# DARK MATTER <br> IN THE UNIVERSE 

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## Outline

o Introduction

0 WIMPs
o Indirect Detection

0 Direct Detection
o Beyond WIMPs

## Dark matter: compelling evidence there are new particles to be discovered



## No shortage of DM candidates

Axions, WIMPS, Neutralinos, Gravitinos, Axinos, Sneutrinos, Kaluza-Klein particles, Heavy Fourth Generation Neutrinos, Mirror particles, superWIMPs, WIMPzillas, Sterile Neutrinos, Light Scalars, Q-Balls, Brane World Dark Matter, Primordial Black Holes, Asymmetric Dark Matter....
....or, perhaps most likely, something we haven't thought of yet....

## Well-motivated DM candidates

- Thermal WIMPs (e.g. SUSY neutralinos)
$\rightarrow$ well motivated theoretically \& good detection prospects
- Axions
$\rightarrow$ motivated by QCD strong CP
- Asymmetric Dark Matter
$\rightarrow$ motivated by $\Omega_{D M} \approx 5 \Omega_{b}$
- Sterile neutrinos
$\rightarrow$ new physics already required in neutrino sector
- DM with only gravitational interactions
$\rightarrow$ Nightmare scenario!


## Thermal Relic Dark Matter

(1) Dark matter initially in thermal equilib:

$$
\chi \chi \leftrightarrow \bar{f} f
$$

(2) Universe cools and the non-relativistic DM is Boltzmann suppressed: $N \sim(m T)^{3 / 2} e^{-m / T}$
(3)"Freeze out" at $\mathrm{m} / \mathrm{T} \approx 20$.

$$
N=\text { constant } \propto \frac{1}{\langle\sigma v\rangle}
$$


$\rightarrow$ Final dark matter abundance proportional to inverse of the annihilation cross section.

## "WIMP Miracle"

* The thermal relic picture sets the "natural scale" for the dark matter annihilation cross section:

$$
\Omega_{D M} \sim 0.2 \text { implies }\langle\sigma v\rangle \sim 2 \times 10^{-26} \mathrm{~cm}^{3} \mathrm{~s}^{-1}
$$

* Suggests electroweak-scale parameters since:

$$
\langle\sigma v\rangle \sim \frac{\alpha^{2}}{(100 \mathrm{GeV})^{2}} \sim 10^{-26} \mathrm{~cm}^{3} \mathrm{~s}^{-1}
$$

$\rightarrow$ 1) A compelling argument, given we have other reason to expect new physics at the $\mathrm{GeV}-\mathrm{TeV}$ scale.
$\rightarrow$ 2) Realistic prospects of detection:

- annihilation signals (indirect detection)
- nuclear recoils (direct detection)
- monojets+missing ET (colliders)


## "WIMPless" Miracle?

Actually, thermal freezeout does not single out the electroweak scale. The relic density simply sets

$$
\Omega_{X} \propto \frac{1}{\langle\sigma v\rangle} \sim \frac{m_{X}^{2}}{g_{X}^{4}}
$$

$\rightarrow$ we can choose any m or g , provided we fix the ratio
Note: Partial wave unitarity bounds the cross section

$$
\left(\sigma_{J}\right)_{\max } v^{\prime} \mathrm{rel}=\frac{4 \pi(2 J+1)}{m_{X}^{2} v^{\prime} \text { rel }}
$$

Griest \& Kamionkowski
$\rightarrow$ rules out thermal relic DM for very large masses.

$$
\langle\sigma v\rangle=\langle\sigma v\rangle_{\text {thermal }} \Rightarrow m_{\chi}<300 \mathrm{TeV}
$$

## Complementary ways to probe (non-gravitational) DM interactions



## Indirect detection

Search for DM annihilation (or decay) products from regions where DM density is high and (ideally) backgrounds are low:


## Indirect detection with MW dwarfs

## Dwarf spheroidal galaxies

Gammaray
Space Telescope
$\checkmark$ dSphs are DM dominated systems (they have very high M/L ratios).
$\checkmark$ Many dSphs are closer than 100 kpc to the Galactic Centre.
$\checkmark$ Low background
Negligible astrophysical backgrounds
$\rightarrow$ robust limits


## Fermi dwarf results

Fermi-LAT, Ackermann et al, arXiv:1310.0828,


## Cluster limits

Galaxy clusters have large DM density and low background $>$ Good for indirect detection

Uncertainties:
$>$ DM density profile
$>$ Existence of sub-halos (clumpiness boost factor)
$>$ Gamma emission induced by cosmic rays

## Very strong limits! Rule out

 some proposed signals!

Figure 9. Upper limits for the DM annihilation cross-section in the $\tau^{+} \tau^{-}$channel. Line styles are as in Fig. 6, but only the EXT results are shown. The black dashed line is the dwarf galaxy constraint. (Geringer-Sameth \& Koushiappas 2011)

## CMB limits on DM annihilation

- Recombination history of the universe could be modified if DM annihilations inject energy into the photon-baryon plasma.
- Limits depend on:
$\rightarrow$ the fraction of the DM energy absorbed by the plasma $\rightarrow$ typical value $\mathrm{f}=0.2$ (larger for annihilation to electrons)
$\rightarrow$ Velocity dependence of the cross section
$\rightarrow$ If $p$-wave suppressed, annihilation rate is very small
- Currently exclude thermal relics with $m<5 \mathrm{GeV}$


## CMB limits on DM annihilation

Madhavacheril, Sehgal \& Slatyer, arXiv:1310.3815


## CMB limits on DM annihilation



## Indirect detection: positrons?




## DM annihilation signal?

Or maybe pulsars?

Possible contribution from Geminga pulsar.
Yuksel, Kistler and Stanev, PRL 2009

## Gamma ray lines - the smoking gun...

 Fermi Gamma ray line search from $5-300 \mathrm{GeV}$


Fermi 1305.5597
No globally significant line signal
Note: gamma ray lines should be loop suppressed, thus subdominant to continuum gammas.

## Fermi gamma ray line at ~ 130 GeV ?



Weniger 1204.2797, and several other groups.


A surprise! Remember, gamma ray lines are loop suppressed. Official Fermi-LAT analysis with more data found a lower significance.

## Galactic center emission - dark matter?

Extended source of $1-3 \mathrm{GeV}$ gamma ray emission at the Galactic centre is seen in Fermi-LAT data. (Hooper et al. + other groups.)

Spatial distribution consistent with DM distribution
Can be fit by annihilation to

- bbar with 40 GeV DM mass
- tau+tau- with 10 GeV DM mass
with a cross section roughly consistent with a thermal relic.
BUT, unresolved point sources (e.g. millisecond pulsars) can mimic this signal.

And, at these low DM masses, in tension with indirect detection limits from clusters.

## Galactic center emission - dark matter?



Abazajian et al, arXiv:1402.4090

## Comparing limits



Bell, Horiuchi \& Shoemaker, arXiv:1408.xxxx

# Interpreting the GC emission in an asymmetric DM model 



## ADM $\rightarrow$ tau+tau- channel consistent with thermal freezeout

## Direct Detection

## in the summer, DM "wind" moving against wind


in the winter, moving away from wind

## Spin-independent vs Spin-dependent

Spin independent - DM interacts coherently with whole nucleus

- $\mathrm{A}^{2}$ enhancement of cross section

Spin dependent - DM couples to spin of nucleus
(a)Operators for Dirac fermion DM

| Name | Operator | Dimension | $\mathrm{SI} / \mathrm{SD}$ |
| :--- | :---: | :---: | :---: |
| D1 | $\frac{m_{q}}{\Lambda^{3}} \bar{\chi} \chi \bar{q} q$ | 7 | SI |
| D5 | $\frac{1}{\Lambda^{2}} \bar{\chi} \gamma^{\mu} \chi \bar{q} \gamma_{\mu} q$ | 6 | SI |
| D8 | $\frac{1}{\Lambda^{2}} \bar{\chi} \gamma^{\mu} \gamma^{5} \chi \bar{q} \gamma_{\mu} \gamma^{5} q$ | 6 | SD |
| D9 | $\frac{1}{\Lambda^{2}} \bar{\chi} \sigma^{\mu \nu} \chi \bar{q} \sigma_{\mu \nu} q$ | 6 | SD |
| D11 | $\frac{\alpha_{s}^{s}}{\Lambda^{3}} \bar{\chi} \chi G^{\mu \nu} G_{\mu \nu}$ | 7 | SI |

(b)Operators for Complex scalar DM

| Name | Operator | Dimension | SI/SD |
| :--- | :---: | :---: | :---: |
| C1 | $\frac{m_{q}}{\Lambda^{2}} \phi^{\dagger} \phi \bar{q} q$ | 6 | SI |
| C 3 | $\frac{1}{\Lambda^{2}} \phi^{\dagger} \overleftrightarrow{\partial}_{\mu} \phi \bar{q} \gamma^{\mu} q$ | 6 | SI |
| C 5 | $\frac{\alpha_{s}}{\Lambda^{2}} \phi^{\dagger} \phi G^{\mu \nu} G_{\mu \nu}$ | 6 | SI |

Further operators, not shown, have a velocity suppressed WIMP-nucleon cross section

## Direct Detection Results



## Future: "neutrino floor" is irreducible background

Coherent neutrino-nucleus scattering of solar/atmospheric neutrinos!


## Spin-dependent WIMP-nucleon scattering



Xenon100 1301.6620

## Solar WIMPS

- Dark matter accumulates and annihilates in the centre of Sun
- Only neutrinos escape the Sun $\rightarrow$ IceCube, SuperK
- Capture determined by WIMP-nucleon scattering cross section.
$\rightarrow$ capture rate $=$ annihilation rate (in equilibrium)
- Hence probes the same quantity as direct detection experiments
$\rightarrow$ Competitive limits for spin-dependent cross sections



## Spin-dependent WIMP-nucleon scattering

IceCube solar WIMP limits more sensitive than nuclear recoil experiments.


Complementary: Collider, nuclear recoil, and solar wimp searches rely on very different assumptions.

## LHC vs direct detection

## Spin-independent

## Spin-dependent



## DAMA/LIBRA annual modulation



Now $8.9 \sigma$ confidence level

## Is the DAMA signal really DM ?

## Something is modulated

There is strong motivation to check the systematics with an experiment in the southern hemisphere.
$\rightarrow$ True DM signal should have the same modulation phase.
$\rightarrow$ The phase of a background modulation could be expected to change with location (sessional variation of atmosphere, etc).
$\rightarrow$ Various proposed southern hemisphere experiments :

- South pole (DMIce)
- Chile-Argentina (ANDES)
- Australia


## DM in

## Australia!

## Mine identified in Stawell

## (near Melbourne)

Studies to assess suitability of the site are underway.
(E. Barberio et al.)


## Beyond WIMPs

## Leptophilic WIMP?

- Suppose DM couples only to leptons (at tree level)
- Usual direct detection and mono-jet bounds not applicable.
- Even so, this scenario is strongly constrained


Direct detection loop-suppressed, yet still yields strong limits

Collider production via Drell-Yan process

Bell et al 1407.4566.
See also: Kopp 0907.3159
Altmannshofer 1406.1269

## Leptophilic WIMP

Direct detection still requires the new-physics scale to be high
$\rightarrow$ some tension with relic density requirement

Bell et al 1407.4566.




## Sterile neutrino dark matter

o keV sterile neutrinos $\rightarrow$ good candidate for warm DM
o Produced in early universe via active-sterile oscillations
o Exclusion principle prevents arbitrary high density $\rightarrow$ dense galaxies set lower limit on mass (Tremaine-Gunn bound)
o Unstable. Decays produce $x$-ray line.


## Possible signal in Perseus and other galaxy

 clusters.$$
\begin{aligned}
& \mathrm{E}=3.57 \pm 0.02 \mathrm{keV} \\
& \Rightarrow \mathrm{~m}=7.1 \mathrm{keV}
\end{aligned}
$$

Bulbul et al 1402.2301, ApJ


## Caution: many nearby atomic transition lines

## Sterile neutrino DM parameter space



Horiuchi et al 1311.0282

## Asymmetric dark matter

Two birds with one stone:
(i) Relic DM abundance
(ii) baryon-antibaryon asymmetry

* Motivation: $\Omega_{D M} \approx 5 \Omega_{b}$

Assume DM density set by a matter anti-matter asymmetry of the same size as the baryon asymmetry.
then $\mathrm{n}_{\mathrm{DM}} \approx \mathrm{n}_{\mathrm{b}}$ (assuming complete asymmetry)
and $\quad m_{D M} \approx 5 m_{b} \approx 5 \mathrm{GeV}$ (prediction for DM mass)

* ADM replaces $\Omega_{D M} \approx \Omega_{b}$ puzzle, with a $m_{D M} \approx m_{b}$ puzzle


## Asymmetric dark matter

## Requirements:

- Mechanism to simultaneously create $\mathrm{B}($ visible) and $\mathrm{B}($ dark $)$ asymmetries, or create an asymmetry in one sector and communicate it to the other.
- Sufficiently large DM annihilation cross section to annihilate the symmetric part (to leave only particles and no antiparticles).


## Implications:

- Light DM.
- Suppressed indirect detection (nothing to annihilate with)
- Large annihilation cross section means either sizeable couplings with SM particles, or else new light degrees of freedom.


## ADM annihilation cross section

## WIMPs - relic density set by annihilation cross section

ADM - relic density set by asymmetry, provided annihilation cross section is big enough to remove the symmetric part
$\rightarrow$ still need a WIMP-like cross section!
Fractional asymmetry: $\quad r \equiv \frac{n(\bar{\chi})}{n(\chi)}$

$$
r_{\infty} \approx \exp \left[-2\left(\frac{\sigma_{0}}{\sigma_{0, \mathrm{WIMP}}}\right)\left(\frac{1-r_{\infty}}{1+r_{\infty}}\right)\right] \xrightarrow{r_{\infty} \ll 1} \exp \left[-2 \sigma_{0} / \sigma_{0, \mathrm{WIMP}}\right]
$$

For $\mathrm{r}_{\infty}<0.1$, require: $\quad \sigma_{0} \gtrsim 1.4 \sigma_{0, \mathrm{WIMP}}$
Graesser et al., arXiv:1103.2771

## ADM - indirect detection limits

Bell, Horiuchi \& Shoemaker, arXiv:1408.xxxx,


## Current limits



## Future limits

## Dark Radiation

o Dark radiation = relativistic dark particles. E.g. dark photons,dark neutrinos, or similar.
o Needed in some models (e.g. many asymmetric DM models)
o This radiation need not have the same temperature as ordinary radiation.
o e.g. models with "dark atoms" have hydrogen-like states formed from two oppositely charged particles interacting via a massless U(1) $)_{\mathrm{D}}$ gauge boson.

The dark sector may have a particle spectrum as rich as the visible sector

- Dark Radiation leaves imprint in CMB (usually discussed in term of "effective number of neutrinos")
- CMB accommodates extra radiation:


Archidiacono et al, 1307.0637


- Can be more complicated: dark matter-dark radiation coupling leads to "dark acoustic oscillations"


## Dark Matter Self Interactions

Dark matter should not strongly self interact.

- The Bullet Cluster
- Halo shapes (self interactions make galaxies too spherical)
* But some amount of self interaction is usually expected.

This is ok, and maybe even be desired:
$\rightarrow$ helps to alleviate the CDM problem of too much structure on small scales. However, there are other solutions to this problem, including warm dark matter, decaying dark matter, ...

## Outlook

o WIMP...is this idea compelling, or are we searching under the lamp post?
o If DM is a conventional WIMP, discovery should be close!
o ADM...is the similarity of the dark and visible matter densities an important clue, or just a red herring?

O Direct detection, indirect detection, colliders, solar WIMP searches, cosmological probe...all have good sensitivity and provide complementary information.
o We should remember that many dark-sector models have a rich spectrum of new particles. Indeed, DM may be multi-component.

## Extra slides

## Future: "neutrino floor" is irreducible background

Coherent neutrino-nucleus scattering of solar/atmospheric neutrinos!



## Cushman et al arXiv:1310.8327

## Linking dark matter and baryogensis

## Connect (i) Relic DM abundance

(ii) baryon-antibaryon asymmetry

Various ideas: Asymmetric dark matter, WIMPy baryogensis, Baryomorphosis, DM assimilation, ......

|  | Asymmetric <br> dark matter | WIMPy <br> baryogenesis |
| :--- | :---: | :--- |
| WIMP miracle | $\mathbf{x}$ | $\checkmark$ |
| Explain $\Omega_{\mathrm{DM}} \approx \Omega_{\mathrm{b}}$ | $\checkmark$ | $\mathbf{x}$ |

ADM: Many papers! See reviews by Petraki and Volkas 1305.4939 and Zurek 1308.0388.
WIMPy baryogenesis: Cui, Randall and Shuve, 1112.2704; Bernal et al., 1210.0094,
Bernal et al., 1307.6878; Kumar \& Stengel, 1309.1145
Baryomorphosis: McDonald 1009.3227 Dark matter assimilation: D’Eramo et al., 1111.5615

## WIMPy baryogensis

Require WIMP annihilation satisfy the Sakharov conditions
$\rightarrow$ a baryon asymmetry can be generated from DM annihilations


## Asymmetry in exotic antibaryons, which decay to SM baryons

## dSph radio limits

Most annihilation channels produce e- and e+
e+, e- loose energy (multiple processes) as they propagate
> Includes synchrotron radiation, at radio wavelengths
$>$ Rates depend upon diffusion assumptions, especially magnetic field strengths


Regis et al, arXiv:1407.4948

## current

## future




Regis et al, arXiv:1407.4948

## Effective operators for DM interactions

Model-independent description of
DM interactions with SM particles:

$$
\begin{aligned}
& L_{E f f}=\frac{1}{\Lambda_{\text {eff }}^{2}} \bar{\chi} \Gamma_{\chi} \chi \bar{q} \Gamma_{q} q \\
& \Gamma_{\chi, q} \in\left\{1, \gamma^{5}, \gamma^{\mu}, \gamma^{\mu} \gamma^{5}, \sigma^{\mu \nu}\right\} .
\end{aligned}
$$

## Advantages:

- model-independent description

Disadvantages:

- breaks down if $q^{2}$ is large or mediators light

| Name | Operator | Coefficient | DD |
| :---: | :---: | :---: | :---: |
| D1 | $[\bar{\chi} \chi][\bar{f} f]$ | $m_{f} \Lambda^{-3}$ | SI |
| D2 | $\left[\bar{\chi} \gamma^{5} \chi\right][\bar{f} f]$ | $i m_{f} \Lambda^{-3}$ | - |
| D3 | $[\bar{\chi} \chi]\left[\bar{f}{ }^{5} f\right]$ | $i m_{f} \Lambda^{-3}$ | - |
| D4 | $\left[\bar{\chi} \gamma^{5} \chi\right]\left[\bar{f} \gamma^{5} f\right]$ | $m_{f} \Lambda^{-3}$ | - |
| D5 | $\left[\bar{\chi} \gamma^{\mu} \chi\right]\left[\bar{f} \gamma_{\mu} f\right]$ | $\Lambda^{-2}$ | SI |
| D6 | $\left[\bar{\chi} \gamma^{\mu} \gamma^{5} \chi\right]\left[\bar{f} \gamma_{\mu} f\right]$ | $\Lambda^{-2}$ | - |
| D7 | $\left[\bar{\chi} \gamma^{\mu} \chi\right]\left[\bar{f} \gamma_{\mu} \gamma^{5} f\right]$ | $\Lambda^{-2}$ | - |
| D8 | $\left[\bar{\chi} \gamma^{\mu} \gamma^{5} \chi\right]\left[\bar{f} \gamma_{\mu} \gamma^{5} f\right]$ | $\Lambda^{-2}$ | SD |
| D9 | $\left[\bar{\chi} \sigma^{\mu \nu} \chi\right]\left[\bar{f} \sigma_{\mu \nu} f\right]$ | $\Lambda^{-2}$ | SD |
| D10 | $\left[\bar{\chi} \sigma^{\mu \nu} \gamma^{5} \chi\right]\left[\bar{f} \sigma_{\mu \nu} f\right]$ | $i \Lambda^{-2}$ | - |
| D11 | $[\bar{\chi} \chi]\left[G_{\mu \nu} G^{\mu \nu}\right]$ | $\alpha_{S} \Lambda^{-3}$ | SI |
| D12 | $\left[\bar{\chi} \gamma^{5} \chi\right]\left[G_{\mu \nu} G^{\mu \nu}\right]$ | $i \alpha_{S} \Lambda^{-3}$ | - |
| D13 | $[\bar{\chi} \chi]\left[G_{\mu \nu} \tilde{G}^{\mu \nu}\right]$ | $i \alpha_{S} \Lambda^{-3}$ | - |
| D14 | $\left[\bar{\chi} \gamma^{5} \chi\right]\left[G_{\mu \nu} \tilde{G}^{\mu \nu}\right]$ | $\alpha_{S} \Lambda^{-3}$ | - |

## Strong bounds on EFT operators!

Bounds on EFT operators are becoming quite constraining!

* Relic density
$\rightarrow$ upper limit on $\Lambda_{\text {eff }}$ (to prevent over-closure)
* Direct detection, collider, and indirect detection
$\rightarrow$ lower limits on $\wedge_{\text {eff }}$ (no signals)
For many operators, these limits are approaching!
If the EFT description is relevant for DM, we should see a signal soon!


## Dark Matter at the LHC

$\square$ The dominant DM production process is invisible (DM stable, weakly interacting) : $\bar{q} q \rightarrow \chi \chi$

Need visible particles in the final state, to recoil against missing transverse energy

$$
\bar{q} q \rightarrow \chi \chi+\text { single SM particle }
$$

Mono-X process in which DM is visible as a high pT state + missing ET
$\rightarrow$ Mono-jet, mono-photon, mono-Z, mono-W, mono-Higgs


## Beyond an EFT:

## t-channel scalar

 mediator
(a1)
H.An et al, 1308.0592

See also:
Chang et al. , 1307.8120
Bai \& Berger, 1308.0612
DiFranzo et al., 1308.2679

## Mediator pair production

