- $\blacksquare$  Small width for signal due to the resolution on  $y_{\gamma\gamma}$  and  $\xi_{1,2}$



Proton timing will be measured by forward detectors

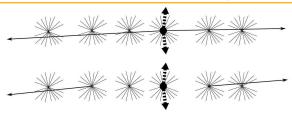






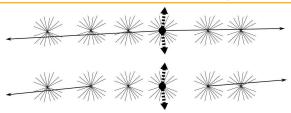
- Proton timing will be measured by forward detectors
  - 10 ps resolution assumed → proton vertex constrained within 2.1 milimeters



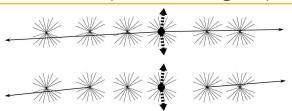


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  - Resolution driven by forward timing detectors



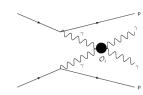
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  - Requirement of 1 converted  $\gamma \rightarrow <$  1 mm resolution on the  $\gamma$  vertex
  - Resolution driven by forward timing detectors
- $\blacksquare$  additional background rejection factor of 40 at  $\mu=50$
- No need to use it for us, robustness of the analysis
- can be used for unknown backgrounds (beam-induced)



### Event selection summary

#### Kinematic cuts

- $p_{T1}^{\gamma} > 200 \text{ GeV}, p_{T2}^{\gamma} > 100 \text{ GeV}$
- $2 m_{\gamma\gamma} > 600 \text{ GeV}$

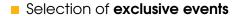


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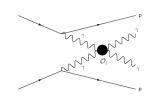


$$p_{T1}^{\gamma} > 200 \text{ GeV}, p_{T2}^{\gamma} > 100 \text{ GeV}$$

$$m_{\gamma\gamma} > 600 \text{ GeV}$$



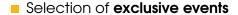
- $\frac{p_{72}}{p_{71}} > 0.95$
- $|\Delta \Phi| > \pi 0.01$



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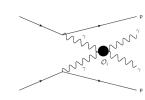
- $p_{T1}^{\gamma} > 200 \text{ GeV}, p_{T2}^{\gamma} > 100 \text{ GeV}$
- $m_{\gamma\gamma} > 600 \text{ GeV}$



- $\frac{p_{12}}{p_{11}} > 0.95$
- $|\Delta \Phi| > \pi 0.01$



- $m_{pp}^{miss} = m_{\gamma\gamma} \pm 3\%$
- $|y_{\gamma\gamma} y_{DD}| < 0.03$ with  $y_{pp} = (0.5 * ln(\frac{\xi_1}{\xi_2}))$
- 3 Possible proton timing measurement with forward detectors



## Expected events for $\zeta_1^{\gamma} = 2 \cdot 10^{-13} \cdot \text{GeV}^{-4}$

 $\sqrt{s}$  = 14 TeV, L = 300 fb<sup>-1</sup>, at least one converted  $\gamma$ 

Cut / Process	Signal	Excl.	DPE	e <sup>+</sup> e <sup>-</sup> ,dijet + pu	$\gamma\gamma$ + pu
$0.015 < \xi < 0.15, p_{\text{T1/2}} > 50 \text{GeV}$	20.8	3.7	48.2	2.8 · 10 <sup>4</sup>	1.0 · 10 <sup>5</sup>
$p_{\rm T1} > 200 {\rm GeV}, p_{\rm T2} > 100 {\rm GeV}$	17.6	0.2	0.2	1.6	2968
$m_{\gamma\gamma} > 600$ GeV	16.6	0.1	0	0.2	1023
$p_{\rm T2}/p_{\rm T1} > 0.95,  \Delta\phi  > \pi - 0.01$	16.2	0.1	0	0	80.2
$\sqrt{\xi_1\xi_2s}=m_{\gamma\gamma}\pm 3\%$	15.7	0.1	0	0	2.8
$ y_{\gamma\gamma} - y_{pp}  < 0.03$	15.1	0.1	0	0	0

- Signal selection efficiency > 70% (after preselection)
  - Acceptance increased by a factor 3-4 when adding unconverted photons (with EM "pointing")
- - 1.5 background events expected at  $\mu$  = 200 Robust analysis, very good background control
  - proton time-of-flight **not used**Possible additional rejection factor of 40 at  $\mu$  = 50





# Final discovery (5 $\sigma$ ) and exclusion (95% CL) sensitivities on $\zeta_1$ and $\zeta_2$



S. Fichet, G. von Gersdorff, O. Kepka, B. Lenzi, C. Royon, M. Saimpert, arXiv:1312.5153 (2013)

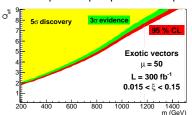
Luminosity	300 fb <sup>-1</sup>	300 fb <sup>-1</sup>	300 fb <sup>-1</sup>	3000 fb <sup>-1</sup>		
pile-up (μ)	pile-up (μ) 50		50	200		
coupling	$\geq$ 1 conv. $\gamma$	$\geq$ 1 conv. $\gamma$	all $\gamma$	all $\gamma$		
(GeV <sup>-4</sup> )	5 σ	95% CL	95% CL	95% CL		
ζ <sub>1</sub> f.f.	$\zeta_1$ f.f. $1 \cdot 10^{-13}$		5 · 10 <sup>-14</sup>	2.5 · 10 <sup>-14</sup>		
$\zeta_1$ no f.f.	$3.5 \cdot 10^{-14}$	2.5 · 10 <sup>-14</sup>	1.5 · 10 <sup>-14</sup>	7 · 10 <sup>-15</sup>		
ζ <sub>2</sub> f.f.			1 · 10 <sup>-13</sup>	4.5 · 10 <sup>-14</sup>		
$\zeta_2$ no f.f.	$7.5 \cdot 10^{-14}$	5.5 · 10 <sup>-14</sup>	3 · 10 <sup>-14</sup>	1.5 · 10 <sup>-14</sup>		

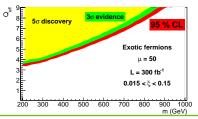
 KK gravitons and strongly-coupled dilaton/radion (warped extra-dimensions) can be discovered in the multi-TeV range

# **Full amplitude** computation for generic heavy charged fermions/vectors contributions (preliminary)



- The existence of new heavy charged particles will enhance the  $\gamma\gamma\gamma\gamma$  coupling at high mass via loops
- This enhancement can be parametrized by **only the** mass and the effective charge  $Q_{eff} = Q.N^{1/4}$ , N multiplicity
- Generic implementation for fermions and vectors implemented in FPMC
- Paper in preparation, preliminary M-Q<sub>eff</sub> exclusion plane





# **Full amplitude** computation for generic heavy charged fermions/vectors contributions (preliminary)



Link full amplitude - effective field theory

$$\zeta_i^{\gamma} = c_i^s Q_{\rm eff}^4 m^{-4} \alpha_{\rm em}^2$$
,  $c_i \simeq$  0.01 (0.1) for fermions (vectors)

Typical sensitivity with the full amplitude calculation

$$M = 800 \text{ GeV}$$
,  $Q_{eff} > 7$  (4) for fermions (vectors)

- Gives a coupling of  $\simeq 10^{-15}$  in terms of  $\zeta_i$
- Corresponds to the sensitivity we had using the effective field theory → successful cross-check of the method

#### Conclusion



- Forward proton tagging at the LHC seems promising to probe anomalous Quartic Gauge Couplings
  - proton tagging associated by high energy object detections in the central EM calorimeter allow to completely suppress the background
  - $WW\gamma\gamma$  and  $ZZ\gamma\gamma$  couplings already studied with positive outputs (improvement by a factor > **100**)
  - $\gamma\gamma\gamma\gamma$  coupling: sensitivities around  $10^{-13} 10^{-14}$  GeV<sup>-4</sup>, down to  $7 \cdot 10^{-15}$  GeV<sup>-4</sup>  $\rightarrow$  allows to probe directly a large panel of new physics models no previous constraints from collider experiments
- γγγγ coupling: a channel probing exotic heavy charged vectors/fermions (even scalars, sensitivity is smaller) in a completely model-independent way
  - sensitive for vectors (fermions) up to 1400 (920) GeV

#### Conclusion



- Effective field theory:  $5\sigma$  discovery with less luminosity (1 fb<sup>-1</sup>, 10 fb<sup>-1</sup>, 50 fb<sup>-1</sup>):  $7 \cdot 10^{-13}$ ,  $2 \cdot 10^{-13}$ ,  $9 \cdot 10^{-14}$  GeV<sup>-4</sup>
- $\sim \gamma \gamma \gamma \gamma$  paper submitted : arXiv:1312.5153
- More detailed paper including the full amplitude calculations for loop contributions in preparation



# Probing anomalous quartic gauge couplings using proton tagging at the Large Hadron Collider

Back-up

Matthias Saimpert<sup>1</sup>
E. Chapon, S. Fichet, G. von Gersdorff,
O. Kepka, B. Lenzi, C. Royon

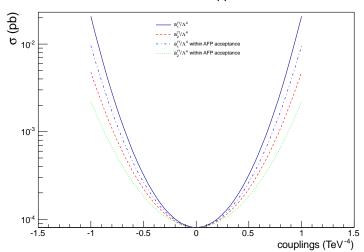
<sup>1</sup>CEA Saclay - Irfu/SPP

16/04/2014

## Integrated total cross-section against couplings

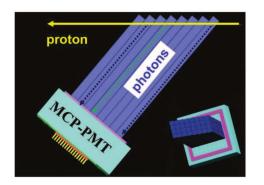


#### Form factor applied



# Forward timing detectors: inefficiencies due to pile-up protons





	Inefficiencies - 2mm bar detector										
Ì	Bar	1	2	3	4	5	6	7	8	9	10
ĺ	$\mu = 50$	0.129	0.085	0.067	0.057	0.049	0.046	0.043	0.040	0.036	0.011
ĺ	$\mu = 100$	0.185	0.122	0.097	0.082	0.071	0.066	0.062	0.057	0.051	0.016

M. Saimpert. Search for new states of matter wih the ATLAS experiment at the LHC, Master Thesis MINES ParisTech (2013)



#### The BSM amplitudes



- Loops of spin 0,1/2, 1 new electric particles contribute to  $4\gamma$  . Because all vertices are fixed by gauge invariance, the NP contributions depend only on spin, mass and
- For example in the effective theory limit :  $\zeta_i^{\gamma} = \alpha_{em}^2 Q^4 m^{-4} N c_{i,s}$

$$c_{1,s} = \begin{cases} \frac{1}{288} & s = 0 \\ -\frac{1}{36} & s = \frac{1}{2} \\ -\frac{5}{22} & s = 1 \end{cases}, \quad c_{2,s} = \begin{cases} \frac{1}{360} & s = 0 \\ \frac{7}{90} & s = \frac{1}{2} \\ \frac{27}{40} & s = 1 \end{cases}$$
 Scalar loops are smaller!

- Full amplitudes for fermions and vectors are now implemented in FPMC.
- Amplitudes get enhanced near the threshold



## $\gamma\gamma\gamma\gamma$ full amplitude calculation (S. Fichet)



#### The SM background



- All electric particles of the SM contribute : leptons, quarks and W bosons
- The imaginary part of certain W helicity amplitudes grows with the energy, while the fermion ampliudes are finite. Background is dominated by the W loop
- When the new particle is real, it interfers with the W loop.
  - On-shell NP signal enhanced by SM interference
- All SM background amplitudes are implemented in FPMC (+ swiches to separately turn off them)
- One can check that SM fermions contributions are negligible.
  - Keeping only the W loop provides a huge gain of CPU time!