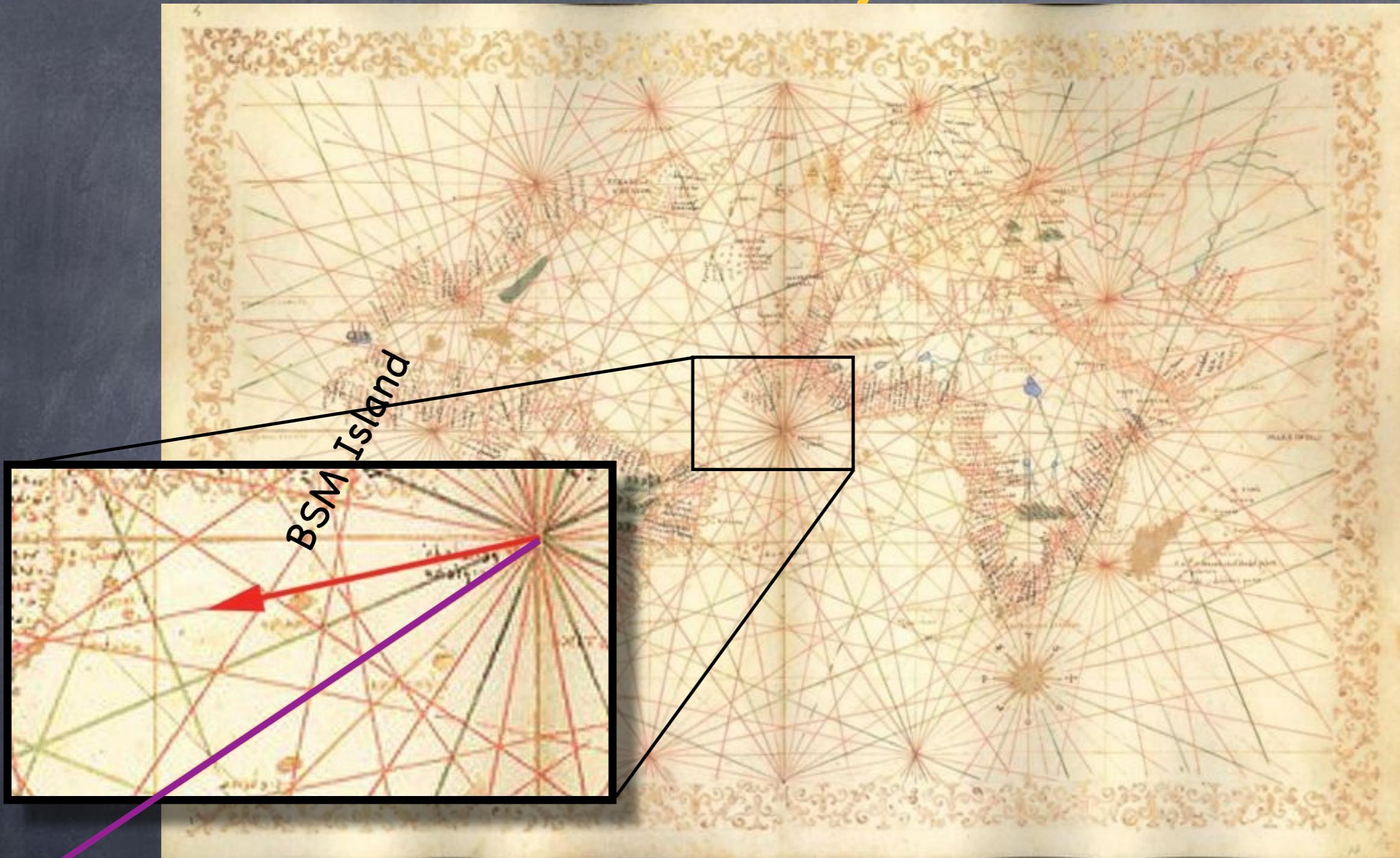


BSM Primary Effects



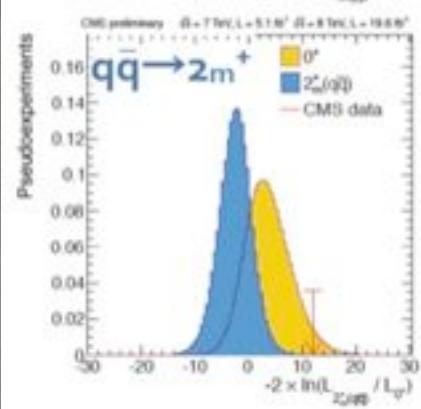
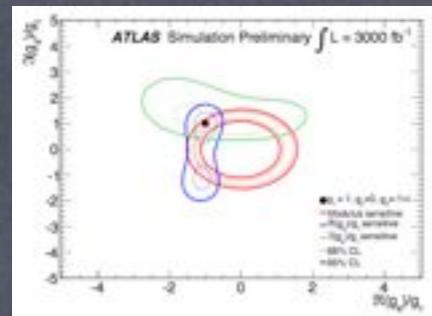
Francesco Riva (EPFL - Lausanne)

In Collaboration with:

Pomarol, Gupta, Liu, Falkowski, Sanz, Masso, Espinosa, Elias-Miro, Biekötter, Knochel, Krä
(1308.2803 ,1308.1879, 1405.0181, 1406.7320, XXX)

Motivation

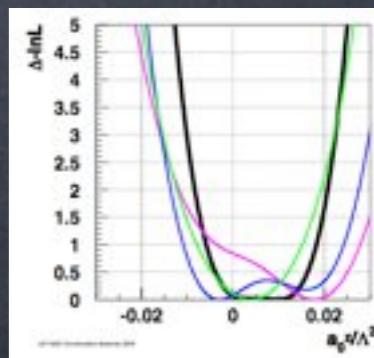
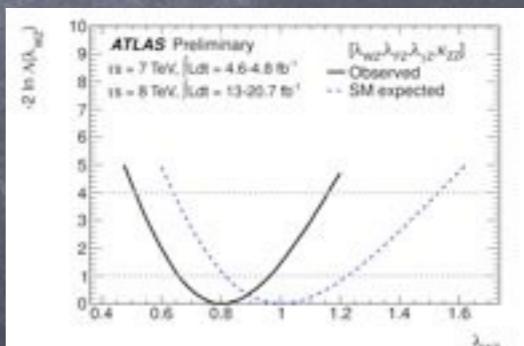
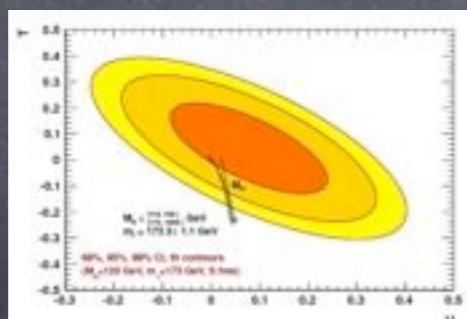
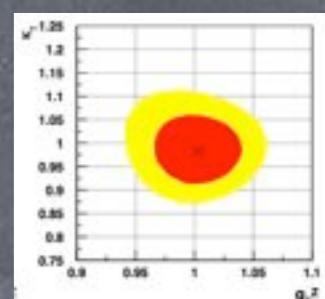
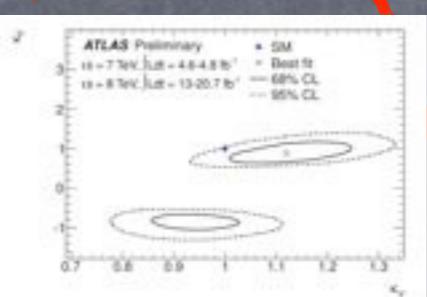
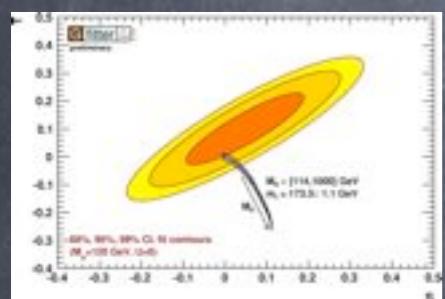
Searches for New Physics



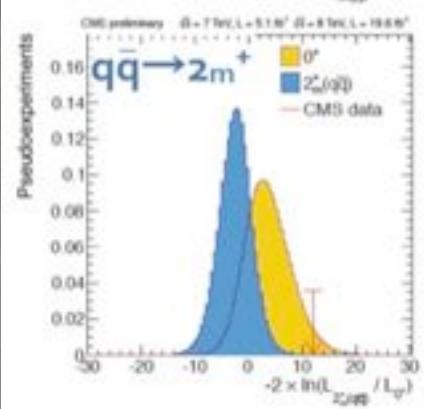
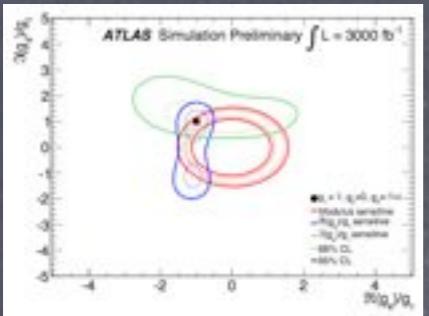
$$\mathcal{L}^{SM}$$

Direct

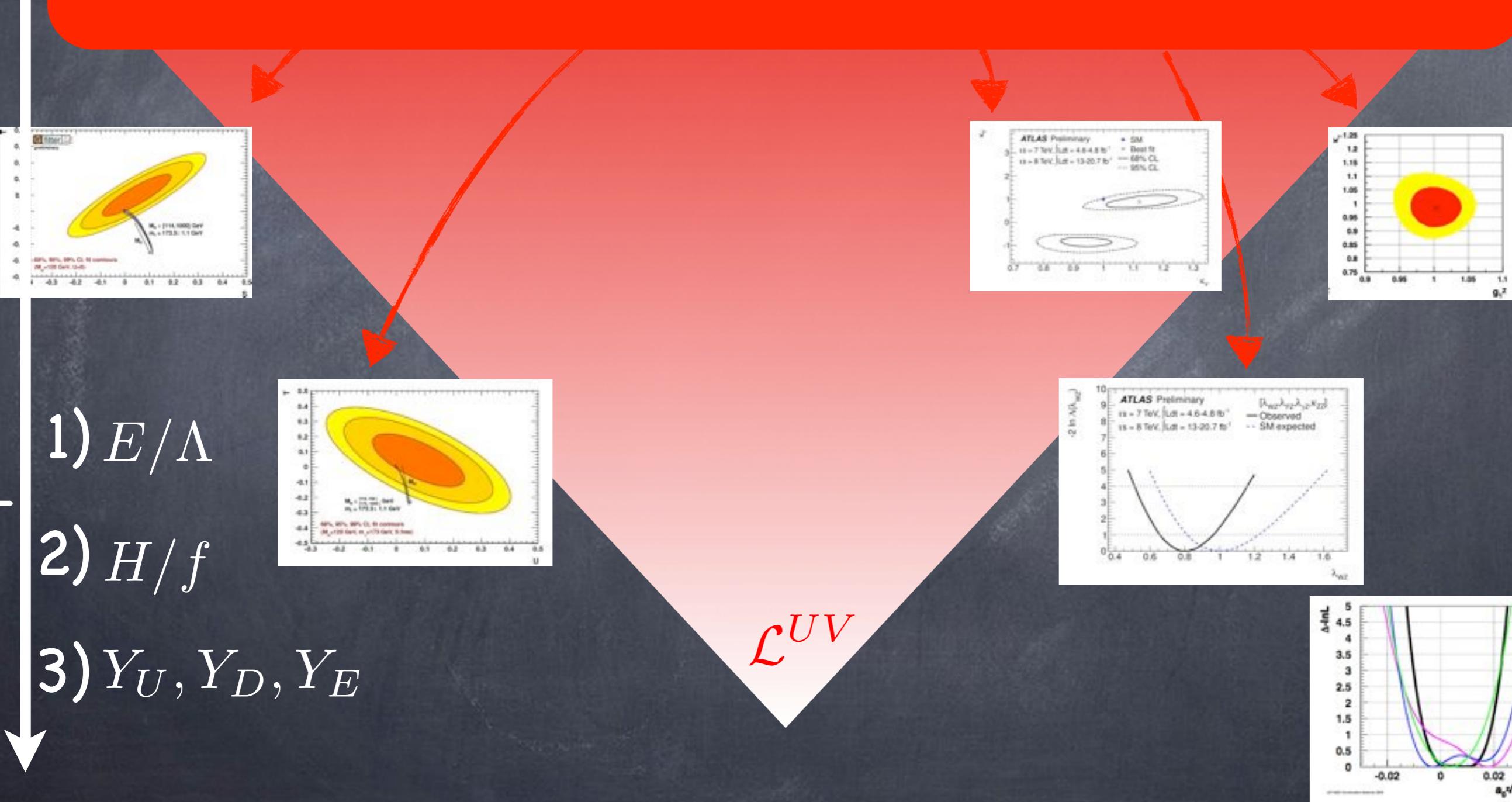
Precision



Motivation



\mathcal{L}^{SM}



Motivation

$$\mathcal{L}^{SM} \equiv$$

1) No direct findings: $M_{new}^i \sim \Lambda \gg gv$

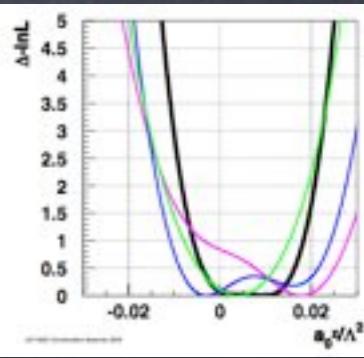
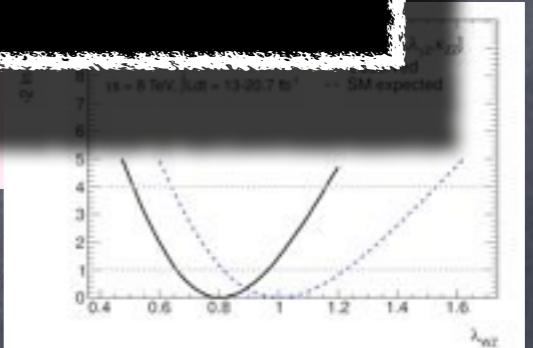
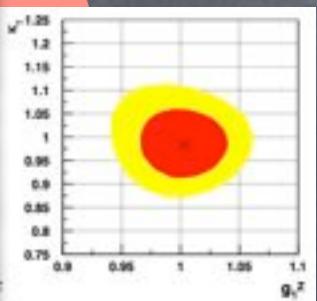
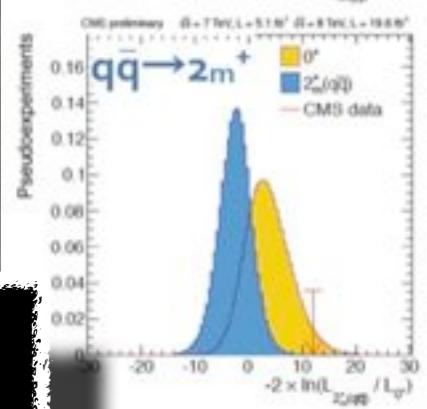
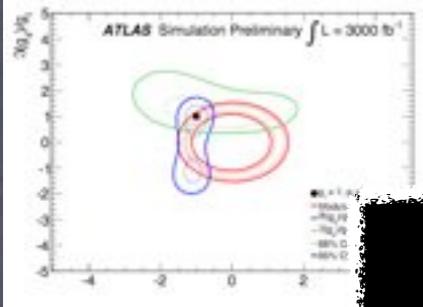
→ Expansion in E/Λ

1) E/Λ

2) H/f

3) Y_U, Y_D, Y_E

$$\mathcal{L}^{UV}$$



Motivation

Expansion

$$1) E/\Lambda$$

$$2) H/f$$

$$3) Y_U, Y_D, Y_E$$

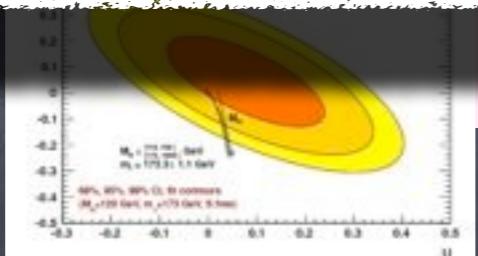
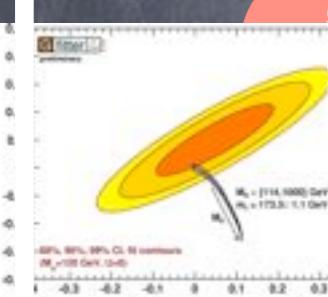
$$\mathcal{L}^{SM} \equiv$$

2) Higgs is excitation around EWSB vacuum

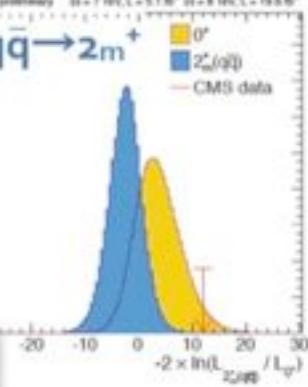
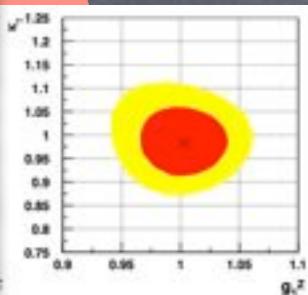
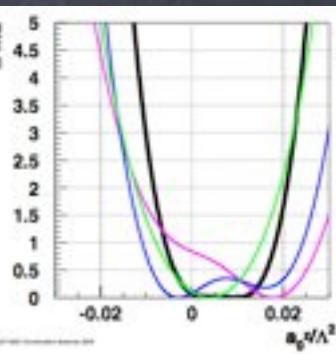
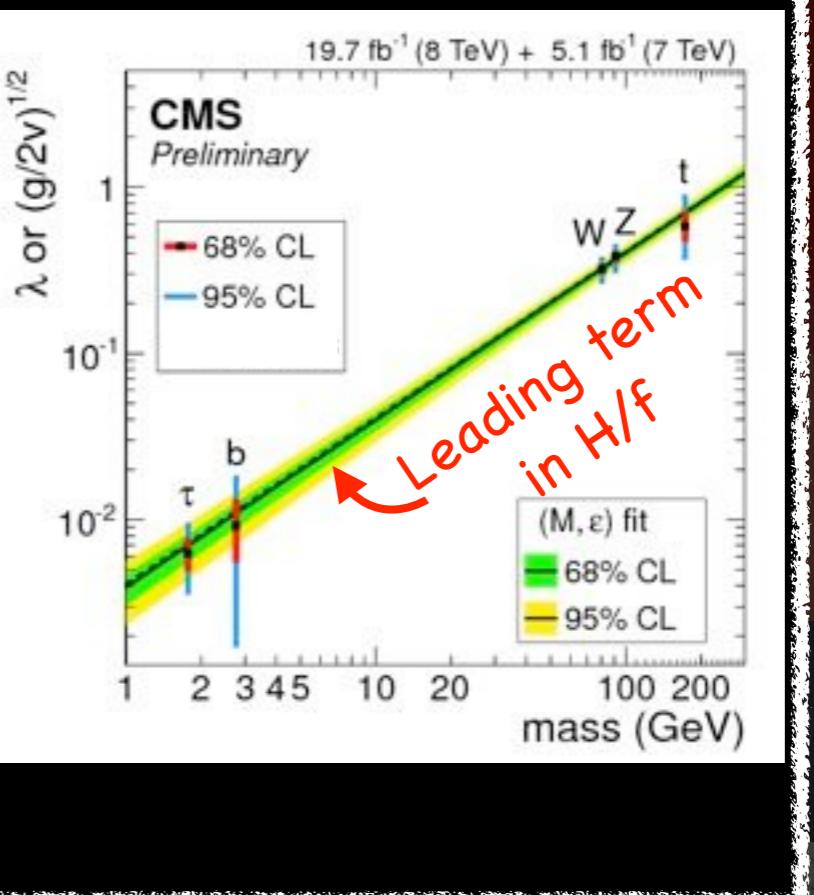
→ Expansion in H/f

$$(f \equiv \Lambda/g_*)$$

$$v + h$$



$$\mathcal{L}^{UV}$$



Motivation

3) Minimal Flavor Violation

→ Expansion in Y_U, Y_D, Y_E

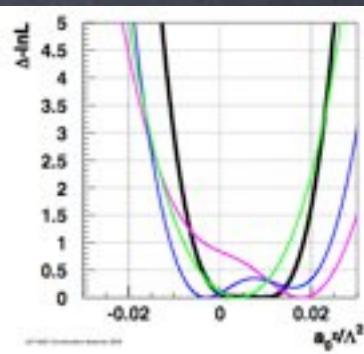
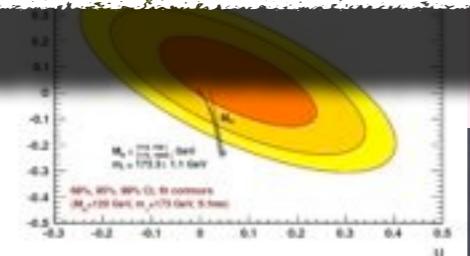
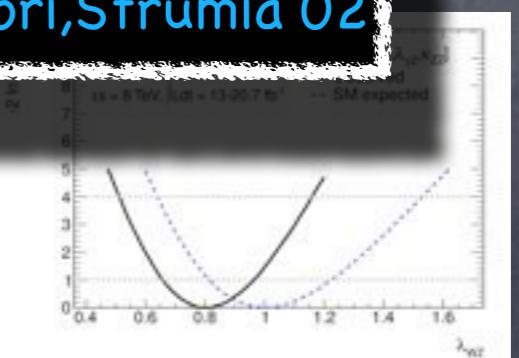
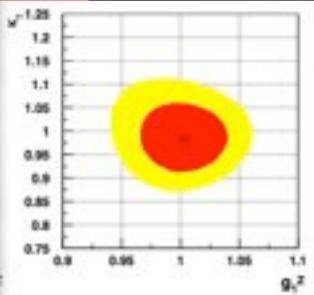
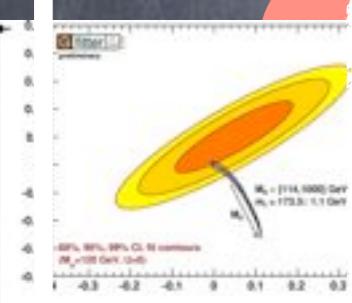
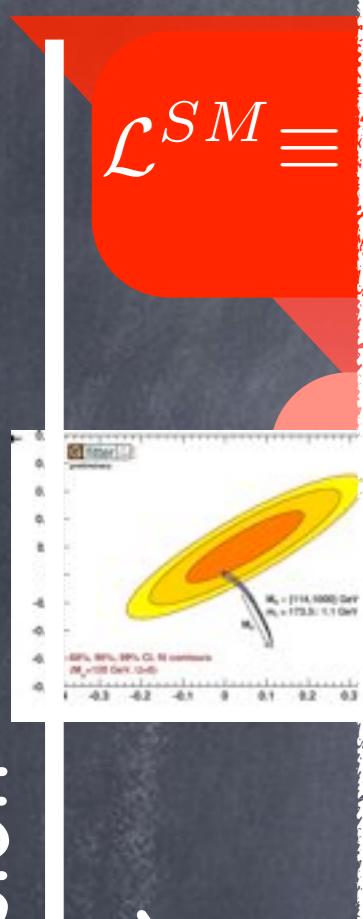
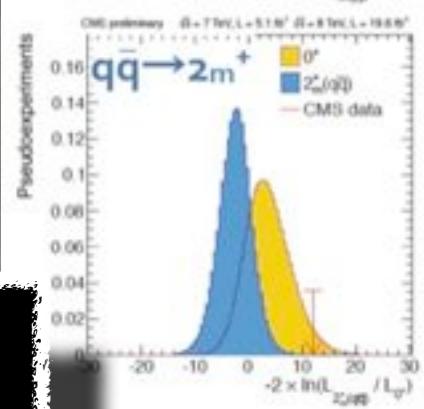
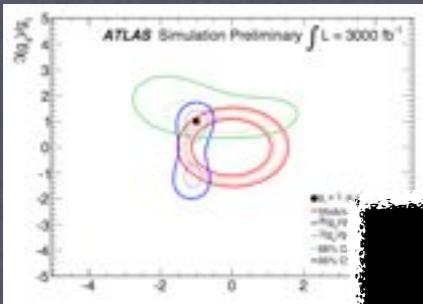
D'ambrogio, Giudice,
Isidori, Strumia'02

1) E/Λ

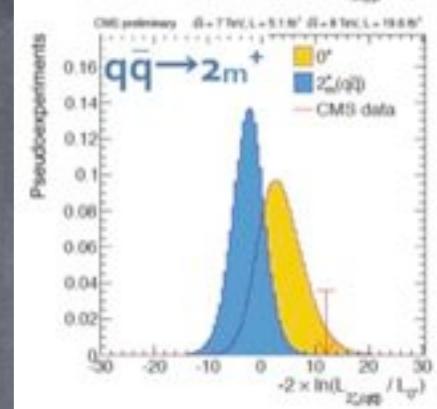
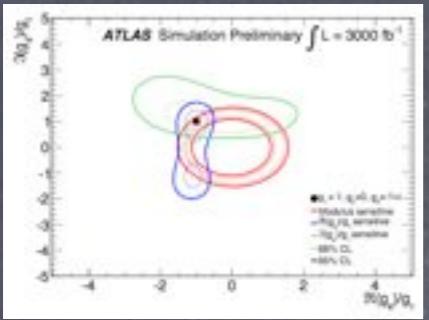
2) H/f

3) Y_U, Y_D, Y_E

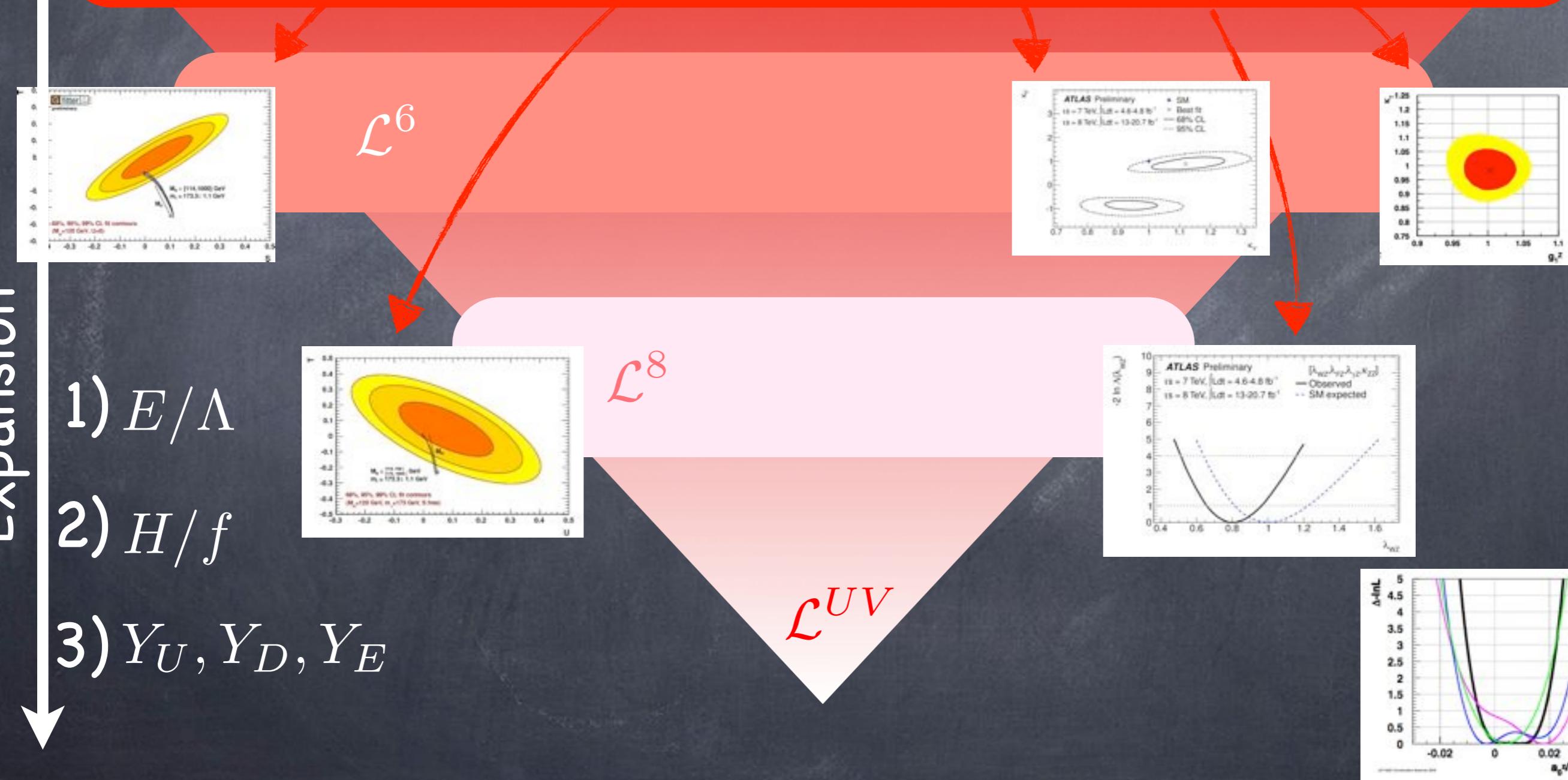
\mathcal{L}^{UV}



Motivation



$$\mathcal{L}^{SM} \equiv \mathcal{L}^4$$



Motivation

$$\mathcal{L}_{\text{eff}} = \frac{\Lambda^4}{g_*^2} \mathcal{L} \left(\frac{D_\mu}{\Lambda}, \frac{g_* H}{\Lambda}, \frac{g_* f_{L,R}}{\Lambda^{3/2}}, \frac{g F_{\mu\nu}}{\Lambda^2} \right) \simeq \mathcal{L}_4 + \mathcal{L}_6 + \dots$$

$$\xrightarrow{\text{SM}} \sum_i \frac{c_i}{\Lambda^2} O_i$$

Buchmuller,Wyler'86;
Giudice,Grojean,Pomarol,Rattazzi'07
Grzadkowski,Iskrzynski,Misiak,Rosiek'10

$$\mathcal{L}^{SM} \equiv \mathcal{L}^4$$

$$\mathcal{L}^{BSM} \simeq \mathcal{L}^6$$

- Parameters: 17
- Accidental relations
- Accidental relations ?

(due to d=4 Lagrangian)

e.g. $m_W = m_Z \cos \theta_W$

$$g_{h\bar{f}f} = m_f/v$$

This Talk: HIGGS PHYSICS
(one family, CP conserving)

Are there relations between different H-observables?

Parameters for BSM: Higgs-only

Higgs Physics Only

v	\leftarrow	$\mathcal{O}_r = H ^2 (D_\mu H)^\dagger (D^\mu H)$
m_d	\leftarrow	$\mathcal{O}_{y_d} = y_d H ^2 \bar{Q}_L H d_R$
m_e	\leftarrow	$\mathcal{O}_{y_e} = y_e H ^2 \bar{L}_L H e_R$
m_u	\leftarrow	$\mathcal{O}_{y_u} = y_u H ^2 \bar{Q}_L \tilde{H} u_R$
$\langle h \rangle$	\leftarrow	$\mathcal{O}_{GG} = \frac{g_s^2}{4} H ^2 G_{\mu\nu}^A G^{A\mu\nu}$
g'	\leftarrow	$\mathcal{O}_{BB} = \frac{g'^2}{4} H ^2 B_{\mu\nu} B^{\mu\nu}$
g	\leftarrow	$\mathcal{O}_{WW} = \frac{g^2}{4} H ^2 W_{\mu\nu}^a W^{a\mu\nu}$
m_h	\leftarrow	$\mathcal{O}_6 = \lambda H ^6$

In the vacuum $\langle h \rangle = v$, operators $|H|^2 \times \mathcal{L}_{SM}$ only redefine SM parameters! \rightarrow Observable only in Higgs physics!

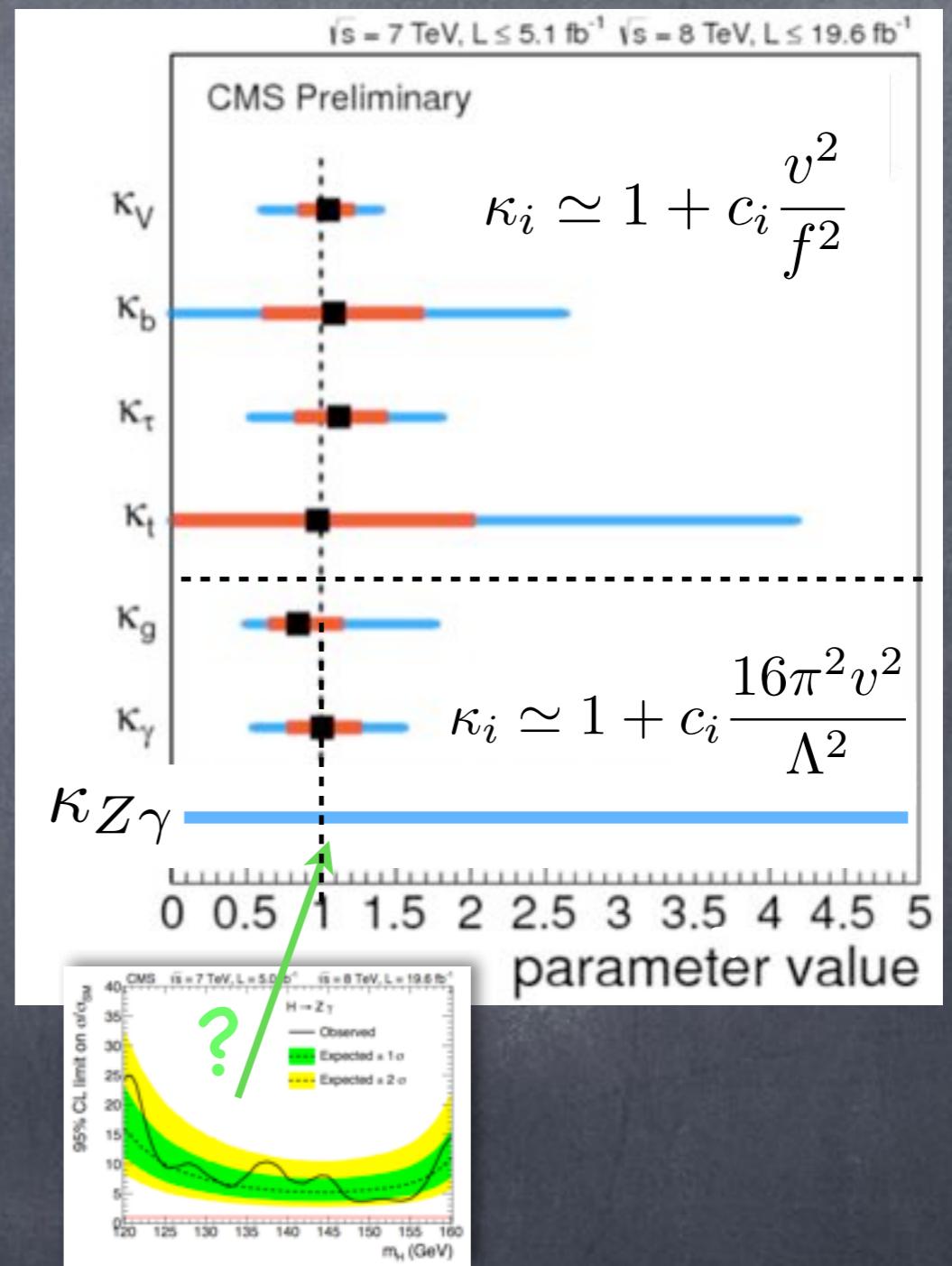
$$\frac{1}{g_s^2} G_{\mu\nu} G^{\mu\nu} + \frac{|H|^2}{\Lambda^2} G_{\mu\nu} G^{\mu\nu} = \left(\frac{1}{g_s^2} + \frac{v^2}{\Lambda^2} \right) G_{\mu\nu} G^{\mu\nu} + h \frac{2v}{\Lambda^2} G_{\mu\nu} G^{\mu\nu} + \dots$$

Parameters for BSM: Higgs-only

Higgs Physics Only

v	\leftarrow	$\mathcal{O}_r = H ^2 (D_\mu H)^\dagger (D^\mu H)$	\rightarrow
m_d	\leftarrow	$\mathcal{O}_{y_d} = y_d H ^2 \bar{Q}_L H d_R$	\rightarrow
m_e	\leftarrow	$\mathcal{O}_{y_e} = y_e H ^2 \bar{L}_L H e_R$	\rightarrow
$ m_u $	\leftarrow	$\mathcal{O}_{y_u} = y_u H ^2 \bar{Q}_L \tilde{H} u_R$	\rightarrow
$\langle h \rangle$	\leftarrow	$\mathcal{O}_{GG} = \frac{g_s^2}{4} H ^2 G_{\mu\nu}^A G^{A\mu\nu}$	\rightarrow
g'	\leftarrow	$\mathcal{O}_{BB} = \frac{g'^2}{4} H ^2 B_{\mu\nu} B^{\mu\nu}$	\rightarrow
g	\leftarrow	$\mathcal{O}_{WW} = \frac{g^2}{4} H ^2 W_{\mu\nu}^a W^{a\mu\nu}$	\times
m_h	\leftarrow	$\mathcal{O}_6 = \lambda H ^6$	\times

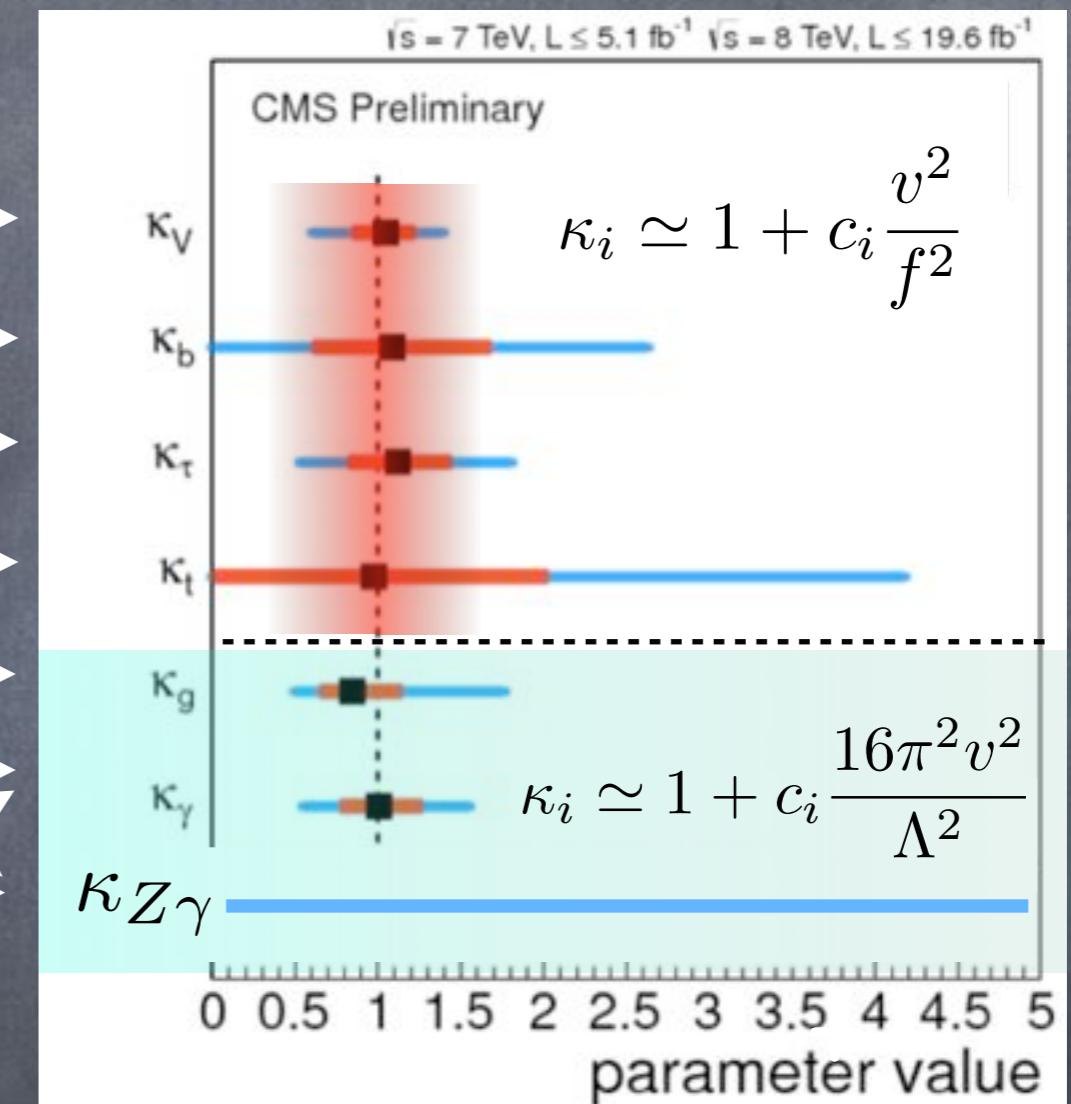
$h^3?$



Parameters for BSM: Higgs-only

Higgs Physics Only

v	\leftarrow	$\mathcal{O}_r = H ^2 (D_\mu H)^\dagger (D^\mu H)$	\rightarrow
m_d	\leftarrow	$\mathcal{O}_{y_d} = y_d H ^2 \bar{Q}_L H d_R$	\rightarrow
m_e	\leftarrow	$\mathcal{O}_{y_e} = y_e H ^2 \bar{L}_L H e_R$	\rightarrow
$ m_u $	\leftarrow	$\mathcal{O}_{y_u} = y_u H ^2 \bar{Q}_L \tilde{H} u_R$	\rightarrow
$\langle h \rangle$	\leftarrow	$\mathcal{O}_{GG} = \frac{g_s^2}{4} H ^2 G_{\mu\nu}^A G^{A\mu\nu}$	\rightarrow
g'	\leftarrow	$\mathcal{O}_{BB} = \frac{g'^2}{4} H ^2 B_{\mu\nu} B^{\mu\nu}$	\rightarrow
g	\leftarrow	$\mathcal{O}_{WW} = \frac{g^2}{4} H ^2 W_{\mu\nu}^a W^{a\mu\nu}$	\times
m_h	\leftarrow	$\mathcal{O}_6 = \lambda H ^6$	\times



Is the EFT expansion justified by these constraints?

$$c_y \frac{v^2}{f^2} \ll 1$$

$$c_{GG} \frac{m_h^2}{\Lambda^2} \ll 1$$

Parameters for BSM: Higgs+EW

Higgs Physics Only

$$\begin{aligned}\mathcal{O}_r &= |H|^2 (D_\mu H)^\dagger (D^\mu H) \\ \mathcal{O}_{y_d} &= y_d |H|^2 \bar{Q}_L H d_R \\ \mathcal{O}_{y_e} &= y_e |H|^2 \bar{L}_L H e_R \\ \mathcal{O}_{y_u} &= y_u |H|^2 \bar{Q}_L \tilde{H} u_R \\ \mathcal{O}_{GG} &= \frac{g_s^2}{4} |H|^2 G_{\mu\nu}^A G^{A\mu\nu} \\ \mathcal{O}_{BB} &= \frac{g'^2}{4} |H|^2 B_{\mu\nu} B^{\mu\nu} \\ \mathcal{O}_{WW} &= \frac{g^2}{4} |H|^2 W_{\mu\nu}^a W^{a\mu\nu} \\ \mathcal{O}_6 &= \lambda |H|^6\end{aligned}$$

EW and Higgs physics

$$\begin{aligned}\mathcal{O}_{WB} &= \frac{gg'}{4} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu} \\ \mathcal{O}_T &= \frac{1}{2} \left(H^\dagger \overset{\leftrightarrow}{D}_\mu H \right)^2 \\ \mathcal{O}_R^u &= (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{u}_R \gamma^\mu u_R) \\ \mathcal{O}_R^d &= (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{d}_R \gamma^\mu d_R) \\ \mathcal{O}_R^e &= (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{e}_R \gamma^\mu e_R) \\ \mathcal{O}_L^q &= (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \gamma^\mu Q_L) \\ \mathcal{O}_L^{(3)q} &= (i H^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \sigma^a \gamma^\mu Q_L) \\ \mathcal{O}_L &= (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{L}_L \gamma^\mu L_L) \\ \mathcal{O}_L^{(3)} &= (i H^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H) (\bar{L}_L \sigma^a \gamma^\mu L_L)\end{aligned}$$

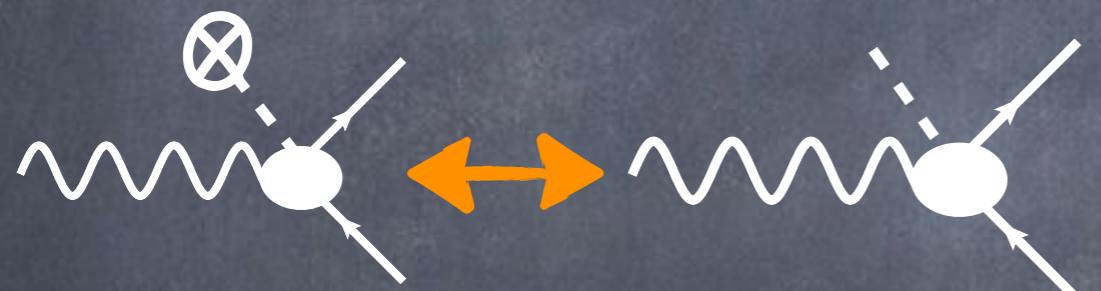
Parameters for BSM: Higgs+EW (see Adam's talk)

In the vacuum $\langle h \rangle = v$, these operators can be measured!

7 of these operators modify:

$$Z\bar{\nu}\nu \ Z\bar{e}_L e_L \ Z\bar{e}_R e_R$$

$$Z\bar{u}_L u_L \ Z\bar{u}_R u_R \ Z\bar{d}_L d_L \ Z\bar{d}_R d_R$$



Constrained by LEP1* $\sim 1/1000$!

Impact of these operators in H-physics is irrelevant

EW and Higgs physics

$$\mathcal{O}_{WB} = \frac{gg'}{4}(H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$$

$$\mathcal{O}_T = \frac{1}{2} \left(H^\dagger \overset{\leftrightarrow}{D}_\mu H \right)^2$$

$$\mathcal{O}_R^u = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{u}_R \gamma^\mu u_R)$$

$$\mathcal{O}_R^d = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{d}_R \gamma^\mu d_R)$$

$$\mathcal{O}_R^e = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{e}_R \gamma^\mu e_R)$$

$$\mathcal{O}_L^q = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{Q}_L \gamma^\mu Q_L)$$

$$\mathcal{O}_L^{(3)q} = (iH^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H)(\bar{Q}_L \sigma^a \gamma^\mu Q_L)$$

$$\mathcal{O}_L^{(3)} = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{L}_L \gamma^\mu L_L)$$

$$\mathcal{O}_L^{(3)} = (iH^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H)(\bar{L}_L \sigma^a \gamma^\mu L_L)$$

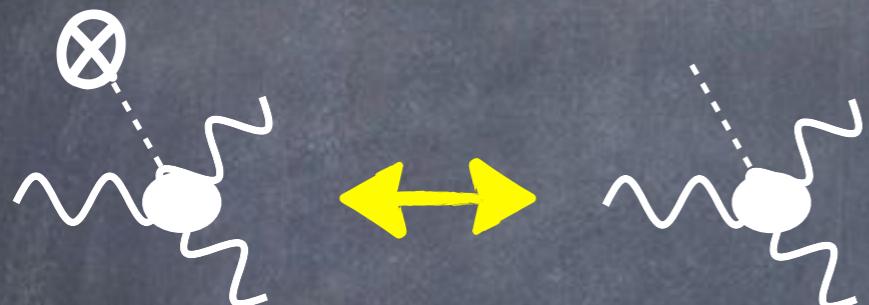
(Gupta),Pomarol,FR'13-14; Falkowski,FR,Sanz'to appear

*= if α, m_Z, m_W are used as input parameters, no other dim-6 operators affect LEP1 measurements!

Parameters for BSM: Higgs+EW (see Adam's talk)

In the vacuum $\langle h \rangle = v$, these operators can be measured!

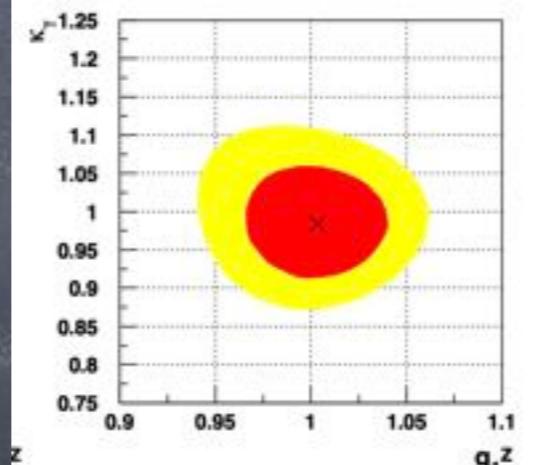
② of these modify TGCs:



$$g_Z^1 \quad \kappa_\gamma$$

Hagiwara,Hikasa,
Peccei,Zeppenfeld'87

LEP2($e e \rightarrow W W$)
constrained* $\sim 1/100$



EW and Higgs physics

$\mathcal{O}_{WB} = \frac{gg'}{4}(H^\dagger \sigma^a H)W_{\mu\nu}^a B^{\mu\nu}$
$\mathcal{O}_T = \frac{1}{2} \left(H^\dagger \overset{\leftrightarrow}{D}_\mu H \right)^2$
$\mathcal{O}_R^u = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{u}_R \gamma^\mu u_R)$
$\mathcal{O}_R^d = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{d}_R \gamma^\mu d_R)$
$\mathcal{O}_R^e = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{e}_R \gamma^\mu e_R)$
$\mathcal{O}_L^q = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{Q}_L \gamma^\mu Q_L)$
$\mathcal{O}_L^{(3)q} = (iH^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H)(\bar{Q}_L \sigma^a \gamma^\mu Q_L)$
$\mathcal{O}_L = (iH^\dagger \overset{\leftrightarrow}{D}_\mu H)(\bar{L}_L \gamma^\mu L_L)$
$\mathcal{O}_L^{(3)} = (iH^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H)(\bar{L}_L \sigma^a \gamma^\mu L_L)$

→ We can include these 2 combinations in H-physics studies
(but recall connection with TGC!)

*= Non-Higgs operator $g\epsilon_{abc}W_\mu^{a\nu}W_{\nu\rho}^bW^{c\rho\mu}$ can interfere with extraction of bounds
(see backup slides)

Small Summary: Parameters

$\mathcal{O}_r = H ^2 (D_\mu H)^\dagger (D^\mu H)$
$\mathcal{O}_{y_d} = y_d H ^2 \bar{Q}_L H d_R$
$\mathcal{O}_{y_e} = y_e H ^2 \bar{L}_L H e_R$
$\mathcal{O}_{y_u} = y_u H ^2 \bar{Q}_L \tilde{H} u_R$
$\mathcal{O}_{GG} = \frac{g_s^2}{4} H ^2 G_{\mu\nu}^A G^{A\mu\nu}$
$\mathcal{O}_{BB} = \frac{g'^2}{4} H ^2 B_{\mu\nu} B^{\mu\nu}$
$\mathcal{O}_{WW} = \frac{g^2}{4} H ^2 W_{\mu\nu}^a W^{a\mu\nu}$
$\mathcal{O}_6 = \lambda H ^6$



$\kappa_V, \kappa_b, \kappa_\tau, \kappa_t, \kappa_G, \kappa_{\gamma\gamma}, \kappa_{Z\gamma}, \kappa_{h^3}$

$\mathcal{O}_{WB} = \frac{gg'}{4} (H^\dagger \sigma^a H) W_{\mu\nu}^a B^{\mu\nu}$
$\mathcal{O}_T = \frac{1}{2} \left(H^\dagger \overset{\leftrightarrow}{D}_\mu H \right)^2$
$\mathcal{O}_R^u = (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{u}_R \gamma^\mu u_R)$
$\mathcal{O}_R^d = (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{d}_R \gamma^\mu d_R)$
$\mathcal{O}_R^e = (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{e}_R \gamma^\mu e_R)$
$\mathcal{O}_L^q = (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \gamma^\mu Q_L)$
$\mathcal{O}_L^{(3)q} = (i H^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H) (\bar{Q}_L \sigma^a \gamma^\mu Q_L)$
$\mathcal{O}_L^L = (i H^\dagger \overset{\leftrightarrow}{D}_\mu H) (\bar{L}_L \gamma^\mu L_L)$
$\mathcal{O}_L^{(3)L} = (i H^\dagger \sigma^a \overset{\leftrightarrow}{D}_\mu H) (\bar{L}_L \sigma^a \gamma^\mu L_L)$



g_Z^1, κ_γ

$\delta g_{ZeL}, \delta g_{ZeR}, \delta g_{Z\nu}, \delta g_{ZuL}, \delta g_{ZdL}, \delta g_{ZuR}, \delta g_{ZdR}$

Might as well use these as parameters, to keep relations between observables manifest!

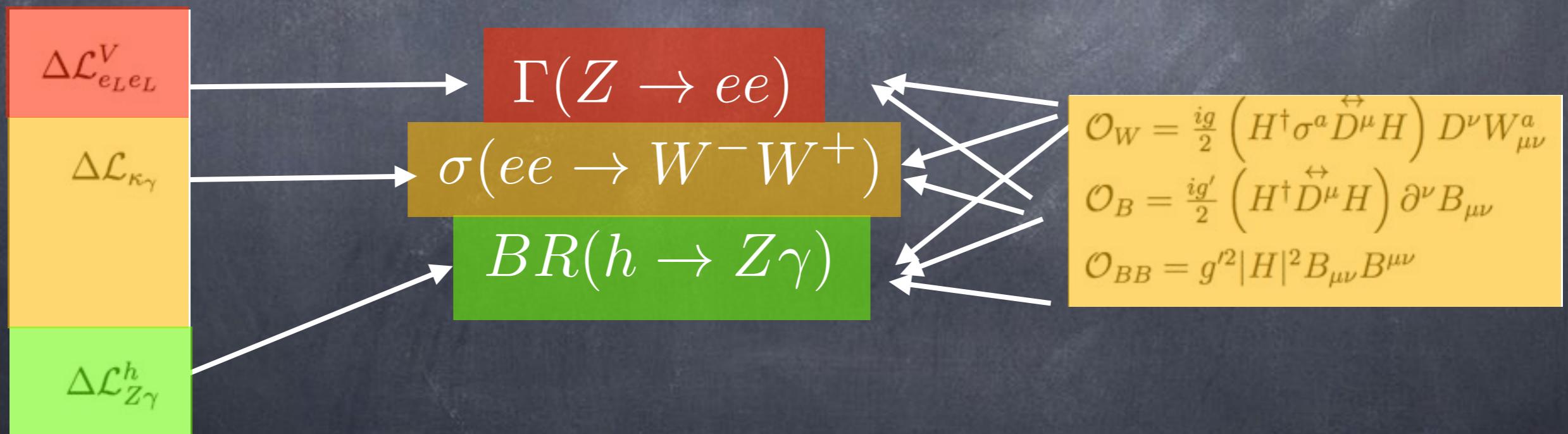
Parameters \rightarrow Relations

“BSM Primaries”
Parametrization

- Mass eigenstate basis
- 1 to 1 with best experiments
- No theoretical correlation
(orthogonal to other experiments)

Usual Operator
Parametrization

- Gauge invariance manifest
- Physics unclear
- Large theo. correlations



Parameters -> Relations

“BSM Primaries”
Parametrization

- Mass eigenstate basis
- 1 to 1 with best experiments
- No theoretical correlation

(orthogonal to other experiments)

Usual Operator
Parametrization

- Gauge invariance manifest
- Physics unclear
- Large theo. correlations

Z and W couplings
related also at dim-6
(and related to $hVff$ from $h=\hat{h}+v$)

$$\Delta \mathcal{L}_{e_L e_L}^Z = \delta g_{eL}^Z \frac{h^2}{v^2} \left[Z^\mu \bar{e}_L \gamma_\mu e_L - \frac{c_{\theta_W}}{\sqrt{2}} (W^{+\mu} \bar{\nu}_L \gamma_\mu e_L + \text{h.c.}) \right]$$

$$\begin{aligned} \Delta \mathcal{L}_{\kappa_\gamma} = & \frac{\delta \kappa_\gamma}{v^2} \left[ieh^2 (A_{\mu\nu} - t_{\theta_W} Z_{\mu\nu}) W^{+\mu} W^{-\nu} \right. \\ & \left. + Z_\nu \partial_\mu h^2 (t_{\theta_W} A^{\mu\nu} - t_{\theta_W}^2 Z^{\mu\nu}) + \frac{(h^2 - v^2)}{2} \left(t_{\theta_W} Z_{\mu\nu} A^{\mu\nu} + \frac{c_{2\theta_W}}{2c_{\theta_W}^2} Z_{\mu\nu} Z^{\mu\nu} + W_{\mu\nu}^+ W^{-\mu\nu} \right) \right] \end{aligned}$$

$$\Delta \mathcal{L}_{Z\gamma}^h = \kappa_{Z\gamma} \left(\frac{\hat{h}}{v} + \frac{\hat{h}^2}{2v^2} \right) \left[t_{\theta_W} A_{\mu\nu} Z^{\mu\nu} + \frac{c_{2\theta_W}}{2c_{\theta_W}^2} Z_{\mu\nu} Z^{\mu\nu} + W_{\mu\nu}^+ W^{-\mu\nu} \right]$$

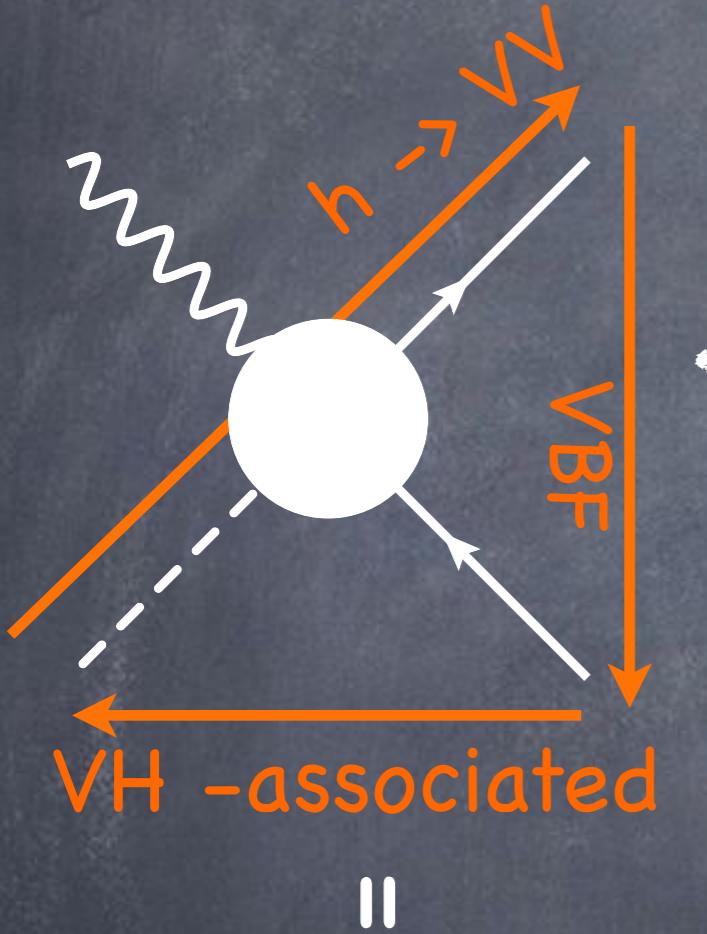
$h \rightarrow Z\gamma$ related to $h \rightarrow WW, ZZ$

Gupta, Pomarol, FR'14

BSM Relations in Higgs physics

Off-shell Higgs physics:

$$\begin{array}{c} \kappa_V, \kappa_b, \kappa_\tau, \kappa_t, \kappa_G, \kappa_{\gamma\gamma}, \kappa_{Z\gamma}, \kappa_{h^3} g_Z^1, \kappa_\gamma \\ \delta g_{ZeL}, \delta g_{ZeR}, \delta g_{Z\nu}, \delta g_{ZuL}, \delta g_{ZdL}, \delta g_{ZuR}, \delta g_{ZdR} \end{array}$$



$$\leftarrow \mathcal{L} = \dots h V^\mu V_\mu + \dots h V^\mu \bar{f} \gamma_\mu f + \dots h V^{\mu\nu} V_{\mu\nu}$$

LEP1

$$a_f^Z = 2\delta g_f^Z - 2\delta g_1^Z (g_f^Z c_{2\theta_W} + e Q s_{2\theta_W}) + 2\delta \kappa_\gamma g' Y \frac{s_{\theta_W}}{c_{\theta_W}^2}$$

TGC

$$\hat{a}_f^Z = 2g_f^Z + \frac{g_f^Z v}{m_Z^2 c_{\theta_W}^2} (\delta g_{VV}^h + \delta g_1^Z e^2 v - \delta \kappa_\gamma g'^2 v) ,$$

$$b_f^Z = 2 \frac{g_f^Z}{c_{\theta_W}^2} (-\delta \kappa_\gamma - \kappa_{Z\gamma} c_{2\theta_W} - 2\kappa_{\gamma\gamma} c_{\theta_W}^2) ,$$

$\hat{b}_f^Z = -2e Q_f t_{\theta_W} \kappa_{Z\gamma}$, Other Higgs Processes

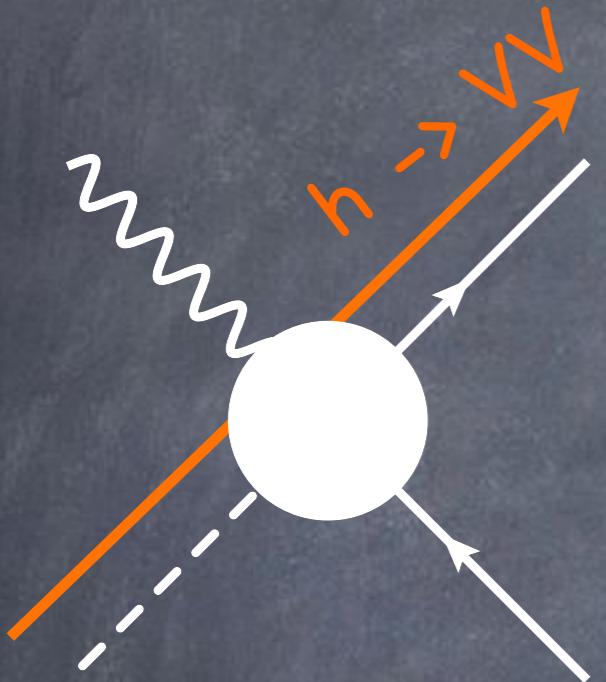
$$\frac{1}{v} \epsilon^{*\mu}(q) J_f^{V\nu}(p) [A_f^V \eta_{\mu\nu} + B_f^V (p \cdot q \eta_{\mu\nu} - p_\mu q_\nu)]$$

$$\begin{aligned} A_f^V &= a_f^V + \hat{a}_f^V \frac{m_V^2}{p^2 - m_V^2} \\ B_f^V &= b_f^V \frac{1}{p^2 - m_V^2} + \hat{b}_f^V \frac{1}{p^2} \end{aligned}$$

BSM Relations for Run 2

Deviations in different. distr. of $h \rightarrow Z\bar{f}f$ or $h \rightarrow W\bar{f}f$

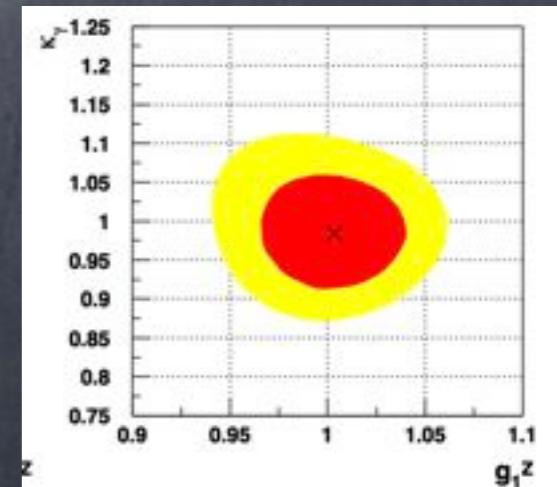
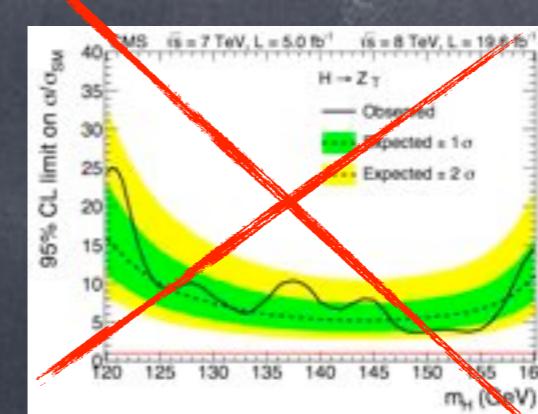
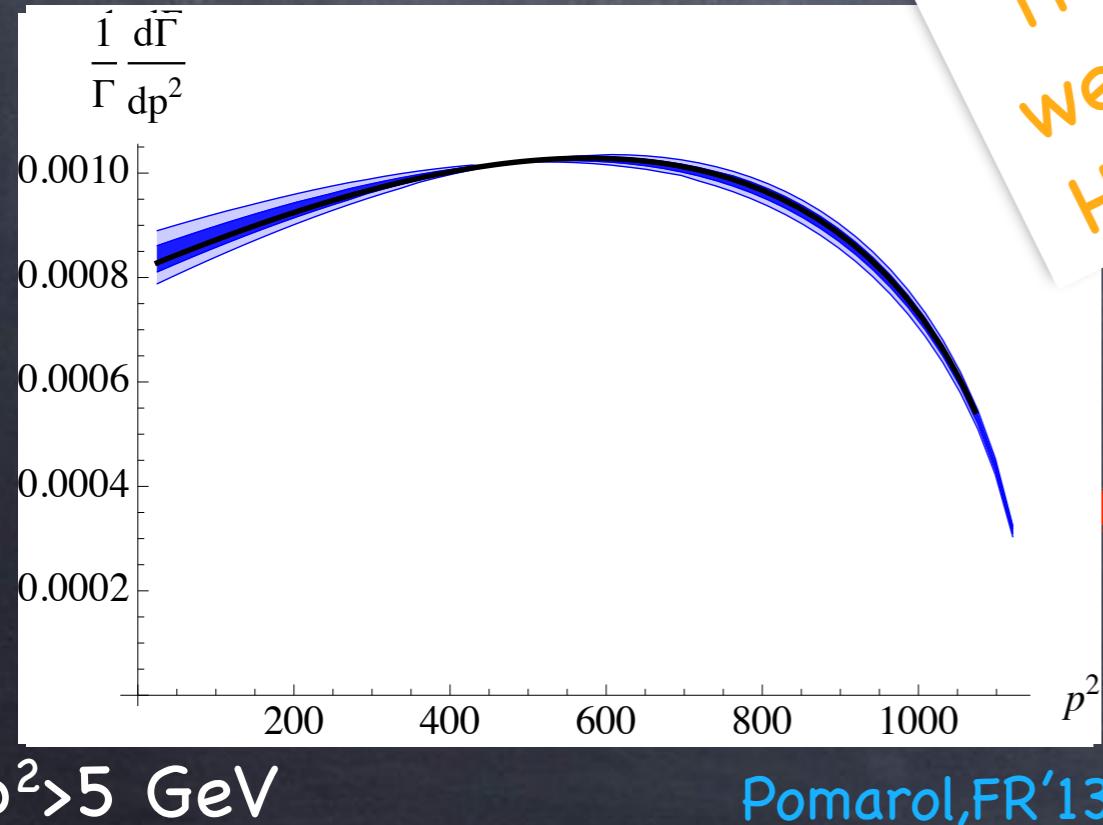
See e.g. Isidori,(Manohar),Trott'13
Falkowski,Vega-Morales'14



LEP 1
~~Related with Zff couplings~~

Related with Triple Gauge Coupling

Related with α/α_{SM}
This is the sensitivity
we are aiming to make
H-physics competitive!

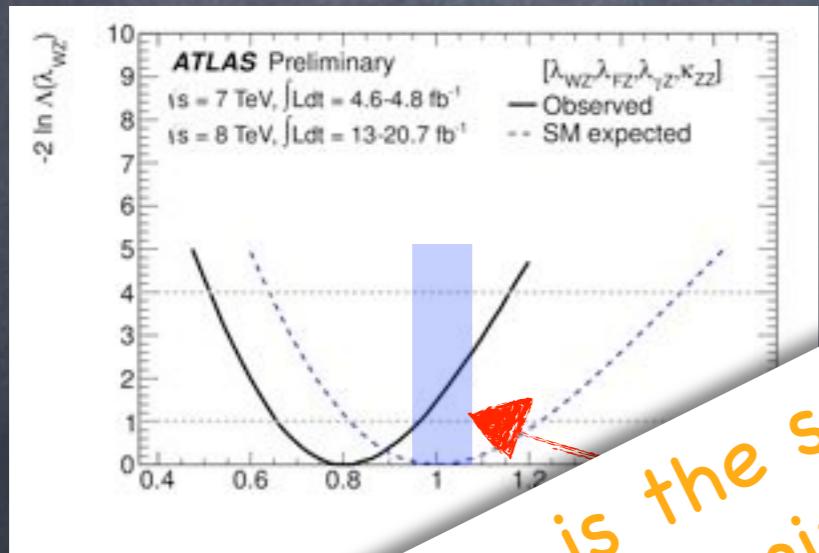


BSM Relations at Run 1

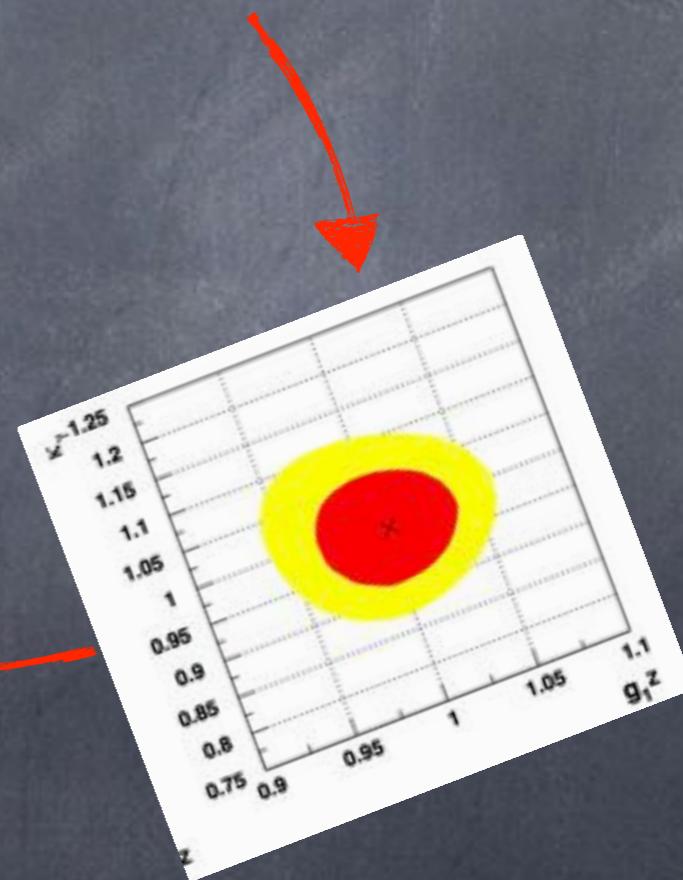
Custodial Symmetry in h decays $h \rightarrow VV^*$ λ_{WZ}

- Off-Shell V → Integrated Decay Width already sensitive
- $m_Z \neq m_W$ → to p-dependence of hVV coupling!

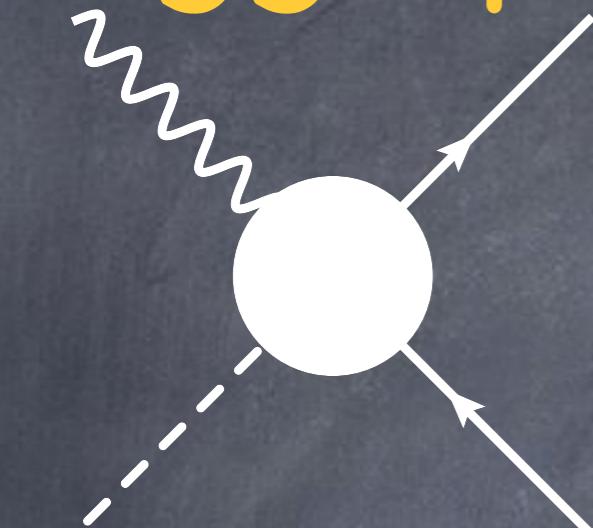
$$\lambda_{WZ}^2 - 1 \simeq 0.6\delta g_1^Z - 0.5\delta\kappa_\gamma - 1.6\kappa_{Z\gamma}$$



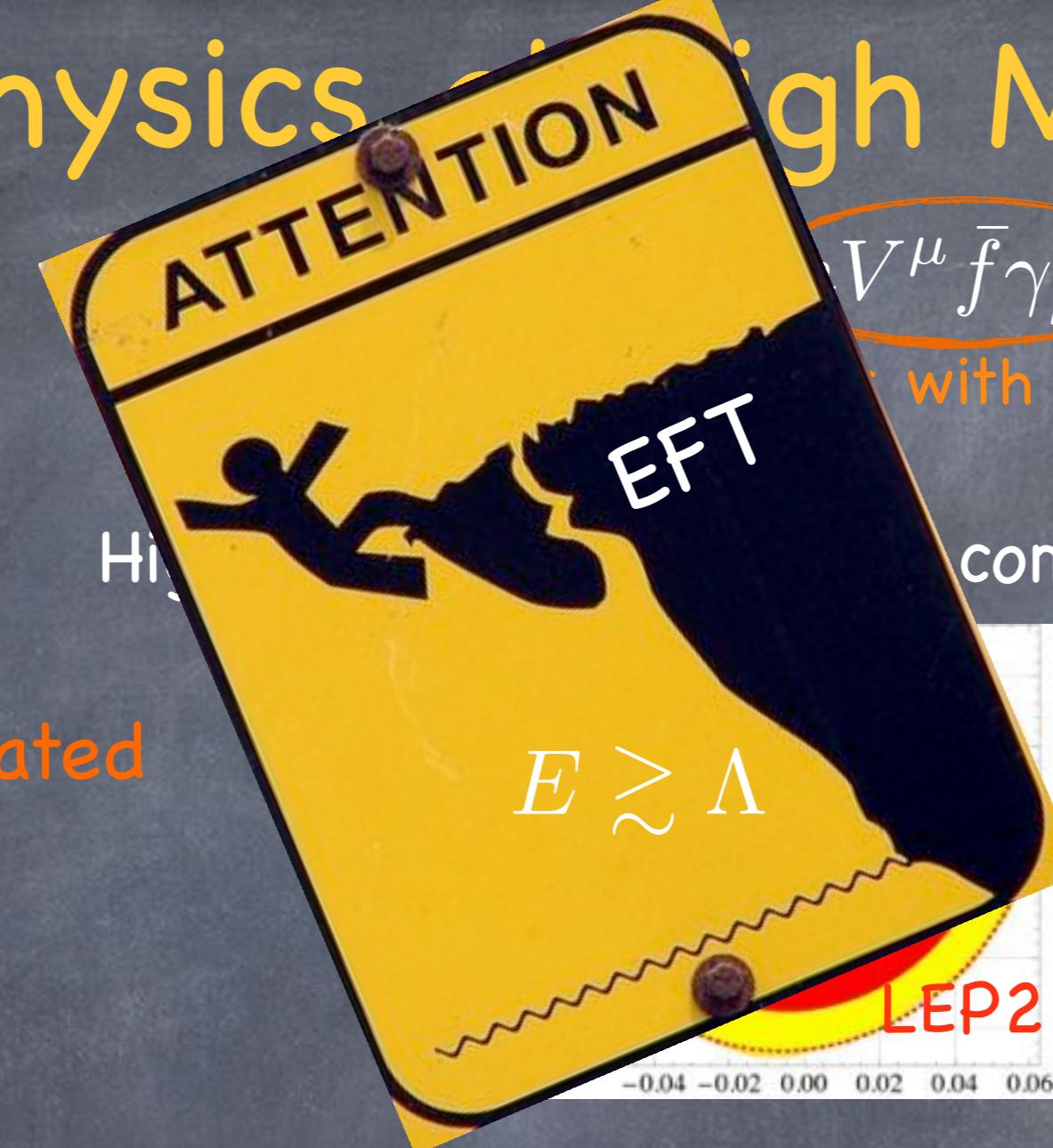
This is the sensitivity
we are aiming to make
H-physics competitive!



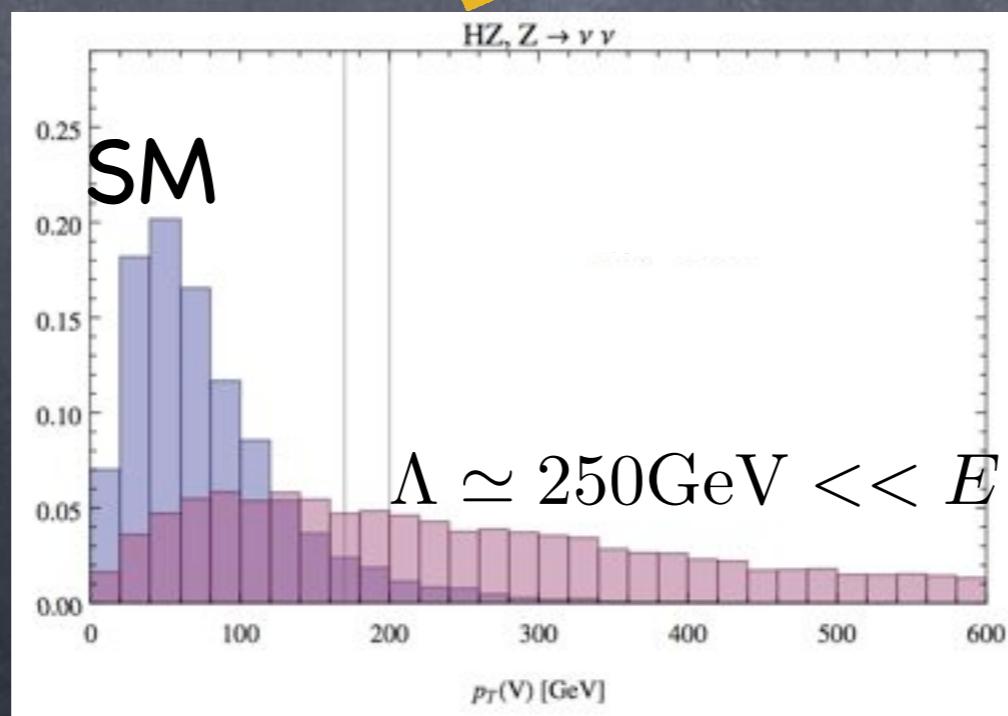
Higgs physics at high Momentum?



$V^* \rightarrow VH$ associated



However:



Only valid for some theories!

Gupta,Pomarol,FR'14;
Biekötter,Knochel,Krämer,Liu,FR '14

$$hV^\mu \bar{f} \gamma_\mu f + \dots h V^{\mu\nu} V_{\mu\nu}$$

with Energy!

compete with TGC at LEP

Isidori,Trott'13

Ellis,Sanz,You'14

Corbett,Eboli,

Gonzalez-Garcia,Fraile'12-13

Biekötter,Knochel,Krämer,Liu,FR '14

Beneke,Boito,Wang'14

Conclusions

- EFT: consistent framework to search for leading BSM effects
 - Results must satisfy $E \ll \Lambda \rightarrow v \ll f$
(not always true, at present)
- Parametrization of BSM for Higgs physics:

7

$$\{\delta g_{ZeL}, \delta g_{ZeR}, \delta g_{Z\nu}, \delta g_{ZuL}, \delta g_{ZdL}, \delta g_{ZuR}, \delta g_{ZdR}\}$$

2

$$\{g_1^Z, \kappa_\gamma\}$$

8

$$\{\kappa_g, \kappa_\gamma, \kappa_V, \kappa_t, \kappa_b, \kappa_\tau, \kappa_{Z\gamma}, \kappa_{h^3}\}$$

→ Basis independent relations between EW and Higgs observables

LEP2

2 Parameter fit

Parameter	68% C.L.	95% C.L.	Correlations
g_1^Z	$1.004^{+0.024}_{-0.025}$	[+0.954, +1.050]	1.00 +0.11
κ_γ	$0.984^{+0.049}_{-0.049}$	[+0.894, +1.084]	+0.11 1.00

Parameter	68% C.L.	95% C.L.	Correlations
Δg_1^Z	$-0.060^{+0.031}_{-0.030}$	[-0.118, +0.002]	1.0 -0.55 -0.41
λ_γ	$0.038^{+0.031}_{-0.032}$	[-0.027, +0.099]	-0.55 1.0 -0.04
$\Delta \kappa_\gamma$	$0.077^{+0.070}_{-0.070}$	[-0.050, +0.218]	-0.41 -0.04 1.0

LEP2 - Combined

Delphi

