The background of the slide is a composite image. On the left, there is a photograph of a nebula with vibrant orange and yellow gases against a dark background. In the center, a long, diagonal spectrum is displayed, composed of numerous small pixels in various colors (red, orange, yellow, green, blue, purple) that form a distinct arc. On the right side, the Earth is visible, showing its blue oceans and green continents from space.

ISOLDE Nuclear Reaction  
and  
Nuclear Structure Course  
22-25 April 2014

Nuclear Reaction at Intermediate to Relativistic Energies: what data tell us (I)

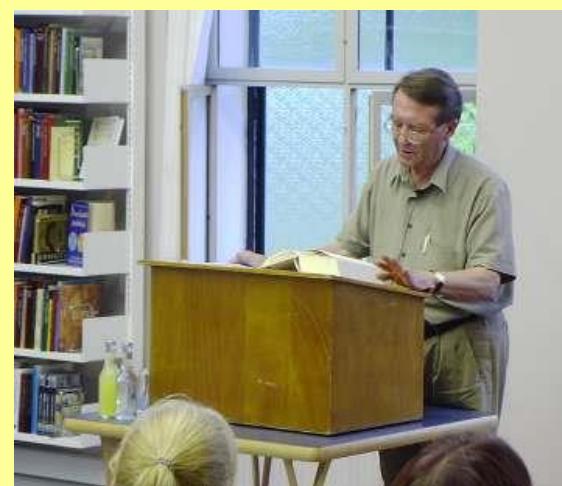
Jim Hamm

I do not know much about anything;  
what I know is already commonplace.  
Wouldn't it be nice to know just one fact,  
for example the proportions of the spots of  
the ladybird.



Jeg kender ikke meget til meget;  
det jeg kender til er i forvejen meget udbredt.  
Tænk at vide bare een bestemt ting,  
for eksempel forholdet mellem mariehønens  
pletter.

Benny Andersen,  
Viden 1965

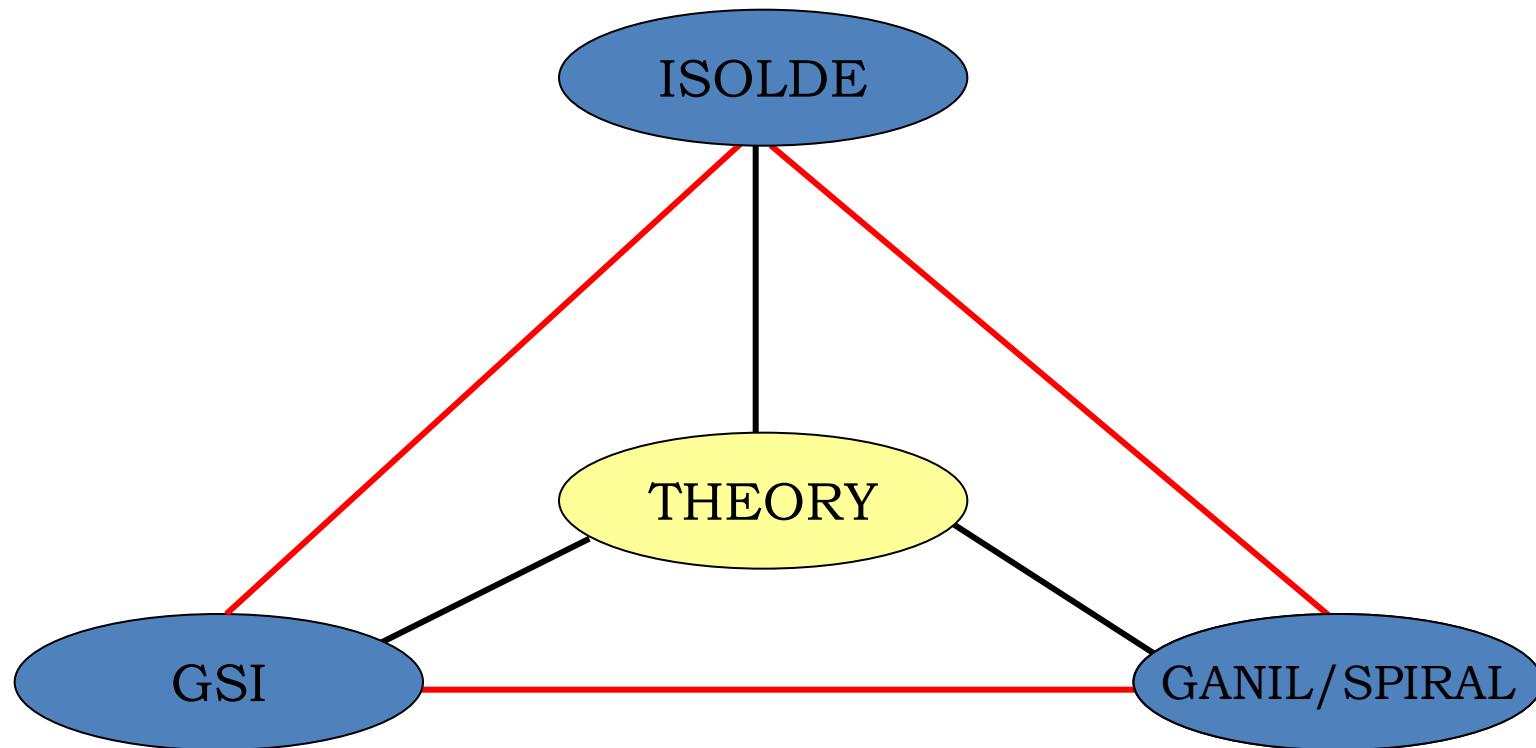
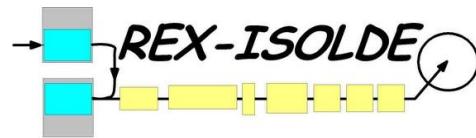


# Physics With Radioactive Beams



Nobel Symposium NS  
152

# HIE-ISOLDE



# Light Nuclei in Nature

Proton Rich

Neutron Rich

Resonances

**Z**



			15F unbound	16F unbound	17F 64.8 s	18F 109.7 m	19F	20F 11 s	21F 4.16 s
	12O unbound	13O 8.58 ms	14O 70.6 s	15O 2.03 m	16O	17O	18O	19O 27.1 s	20O 13.5 s
	11N unbound	12N 20.4 m	13N 20.4 m	14N	15N	16N 7.13 s	17N 4.17 s	18N 0.63 s	19N 329 ms
	9C 125 ms	10C 19.3 s	11C 20.4 m	12C	13C	14C 5730 y	15C 2.45 s	16C 0.747 s	17C 193 ms
	8B 770 ms	9B unbound	10B	11B	12B 20.20 ms	13B 17.33 ms	14B 13.8 ms	15B 10.4 ms	16B unbound
	7Be unbound	8Be unbound	9Be	10Be 1.6 10 <sup>6</sup> y	11Be 13.8 s	12Be 23.6 ms	13Be unbound	14Be 4.35 ms	15Be unbound
	6Li	7Li	8Li 840 ms	9Li 179 ms	10Li unbound	11Li 8.5 ms	12Li unbound	13Li unbound	
	3He	4He	5He unbound	6He 808 ms	7He unbound	8He 119 ms	9He unbound	10He unbound	

1H    2H

3H  
12.323 y

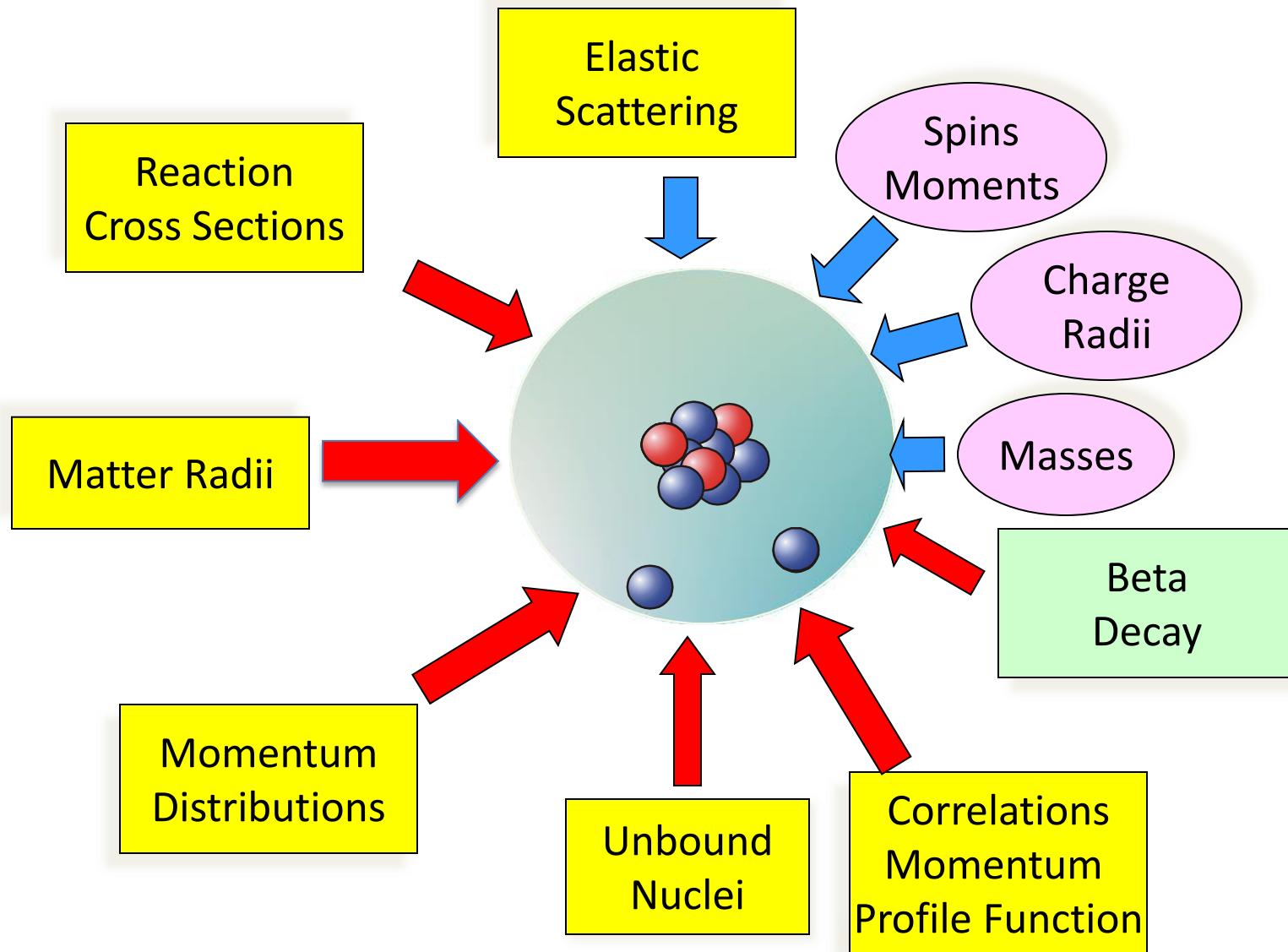
5H  
unbound

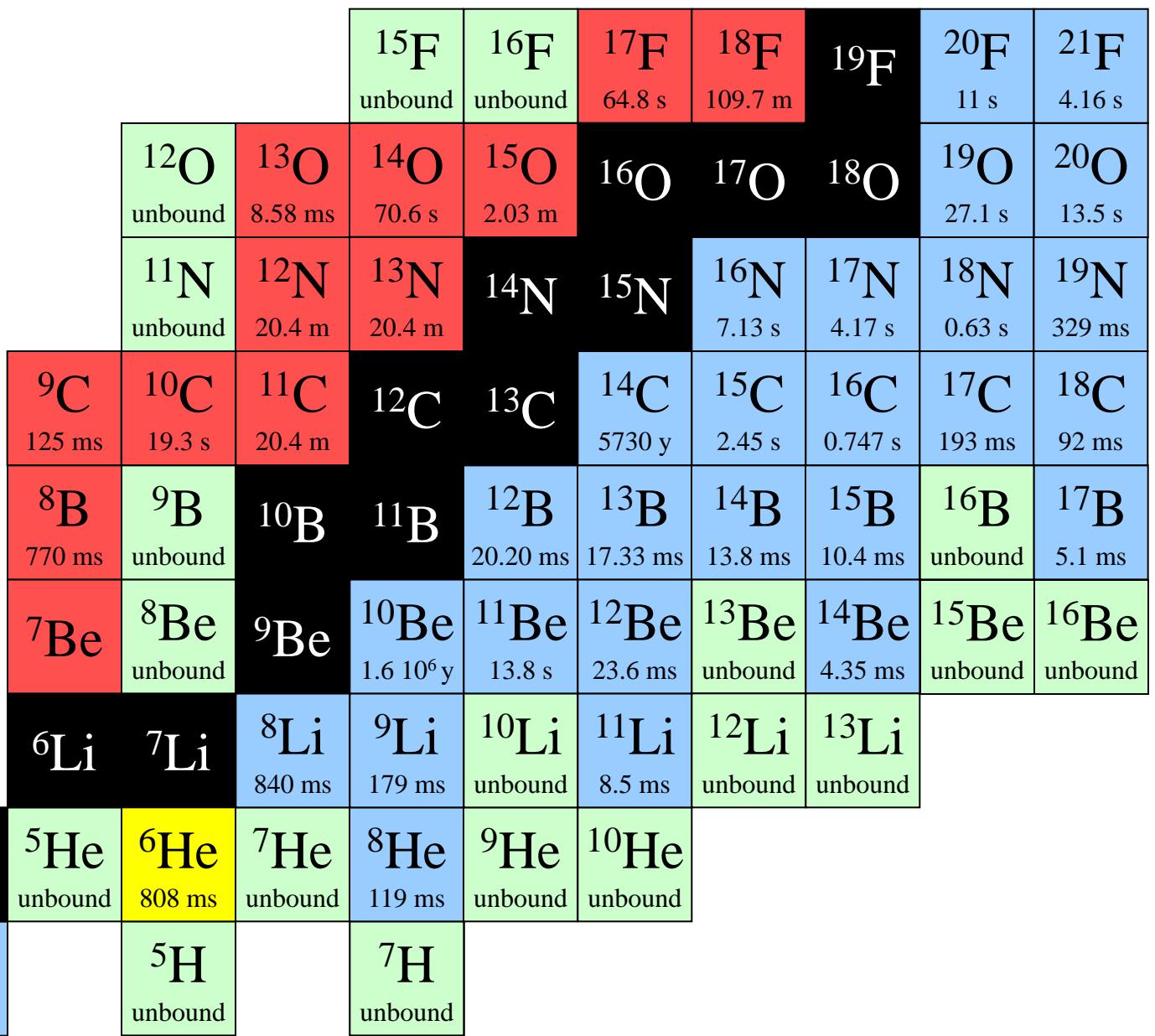
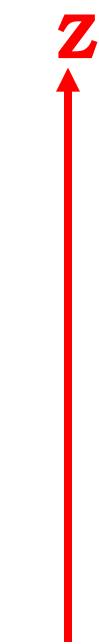
7H  
unbound

n  
10.25 m

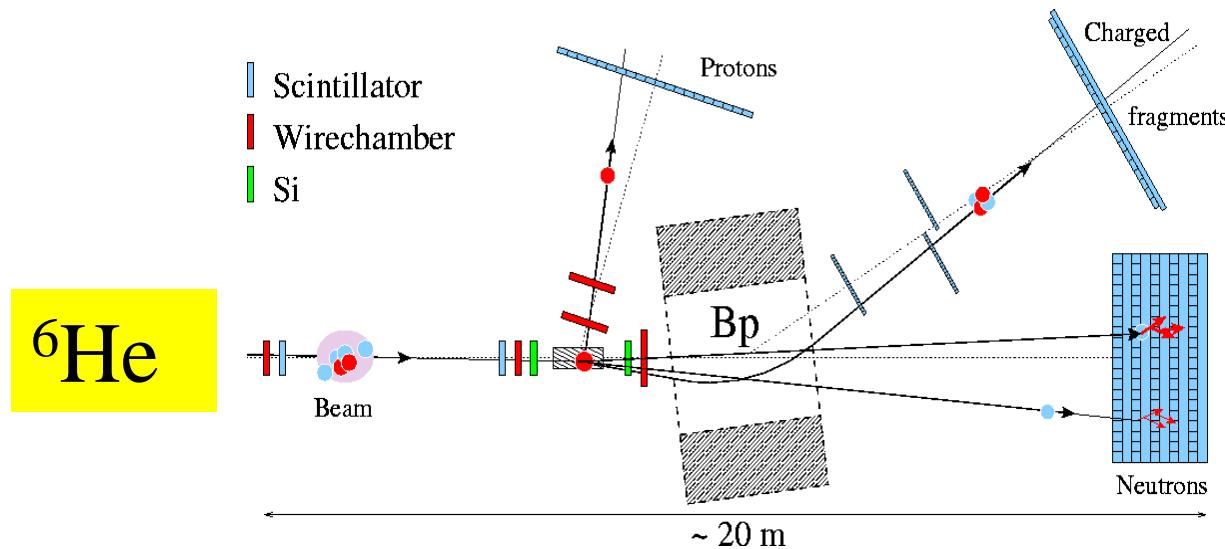
N

# Experimental Studies of Dripline Nuclei





n  
10.25 m



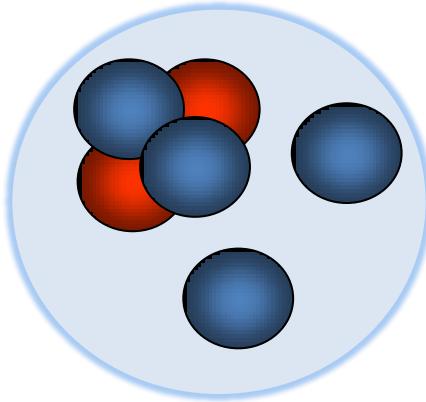
$$\mathbf{p}_{n\alpha} = \frac{m_n m_\alpha}{m_n + m_\alpha} \left[ \frac{\mathbf{p}_{n_2}}{m_n} - \frac{\mathbf{p}_\alpha}{m_\alpha} \right]$$

$$\mathbf{p}_{n(n\alpha)} = \frac{m_n(m_\alpha + m_n)}{2m_n + m_\alpha} \left[ \frac{\mathbf{p}_{n_1}}{m_n} - \frac{(\mathbf{p}_{n_2} + \mathbf{p}_\alpha)}{m_n + m_\alpha} \right]$$

$$\mathbf{p}_{cm} = \mathbf{p}_{n_1} + \mathbf{p}_{n_2} + \mathbf{p}_\alpha \quad (\mathbf{p}_{cm} = 0) \quad \mathbf{p}_{n_1} = -(\mathbf{p}_{n_2} + \mathbf{p}_\alpha)$$

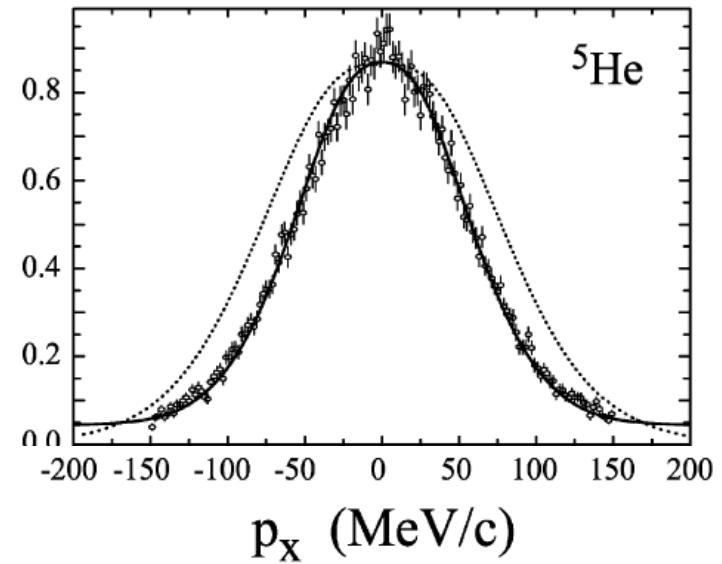
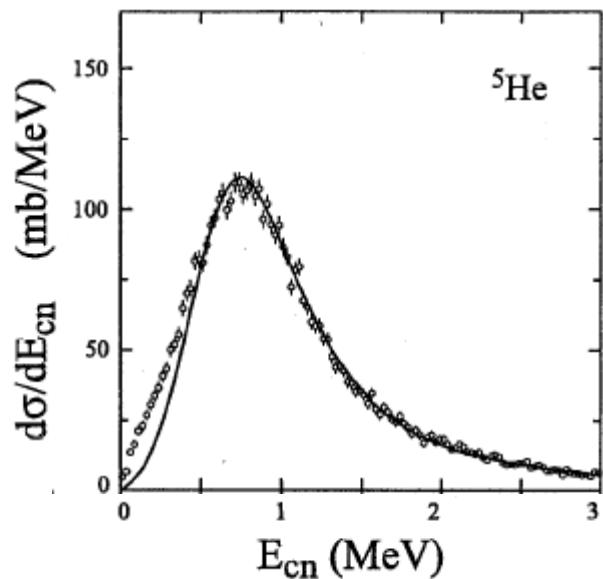
$$E_{fn} = \frac{m_f + m_n}{2m_f m_n} \mathbf{p}_{fn}^2$$

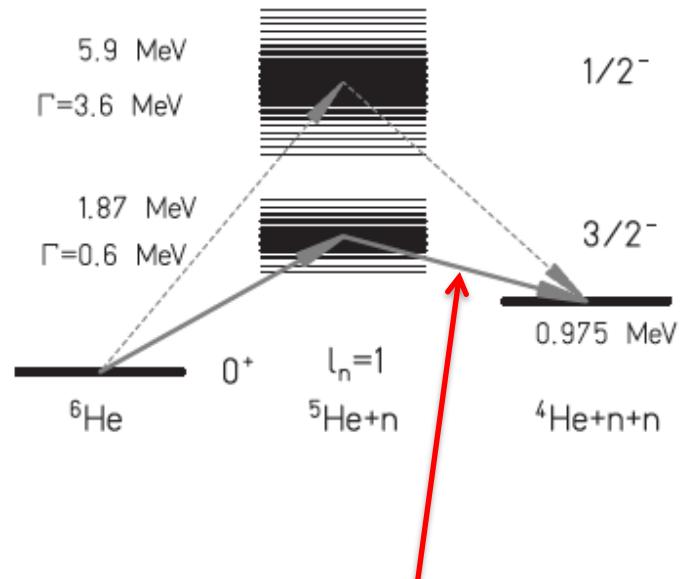
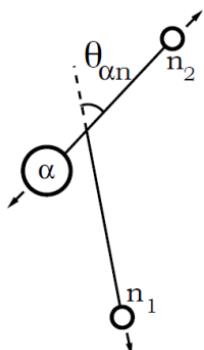
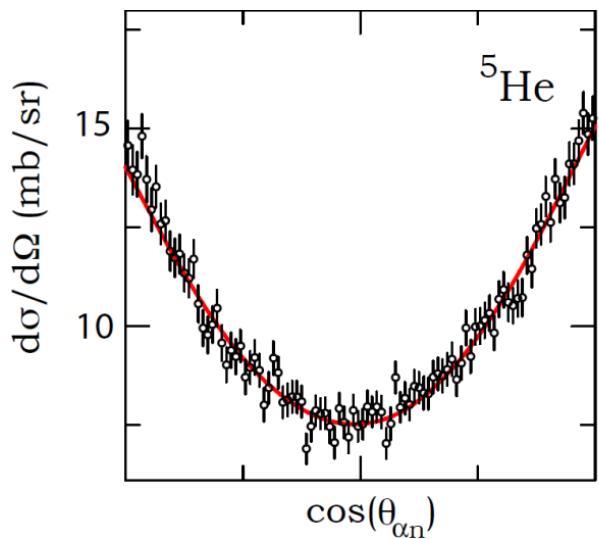
${}^3\text{He}$	${}^4\text{He}$	${}^5\text{He}$	${}^6\text{He}$ 808 ms	${}^7\text{He}$	${}^8\text{He}$ 119 ms	${}^9\text{He}$	${}^{10}\text{He}$
		unbound		unbound		unbound	unbound



$${}^6\text{He} : {}^4\text{He} \ddot{\wedge} (\eta p_{3/2})^2$$

$$\mathbf{p}_{n_1} = -(\mathbf{p}_{n_2} + \mathbf{p}_\alpha)$$

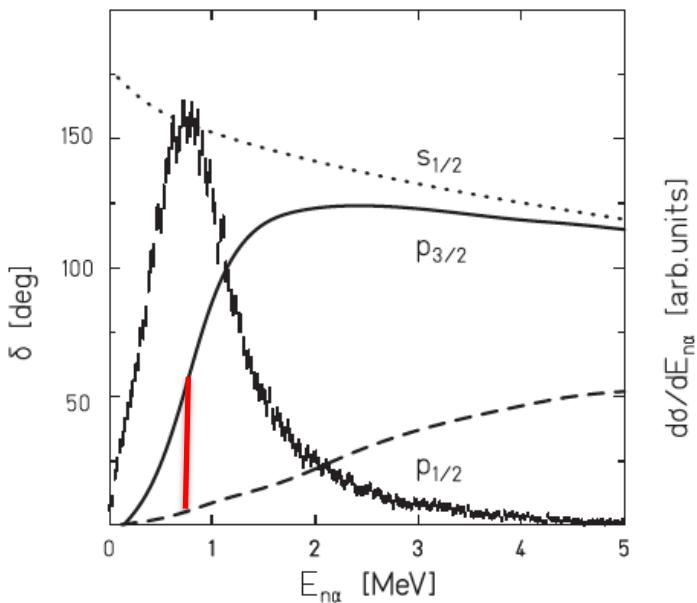




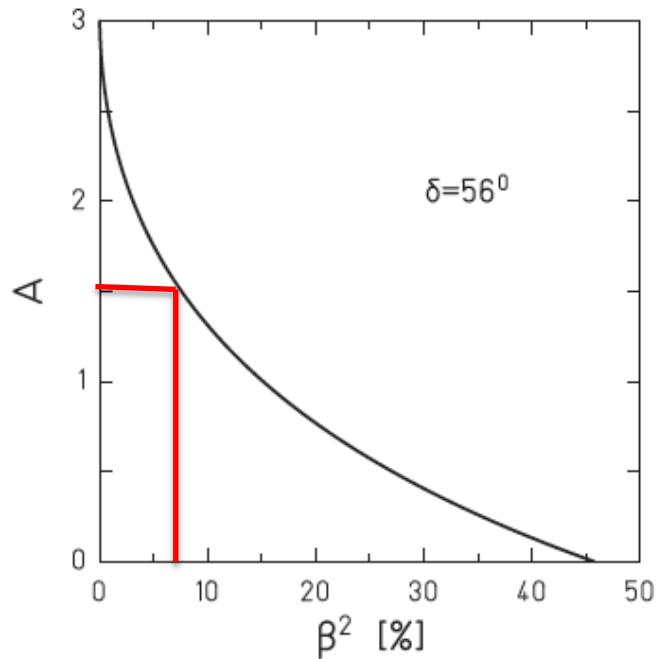
L.V. Chulkov *et al.*, PRL **79** (1997) 201

$$W(\theta_{\alpha n}) = 1 + 1.5(3)\cos^2(\theta_{\alpha n})$$

$$W(\theta_{\alpha n}) = 1 + 3\cos^2(\theta_{\alpha n})$$



$d\sigma/dE_{n\alpha}$  [arb.units]



Morgan, G.L., Walter, R.L.: Phys. Rev. **168**, 1114 (1968)

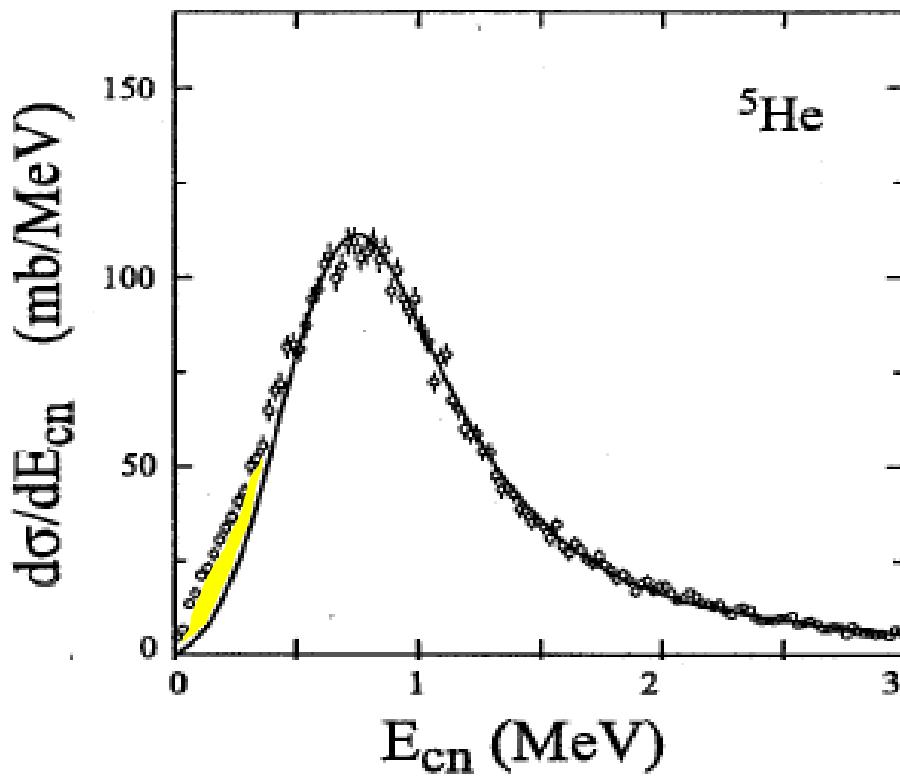
$$\frac{d\sigma}{d \cos(\vartheta_{n\alpha})} \sim \sum_{\nu_1, \nu_2} \left| \alpha f_{\nu_1, \nu_2}(3/2) + \beta e^{i\delta} f_{\nu_1, \nu_2}(1/2) \right|^2$$

$$\alpha^2 + \beta^2 = 1$$

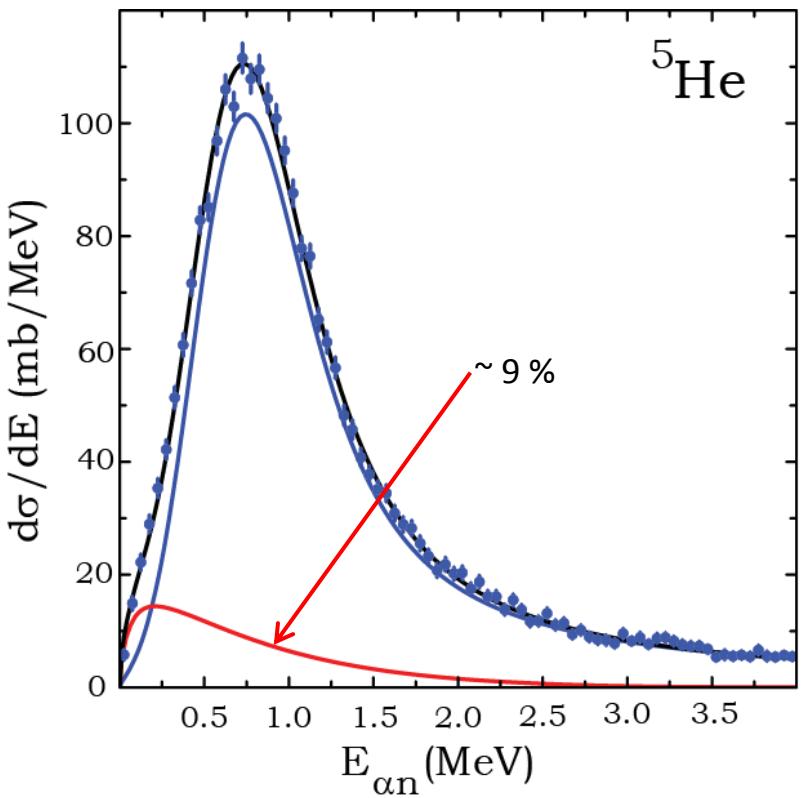
$$W(\theta_{\alpha n}) = 1 + 1.5(3)\cos^2(\theta_{\alpha n})$$

$$(p_{1/2})^2 \sim 7\%$$

L.V. Chulkov, G. Schrieder, Z. Phys **A359** (97) 231



$$d\sigma / dE_{fn} \propto \frac{\Gamma_l(E_{fn})}{[E_r + \Delta_l(E_{fn}) - E_{fn}]^2 + \frac{1}{4} \Gamma_l^2(E_{fn})}$$



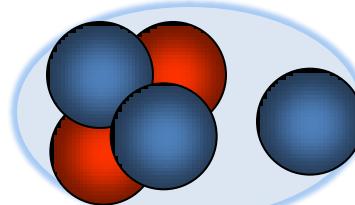
Yu. Aksyutina et al., PLB **679** (2009) 191

$$k = \sqrt{2\mu\varepsilon}$$

$$\varepsilon = 0.94(23) \text{ MeV}$$

$$a = 3.37(38) \text{ fm}$$

$$\varepsilon \approx S_{2n}$$



$$d\sigma / dE_{fn} \propto \frac{\Gamma_l(E_{fn})}{[E_r + \Delta_l(E_{fn}) - E_{fn}]^2 + \frac{1}{4} \Gamma_l^2(E_{fn})}$$

s-wave: Effective-range approximation

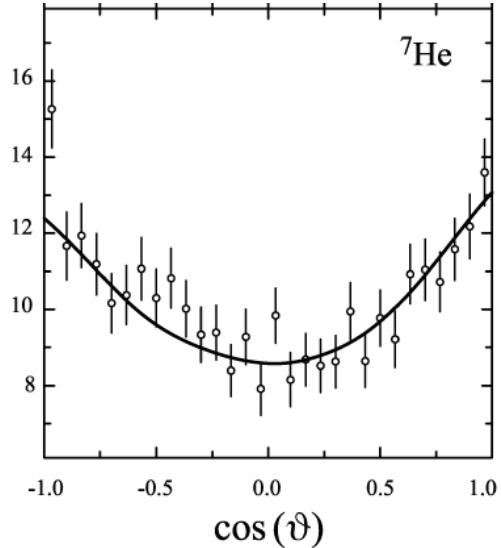
$$\frac{d\sigma}{dE_{fn}} \propto p_{fn} \left[ \frac{1}{k^2 + p_{fn}^{-2}} \right]^2 \left[ \cos(\delta) + \frac{k}{p_{fn}} \sin(\delta) \right]^2$$

$$p_{fn} \cot(\delta) = -\frac{1}{a} + \frac{1}{2} r_0 p_{fn}^{-2} + O(p_{fn}^{-4})$$

Bertch *et al.*, PRC **57** (1998) 1366

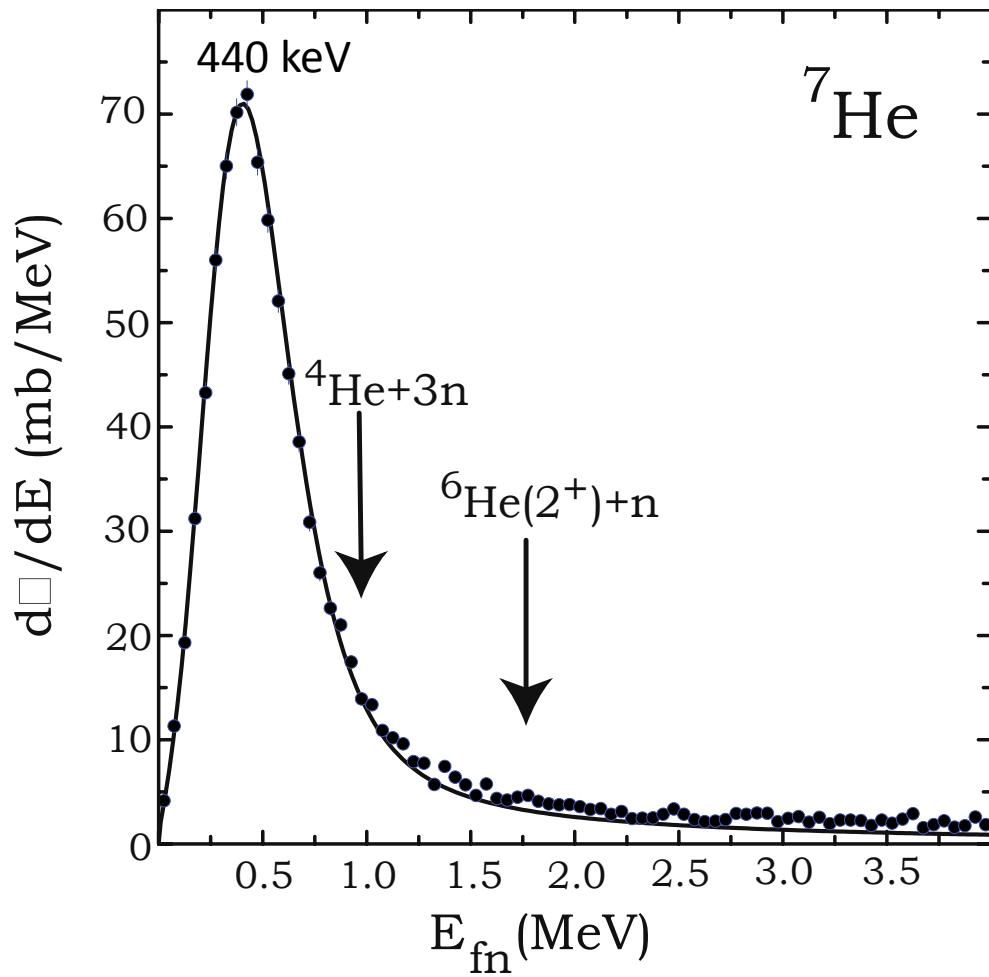
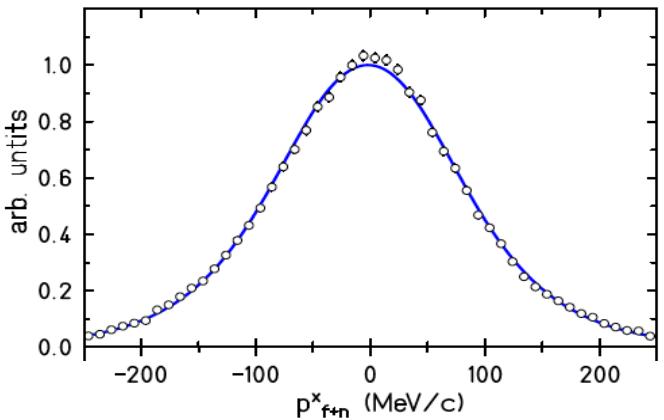
$$\begin{aligned} (p_{3/2})^2 &\sim 86 \% \\ (p_{1/2})^2 &\sim 5 \% \\ (s_{1/2})^2 &\sim 7 \% \end{aligned}$$

B.V. Danilin et al., NPA **632** (98) 383



K. Markenroth *et al.*, NP A**679** (2001) 462

$$W(\vartheta_{^6\text{He}-n}) \sim 1 + 0.7(1) \cos^2(\vartheta_{^6\text{He}-n})$$



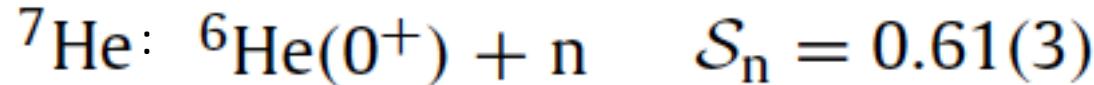
Yu. Aksyutina *et al.*,  
PLB **679** (2009) 191

$$\frac{d\sigma}{dE_{\text{fn}}} \propto \frac{\Gamma_l(E_{\text{fn}})}{[E_r + \Delta_l(E_{\text{fn}}) - E_{\text{fn}}]^2 + \frac{1}{4}\Gamma_l^2(E_{\text{fn}})}$$

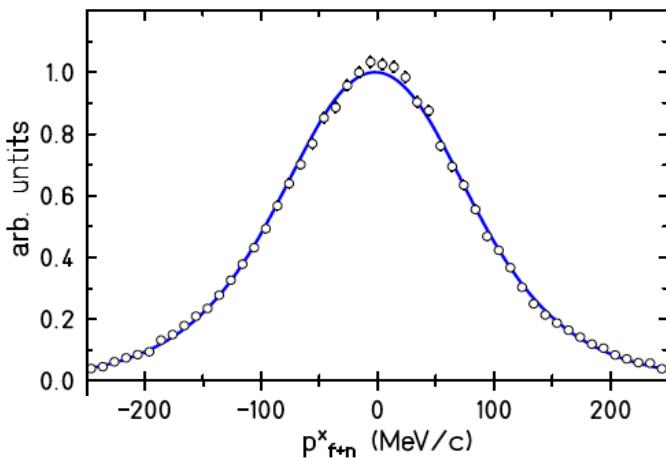
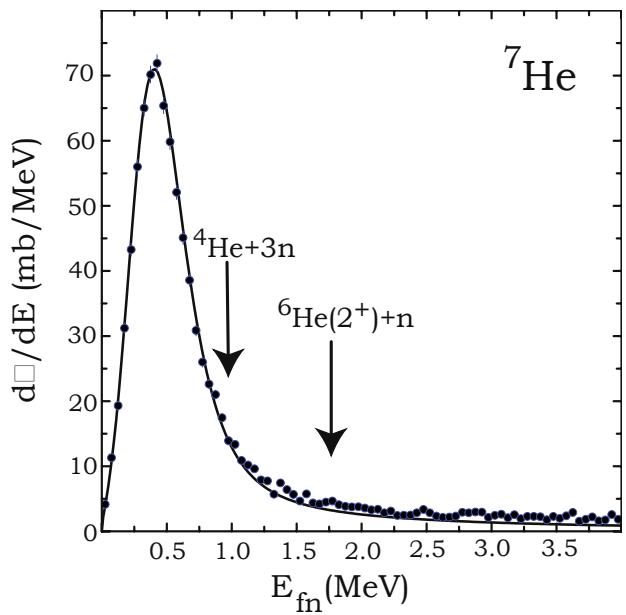
$$\begin{aligned}\Gamma_l &= 2P_l(E_{\text{fn}})\gamma^2 \\ \gamma_{\text{sp}}^2 &= \frac{2}{3}\hbar^2/\mu R^2 \\ S_n &= \gamma_{\text{obs}}^2/\gamma_{\text{sp}}^2\end{aligned}$$

Yu. Aksyutina *et al.*,  
PLB **679** (2009) 191

$${}^5\text{He} \quad \gamma_{\text{obs}}^2 = \gamma_{\text{sp}}^2$$



→ not a pure single-particle  $p_{3/2}$ -state



## Momentum Profile Function

$$P(E_{fn}) = \sqrt{\left\langle (p_f^x + p_n^x)^2 \right\rangle - \left\langle (p_f^x + p_n^x) \right\rangle^2}$$

## Transverse Momentum Distribution

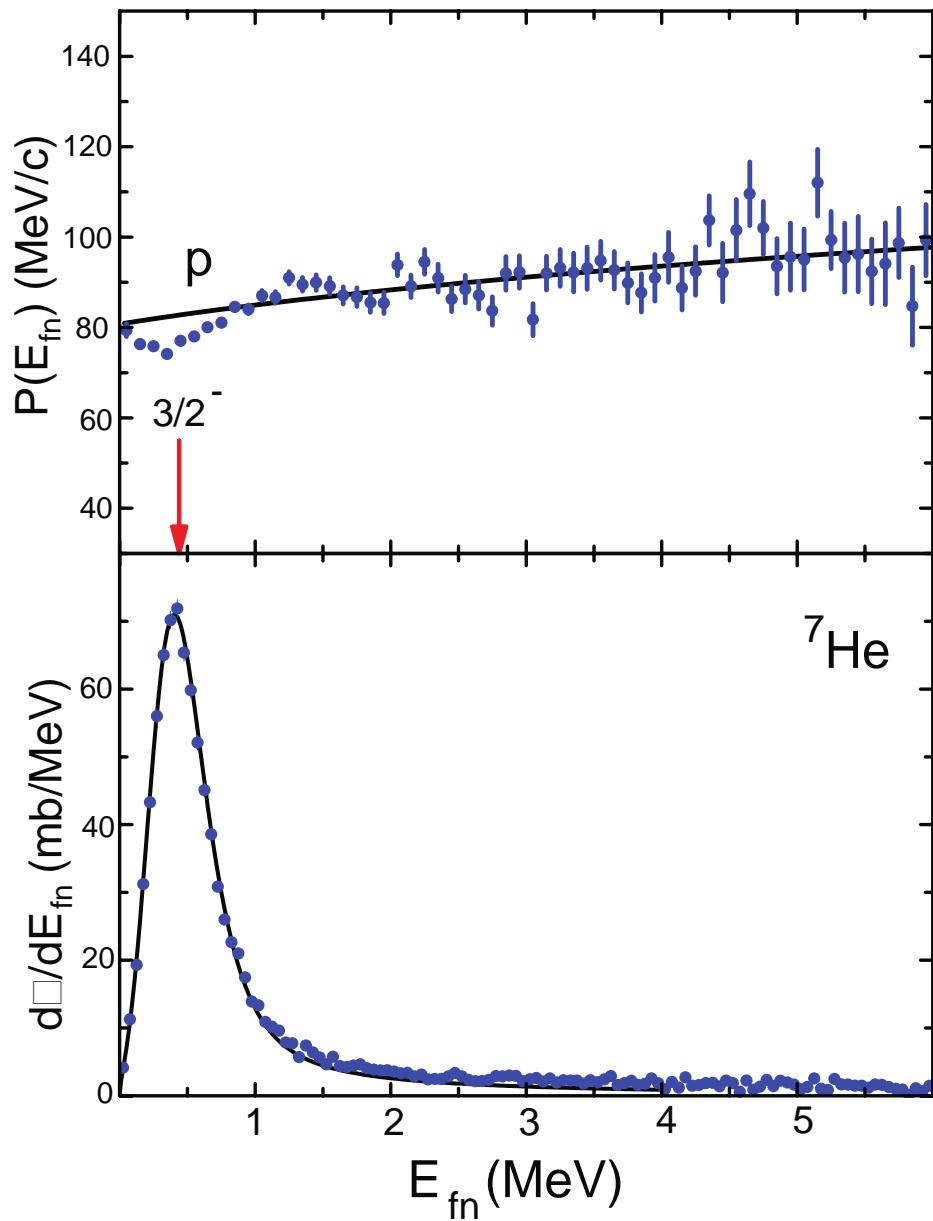
$$\frac{d\sigma_0}{dk_z} = \frac{\sigma_T B^2 \kappa b_{\min}^2}{2\pi} [K_1^2 - K_0^2]$$

$$\begin{aligned} \frac{d\sigma_1}{dk_z} = & \frac{\sigma_T B^2 b_{\min}^2}{2\pi \kappa} \left[ k_z^2 (K_1^2 - K_0^2) + (k_z^2 + \kappa^2) \right. \\ & \left. \times \left( K_2^2 - K_1^2 - \frac{2}{\xi} K_1 K_2 \right) \right] \end{aligned}$$

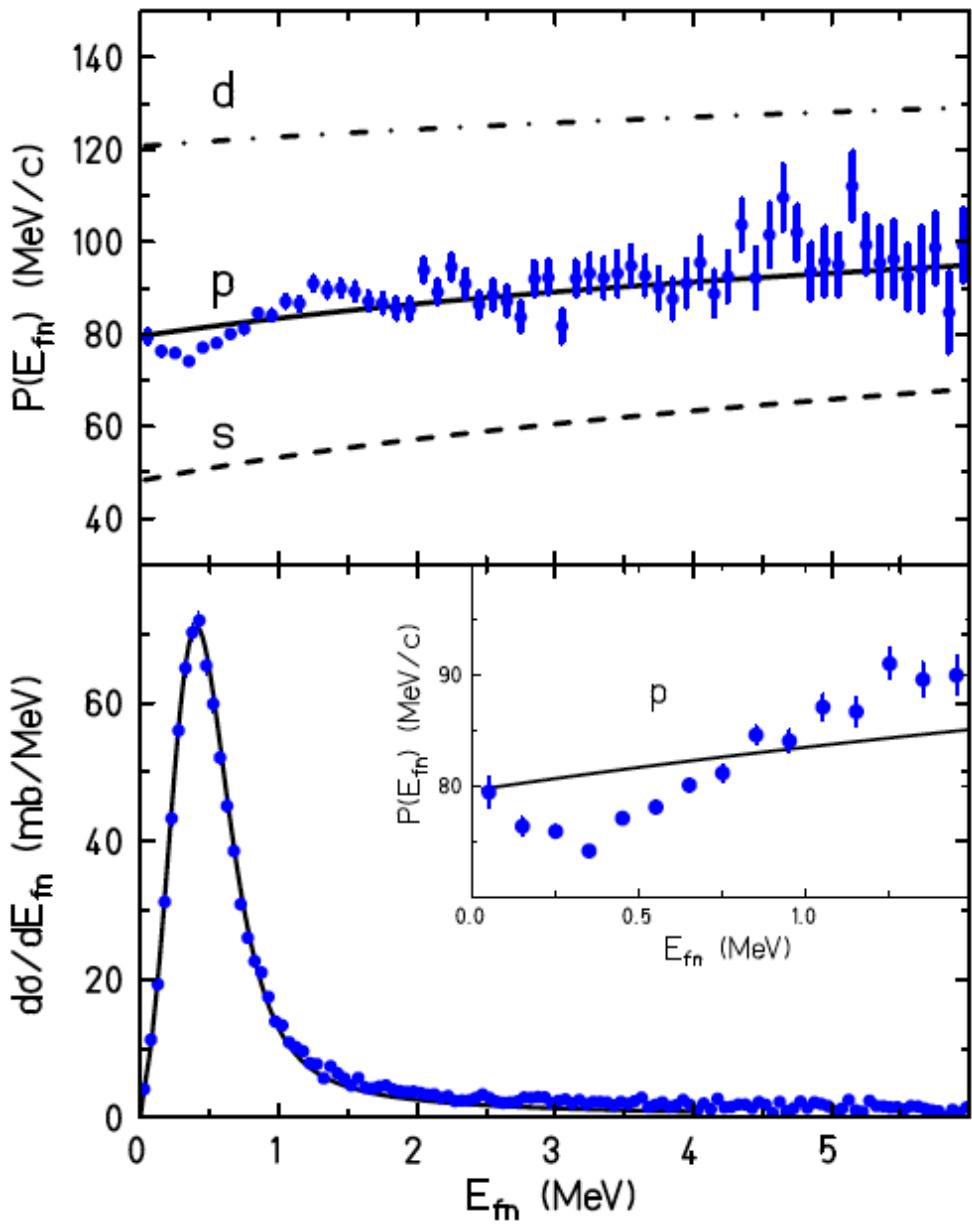
$$\xi = b_{\min} \sqrt{\kappa^2 + k_z^2}$$

$$\kappa = (2\mu S_n)^{1/2}/\hbar$$

P.G. Hansen, PRL **77** (1996) 1016  
D. Basin et al., Phys. Rev. C **57** (1998) 2156



Yu. Aksyutina *et al.*,  
PLB **718** (2013) 1309



$$P_{r.e.}(E_{fn}) = \sqrt{\alpha_s \sigma_s^2 + (1 - \alpha_s) \sigma_p^2}$$

$\uparrow \qquad \qquad \qquad \uparrow$

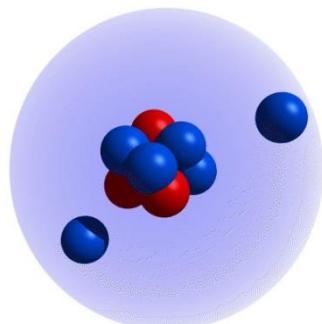
$l = 0 \qquad \qquad l = 1$

↓

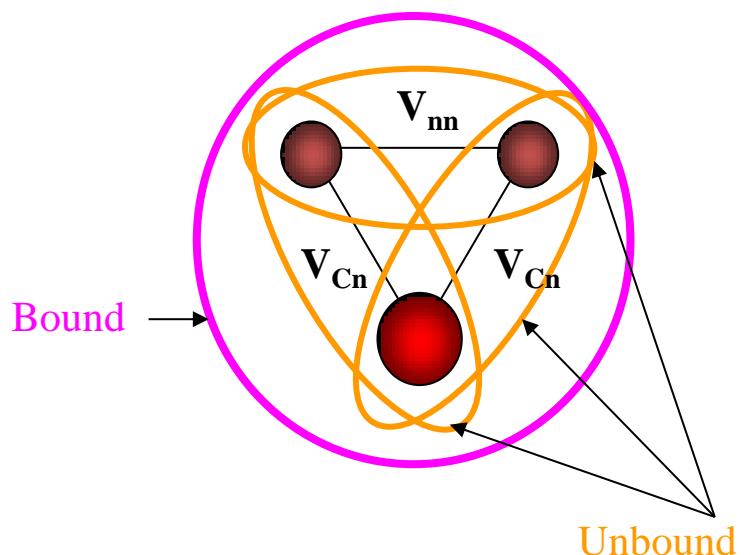
$12\% \ (1s_{1/2})^2$

		$^{11}\text{Be}$ 13.8 s	$^{12}\text{Be}$ 23.6 ms	$^{13}\text{Be}$	$^{14}\text{Be}$ 4.35 ms
	$^9\text{Li}$ 179 ms	$^{10}\text{Li}$	$^{11}\text{Li}$ 8.5 ms		
$^6\text{He}$ 806 ms	$^7\text{He}$	$^8\text{He}$ 119 ms	$^9\text{He}$		

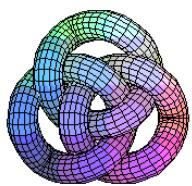
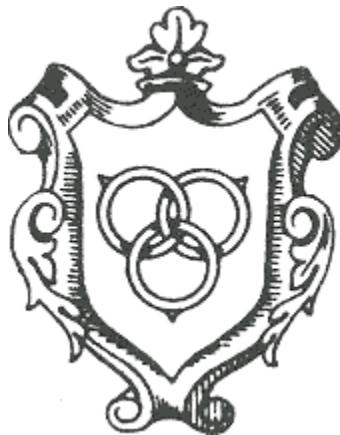
**$^{11}\text{Li}$**   
 $S_n = 369 \text{ keV}$



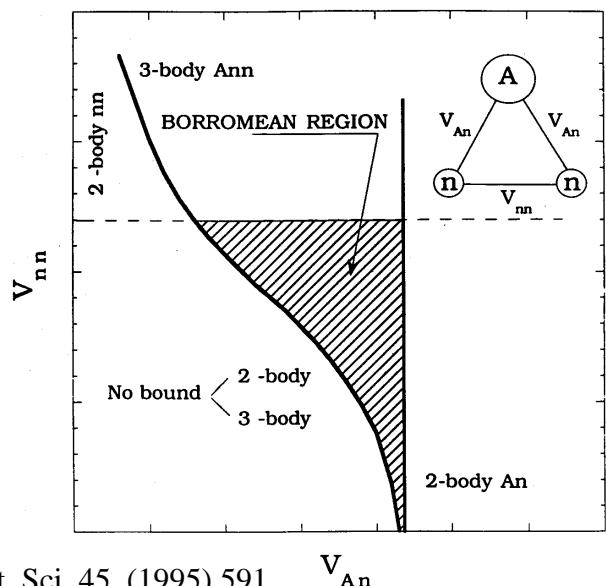
**N=8**



# BORROMEAN

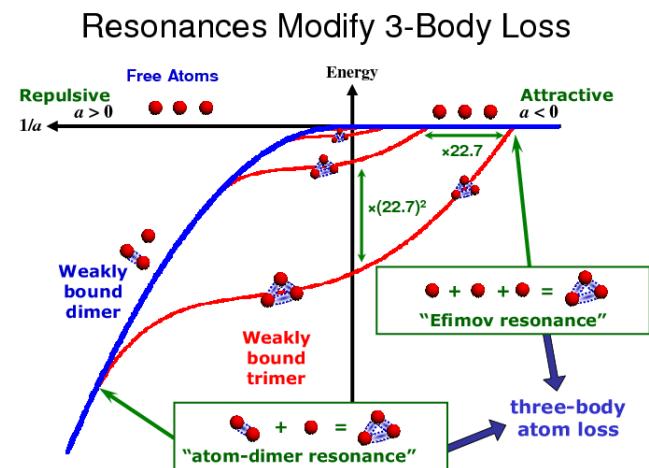


M.V. Zhukov *et al.* Phys. Rep. 231 (1993) 151



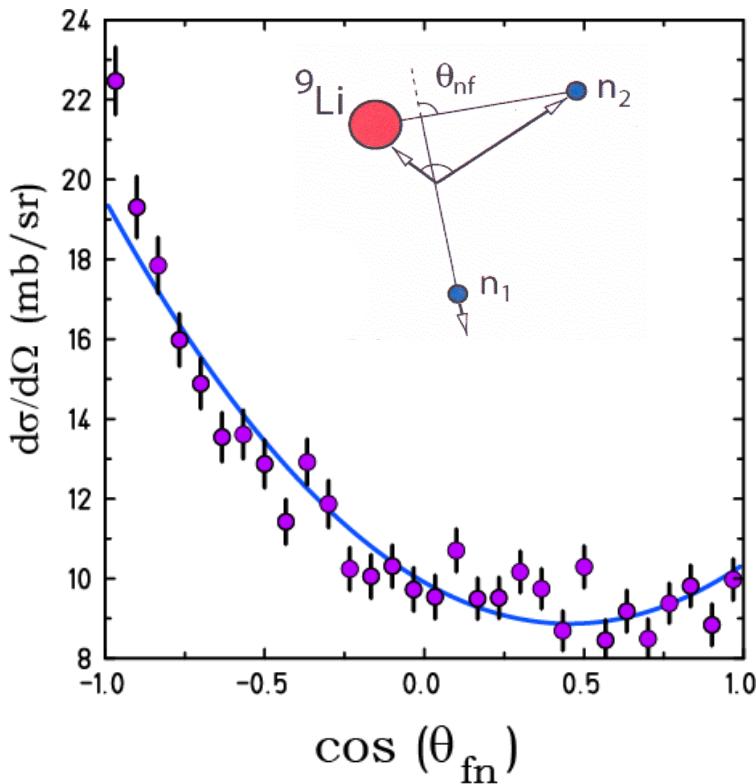
P.G. Hansen *et al.*

Ann.Rev.Nucl.Part. Sci. 45 (1995) 591



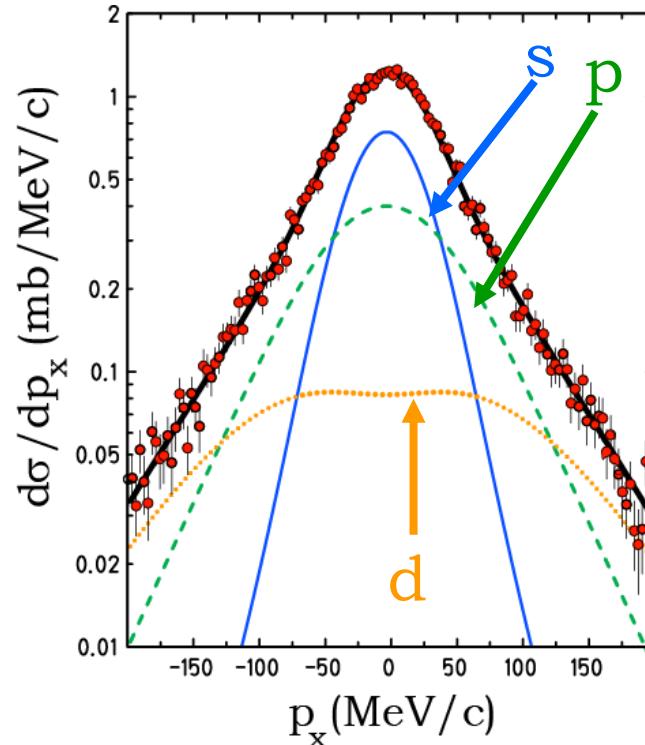
S.E. Pollack *et al.* Science 326 (2009) 1683

# GSI, 287 MeV/u $^{11}\text{Li}$

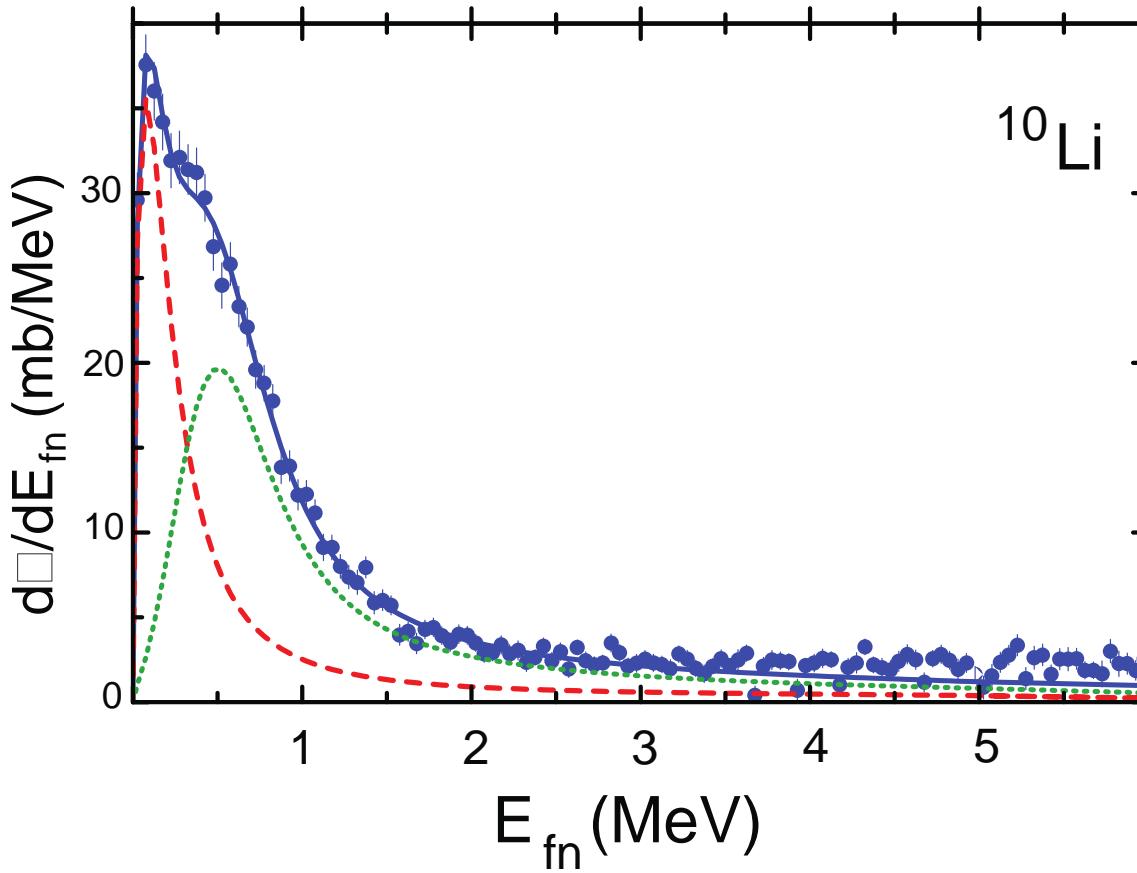


$$W(\theta_{nf}) = 1 - 1.03 \cos(\theta_{nf}) + 1.41 \cos^2(\theta_{nf})$$

$$(1s_{1/2})^2 \quad (0p_{1/2})^2$$

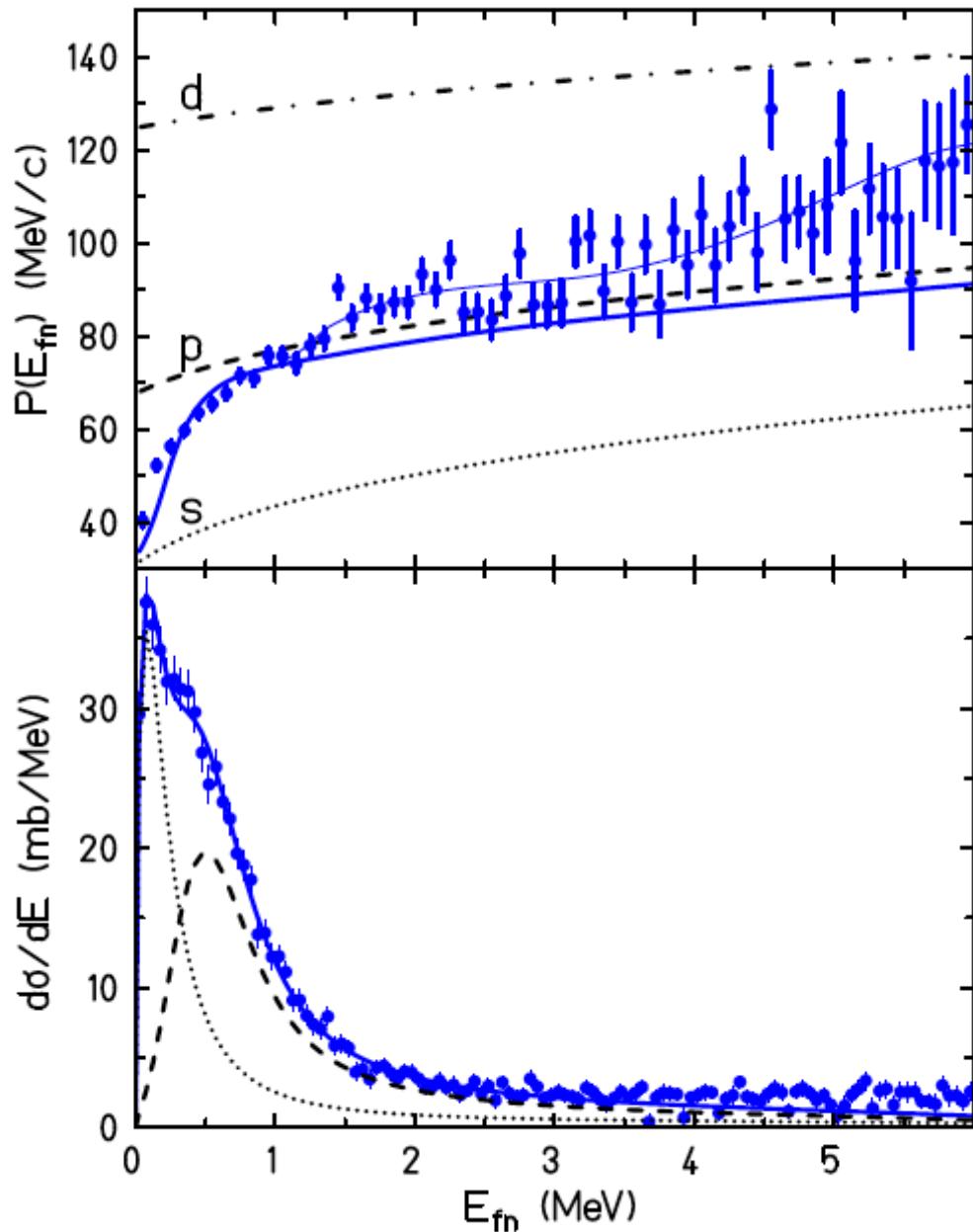


$$\mathbf{p}_{n1} = -(\mathbf{p}_{n2} + \mathbf{p}_{^{9}\text{Li}})$$



Momentum Profile Function

$$P(E_{fn}) = \sqrt{\left\langle (p_f^x + p_n^x)^2 \right\rangle - \left\langle (p_f^x + p_n^x) \right\rangle^2}$$



$$P_{r.e.}(E_{fn}) = \sqrt{\alpha_s \sigma_s^2 + (1 - \alpha_s) \sigma_p^2}$$

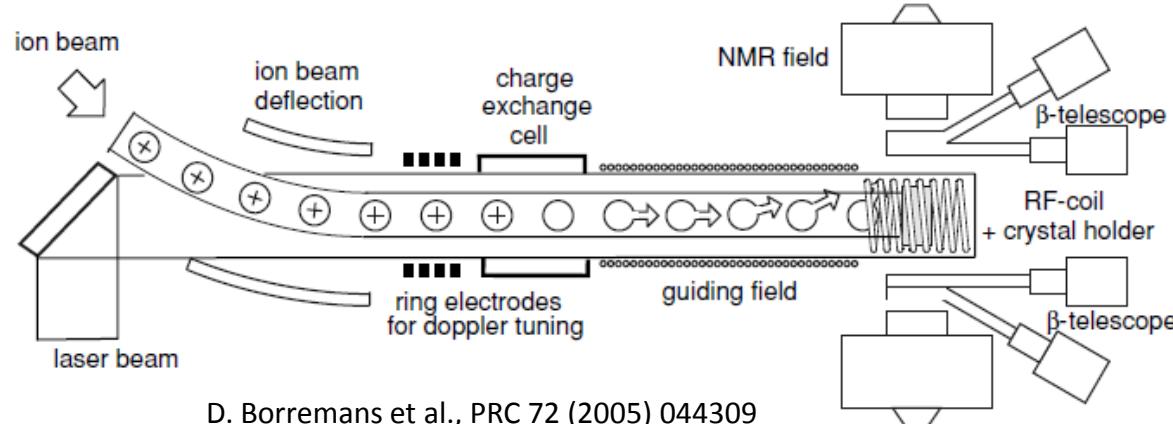
$$\alpha_d = (P_{exp}^2 - P_{r.e.}^2) / (\sigma_d^2 - P_{r.e.}^2)$$

(0d<sub>5/2</sub>)<sup>2</sup>  
11(2) %

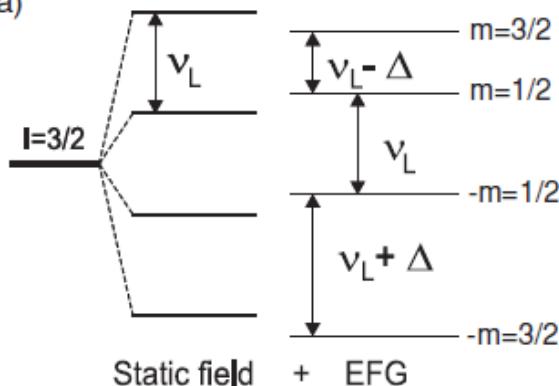
Aksyutina *et al.*  
PLB **718** (2013) 1309

${}^6\text{Li}$	${}^7\text{Li}$	${}^8\text{Li}$	${}^9\text{Li}$	${}^{10}\text{Li}$	${}^{11}\text{Li}$	${}^{12}\text{Li}$	${}^{13}\text{Li}$
		840 ms	179 ms	unbound	8.5 ms	unbound	unbound

$|\pi=3/2^-$        $|\pi=3/2^-$



(a)



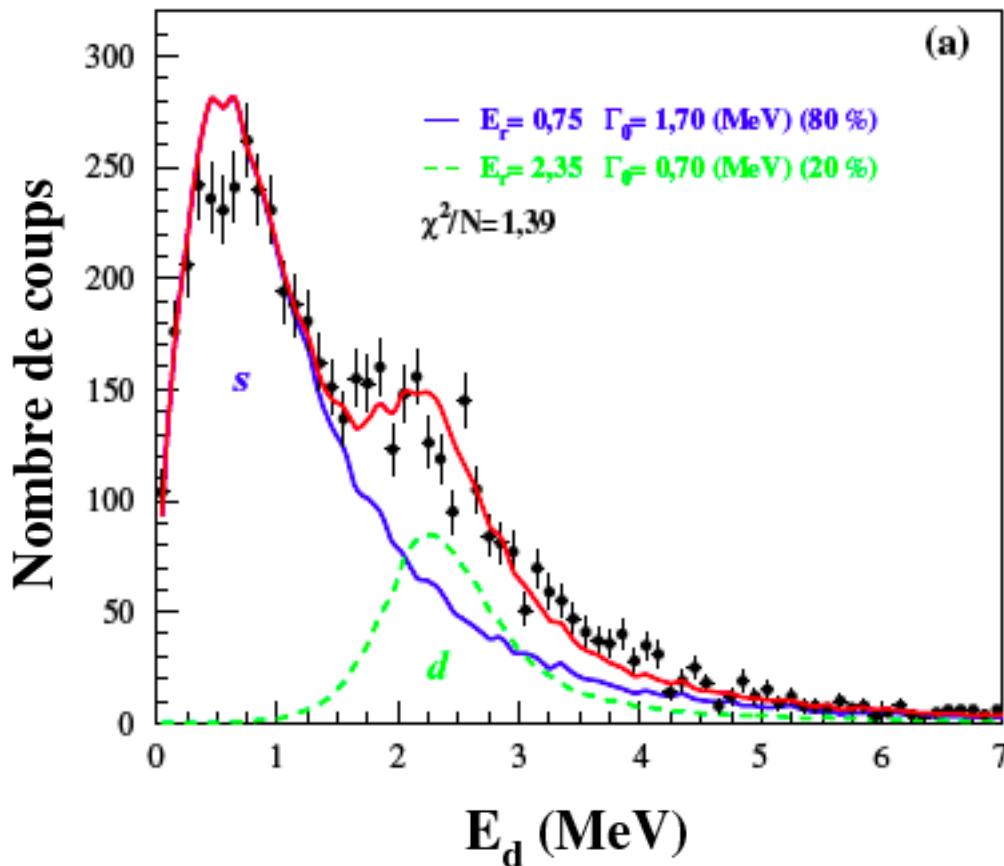
R. Neugart et al., PRL 101(2008) 132502

$|Q({}^{11}\text{Li})/Q({}^9\text{Li})| = 1.088(15)$

$^7\text{Be}$	$^8\text{Be}$	$^9\text{Be}$	$^{10}\text{Be}$	$^{11}\text{Be}$	$^{12}\text{Be}$	$^{13}\text{Be}$	$^{14}\text{Be}$	$^{15}\text{Be}$	$^{16}\text{Be}$
unbound			1.6 $10^6$ y	13.8 s	23.6 ms	unbound	4.35 ms	unbound	unbound

# GANIL

41 MeV/u  $^{14}\text{B}$ , C target



$^{16}\text{Be}$   
Spyrou *et al.*,  
PRL 108(2012) 102501

$^{15}\text{Be}$   
Snyder *et al.*,  
PRC 88(2013) 031303

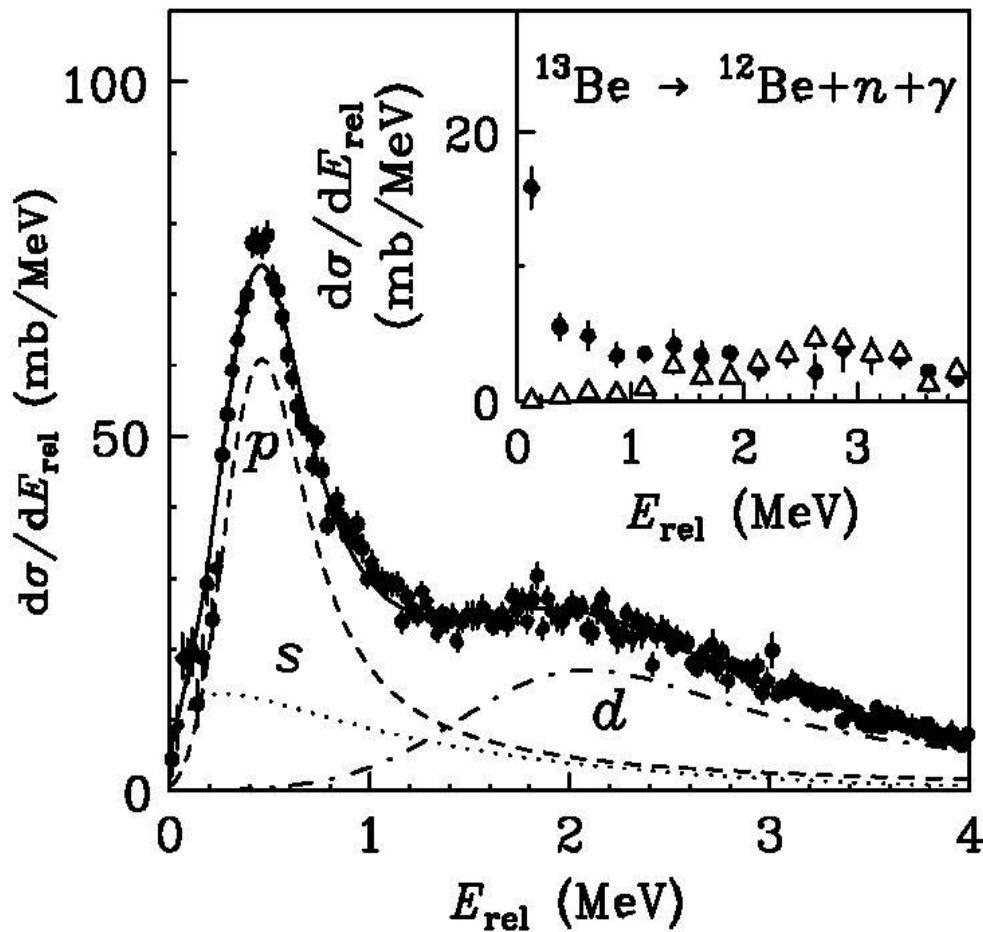
$^{14}\text{Be}$	$^{15}\text{Be}$	$^{16}\text{Be}$
4.35 ms	unbound	unbound

$^{17}\text{B}$   
5.1 ms

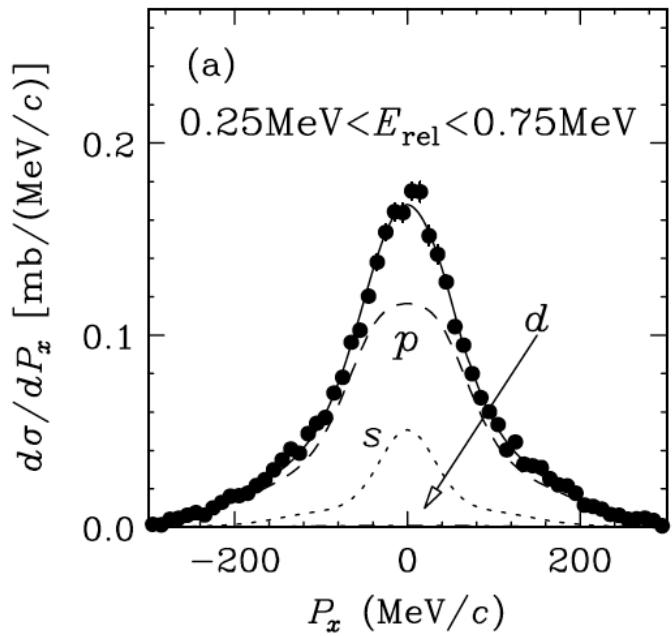
$^{14}\text{B}$   
 $^{13}\text{B}(3/2^-) \otimes (sd)$

**RIKEN**

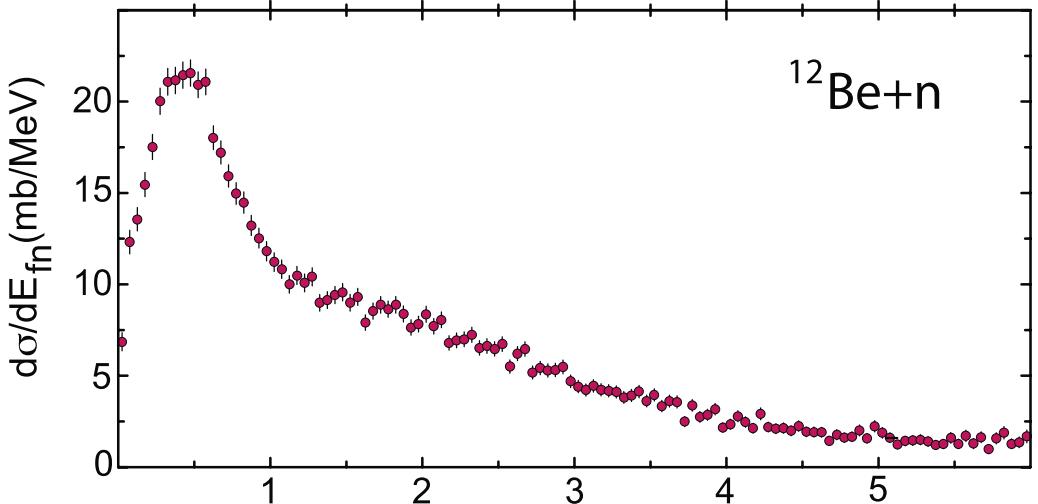
69 MeV/u  $^{14}\text{Be}$ , Lq H target



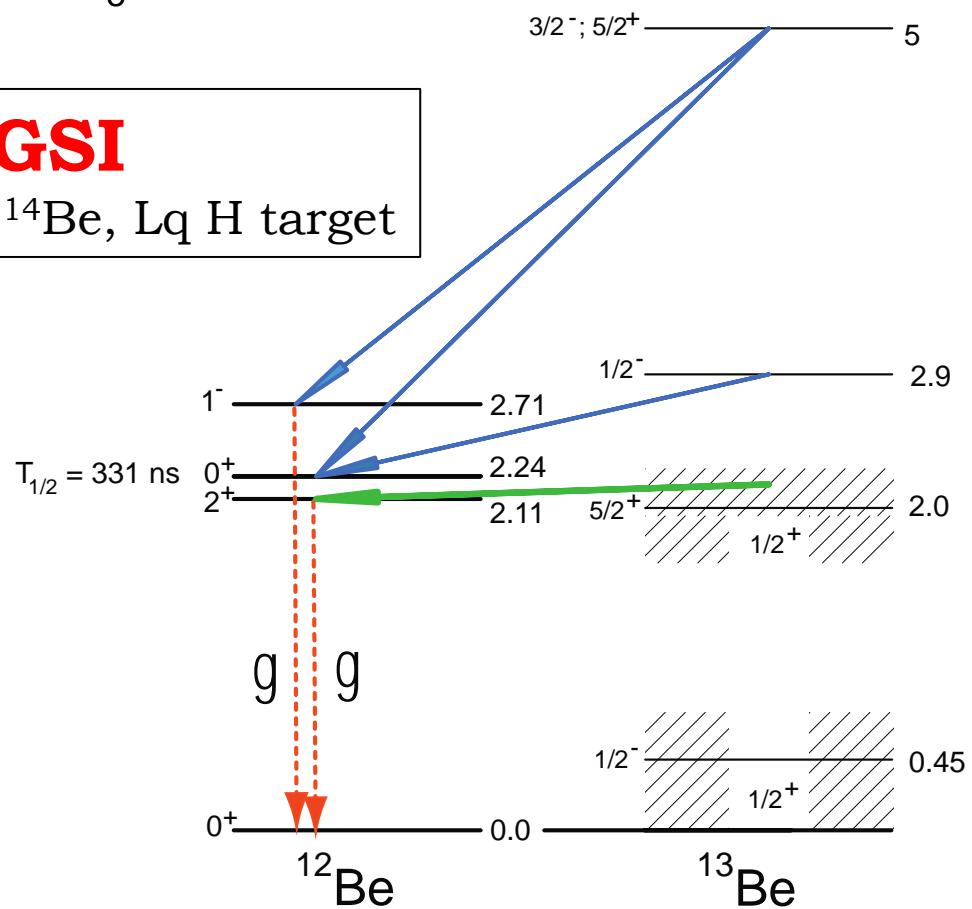
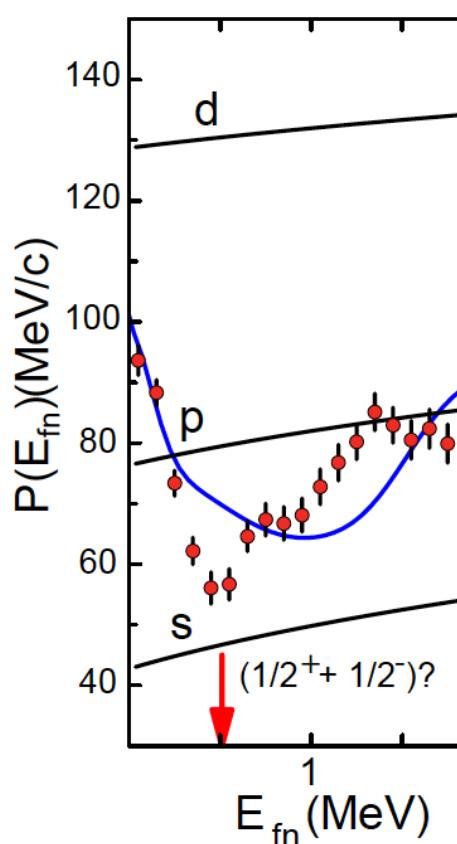
Y. Kondo *et al.*, PLB **690** (2010) 245

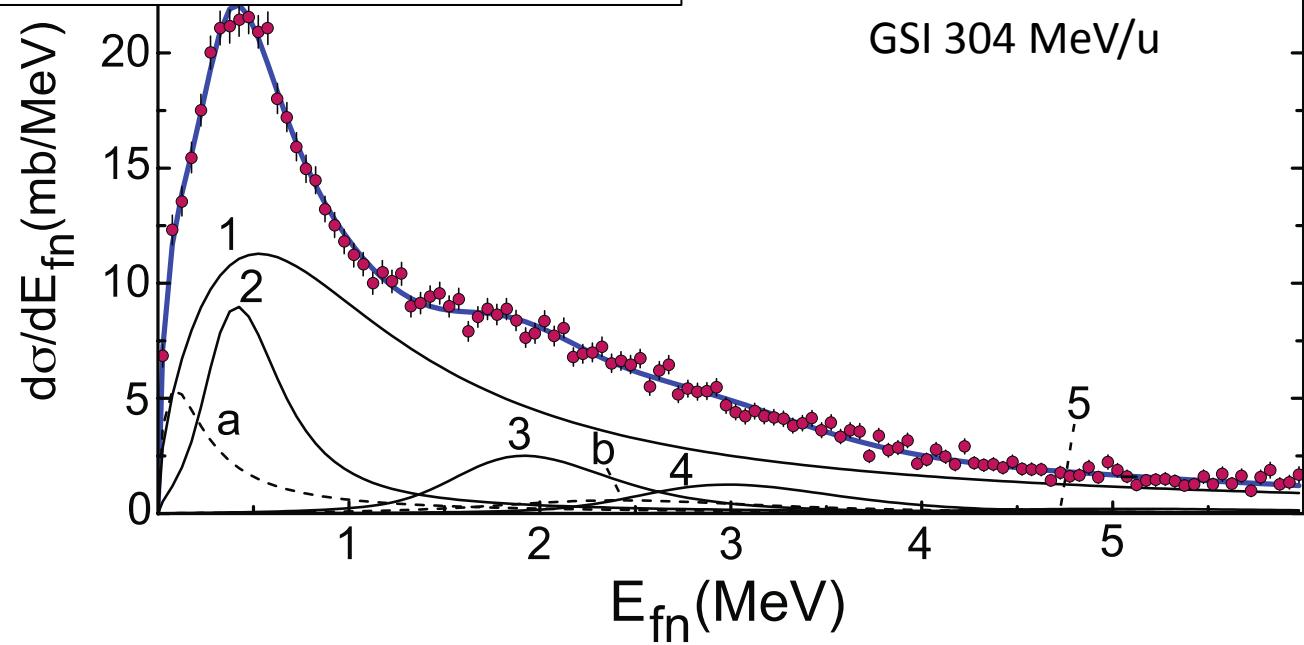
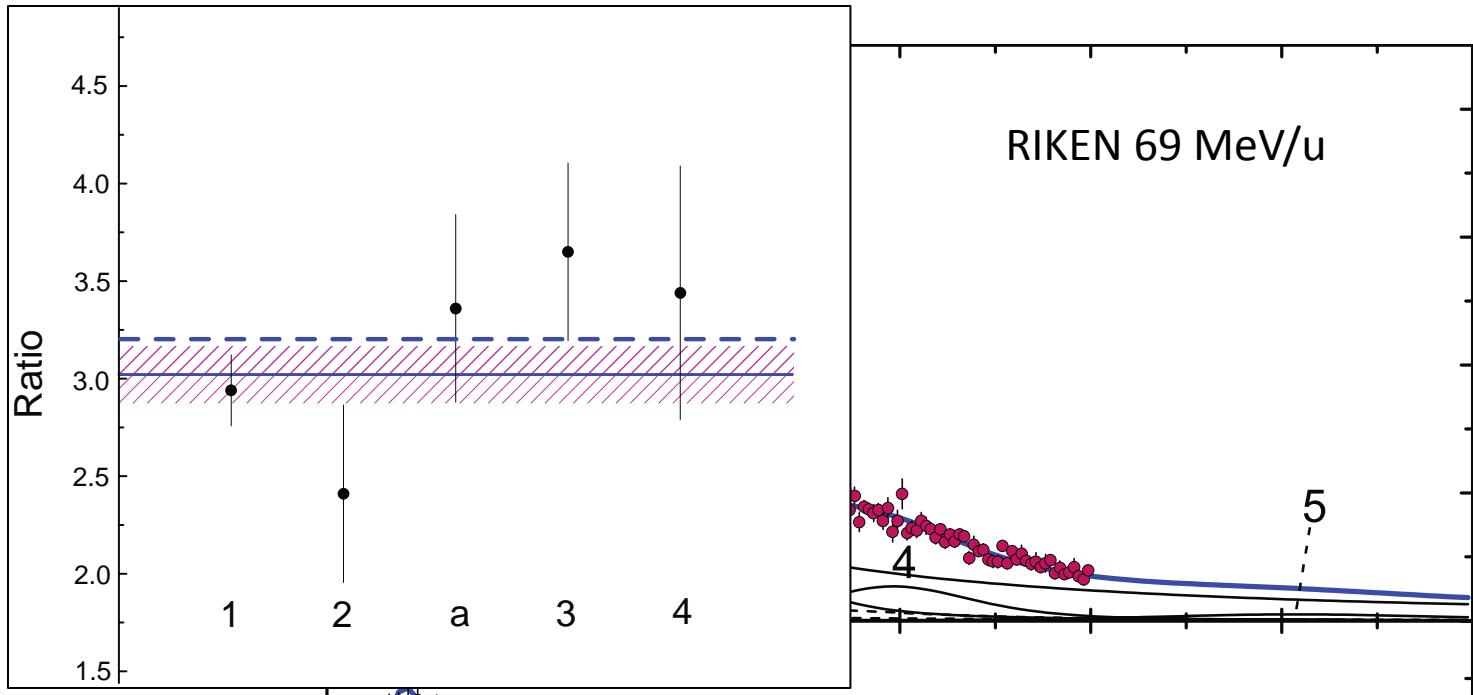


*n - g coincidences*  
2.11 MeV and 2.71 MeV



Yu. Aksyutina *et al.*,  
Phys. Rev. C **87** (2013) 064316





$$\frac{S_{69}}{S_{304}} = 3.2$$

# INTERFERENCE ?

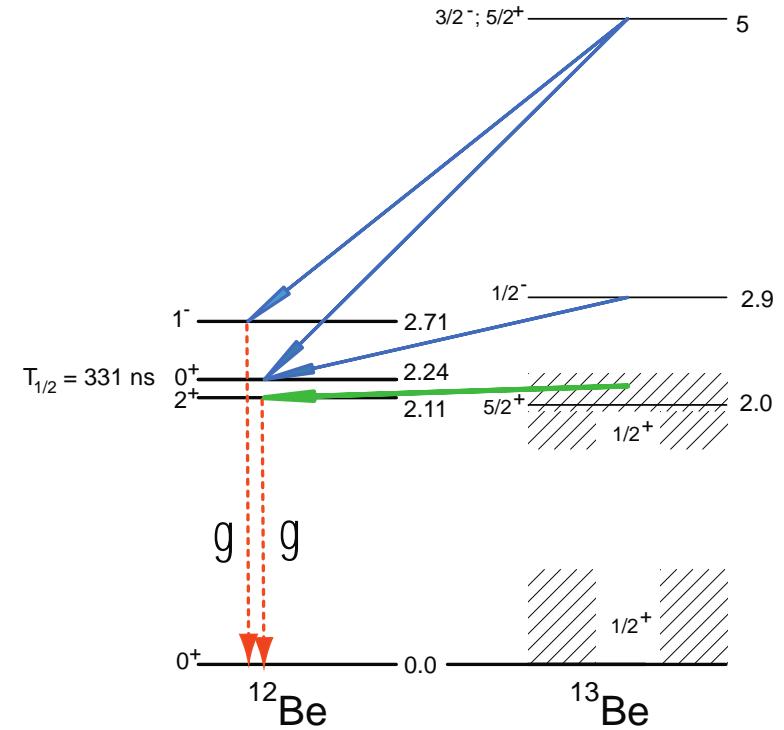
$$^{12}\text{Be}(0^+) = \alpha [{}^{10}\text{Be} \otimes (1s_{1/2})^2] + \beta [{}^{10}\text{Be} \otimes (0p_{1/2})^2] \\ + \gamma [{}^{10}\text{Be} \otimes (0d_{5/2})^2],$$

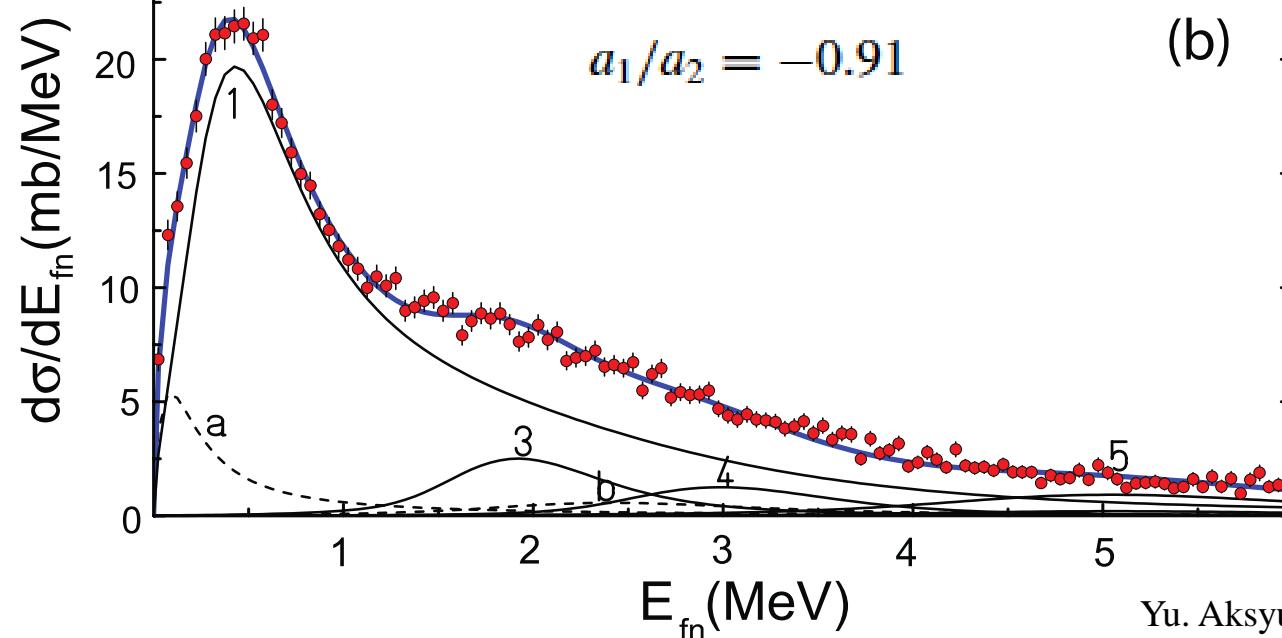
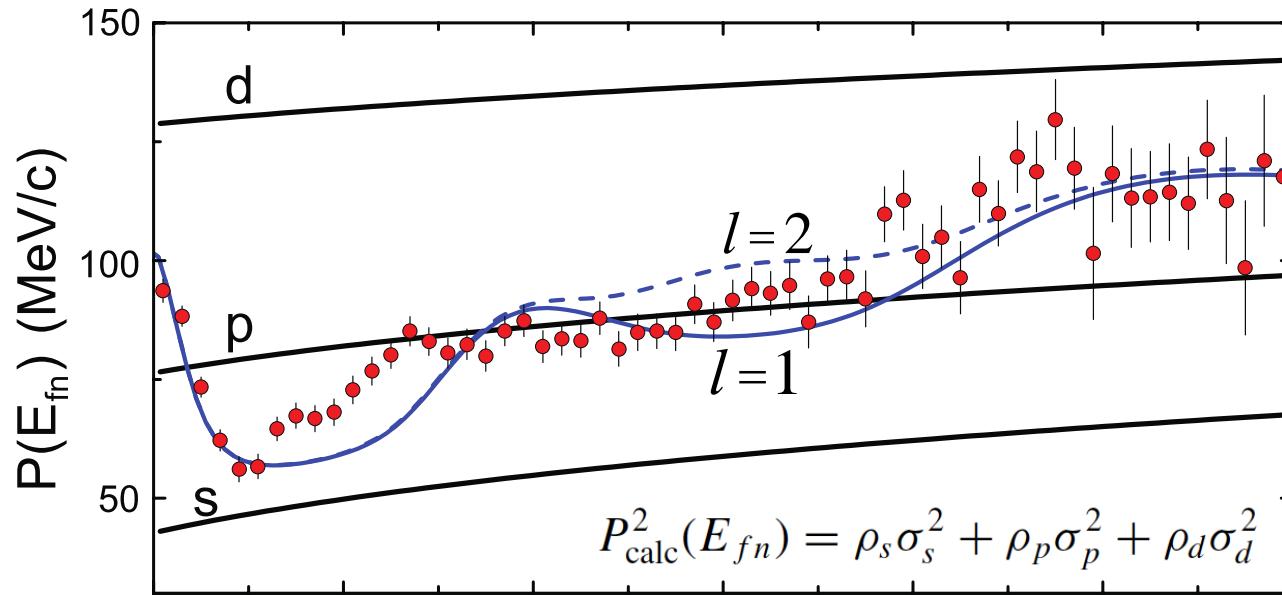
$$\boxed{\alpha^2 = 0.35, \beta^2 = 0.31, \gamma^2 = 0.34}$$

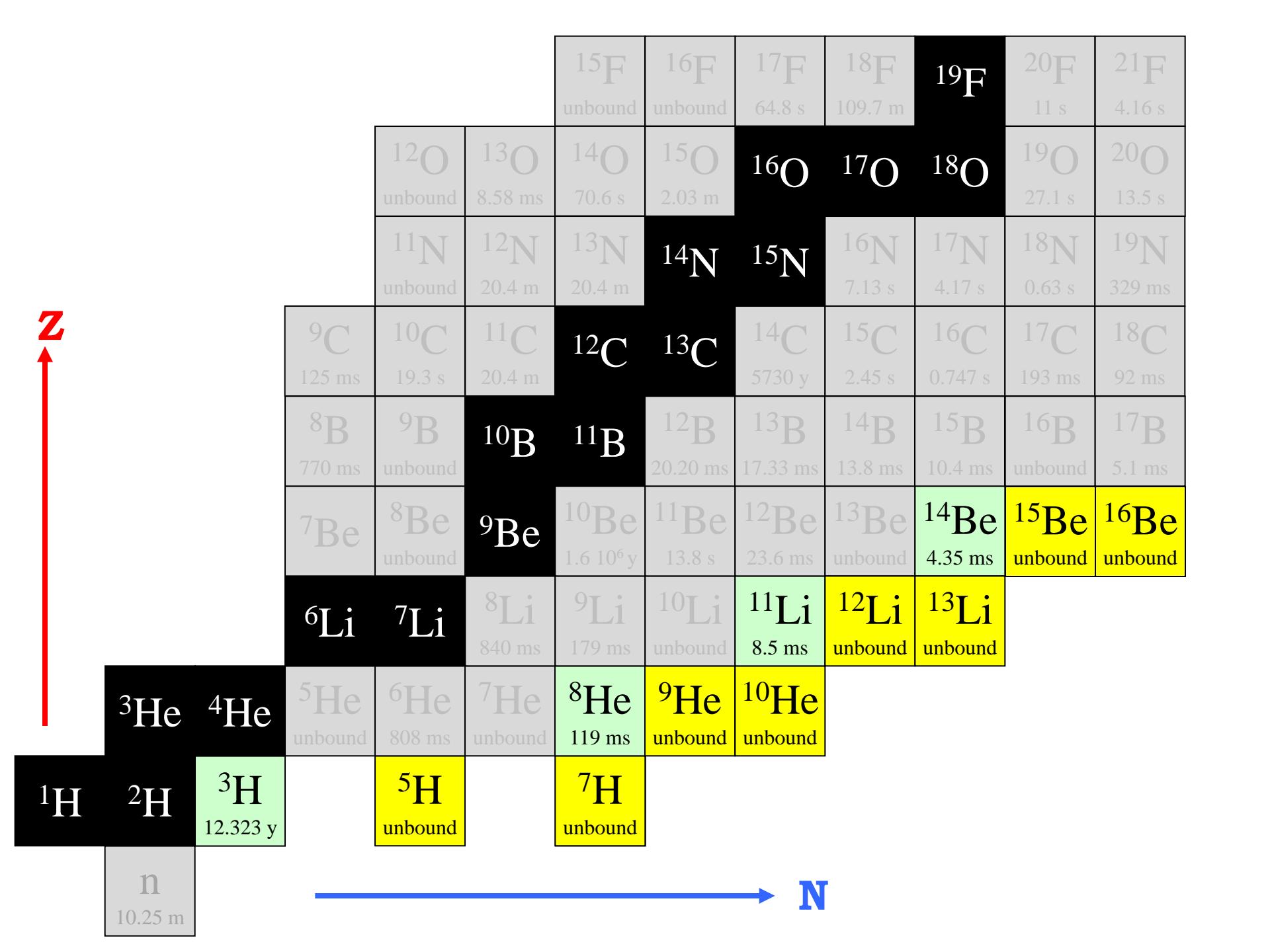
$$^{13}\text{Be}(1/2_2^+) = \mu [{}^{10}\text{Be} \otimes (0p_{1/2})^2 \otimes (1s_{1/2})] \\ - \lambda [{}^{10}\text{Be} \otimes (0d_{5/2})^2 \otimes (1s_{1/2})]$$

$$^{13}\text{Be}(1/2_1^+) = \lambda [{}^{10}\text{Be} \otimes (0p_{1/2})^2 \otimes (1s_{1/2})] \\ + \mu [{}^{10}\text{Be} \otimes (0d_{5/2})^2 \otimes (1s_{1/2})].$$

$$\boxed{d\sigma/dE_{fn} \propto \sqrt{E_{fn}} \left| \frac{a_1}{E_1 - E_{fn} - i\Gamma_1/2} + \frac{a_2}{E_2 - E_{fn} - i\Gamma_2/2} \right|^2},$$







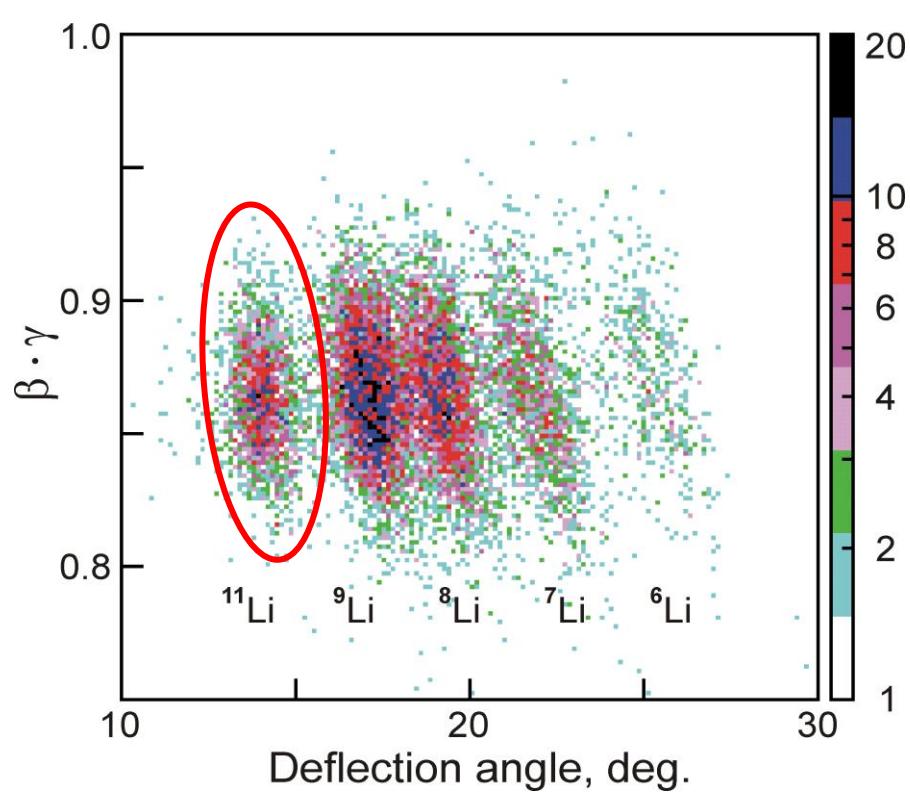
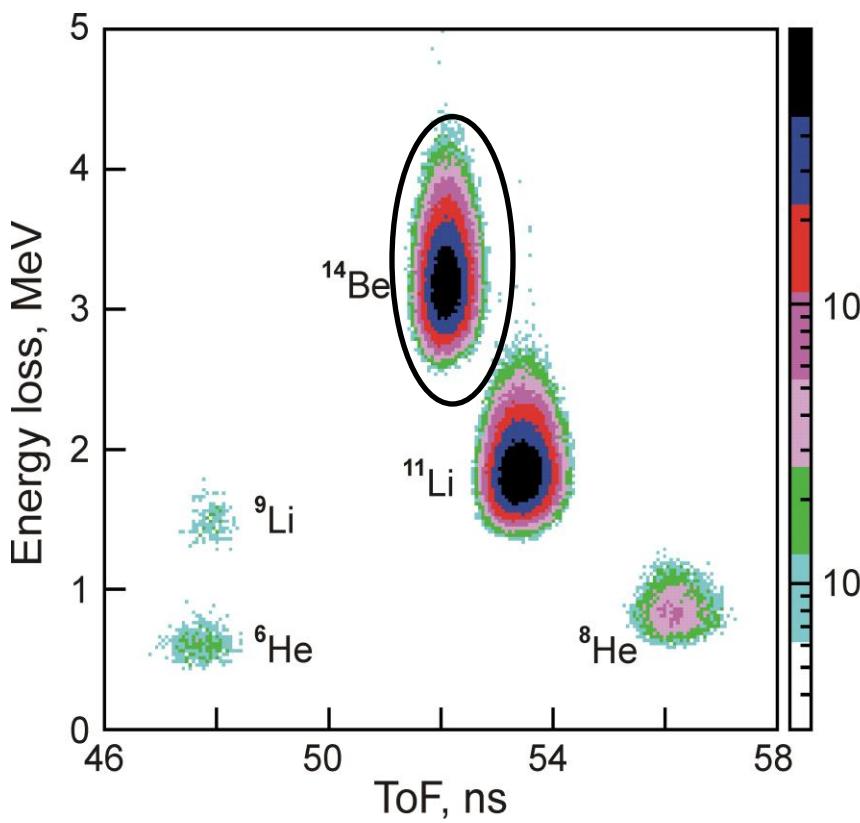
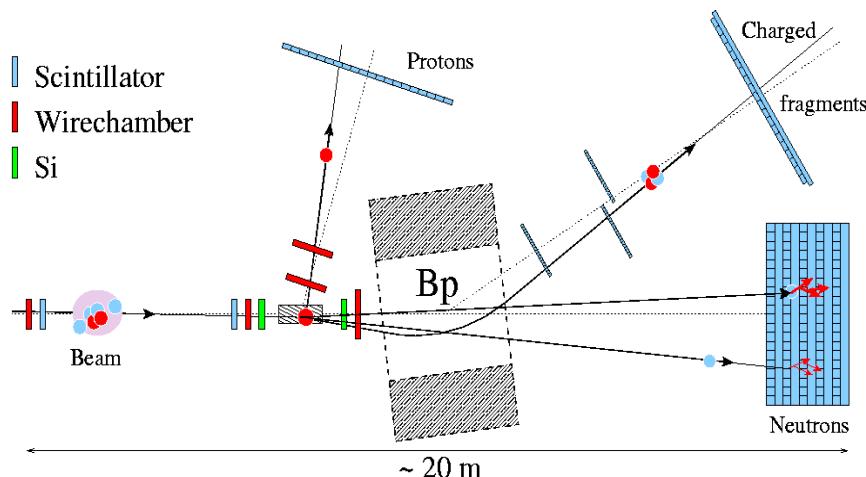
## Dripline nuclei as stepping stones towards the unbound

$^7\text{Be}$	$^8\text{Be}$ unbound	$^9\text{Be}$	$^{10}\text{Be}$ $1.6 \cdot 10^6 \text{ y}$	$^{11}\text{Be}$ $13.8 \text{ s}$	$^{12}\text{Be}$ $23.6 \text{ ms}$	$^{13}\text{Be}$ unbound	$^{14}\text{Be}$ $4.35 \text{ ms}$
$^6\text{Li}$	$^7\text{Li}$	$^8\text{Li}$ $840 \text{ ms}$	$^9\text{Li}$ $179 \text{ ms}$	$^{10}\text{Li}$ unbound	$^{11}\text{Li}$ $8.5 \text{ ms}$	$^{12}\text{Li}$ unbound	$^{13}\text{Li}$ unbound
$^4\text{He}$	$^5\text{He}$ unbound	$^6\text{He}$ $808 \text{ ms}$	$^7\text{He}$ unbound	$^8\text{He}$ $119 \text{ ms}$	$^9\text{He}$ unbound	$^{10}\text{He}$ unbound	
				$^7\text{H}$ unbound			

$^{11}\text{Li}$ :      **Neutron knock-out reaction**  $\rightarrow ^{10}\text{Li}$   
**Proton knock-out reactions**  $\rightarrow ^{9,10}\text{He}$

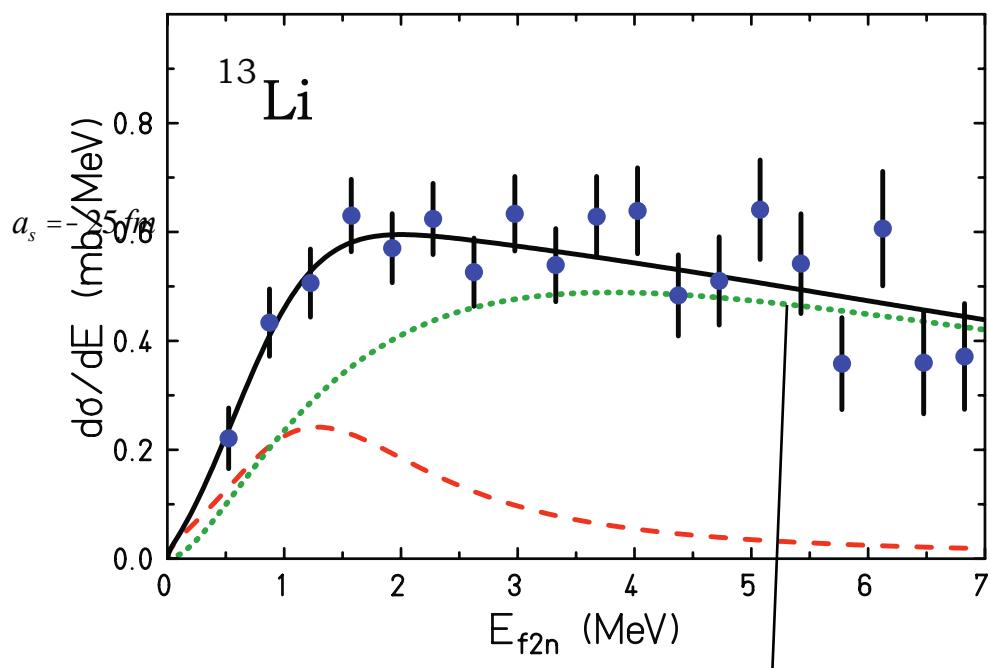
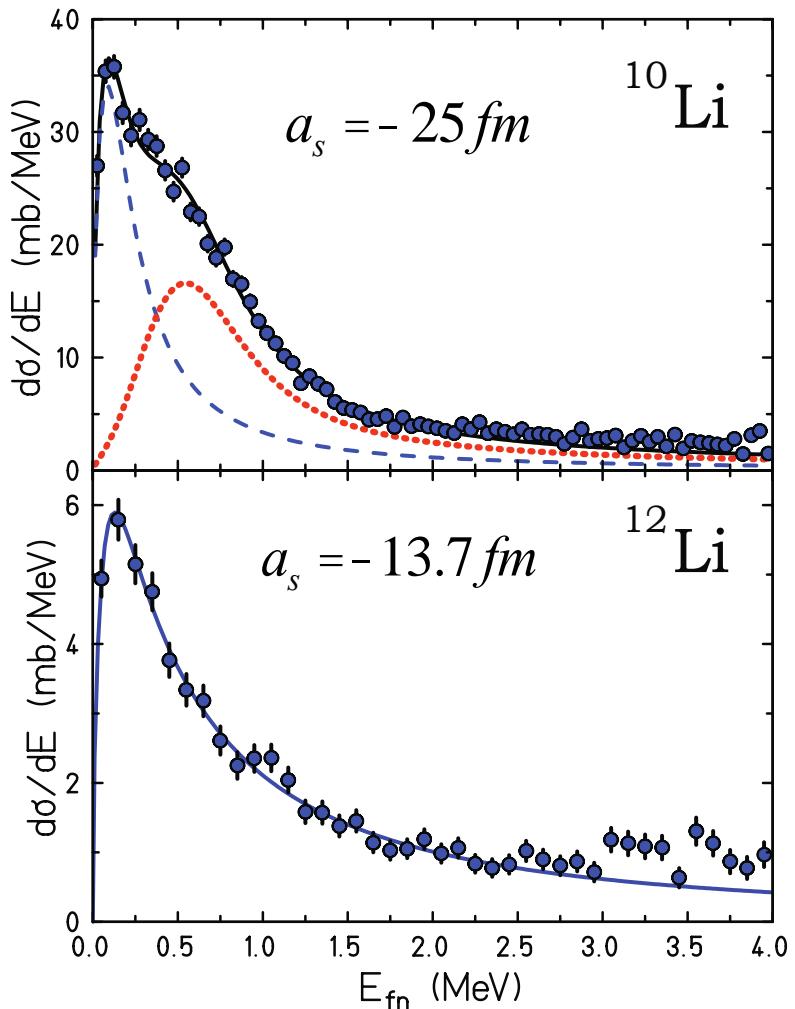
$^{14}\text{Be}$ :      **Neutron knock-out reaction**  $\rightarrow ^{13}\text{Be}$   
**Proton knock-out reactions**  $\rightarrow ^{12,13}\text{Li}$

$^8\text{He}$ :      **Neutron knock-out reaction**  $\rightarrow ^7\text{He}$   
**Proton knock-out reaction**  $\rightarrow (^7\text{H})$



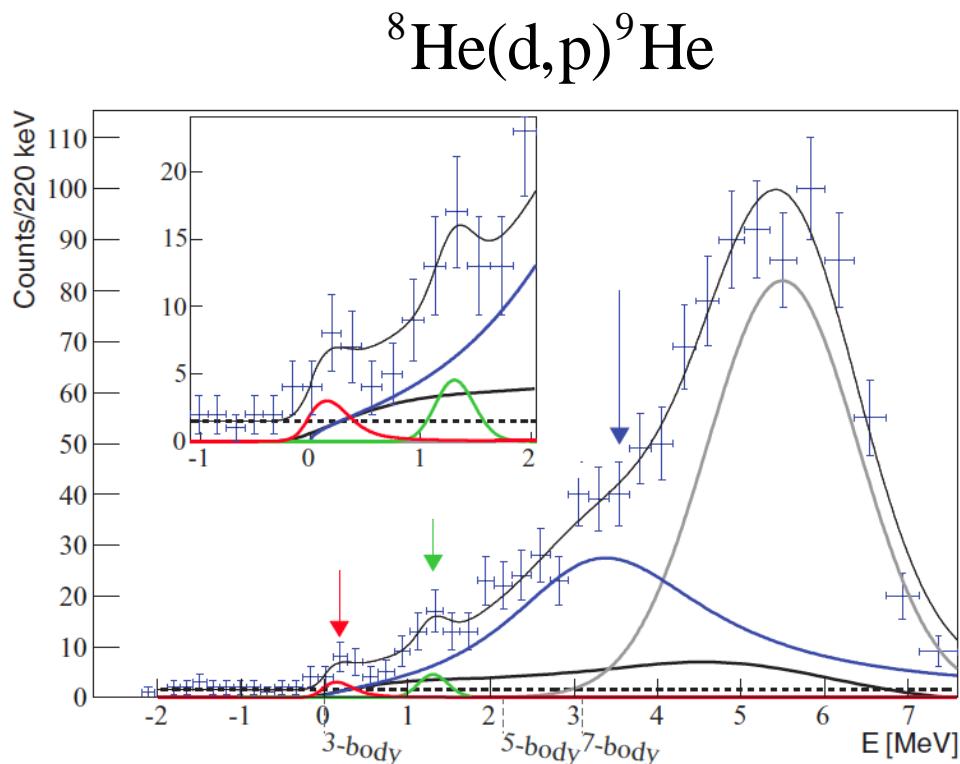
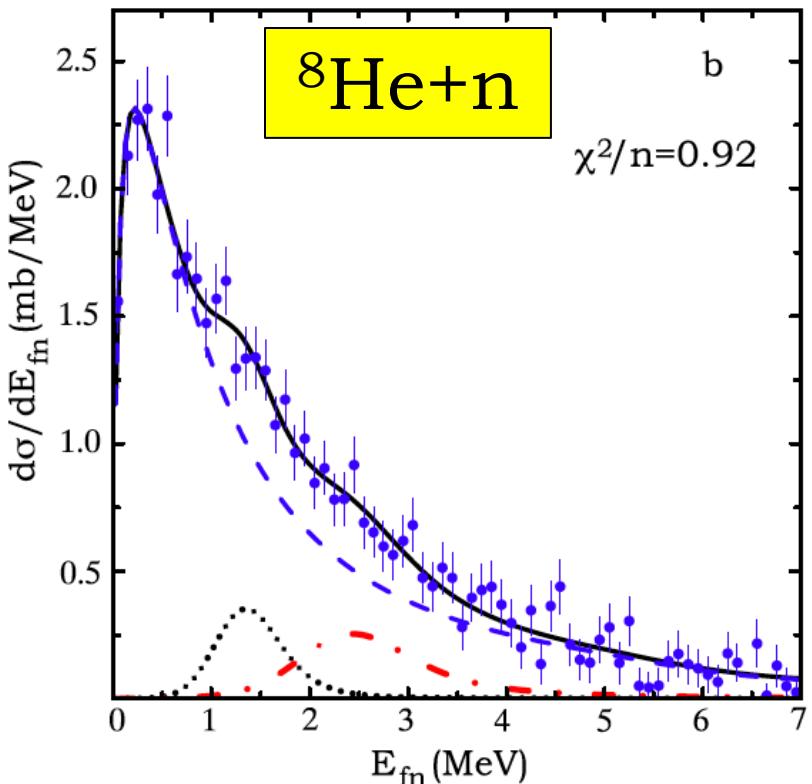
${}^6\text{Li}$	${}^7\text{Li}$	${}^8\text{Li}$	${}^9\text{Li}$	${}^{10}\text{Li}$	${}^{11}\text{Li}$	${}^{12}\text{Li}$	${}^{13}\text{Li}$
		840 ms	179 ms	unbound	8.5 ms	unbound	unbound

Yu. Aksyutina et al., Phys. Lett. **B 666** (2008) 430



Forssén et al. Nucl. Phys. **A 673**(2000) 143

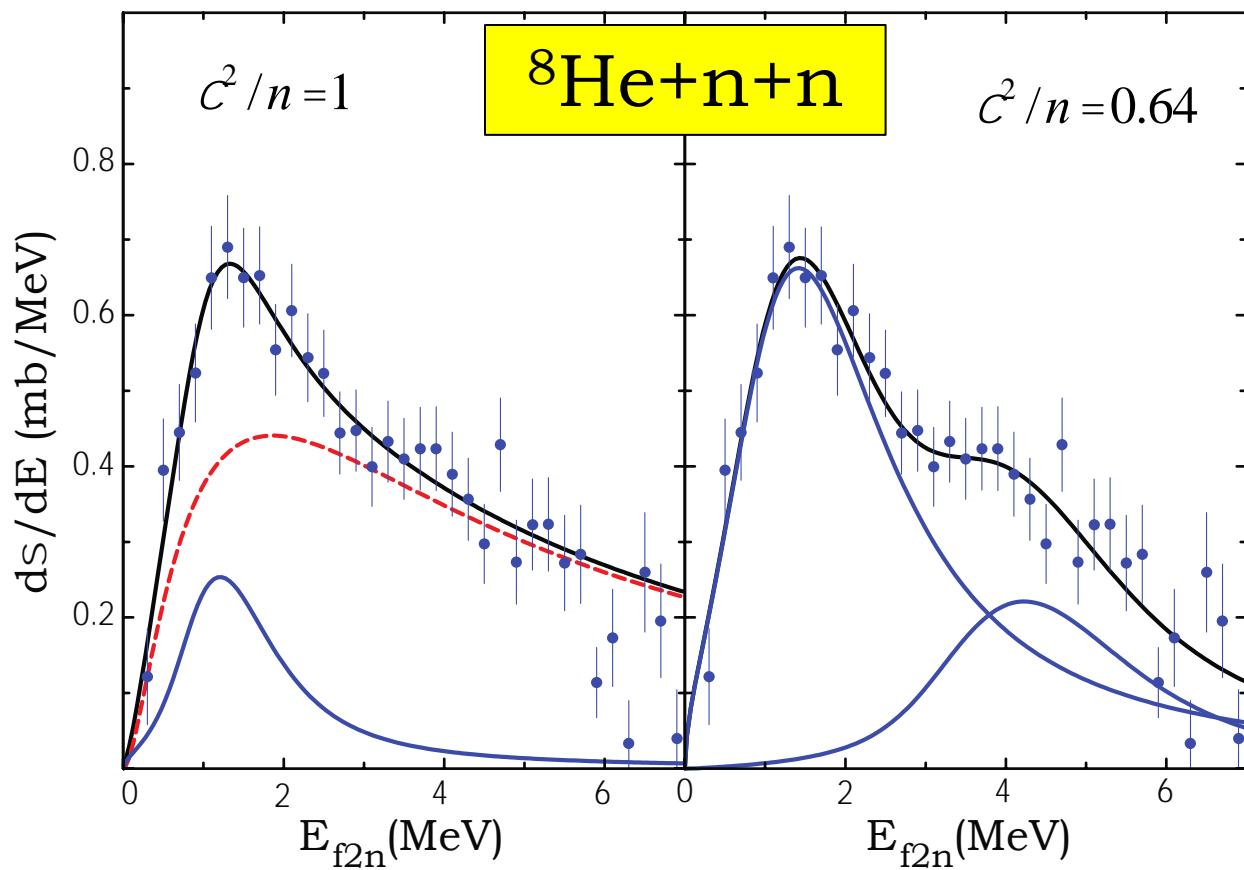
$^3\text{He}$	$^4\text{He}$	$^5\text{He}$ unbound	$^6\text{He}$ 808 ms	$^7\text{He}$ unbound	$^8\text{He}$ 119 ms	$^9\text{He}$ unbound	$^{10}\text{He}$ unbound
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Johansson et al., Nucl. Phys. **A842** (2010) 15

Al Kalanee et al., PRC **88** (2013) 034301

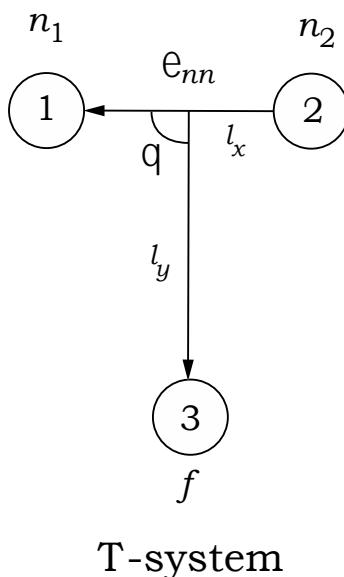
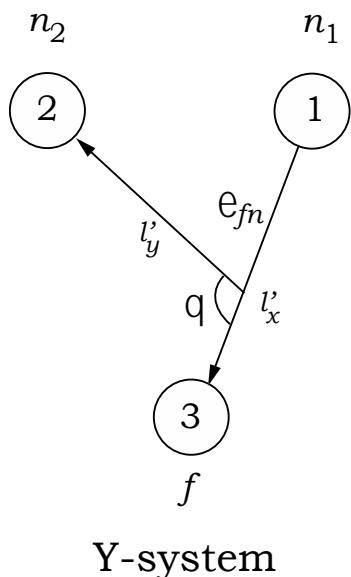
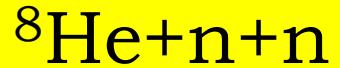
$^3\text{He}$	$^4\text{He}$	$^5\text{He}$ unbound	$^6\text{He}$ 808 ms	$^7\text{He}$ unbound	$^8\text{He}$ 119 ms	$^9\text{He}$ unbound	$^{10}\text{He}$ unbound
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Johansson et al., Nucl. Phys. **A842** (2010) 15

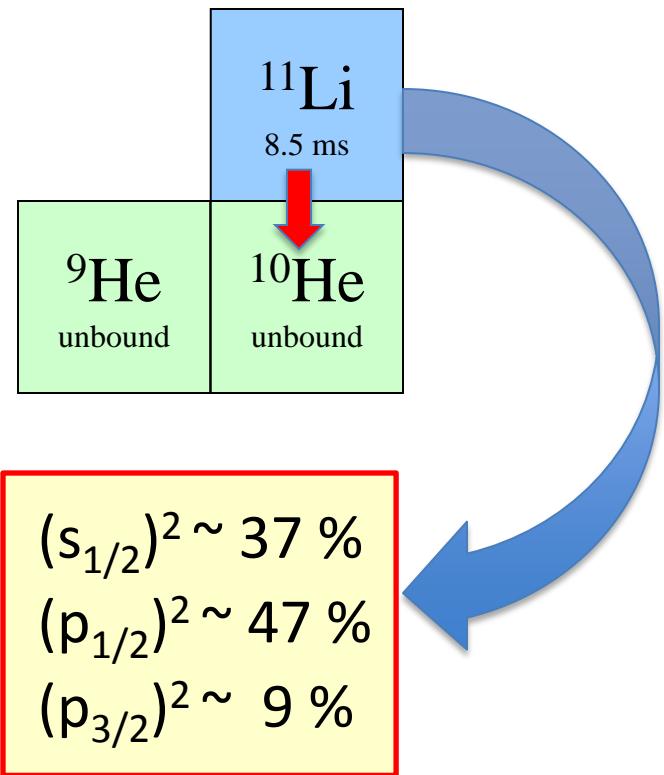
$$\sqrt{2C^2} - \sqrt{2n - 1} = -1.56$$

Neyman-Pearson test

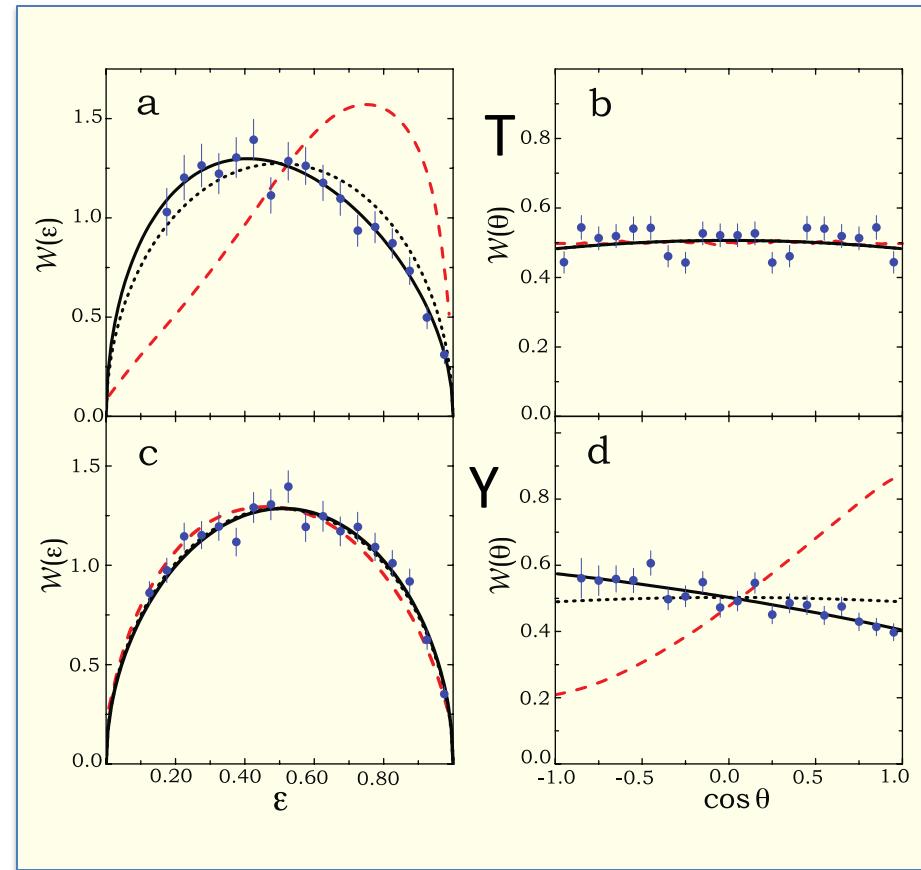
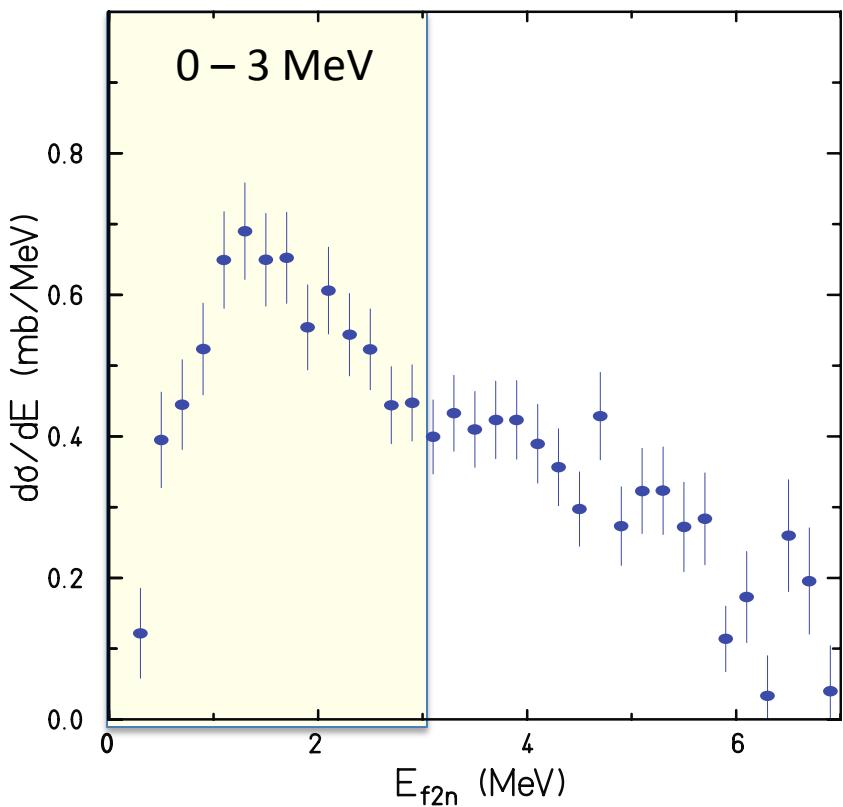


$$\varepsilon_{fn} = E_{fn}/E$$

$$\varepsilon_{nn} = E_{nn}/E$$



# Energy and angular correlations

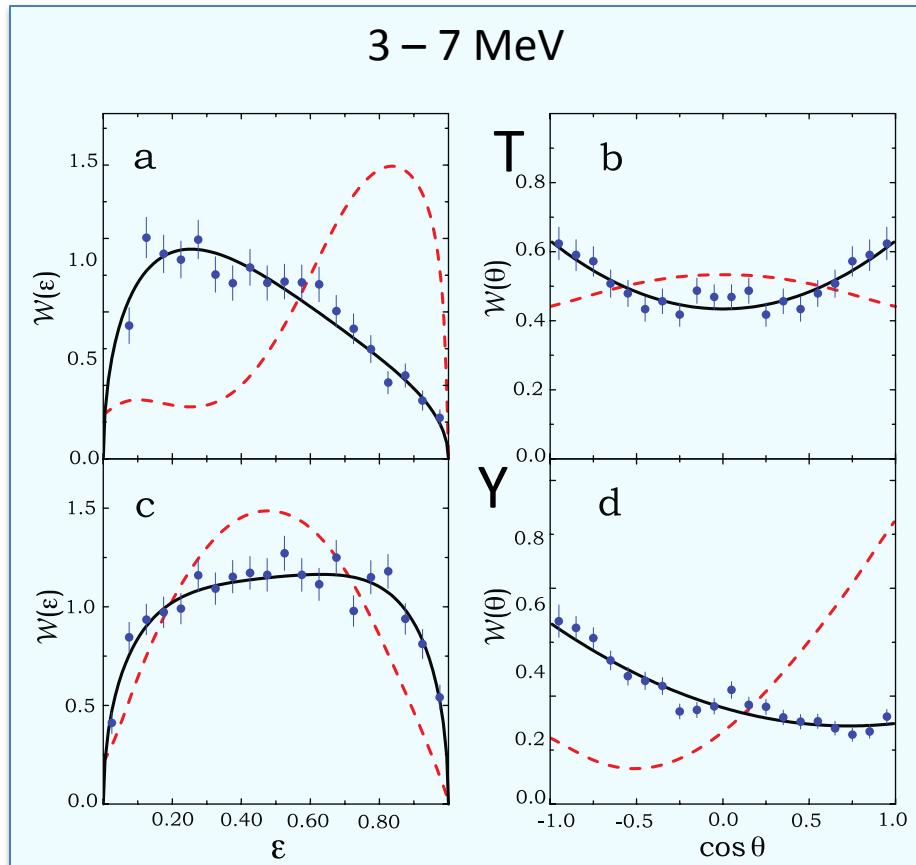
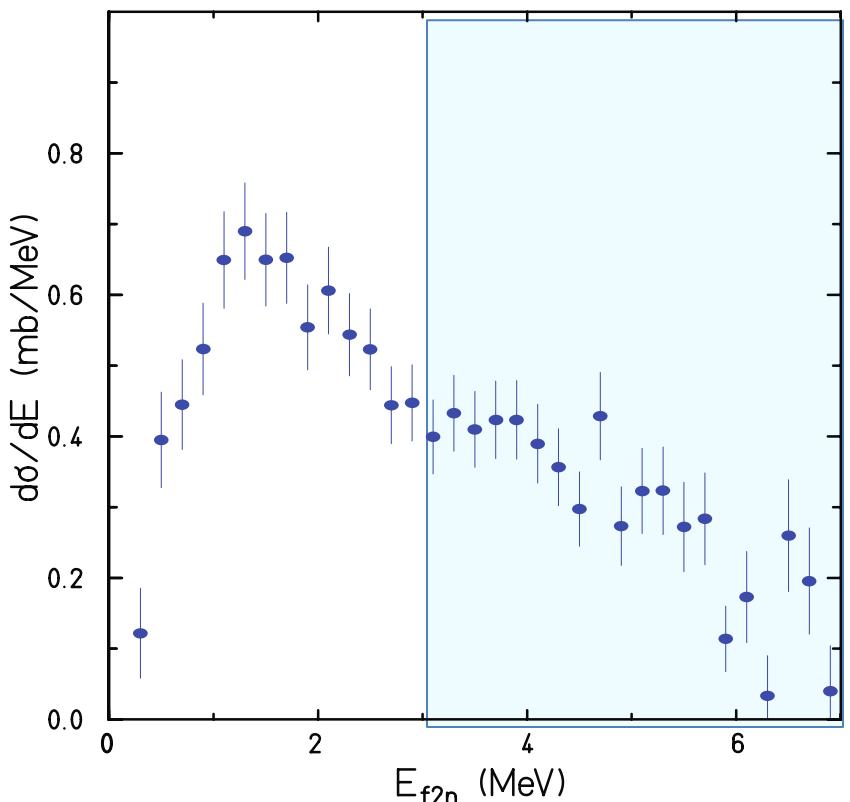


Restricted series of hyperspherical harmonics assuming  
 $I^\pi = 0^+$  and  $K \leq 2$

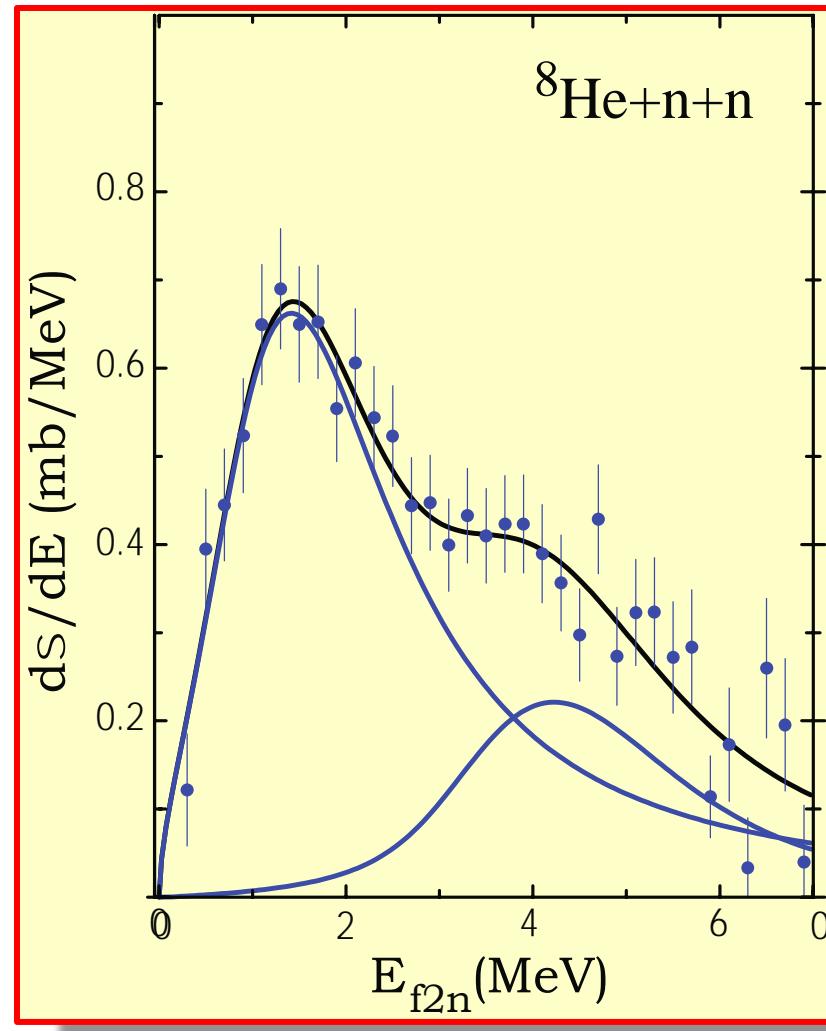
$$\mathcal{W}_{0+}(\varepsilon, \cos \theta)$$

$$\begin{aligned}
 &= \frac{4}{\pi} \sqrt{\varepsilon(1-\varepsilon)} (|A_{00}^{000}|^2 + 4|A_{00}^{000}|[(1-2\varepsilon)|A_{00}^{200}| \cos \varphi_{00}^{200} \\
 &\quad - 2\sqrt{2}\varepsilon(1-\varepsilon)(1-3\cos^2\theta)|A_{22}^{400}| \cos \varphi_{22}^{400}] + 4[(1-2\varepsilon)|A_{00}^{200}|^2 \\
 &\quad - 4\sqrt{2}\varepsilon(1-\varepsilon)(3-2\varepsilon)(1-3\cos^2\theta)|A_{00}^{200}||A_{22}^{400}| \cos(\varphi_{00}^{200} - \varphi_{22}^{400}) \\
 &\quad + 2\varepsilon(1-\varepsilon)((1-\cos^2\theta)|A_{11}^{211}|^2 + 4\varepsilon(1-\varepsilon)(1-3\cos^2\theta)^2 |A_{22}^{400}|^2)])
 \end{aligned}$$

# Energy and angular correlations

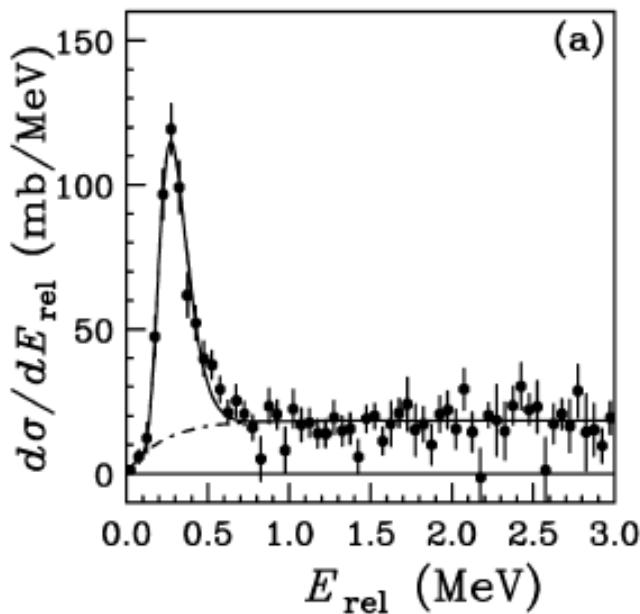


Restricted series of hyperspherical  
harmonics assuming  
 $I^\pi = 2^+$  and  $K \leq 4$



${}^8He + n + n$   
Ground state  $I^\rho = 0^+$   
Excited state  $I^\rho = 2^+$

${}^7\text{Be}$	${}^8\text{Be}$	${}^9\text{Be}$	${}^{10}\text{Be}$	${}^{11}\text{Be}$	${}^{12}\text{Be}$	${}^{13}\text{Be}$	${}^{14}\text{Be}$	${}^{15}\text{Be}$	${}^{16}\text{Be}$
unbound			1.6 $10^6$ y	13.8 s	23.6 ms	unbound	4.35 ms	unbound	unbound

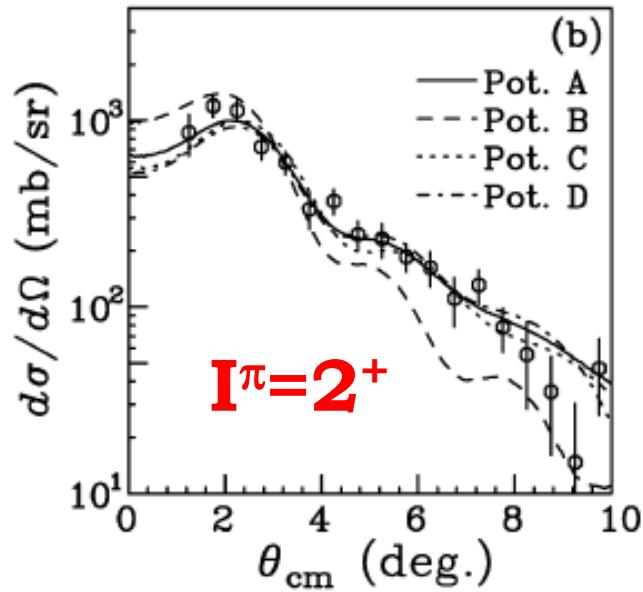


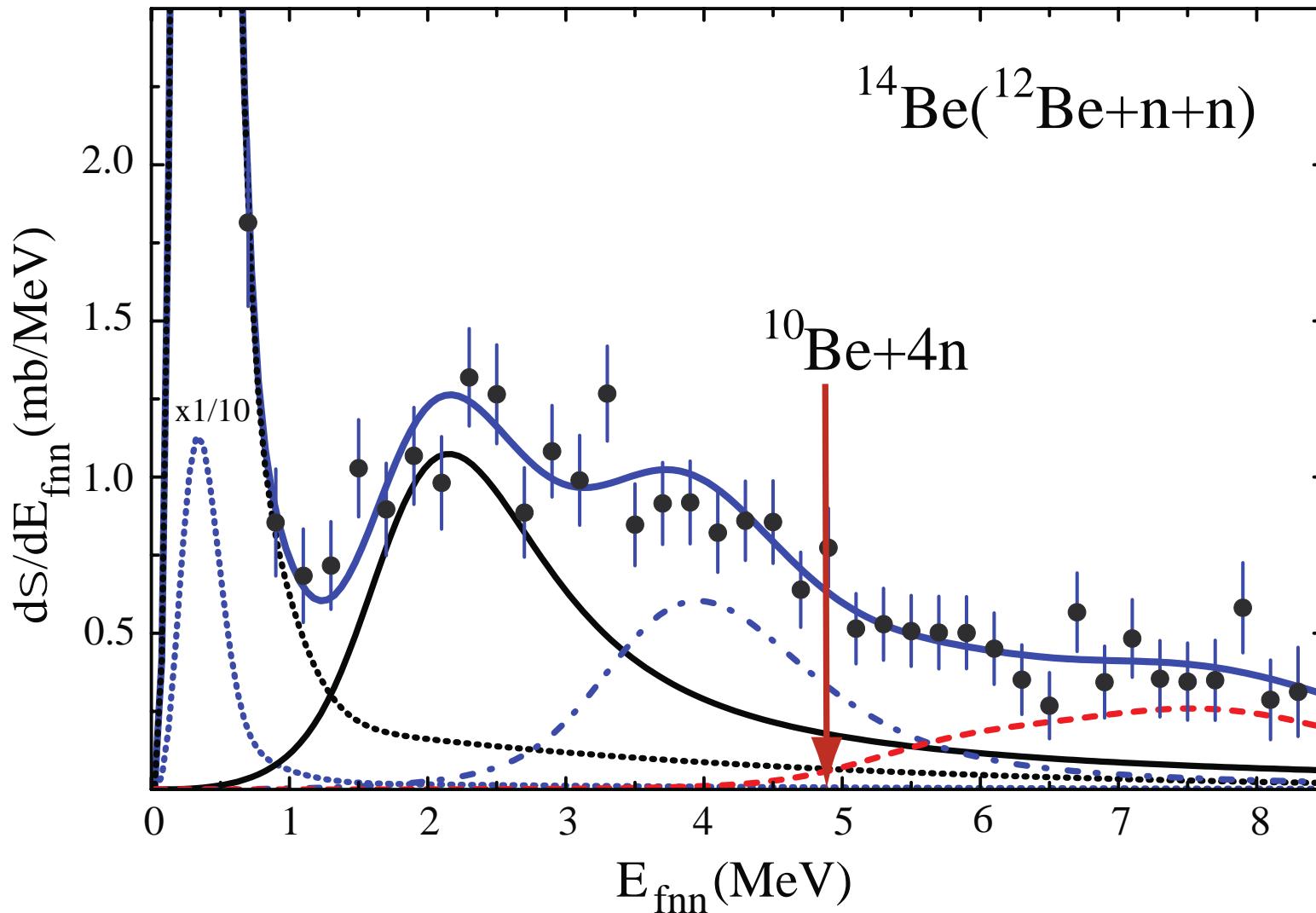
**RIKEN**  
 68 MeV/u  ${}^{14}\text{Be}$ ,  
 ${}^{12}\text{C}$  target

Sugimoto *et al.*,  
 Phys. Lett. **B654** (2007) 160

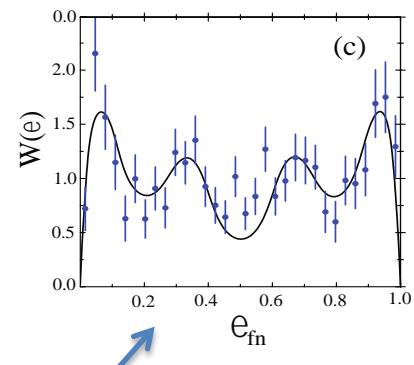
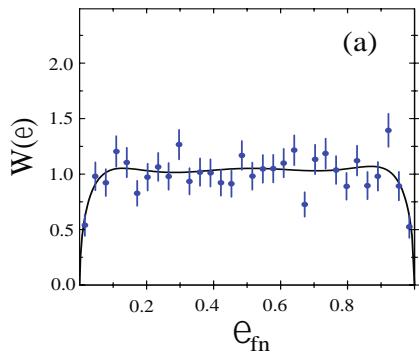
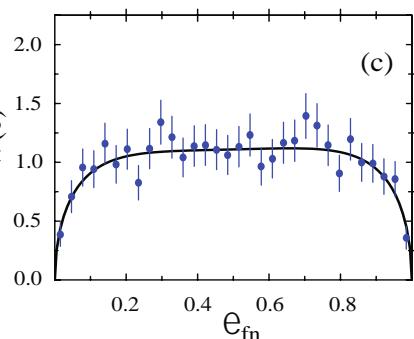
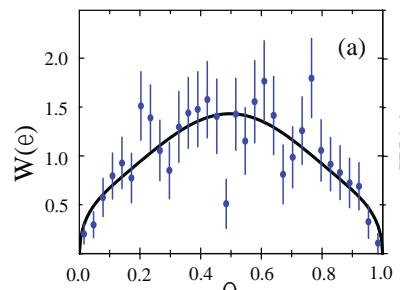
$$2^+ \frac{1.54(13) \text{ MeV}}{12\text{Be} + 2\text{n}}$$

$$0^+ \frac{}{{}^{14}\text{Be}}$$

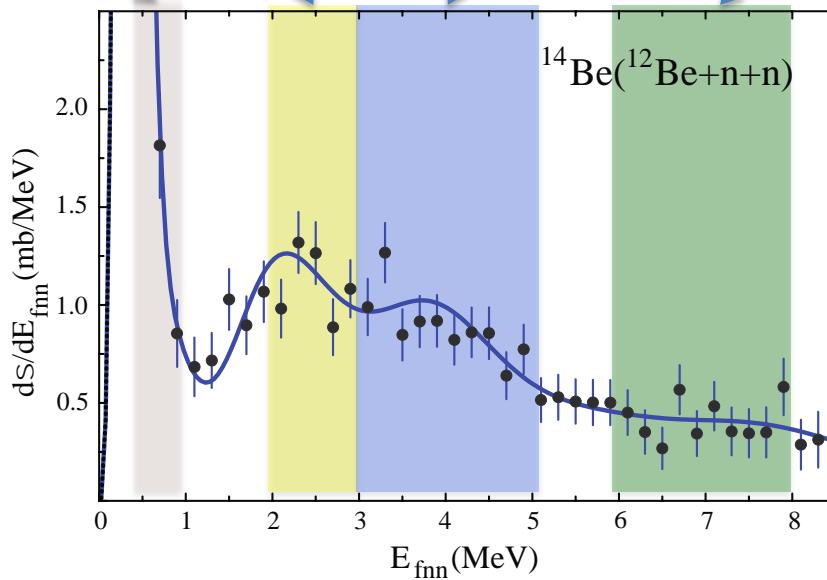




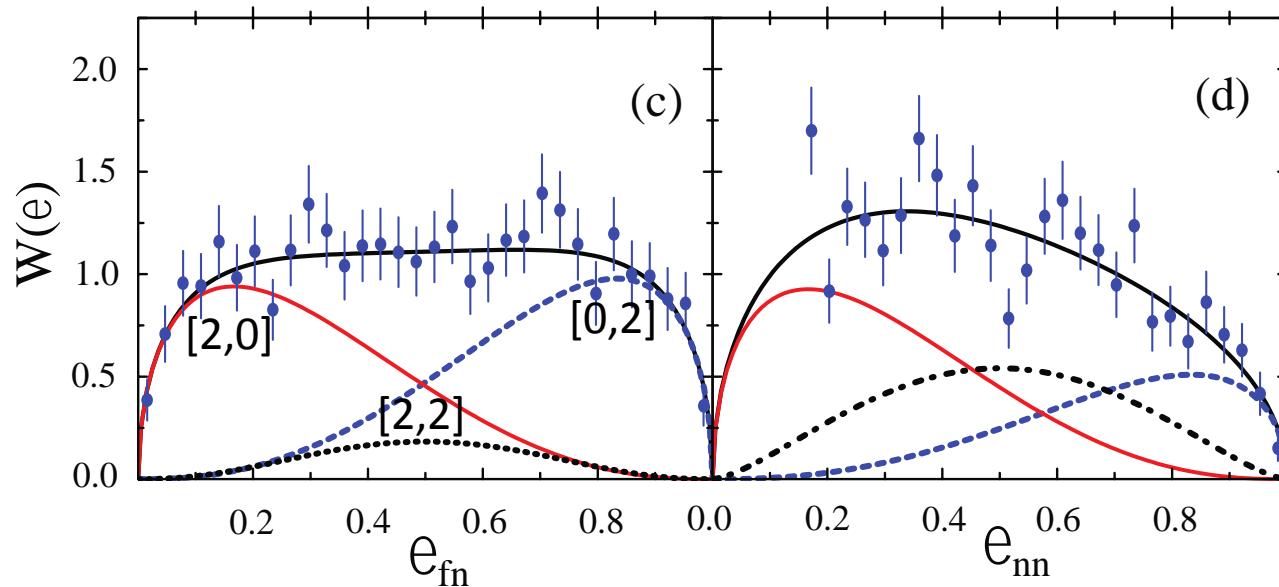
Yu. Aksyutina et al.,  
PRL 111, 242501 (2013)



$$\varepsilon_{fn} = E_{fn}/E$$



$2 \text{ MeV} < E_{\text{fnn}} < 3 \text{ MeV}$



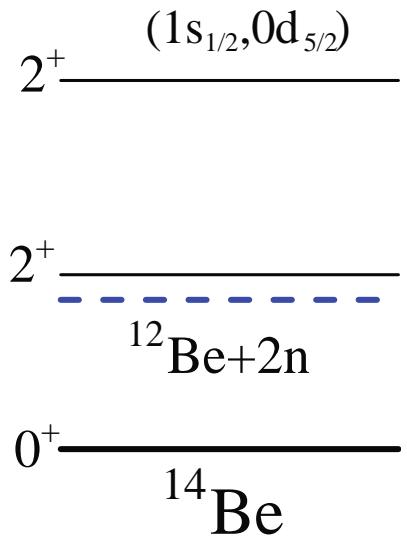
L.M. Delves, Nucl. Phys. **20** (1960) 275

$$W(\epsilon) \sim \sum_i A_i^2 \epsilon^{l_x^i + \frac{1}{2}} (1 - \epsilon)^{l_y^i + \frac{1}{2}}$$

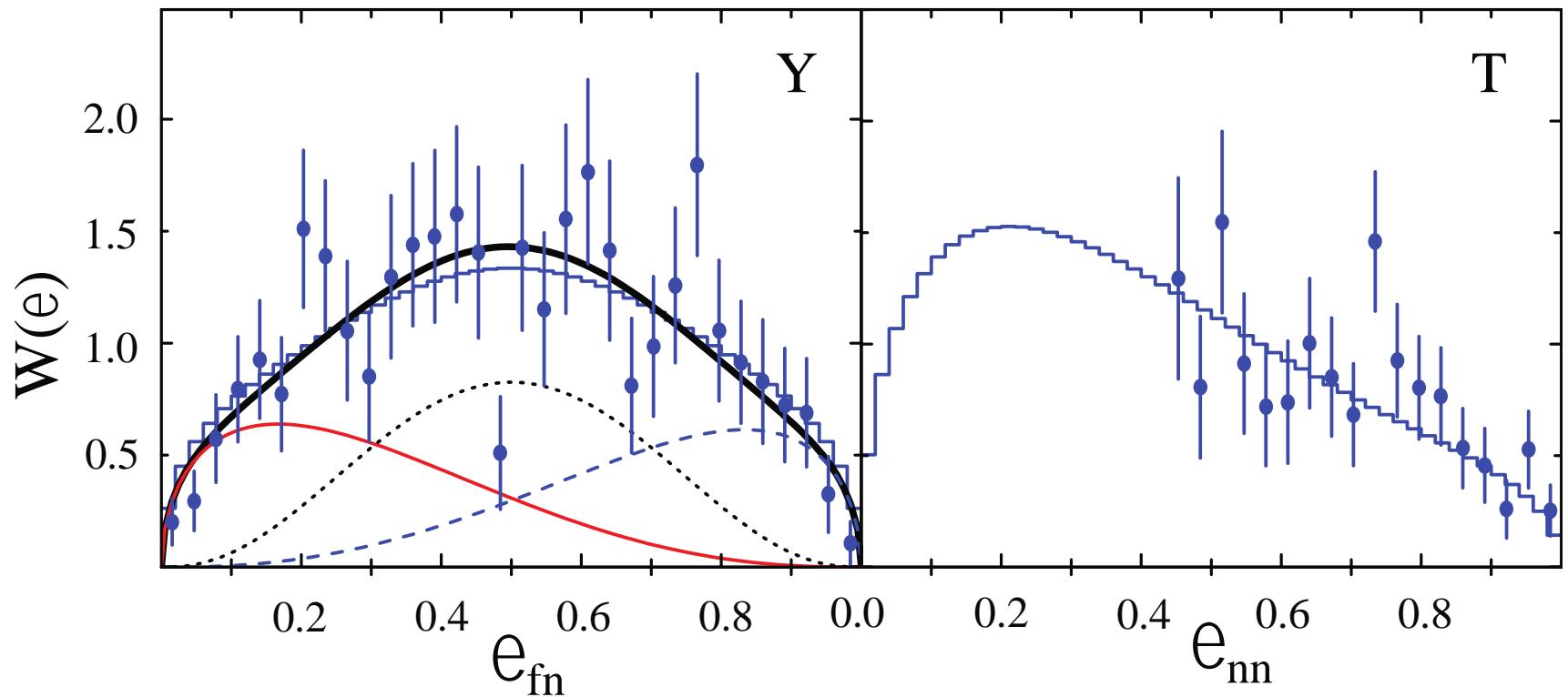
$E^* = 3.54(16) \text{ MeV}, \Gamma = 1.5 \text{ MeV}$

$I^\pi = 2^+$

$(1s_{1/2}, 0d_{5/2}) \approx 90 \%$



$0.5 \text{ MeV} < E_{\text{fnn}} < 1.0 \text{ MeV}$



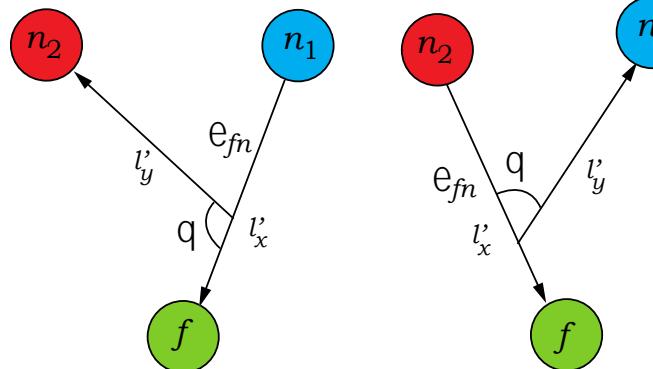
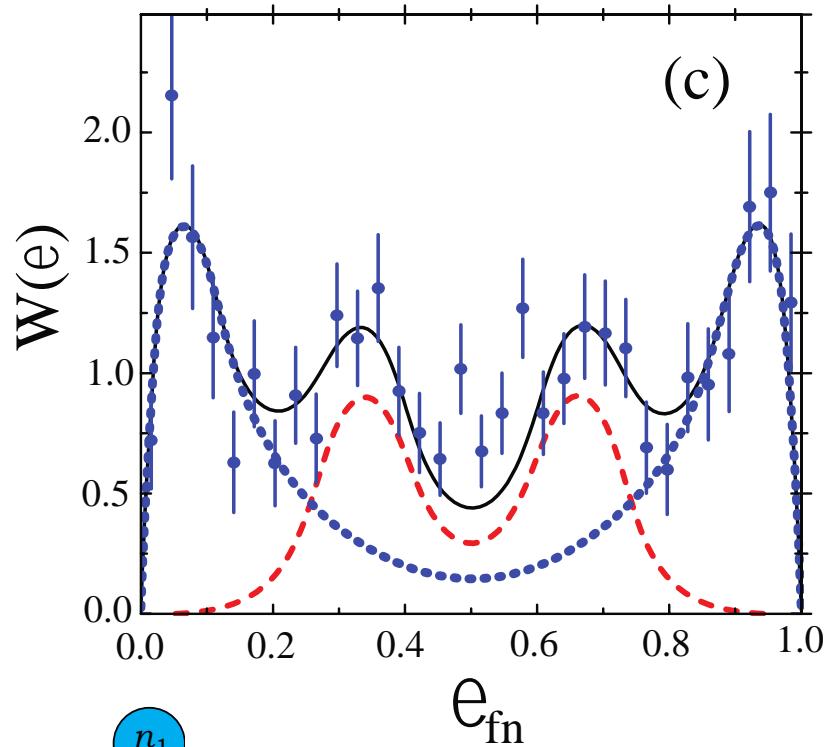
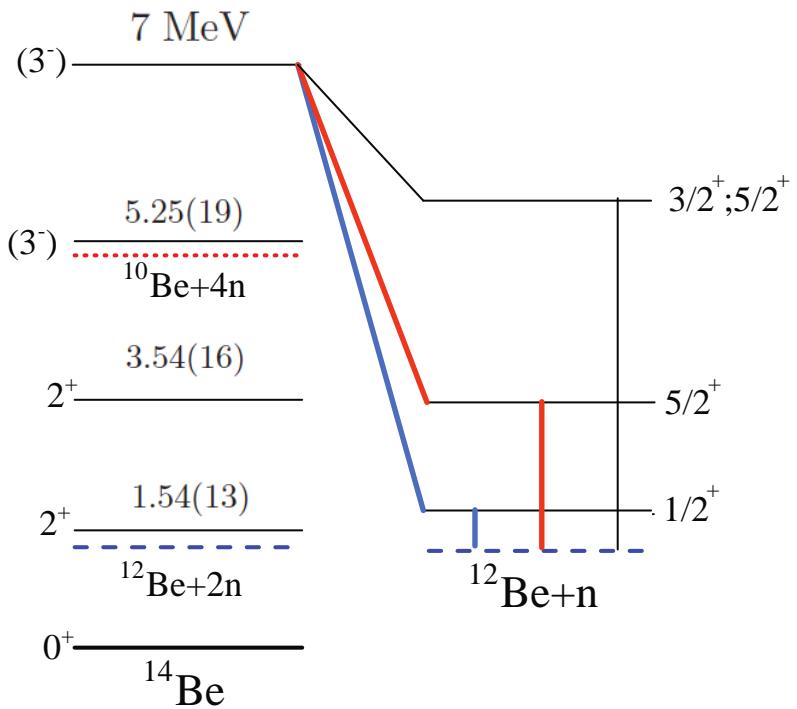
$$^{12}\text{Be}(0^+) = {}^{10}\text{Be} \otimes \left[ \color{red}{\alpha} (1s_{1/2})_{I=0}^2 + \color{blue}{\beta} (0p_{1/2})_{I=0}^2 + \color{green}{\gamma} (0d_{5/2})_{I=0}^2 \right]$$

$$^{14}\text{Be}(2^+) = {}^{10}\text{Be} \otimes \left[ \left\{ \begin{array}{c} \alpha_1 (1s_{1/2})_{I=0}^2 \\ \alpha_2 (0p_{1/2})_{I=0}^2 \\ \alpha_3 (0d_{5/2})_{I=0}^2 \end{array} \right\} (0d_{5/2})_{I=2}^2 + \left\{ \begin{array}{c} \beta_1 (0p_{1/2})_{I=0}^2 \\ \beta_2 (0d_{5/2})_{I=0}^2 \end{array} \right\} (1s_{1/2}, 0d_{5/2})_{I=2} \right]$$

$$\alpha_1\alpha+\alpha_2\beta+\alpha_3\gamma \quad l=2$$

$$\beta_1\beta+\beta_2\gamma \qquad l=0,2$$

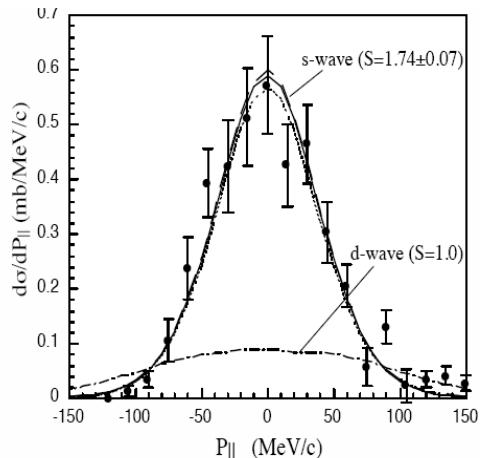
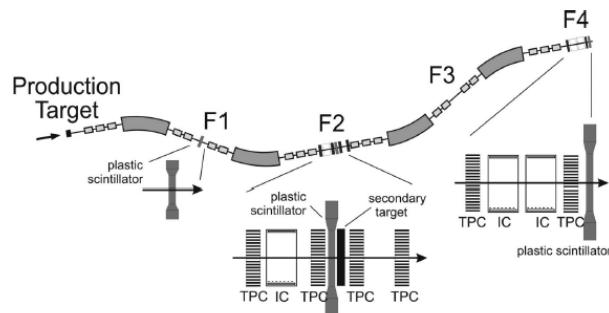
# Sequential Decay



$^{20}\text{F}$ 11 s	$^{21}\text{F}$ 4.16 s	$^{22}\text{F}$ 4.236 s	$^{23}\text{F}$ 2.23 s	$^{24}\text{F}$ 0.34 s	$^{25}\text{F}$ 50 ms	$^{26}\text{F}$ 10.2 ms	$^{27}\text{F}$ 4.29 ms
$^{19}\text{O}$ 27.1 s	$^{20}\text{O}$ 13.5 s	$^{21}\text{O}$ 3.4 s	$^{22}\text{O}$ 2.25 s	$^{23}\text{O}$ 82 ms	$^{24}\text{O}$ 61 ms		
$^{18}\text{N}$ 0.63 s	$^{19}\text{N}$ 329 ms	$^{20}\text{N}$ 142 ms	$^{21}\text{N}$ 95 ms	$^{22}\text{N}$ 24 ms	$^{23}\text{N}$ 14.5 ms		
$^{17}\text{C}$ 193 ms	$^{18}\text{C}$ 92 ms	$^{19}\text{C}$ 49 ms	$^{20}\text{C}$ 14 ms			$^{22}\text{C}$ 6.2 ms	

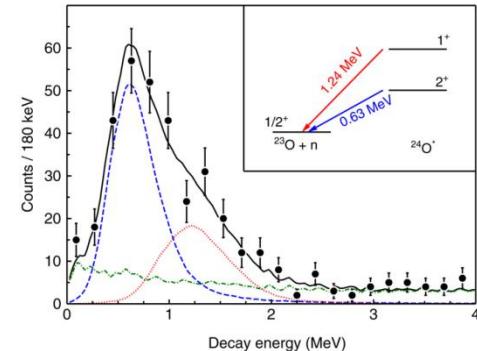
$^{29}\text{F}$ 2.6 ms
$^{28}\text{O}$ unbound

$^{31}\text{F}$   
 $>260$  ns

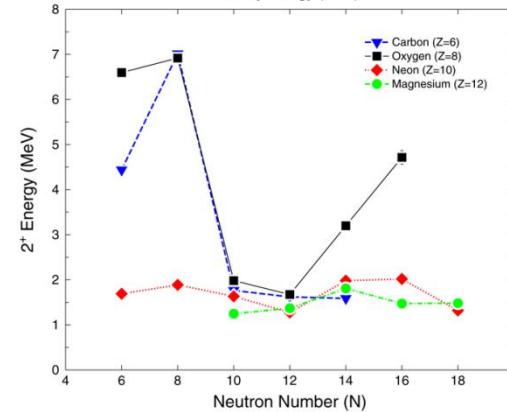


$$S=1.74(19)$$

GSI



MSU



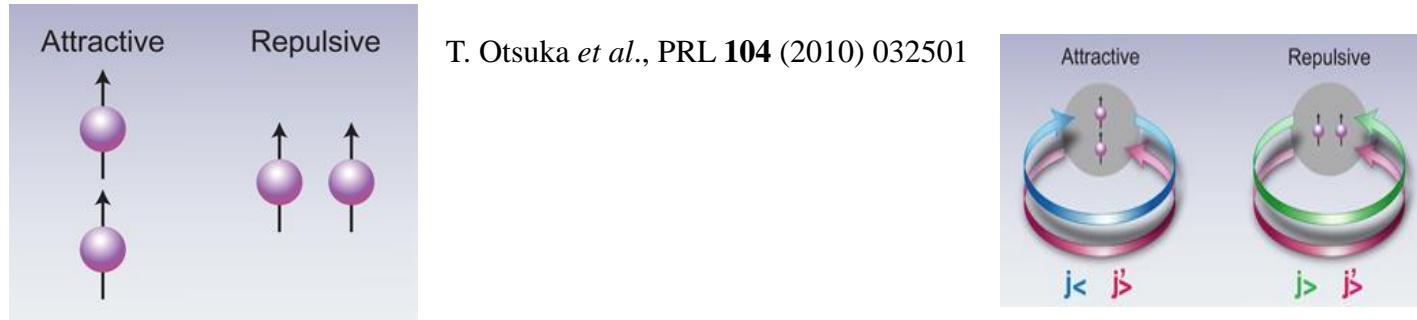
$$2^+: 4.7 \text{ MeV}$$

C.R. Hoffman et al.,  
PLB 672 (2009) 17

# Shell evolution and nuclear forces

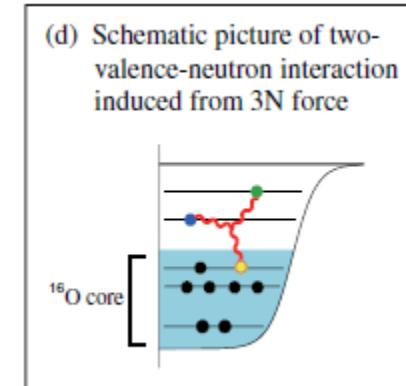
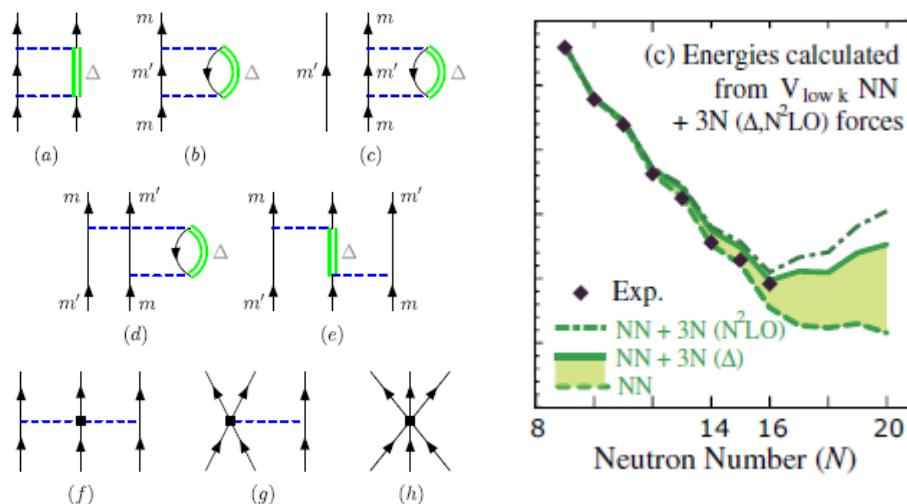
N.A. Smirnova *et al.*, PLB **686** (2010) 109

## Novel Features of Nuclear Forces and Shell Evolution in Exotic Nuclei



## Three-Body Forces and the Limit of Oxygen Isotopes

T. Otsuka *et al.*, PRL **105** (2010) 032501

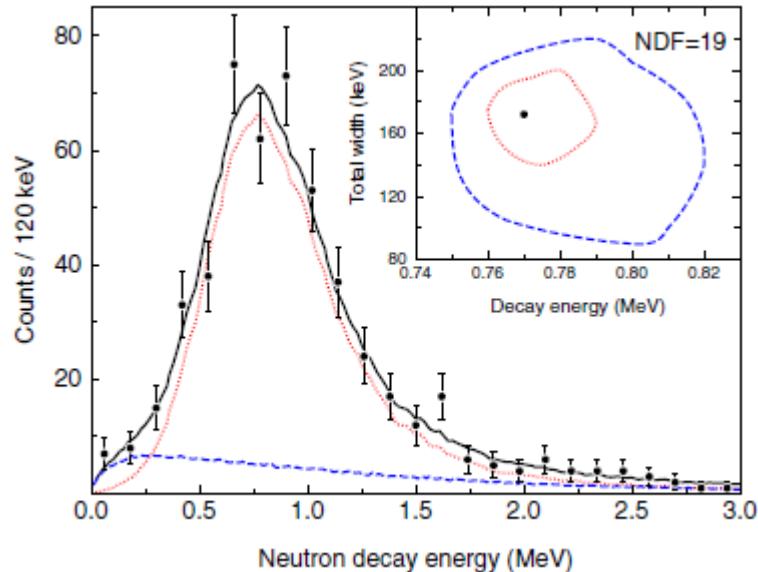


## GSI : Data September 2010

$^{24}\text{F}$	$^{25}\text{F}$	$^{26}\text{F}$	$^{27}\text{F}$	$^{28}\text{F}$	$^{29}\text{F}$	$^{30}\text{F}$	$^{31}\text{F}$
0.34 s	50 ms	10.2 ms	4.9 ms	unbound	2.6 ms	unbound	>260 ns
$^{23}\text{O}$	$^{24}\text{O}$	$^{25}\text{O}$	$^{26}\text{O}$		$^{28}\text{O}$		
82 ms	61 ms	unbound	unbound		unbound		



**MoNA**

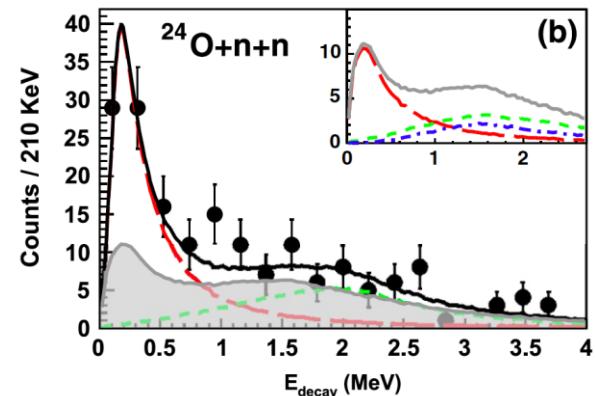
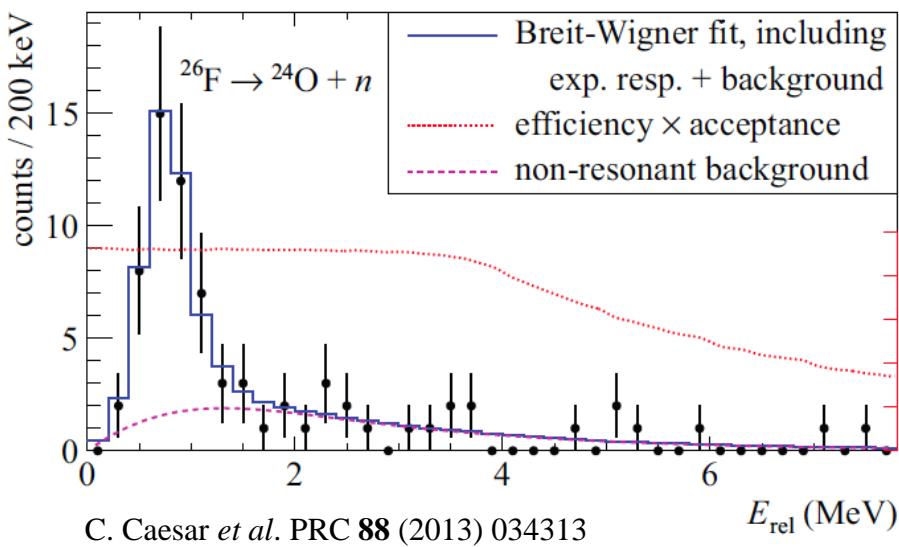


$$E_{\text{decay}} = 770 \text{ keV}, \Gamma = 172 \text{ keV}, I=2$$

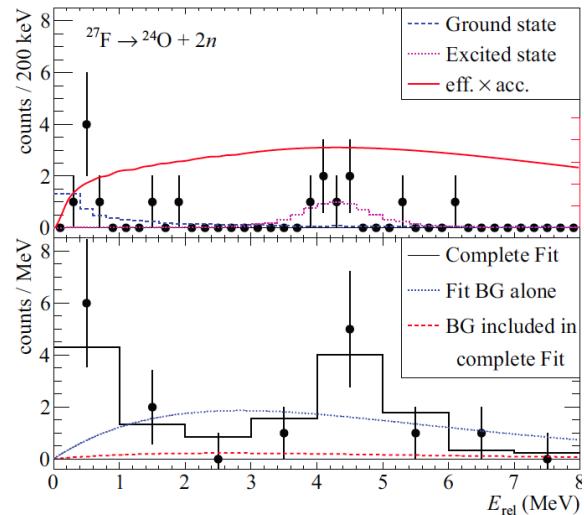
Shell gap  $\nu 1s_{1/2} - \nu 0d_{3/2}$   
4.86(13) MeV

## GSI : Data September 2010

$^{24}\text{F}$ 0.34 s	$^{25}\text{F}$ 50 ms	$^{26}\text{F}$ 10.2 ms	$^{27}\text{F}$ 4.9 ms	$^{28}\text{F}$ unbound	$^{29}\text{F}$ 2.6 ms	$^{30}\text{F}$ unbound	$^{31}\text{F}$ $>260$ ns
$^{23}\text{O}$ 82 ms	$^{24}\text{O}$ 61 ms	$^{25}\text{O}$ unbound	$^{26}\text{O}$ unbound			$^{28}\text{O}$ unbound	



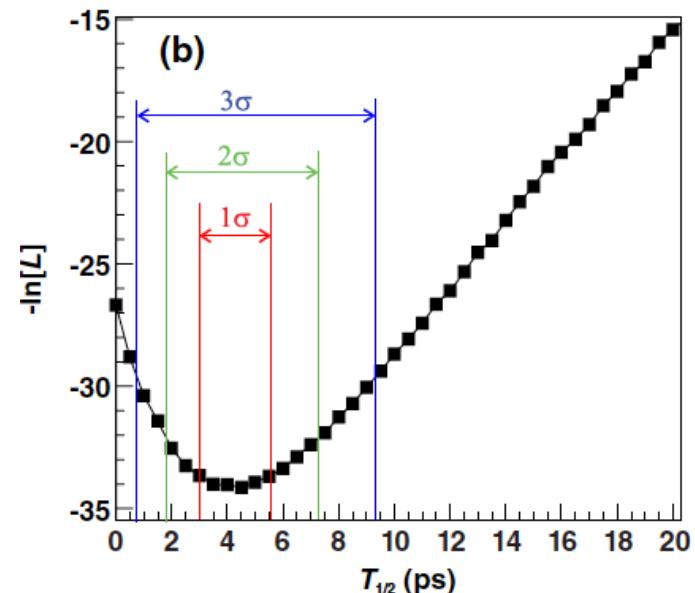
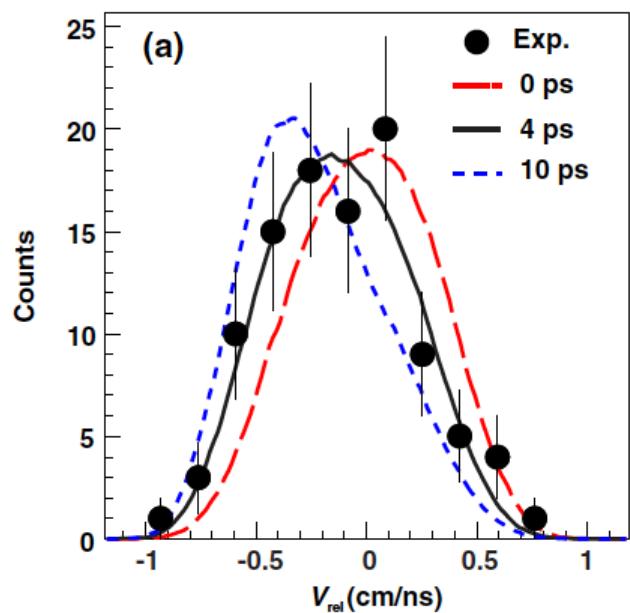
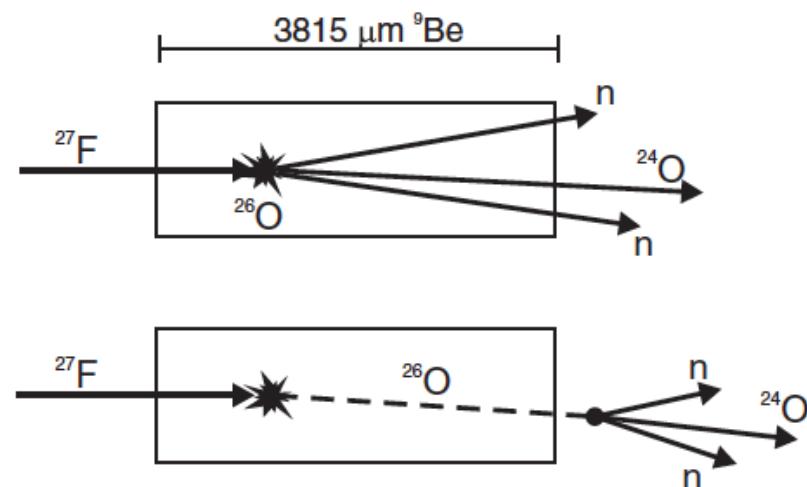
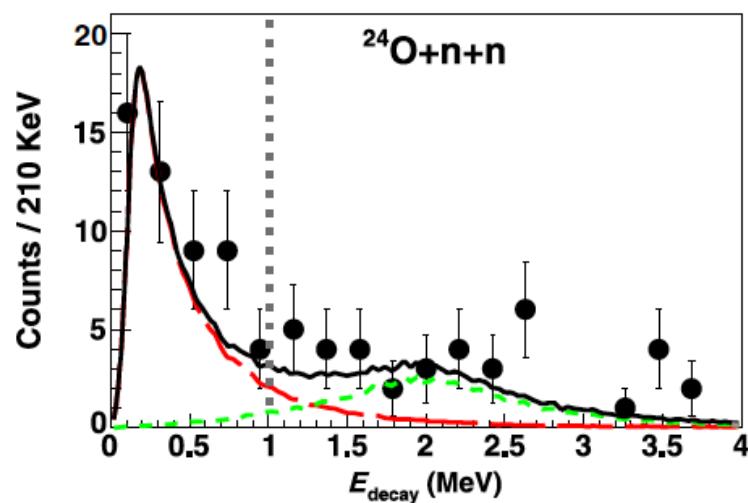
E. Lunderberg *et al.* PRL **108** (2012) 142503



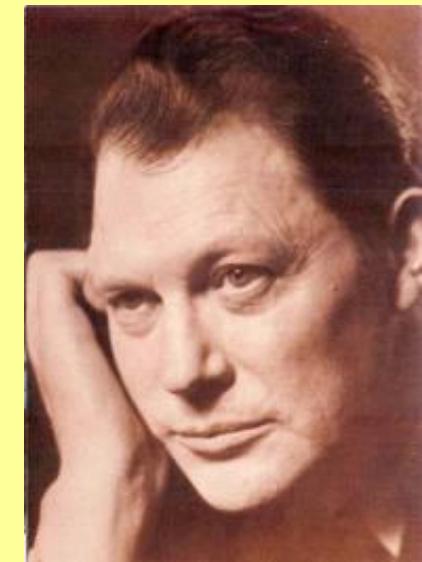
## Two-neutron radioactivity ?

L.V. Grigorenko *et al.* PRL **111** (2013) 042501

Z. Kohley *et al.* PRL **110** (2013) 152501  $\leftarrow$   $4.5^{+1.1}_{-1.5} \pm 3 \text{ ps}$



To view the world in this way  
Is to see it with the eyes of a gambler.  
Our best choice is to cast the die again and again.  
To let it keep rolling in the hope  
That it will give returns high enough  
to inspire new expectations.



Att betrakta världen så är att se  
på den med tärningsögon.  
Vårt bästa val är att kasta tärningen ofta, ofta.  
Att ständigt låta den rulla, att den må falla ut  
med en kvot stor nog  
för en ny förväntan att andas.

Harry Martinson  
Tärningspelet, 1971