## Status of and prospects for W mass and Z mass measurements





On behalf of the **CMS and ATLAS Collaborations** 



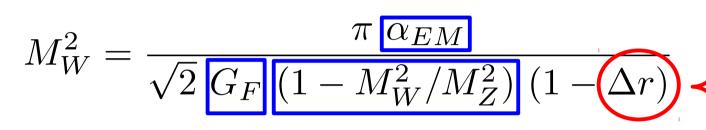
Physics at LHC and beyond Qui-Nhon, Vietnam 10-17 August 2014



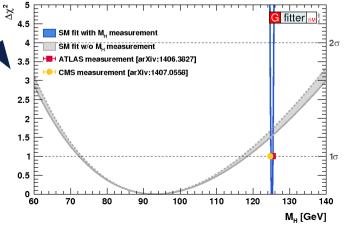
### **Motivation**

The EW sector of the SM, relates  $M_W$  to  $\alpha$ ,  $G_E$ , and  $\sin^2\theta_W$ 

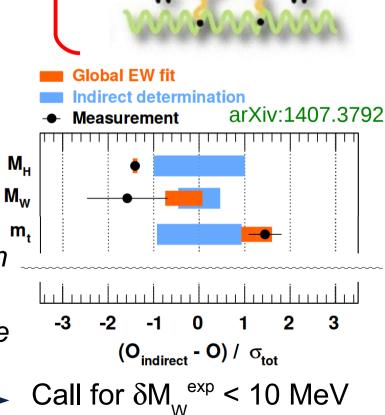
Radiative corrections ∆r are dominated by Top and Higgs loops



- The relation between M<sub>top</sub>, M<sub>H</sub> and M<sub>W</sub>
   provides a stringent test of the SM
- The comparison between the measured M<sub>H</sub> and the predicted M<sub>H</sub> is sensitive to new physics



Indirect determination of  $M_w$  (±8 MeV) is more precise than the experimental measurement



### Status and prospects for W mass measurement

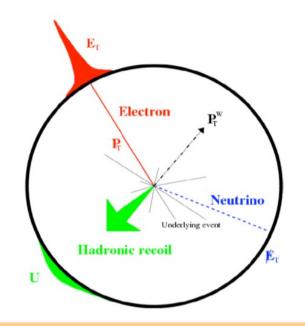
- Methodology for the M<sub>w</sub> extraction
- The Tevatron experience on M<sub>w</sub>
- Measurements of M<sub>w</sub> and M<sub>z</sub> at the LHC
- Theoretical uncertainties for M<sub>w</sub> at the LHC
- ATLAS and CMS measurements to constrain uncertainties on M<sub>w</sub>
- Prospect for M<sub>w</sub> and M<sub>z</sub> measurements at future colliders

### Methodology for the W mass extraction

- Event selection: W leptonic decay  $W \rightarrow I \nu$ , I = e,  $\mu$
- The full kinematic of the W decay cannot be reconstructed, since the longitudinal momentum of the neutrino is unknown

Traditional analyses are based on a template fit extraction from observables sensitive to  $\mathbf{M}_{\mathrm{w}}$ 





Lepton transverse momentum	$p_T^l$
W transverse mass	$M_T = \sqrt{2 \cdot p_T^l p_T^{\nu} \cdot (1 - \cos \Delta \phi(l, \nu))}$
Neutrino transverse momentum (from hadronic recoil)	$p_T^ u$

More sophisticated analysis techniques suggest simultaneous measurements of W and Z observables TS2008-022 Eur. Phys. J. C69 (2010) 379-397

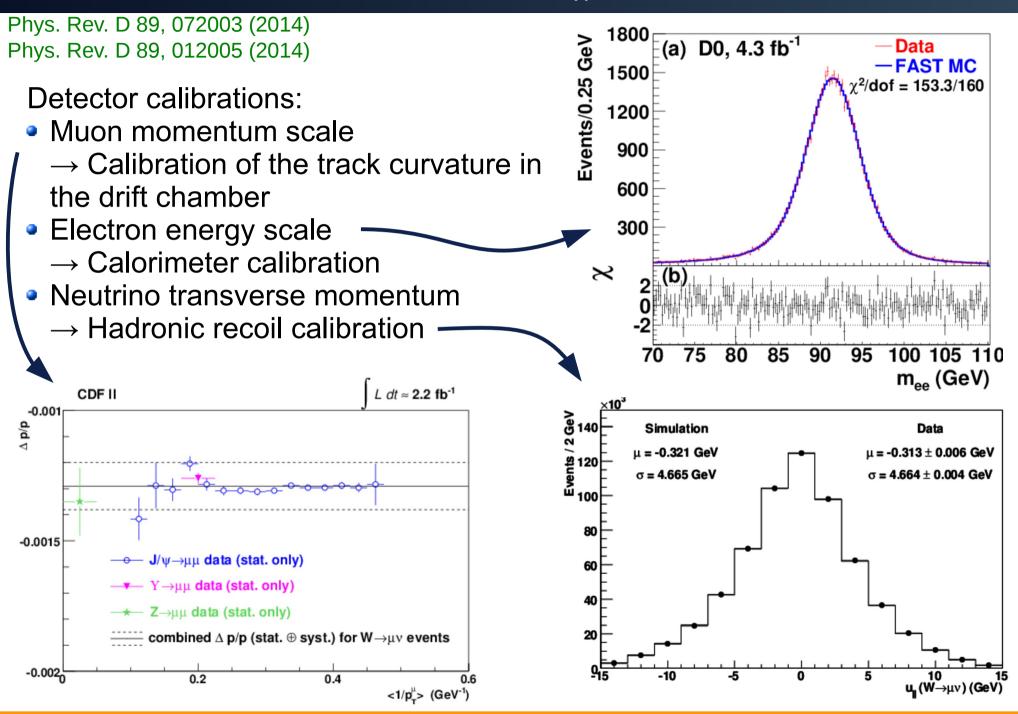


events to constraint both experimental and theory systematics

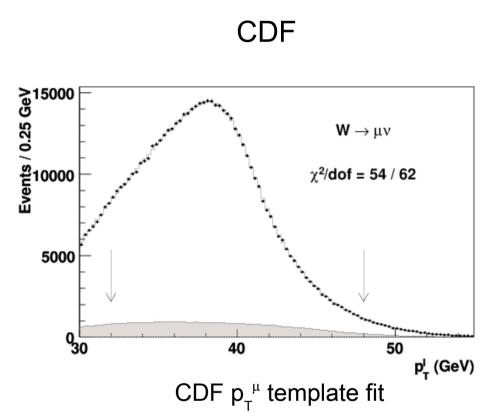
In the same spirit, a common strategy of

template fits analyses is to use  $Z \rightarrow II$ 

## Tevatron measurements of M<sub>w</sub> - Methodology



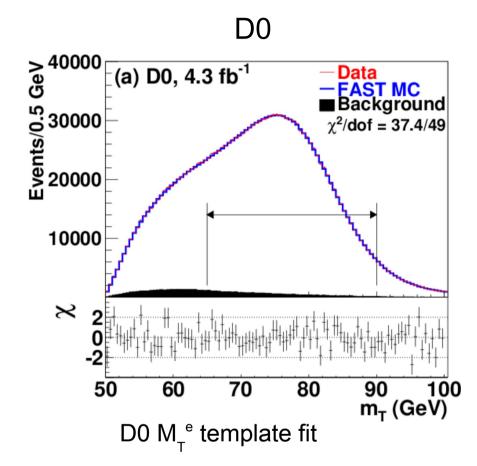
## Tevatron measurements of M<sub>w</sub> - Results



 $M_W$  extracted from  $p_T^{-1}$ ,  $M_T$ ,  $p_T^{-\nu}$  in electron and muon channels  $L = 2.2 \text{ fb}^{-1}$ 

$$M_W = 80.387 \pm 19 \text{ MeV}$$

Phys. Rev. D 89, 072003 (2014)



 $M_{W}$  extracted from  $p_{T}^{-1}$ ,  $M_{T}$  in electron channel  $L = 4.3 (+1) \text{ fb}^{-1}$ 

$$M_W = 80.375 \pm 23 \text{ MeV}$$

Phys. Rev. D 89, 012005 (2014)

### W mass measurement at the LHC

- The M<sub>w</sub> measurement at the LHC follows a strategy similar to the Tevatron
- Important differences:
  - Higher pile-up environment → affect hadronic recoil calibration
  - Potentially larger theoretical uncertainties due to pp instead of pp collisions
  - W<sup>+</sup> and W<sup>-</sup> production is not symmetric → Require a charge dependent analysis

Most precise observables for the M<sub>w</sub> extraction

p <sub>T</sub> I	$M_{T}$
Observable does not depend on hadronic recoil, smaller experimental uncertainty	Depends on hadronic recoil measurement, expected larger experimental uncertainties
Larger theory uncertainty due to higher order QCD, $p_{\scriptscriptstyle T}^{\ \ W}$ modelling, PDF, W polarisation, charm mass	M <sub>T</sub> is quite stable wrt perturbative QCD corrections, smaller PDF uncertainties, smaller non-perturbative QCD uncertainties



M<sub>w</sub> extraction is likely to be limited by theoretical uncertainties



Final balance between theory and expuncertainties will depend on pile-up mitigation algorithms

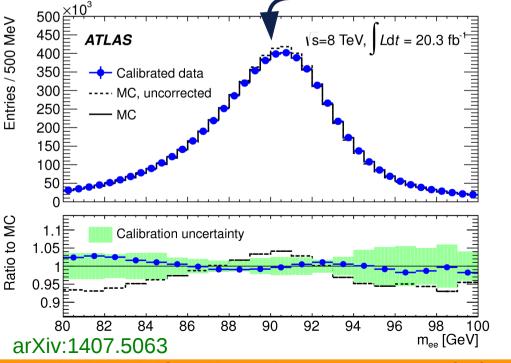
### Z mass measurement at the LHC

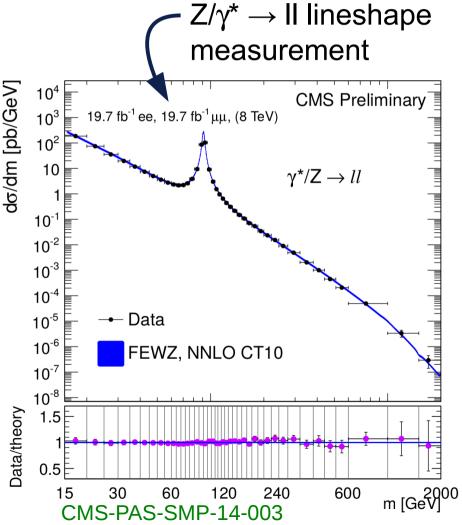
A first step towards the measurement of M<sub>w</sub> at the LHC is the measurement of M<sub>z</sub>

Test the methodology of the  $M_w$  extraction:  $\rightarrow$  Neglect one of the two leptons, extract  $M_{_{7}}$  from  $p_{_{T}}^{^{-1}}$  and hadronic recoil

The lepton energy scale is calibrated by comparing the reconstructed  $M_z$  to the LEP measurement of  $M_z$ 

Electron calibration from  $Z \rightarrow ee$  invariant mass





## Theory systematics for M<sub>w</sub> at the LHC



Why theory systematic uncertainties are crucial for the W mass measurement at the LHC?

W production at the LHC is relatively cheap 7(8) TeV, 5(20) fb<sup>-1</sup>  $\rightarrow$  ~15(75)x10<sup>6</sup> in W  $\rightarrow$  Iv (I = e,  $\mu$ )



- Statistics is not an issue
   → expected ~6(2) MeV
   statistical uncertainty on M<sub>w</sub>
- Main experimental systematic uncertainties
  - Lepton energy scale, resolution, efficiency
- hadronic recoil scale and resolution Estimates depend on the analysis strategy: between 5-15 MeV

Eur.Phys.J.C57:627-651,2008 TS2008-022

Theory uncertainties at the Tevatron: 10-15 MeV Expected similar or higher uncertainties at the LHC

## Constrain M<sub>w</sub> theory uncertainties at the LHC

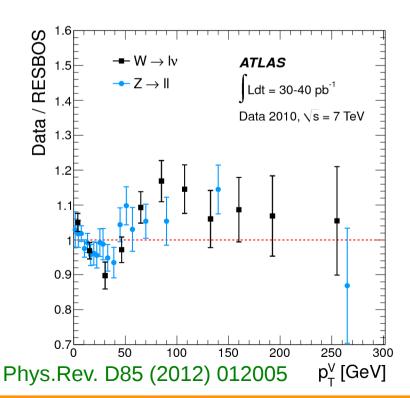
ATLAS and CMS are performing measurements of alternative W, Z observables to control the theoretical models and reduce the uncertainties on the measurement of  $\mathbf{M}_{\mathrm{W}}$ 

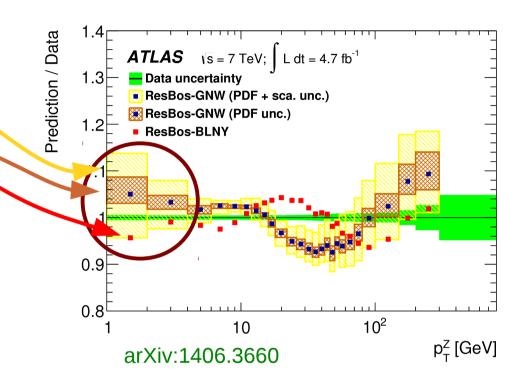
Theory uncertainties	Measurements which can provide constraints to the theoretical models
p <sub>T</sub> <sup>W</sup> modelling	$p_T^{W}, p_T^{Z}$
PDF	W asymmetry, Z rapidity, W + charm
\\/ nologization	Angular coefficients in W, Z
W polarization	production

## p<sub>T</sub><sup>W</sup> and p<sub>T</sub><sup>Z</sup> modelling

Theory uncertainties at low W, Z p<sub>T</sub>

- Factorization, renormalization, resummation or PS scales
- PDF, low x gluon PDF
- Non perturbative QCD parameters.
- Strong-interaction coupling  $\alpha_s(M_z)$
- Heavy flavour masses, HF thresholds in the PDF





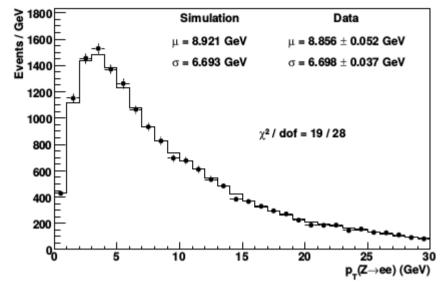
- Measurement of p<sub>T</sub><sup>W</sup> is less precise in the low p<sub>T</sub> region
- However it provides an important cross check

## Constraining p<sub>T</sub><sup>W</sup> and p<sub>T</sub><sup>Z</sup> theory uncertainties

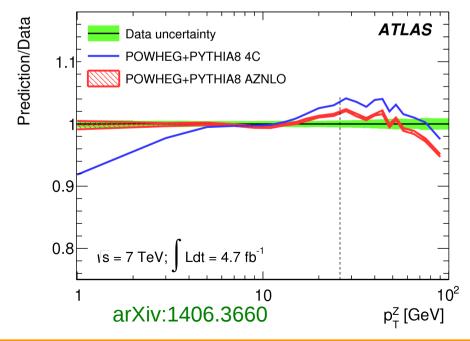
#### CDF W mass measurement

- RESBOS used to model p<sub>T</sub><sup>W,Z</sup>
- Fit non-pQCD parameter  $\mathbf{g_2}$  and  $\alpha_{_{\! S}}$  to  $\mathbf{p_{_{\! T}}}^{_{\! Z}}$  data
- Uncertainties propagated to M<sub>w</sub>
   3, 9, 4, MeV for M<sub>T</sub>, p<sub>T</sub><sup>T</sup>, p<sub>T</sub><sup>V</sup>

$$S = \left[g_1 - g_2 \log \left(\frac{\sqrt{\hat{s}}}{2Q_0}\right) - g_1 g_3 \log(100x_j x_k)\right] b^2$$



Phys. Rev. D 89, 072003 (2014)

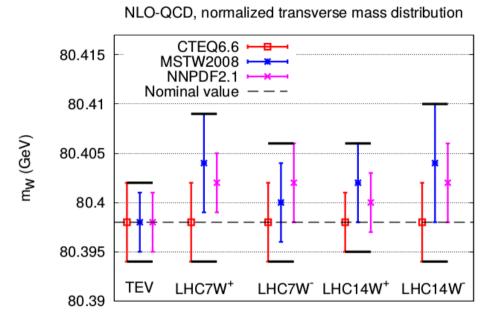


#### **ATLAS**

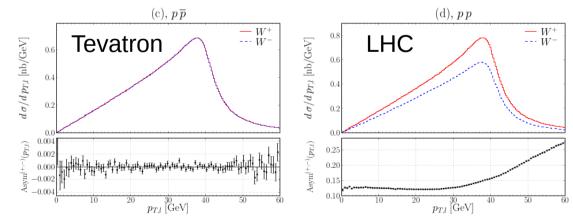
- POWHEG+PYTHIA8 model
- Fit non-pQCD parameters primordial k<sub>T</sub> and ISR cut-off to p<sub>T</sub><sup>Z</sup> data

### PDF uncertainties

PDF uncertainties on M<sub>W</sub> extraction from M<sub>T</sub> (DYNNLO) Phys.Rev.D83:113008,2011 arXiv:1309.1311



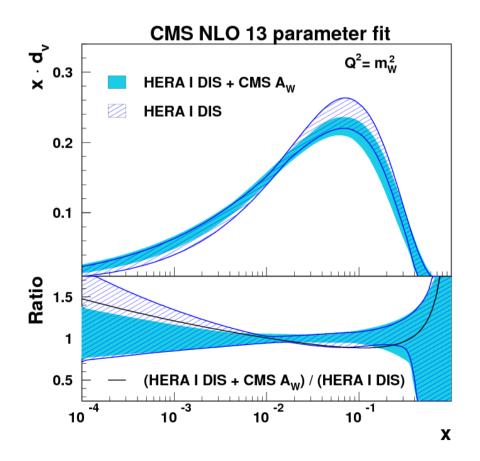
- Normalised distribution are essential to reduce PDF uncertainties
- PDF uncertainties are expected to be larger at the LHC than at the Tevatron, and different between W<sup>+</sup> and W<sup>-</sup>
- Estimated small  $\alpha_{_{\rm S}}$  and  $m_{_{\rm C}}$  uncertainties
- pp collision → no symmetry between W<sup>+</sup> and W<sup>-</sup> lepton p<sub>T</sub> due to u<sub>v</sub> ≠ d<sub>v</sub> and W polarization → Larger PDF uncertainties at LHC than Tevatron
- At the LHC, a charge-dependent separate analysis of W<sup>+</sup> and W<sup>-</sup> helps to reduce PDF uncertainties

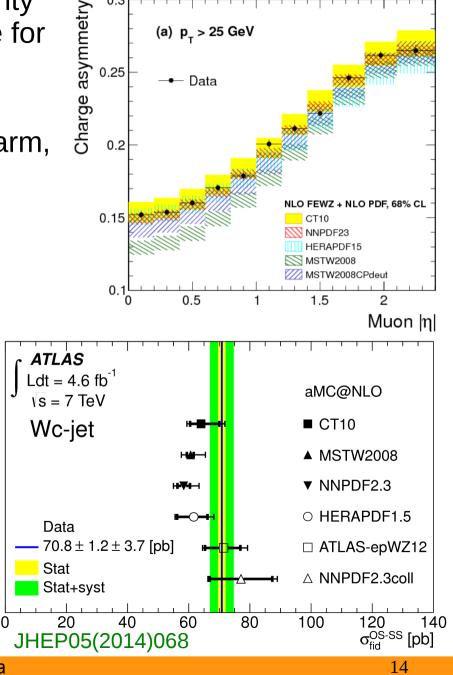


Eur. Phys. J. C69 (2010) 379-397

### PDF uncertainties

- Crucial to reduce the valence PDF uncertainty
- ATLAS and CMS have provided several measurements of W, Z inclusive and W+charm, which can constrain PDF uncertainties





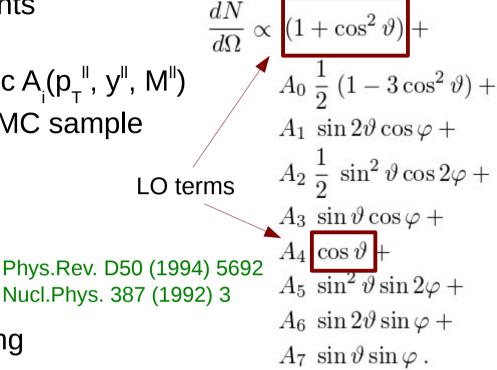
arXiv:1312.6283

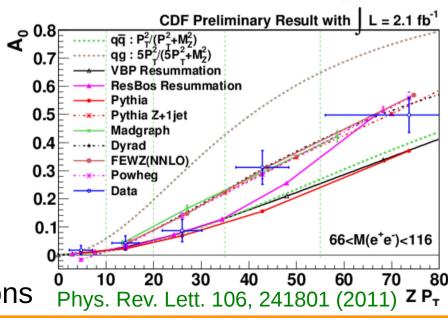
### W, Z polarization

- Set of 8 observables: angular coefficients
   A<sub>i</sub> → ratio of helicity cross sections
- A<sub>i</sub> are functions of the leptons kinematic  $A_i(p_T^{\parallel}, y^{\parallel}, M^{\parallel})$
- A<sub>i</sub> coefficients can be calculated form MC sample with moments method

$$\langle m \rangle = \frac{\int d\sigma(p_T,y,\theta,\phi) \ m \ d\cos\theta \ d\phi}{\int d\sigma(p_T,y,\theta,\phi) \ d\cos\theta \ d\phi}$$

- A<sub>i</sub> can be measured precisely for Z,
   the W measurement is more challenging
  - Related to boson polarization, V-A coupling
  - Provide insight into QCD and EW dynamics
  - Stringent test of predictions and MC generators
- A<sub>0</sub>-A<sub>4</sub> coefficients measured at CDF
- Precise measurements at the LHC of A<sub>0</sub>-A<sub>7</sub>
   can discriminate between different predictions





## Prospects for M<sub>w</sub> at ILC and TLEP

M<sub>w</sub> can be measured at e<sup>+</sup>e<sup>-</sup> colliders through an energy scan of the WW production threshold

Near threshold, the WW cross section is proportional to the non-relativistic W velocity

$$\sigma(WW) \propto \beta_W$$

arXiv:1306.6352

ILC Giga-Z program

- Energy scan 160 to 170 GeV
- $\delta M_W = 6-7 \text{ MeV}$

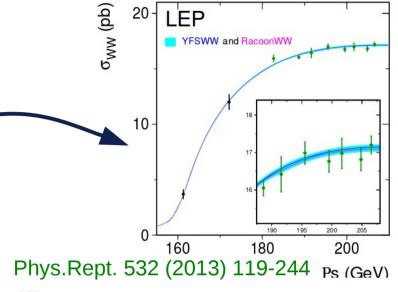
JHEP 1401 (2014) 164

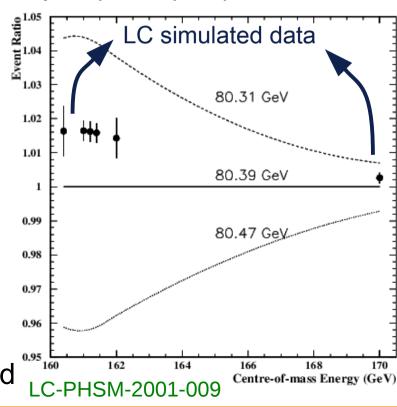
TLEP OkuW program

- $\delta M_{\rm W}$  = 0.5 MeV
  - → dominated by statistical uncertainty

Dominant theory uncertainties

- Initial state QED corrections
- Parametrization of cross section near threshold





## Prospects for M<sub>2</sub> at TLEP

TeraZ program: 7 x 10<sup>11</sup> visible Z in one year of data taking at 91 GeV

M<sub>2</sub> extracted from Z lineshape

Statistical uncertainty: 5 keV

Systematic uncertainty: 100 keV

→ dominated by beam energy calibration

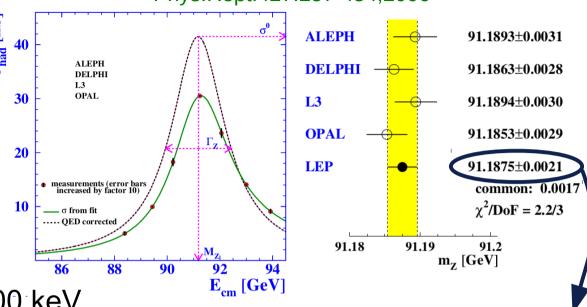
Dominant theory uncertainties

Initial state QED corrections ≤ 100 keV

- Additional lepton pairs ≤ 300 keV
- Lineshape parametrization ≤ 100 keV
  - → Need to revisit theory predictions

 $173200 \pm 900$ 

Phys.Rept.427:257-454,2006



Factor of 20 improvement wrt LEP M<sub>2</sub> measurement

Quantity	Physics	Present	Measured	Statistical	Systematic	Key	Challenge
		precision	from	uncertainty	uncertainty		
m <sub>Z</sub> (keV)	Input	$91187500 \pm 2100$	Z Line shape scan	5 (6) keV	< 100 keV	$E_{\text{beam}}$ calibration	QED corrections
Γ <sub>Z</sub> (keV)	$\Delta \rho \text{ (not } \Delta \alpha_{had})$	$2495200 \pm 2300$	Z Line shape scan	8 (10) keV	< 100 keV	$E_{\text{beam}}$ calibration	QED corrections
$R_{\ell}$	$\alpha_{\rm s}, \delta_{\rm b}$	$20.767 \pm 0.025$	Z Peak	0.00010 (12)	< 0.001	Statistics	QED corrections
$N_{\nu}$	PMNS Unitarity,	$2.984 \pm 0.008$	Z Peak	0.00008 (10)	< 0.004		Bhabha scat.
$N_{\nu}$	and sterile $\nu$ 's	$2.92 \pm 0.05$	Zγ, 161 GeV	0.0010 (12)	< 0.001	Statistics	
$R_{\rm b}$	$\delta_{\mathrm{b}}$	$0.21629 \pm 0.00066$	Z Peak	0.000003 (4)	< 0.000060	Statistics, small IP	Hemisphere correlations
$A_{LR}$	$\Delta \rho$ , $\epsilon_3$ , $\Delta \alpha_{had}$	$0.1514 \pm 0.0022$	Z peak, polarized	0.000015 (18)	< 0.000015	4 bunch scheme, 2exp	Design experiment
mw (MeV)	$\Delta \rho$ , $\epsilon_3$ , $\epsilon_2$ , $\Delta \alpha_{had}$	$80385 \pm 15$	WW threshold scan	0.3 (0.4)MeV	< 0.5 MeV	$E_{beam}$ , Statistics	QED corrections

Input

JHEP 1401 (2014) 164

 $m_{\text{top}} (\text{MeV})$ 

tt threshold scan

10 (12) MeV

< 10 MeV

Statistics

Theory interpretation

## Summary

- The indirect determination of M<sub>w</sub> from the global fit of the SM parameters calls for a target precision of 10 MeV on the measurement of M<sub>w</sub>
- For the measurement of the W mass at the LHC it is necessary to control and reduce many sources of experimental and theoretical uncertainties
- A first step towards the measurement of M<sub>w</sub> is the lepton calibration, and the validation of the analysis techniques through the measurement of M<sub>z</sub>
- ATLAS and CMS have an extensive program of W, Z measurements which can help constraining many of the theoretical uncertainties
- Prospects at future e<sup>+</sup>e<sup>-</sup> colliders ILC (TLEP) for approximately a factor of 2 (20) improvement on  $\delta M_W$  and  $\delta M_Z$

# **BACKUP**

### W mass uncertainties at CDF

$m_T$ fit uncertainties							
Source	$W  o \mu \nu$	$W \rightarrow e v$	Common				
Lepton energy scale	7	10	5				
Lepton energy resolution	1	4	0				
Lepton efficiency	0	0	0				
Lepton tower removal	2	3	2				
Recoil scale	5	5	5				
Recoil resolution	7	7	7				
Backgrounds	3	4	0				
PDFs	10	10	10				
$W$ boson $p_T$	3	3	3				
Photon radiation	4	4	4				
Statistical	16	19	0				
Total	23	26	15				

$p_T^{\nu}$ fit uncertainties							
Source	$W \rightarrow \mu \nu$ $W \rightarrow e \nu$ Correla						
Lepton energy scale	7	10	5				
Lepton energy resolution	1	7	0				
Lepton efficiency	2	3	0				
Lepton tower removal	4	6	4				
Recoil scale	2	2	2				
Recoil resolution	11	11	11				
Backgrounds	6	4	0				
PDFs	11	11	11				
$W$ boson $p_T$	4	4	4				
Photon radiation	4	4	4				
Statistical	22	25	0				
Total	30	33	18				

$p_T^{\ell}$ fit uncertainties							
Source	$W  o \mu \nu$	$W \rightarrow e v$	Common				
Lepton energy scale	7	10	5				
Lepton energy resolution	1	4	0				
Lepton efficiency	1	2	0				
Lepton tower removal	0	0	0				
Recoil scale	6	6	6				
Recoil resolution	5	5	5				
Backgrounds	5	3	0				
PDFs	9	9	9				
$W$ boson $p_T$	9	9	9				
Photon radiation	4	4	4				
Statistical	18	21	0				
Total	25	28	16				

Phys. Rev. D 89, 072003 (2014)

Uncertainties are given in MeV

### W mass uncertainties at D0

Phys. Rev. D 89, 012005 (2014)

Source	Section	$m_T$	$p_T^e$	$E_T$
Experimental				
Electron Energy Scale	$\overline{\mathrm{VIIC4}}$	16	17	16
Electron Energy Resolution	$\overline{ m VIIC 5}$	2	2	3
Electron Shower Model	VC	4	6	7
Electron Energy Loss	$\nabla \mathbf{D}$	4	4	4
Recoil Model	VIID3	5	6	14
Electron Efficiencies	VIIB10	1	3	5
Backgrounds	VIII	2	2	2
$\sum$ (Experimental)		18	20	24
W Production and Decay Model				
PDF	VIC	11	11	14
QED	$\overline{ ext{VIB}}$	7	7	9
Boson $p_T$	$\overline{\text{VIA}}$	2	5	2
$\sum$ (Model)		13	14	17
Systematic Uncertainty (Experimental and Model)		22	24	29
W Boson Statistics	IX	13	14	15
Total Uncertainty		26	28	33

Uncertainties are given in MeV

## PDF uncertainties on M<sub>T</sub>

	CTEQ6.6		MSTW2008		NNPDF2.1		
	$m_W \pm \delta_{ m pdf}$	$ \langle \chi^2 \rangle $	$m_W \pm \delta_{ ext{pdf}}$	$ \langle \chi^2 \rangle $	$m_W \pm \delta_{ ext{pdf}}$	$ \langle \chi^2 \rangle $	$\delta_{ m pdf}^{ m tot}$
Tevatron, $W^{\pm}$	$80.398 \pm 0.004$	1.42	$80.398 \pm 0.003$	1.42	$80.398 \pm 0.003$	1.30	4
LHC 7 TeV $W^+$	$80.398 \pm 0.004$	1.22	$80.404 \pm 0.005$	1.55	$80.402 \pm 0.003$	1.35	8
LHC 7 TeV $W^-$	$80.398 \pm 0.004$	1.22	$80.400 \pm 0.004$	1.19	$80.402 \pm 0.004$	1.78	6
LHC 14 TeV $W^+$	$80.398 \pm 0.003$	1.34	$80.402 \pm 0.004$	1.48	$80.400 \pm 0.003$	1.41	6
LHC 14 TeV $W^-$	$80.398 \pm 0.004$	1.44	$80.404 \pm 0.006$	1.38	$80.402 \pm 0.004$	1.57	8

Phys.Rev.D83:113008,2011

	Tevatron	LHC7W+	LHC7W-	LHC14W+	LHC14W-
$\alpha_s(m_Z) = 0.118$	80.398	80.400	80.398	80.402	80.400
$\alpha_s(m_Z) = 0.119 \text{ (ref)}$	80.398	80.402	80.402	80.400	80.402
$\alpha_s(m_Z) = 0.120$	80.398	80.400	80.398	80.402	80.402

Table 7: Central value of the fit of  $m_w$  obtained with NNPDF2.1, using PDF sets that differ by the  $\alpha_s(m_Z)$  value, for different colliders and energies. The fit has been done on normalized distributions and using normalized templates, and the distributions have been generated at NLO-QCD with DYNNLO.

$m_W$ (GeV)	Tevatron	LHC7W+	LHC7W-	LHC14W+	LHC14W-
$m_c = 1.414 \text{ (ref)}$	80.398	80.402	80.402	80.400	80.402
$m_c = 1.5$	80.398	80.400	80.398	80.398	80.399
$m_c = 1.6$	80.398	80.400	80.400	80.398	80.399
$m_c = 1.7$	80.396	80.400	80.400	80.396	80.398

Table 8: Central value of the fit of  $m_W$  obtained with NNPDF2.1 sets with different values of  $m_c$  for different colliders and energies. We include the default value in NNPDF2.1,  $m_c^2 = 2 \text{ GeV}^2$  as well.