Rencontres du Vietnam 2014: Physics at LHC and beyond Quy Nhon, 12th August 2014



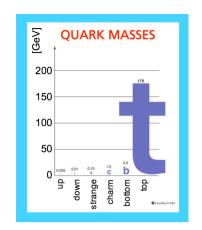
Status of and prospects for top quark mass measurements

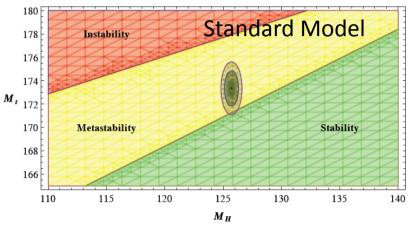
Roberto Tenchini INFN Pisa

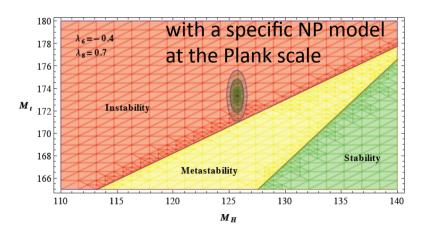
Disclaimer: this is NOT a review talk, the aim is to discuss experimental methods and perspectives. A complete accounting of most recent top mass results is given at http://indico.ific.uv.es/indico/sessionDisplay.py?sessionId=30&confId=2025#20140703

Its highness, the top quark

- Because of the non-decoupling properties of electroweak interactions (Veltman, 1977)
 the top quark gives large contributions to pure EWK radiative corrections (≈G_Fm_t²) hence it is "the quark" for EWK precision physics
- The fate of our universe, i.e. the stability of the EWK vacuum (Degrassi et al., arXiv:1205.6497), might not depend (only) on its mass, nevertheless the top mass value opens a window to new physics up to the Planck scale (Branchina et al., arXiv:1407.4112)



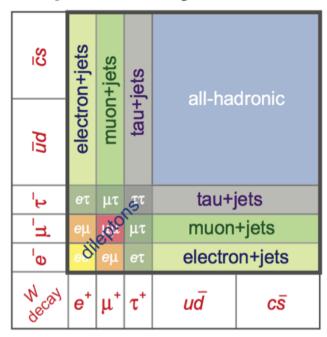




Quick reminder of top properties

- Very high mass \rightarrow very short lifetime: bound states are not formed, almost a free quark $\tau_{top} \approx 0.4 \times 10^{-24} s$
- Hierarchy of CKM elements → only decays to a b quark and a W boson are relevant for this discussion, the decay mode is
 t → Wb
- Dominant production mode at colliders in top pairs, events usually classified and selected according to their W-pair decays
- Top pair decay channels have different rates, different backgrounds, different amount of missing energy (v) providing useful consistency tests

Top Pair Decay Channels



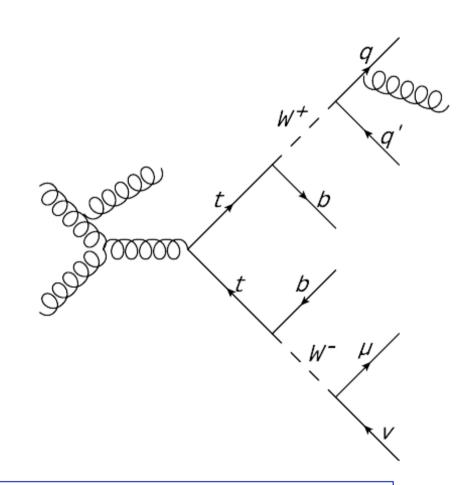
Methods for top mass measurement (1)

- Standard methods at hadron colliders: measure the top mass from the decay products in a specific top pair decay channel
 - from the simplest versions: measure invariant mass of, e.g. three jets in lepton+jets events
 - to the more sophisticated versions: use of the full event information to gain sensitivity, e.g. Matrix Element method
- The standard methods are the most precise with the current statistics
 - they are used in current LHC, Tevatron, World combinations
 - the top mass in EWK fits comes from these methods
- Crucial points for the standard methods
 - accurate calibration of physics objects, in particular Jet Energy Scale: use of kinematic fits for JES calibration in situ, e.g. use the W mass to constraint light quarks jet energy scale (JES) from two-jet invariant mass
 - associate measured objects (jets, leptons, missing E_T) to top candidate: e.g. use b-tagging to choose the right b-jet for the 3-jet combination

Event selection: lepton+jets final state

[example from CMS, TOP-14-001 / JHEP 12 (2012) 105]

- Trigger for isolated muon or electron + jets (pT > 33 GeV)
- Exactly 1 isolated lepton with p_T
 >30 GeV, |η|<2.1 (veto additional isolated e, μ)
- \geq 4 "particle flow" jets (anti-kt, R = 0.5) with $p_T > 30 \text{GeV}, |\eta| < 2.4$
- 2 jets b-tagged among the 4 leading jets
- Composition:
 - 94% ft, 2% W+jets, 3% singletop, 1% other
- 108000 events in 19.5 fb⁻¹ at 8 TeV selected

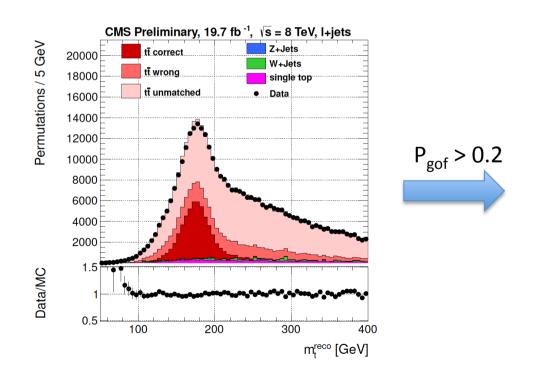


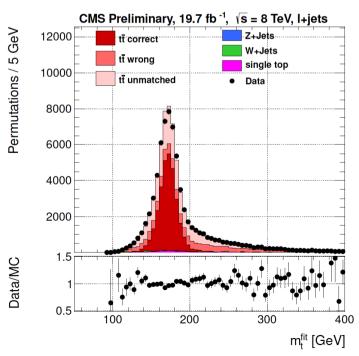
Compare with selections at Tevatron with full statistics: about 2500 events

Event reconstruction

[example from CMS, TOP-14-001 / JHEP 12 (2012) 105]

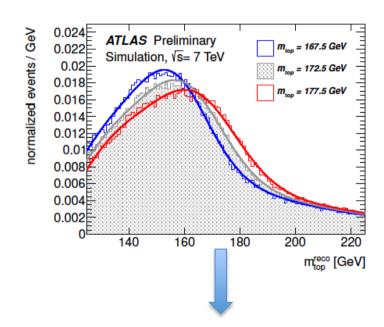
- Assign 4 leading jets to partons from tt decay (obey b-tag)
 - Kinematic fit with constraints: $m_W = 80.4$ GeV, $m_t = m_{tbar}$
 - Weight each permutation by $P_{gof} = exp-(1/2\chi^2)$, select $P_{gof} > 0.2$
- 28750 events in 19.7 fb⁻¹ 2012 data (94% $t\bar{t}$, 44% correct perm.)

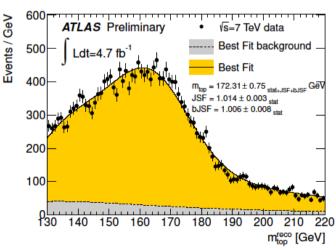




Top mass fitting techniques [example from ATLAS, CONF-2013-046]

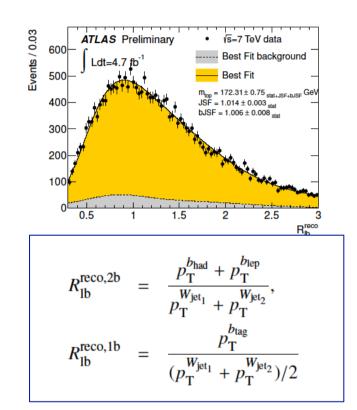
- Invariant mass distributions are distorted by
 - phase space constraints
 - detector resolution
 - wrong particle assignments to jets
 - backgrounds, pileup
 - selection cuts
- Need a MC simulation, tuned to data, to construct templates or probability densities
 - important: at this stage the top mass definition in MC is not too relevant.

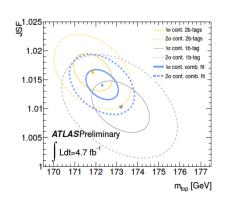


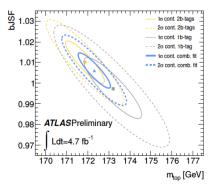


Top mass fitting techniques and JES

- The Jet Energy Scale is the most important source of experimental uncertainties, the W mass constraint is a powerful tool for light quark JES
- Can also find a variable sensitive to b-jet JES and constraint it in situ [ATLAS, CONF-2013-046] in this case b-tagging is used not only for jet classification, but also for JES determination
- Otherwise the simulation is used for b-jet JES, the impact of modeling assumption depends on the jet reconstruction technique



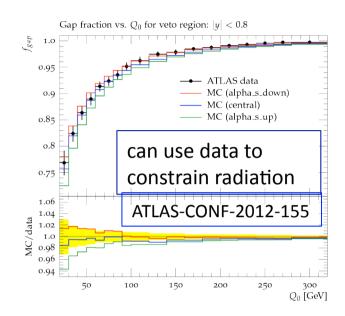


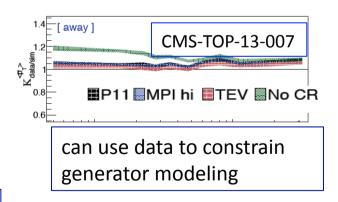


Main sources of systematic uncertainties

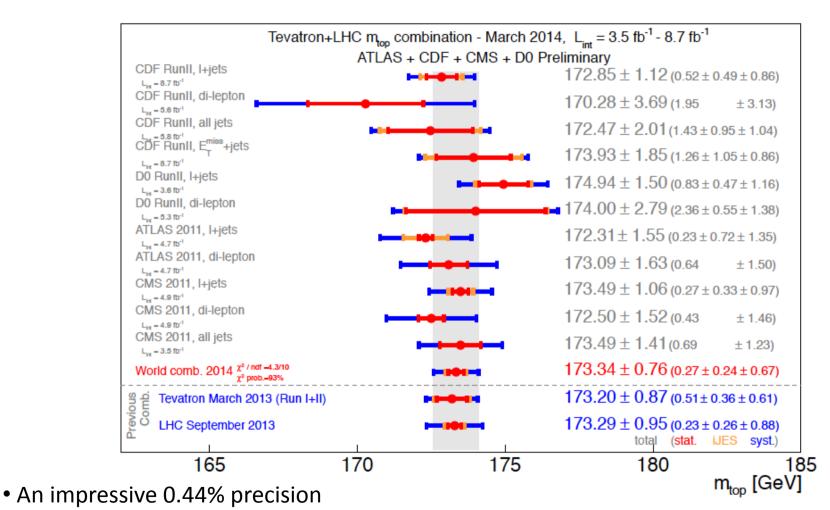
- Jet Energy Scale (depends on technique and jet reco, in situ statistical not included)
 - light jets, detector response [0.2-0.7 GeV]
 - b jets [0.1-0.6 GeV]
- Modeling of gluon radiation [0.3 0.45 GeV]
- Modeling of underlying event [0.1 0.2 GeV]
- Modeling of Colour Reconnection [0.2 0.5 GeV]
- Proton PDF [0.1 0.2 GeV]
- Hadronization, b-fragmentation (included also in JES)
 [0.3 -0.6 GeV]
- b-tagging [0.1 0.8 GeV]
- pileup modeling (included also in JES) (0.1-0.3 GeV)

[The numbers are ranges for illustration only, more details in specific analysis and LHC combination notes]





World combination of m_{top}

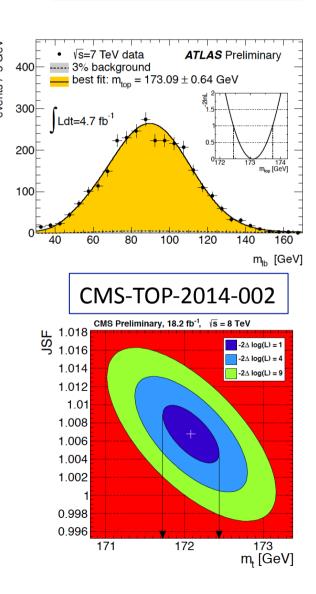


- Some of the most precise measurements non included yet, e.g.
 - D0 full statistics, matrix element method, arXiv:1405.1756, m₊=174.98±0.76
 - CMS l+jets at 8 TeV, L=19.6 fb⁻¹ CMS-TOP-2014-001, m_t =172.04±0.77

A note on the other channels at LHC

- The dilepton and all-hadronic decay channels provide and important cross check, given the difference in colour structure of the final state (next slide).
- The dilepton channel is kinematically underconstrained (2 ν 's), but with low background
- The all-hadronic channel can profit of an accurate in-situ fit of the JES, already providing a result factor 2 better than Tevatron

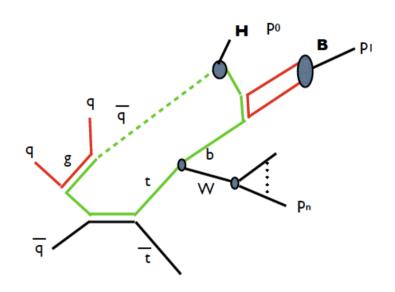
ATLAS-CONF-2013-077



About measuring the top mass from its decay products

- Top is a coloured fermion, it decays before hadronizing, but the b quark from its decay must hadronize
 - there is no way to assign final state
 particles only to the original top, the
 concept is ill-defined as it is the use of a
 pole mass for a coloured particle
 - the effect is expected to be of the order of $\Lambda_{\text{QCD}} \approx 0.2$ GeV but the actual impact depends on the experimental method
 - important to test variables sensitive to the final state definition
 - 2. important to measure the mass with alternative techniques

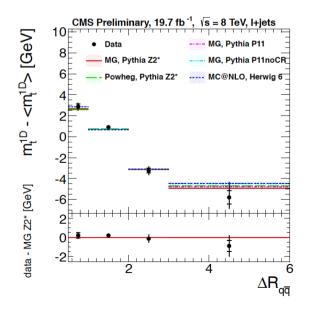
In prospect 1 and 2 will take advantage of the large LHC statistics

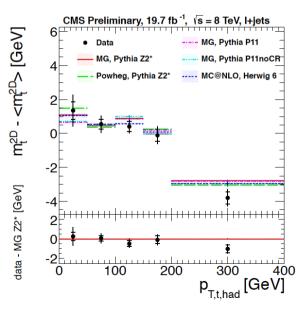


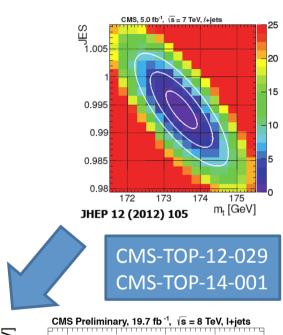
plot courtesy of Michelangelo Mangano

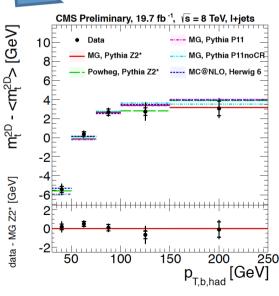
Dependence of Top Mass observable on event kinematics

- test variables sensitive to the final state definition
 - kinematic dependence on final state properly modeled by MC? → 12 kinematic variables checked, related to Color Reconnection, ISR/FRS, b-jet kinematics
 - Good data/MC agreement rules out dramatic effects → need to pursue the study with Run 2 high statistics !!









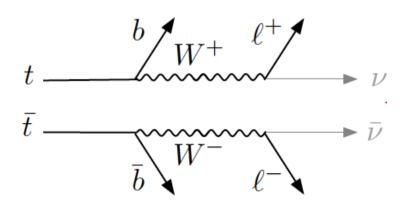
Methods for top mass measurement (2)

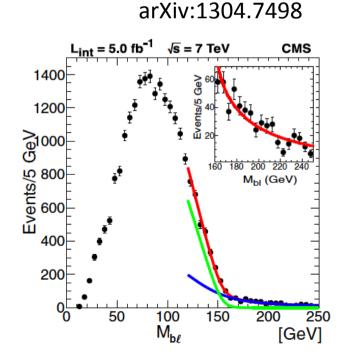
- Given the potential bias in measuring the top mass from its decay products, important to explore alternative techniques, e.g.
 - Measure the decay length (the boost) of B hadrons produced in top decays, the boost is related to the original top mass
 - Select specific channels, for example top with W \rightarrow I ν and B \rightarrow J/ ψ +X decays and measure the three-lepton invariant mass
 - Measure the endpoint of the lepton spectrum or other quantities in top decays
 - Measure the mass from single top events (great potential!)
- <u>Alternative methods have typically larger statistical</u> uncertainties, however at LHC we have large ttbar samples
 - Systematic uncertainties can be controlled with data, again large samples help.
- Another alternative: move away from properties of the decay products
 - extract the top mass from the top cross section

TOP mass from alternative techniques

- Example of a technique already yielding interesting precision:
 Endpoint method
- The shape of the signal can be computed analytically, background data-driven
- Use of MC limited to study underlying assumption: independent decay of two tops (color connections and reconnections violate this assumption)

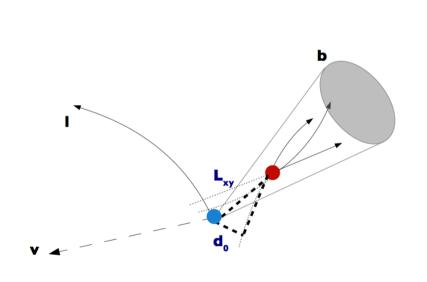
$$M_{\rm t} = 173.9 \pm 0.9 \, ({\rm stat.})^{+1.6}_{-2.0} \, ({\rm syst.}) \, {\rm GeV}$$

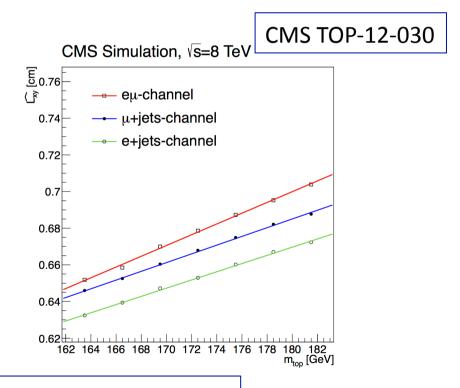




Another example: top mass from the b decay length

 The decay length of b hadrons from top decays is correlated to their boost, i.e. to the top mass



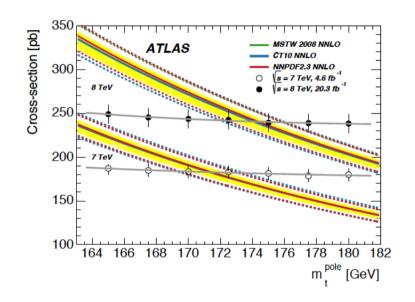


 $m_t = 173.5 \pm 1.5$ stat ± 1.3 syst ± 2.6 p_t(top) GeV

ttbar cross section: mass interpretation

[example from ATLAS, arXiv:1406.5375]

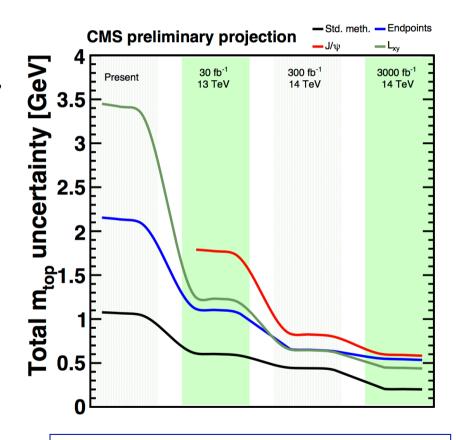
- Measure cross section in the most precise channel: dilepton eµ
- Use b-tagging and double tag method to avoid dependence on btag efficiency
 - interesting by-product: acceptance dependence on m_t is flat because of calibration of the jet acceptance in situ and cancelation with Wt background
- Use recent NNLO calculation of top pair cross section to extract m_t
- The method takes advantage of the excellent luminosity knowledge at LHC (~2%), which is also the long-term experimental limitation, together with the knowledge of the LHC beam energy



$$m_t = 172.9_{-2.6}^{+2.5} \text{ GeV}$$

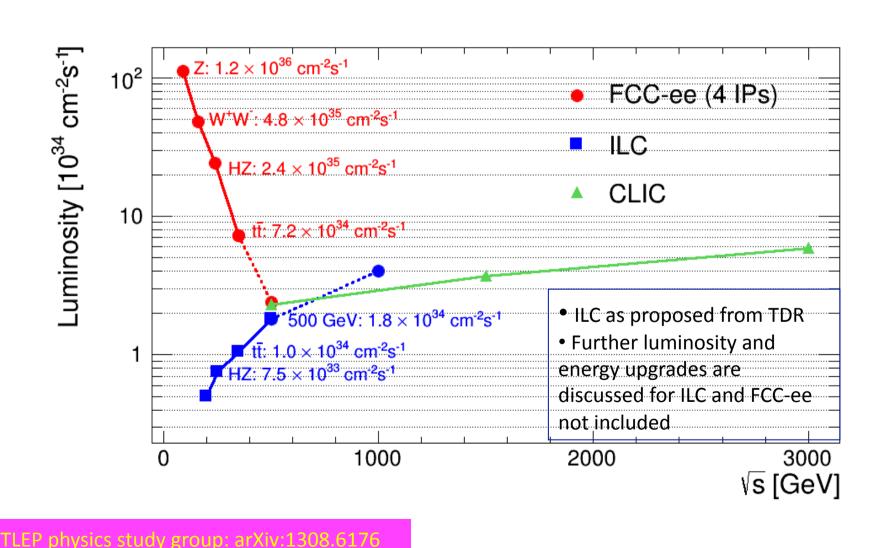
Prospects for top mass at the LHC

- There is potential to improve standard methods, taking advantage of the high statistics for, e.g., in-situ JES calibration, constraining models from differential studies, etc.
- There is even greater potential for alternative methods, most of the current systematic uncertainties can be reduced with higher statistics, e.g. top pt modeling, in-situ JES again
- Improvements on the cross section method are linked to improvements in the luminosity and beam energy uncertainties at LHC
- A optimistic view (maybe realistic give past experience at colliders!) of the evolution in precision is given in the picture



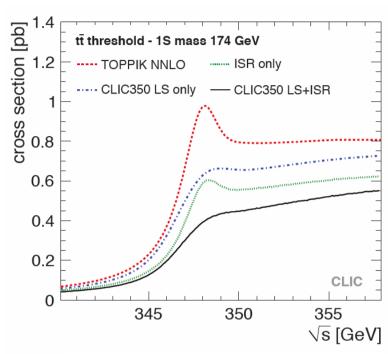
From CMS PAS FTR-13-017, prepared for the "European Strategy for particle physics" discussions

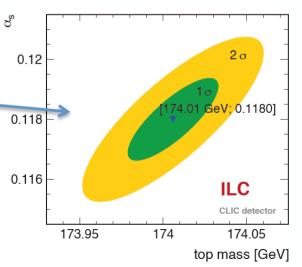
Prospects at future e⁺e⁻ Colliders: Linear and Circular



top mass from threshold scan

- Same conceptual advantage as cross section measurement at LHC, but experimentally dependent only on beam energy spectrum
- Other source of uncertainty is the modification of threshold lineshape related to α_s, top width and top-Higgs Yukawa coupling (with high statistics 2D / 3D fits possible)

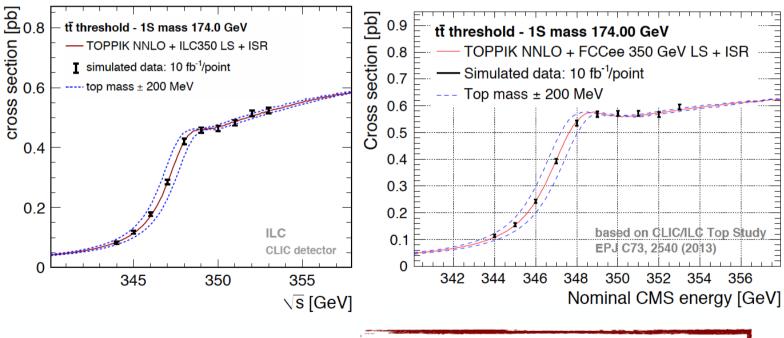




Eur. Phys. J. C73 (2013) 2530

e+e- at 350 GeV: threshold scan

- At ILC statistical uncertainty of 30 MeV with 10 fb-1 scan
- At FCC statistical uncertainty of 10 MeV expected Advantage of a very low level of beamstrahlung at FCC
- Theoretical current uncertainty from higher order QCD contribution ~ 100 MeV
 - Comparing ILC and FCCee assuming identical detector performance



Simulated data points - same integrated luminosity

NB: Assuming unpolarized beams - LC beams can be polarized, increasing cross-sections / reducing backgrounds

From Frank Simon, presented at 7th TLEP-FCC-ee workshop, CERN, June 2014

Conclusions

- The 2014 top mass world average is 173.34±0.76 GeV
- High statistics at LHC offers a unique opportunity of studying the robustness of top mass determination based on top decay products
 - data-based constraints on systematics
 - differential measurements
 - alternative techniques
- We need e+e- colliders to reach precisions in the range 100 MeV → 10 MeV