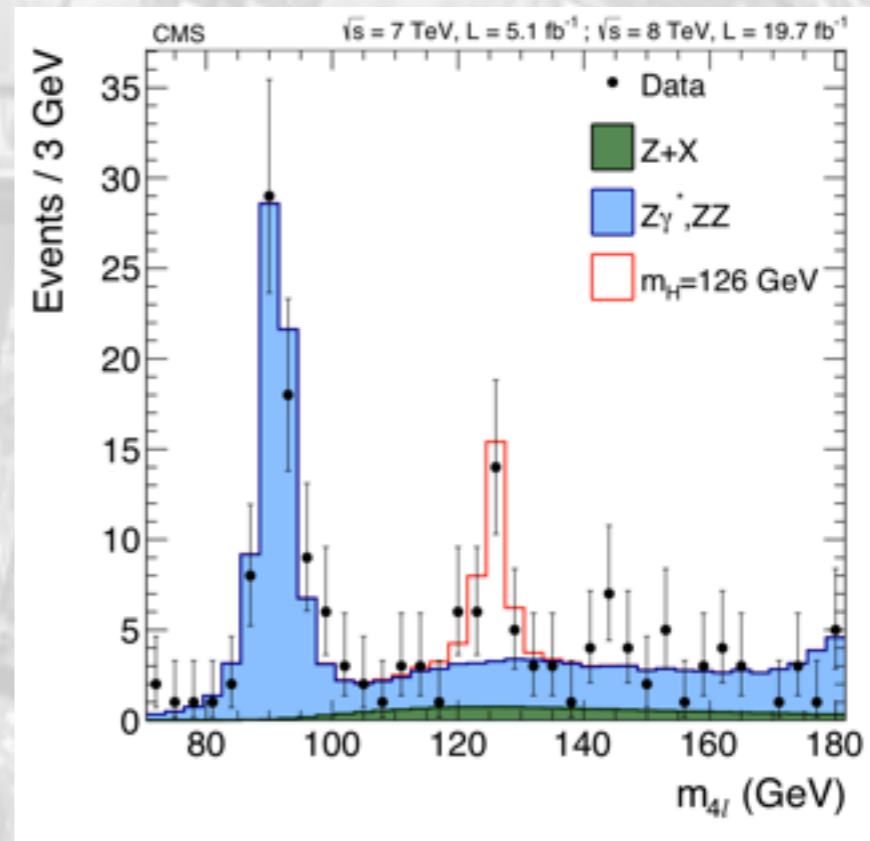
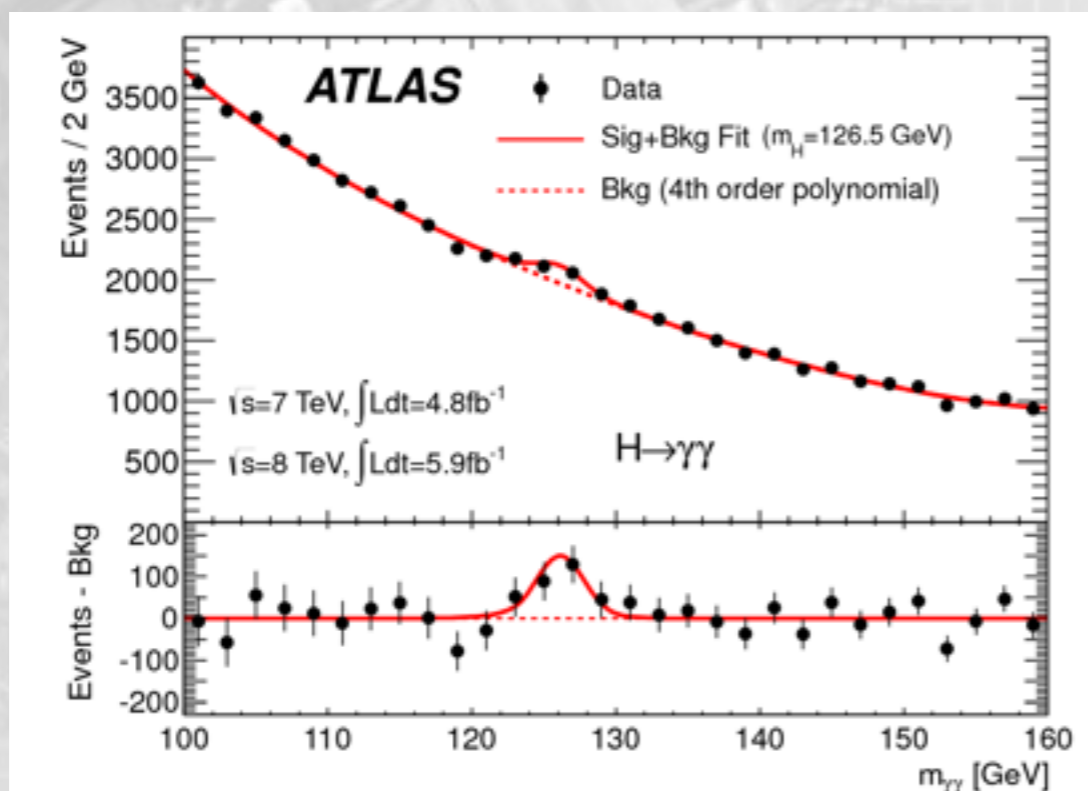




Search for Heavy Resonances at the LHC

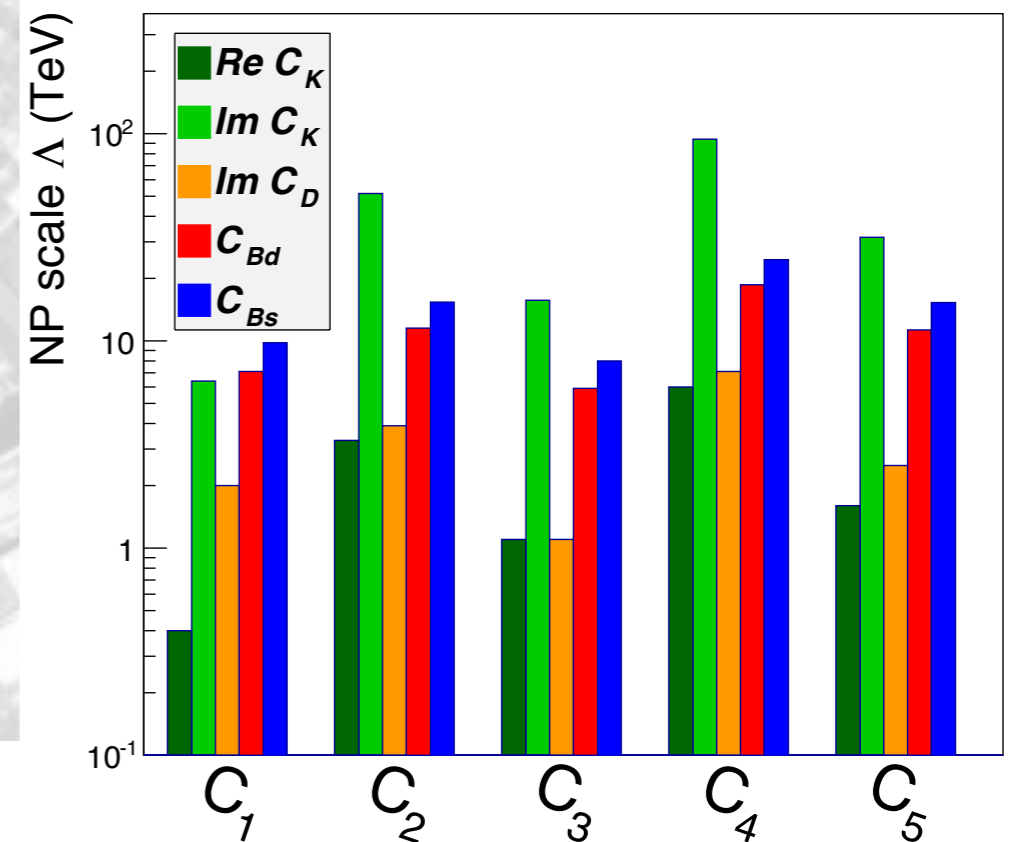
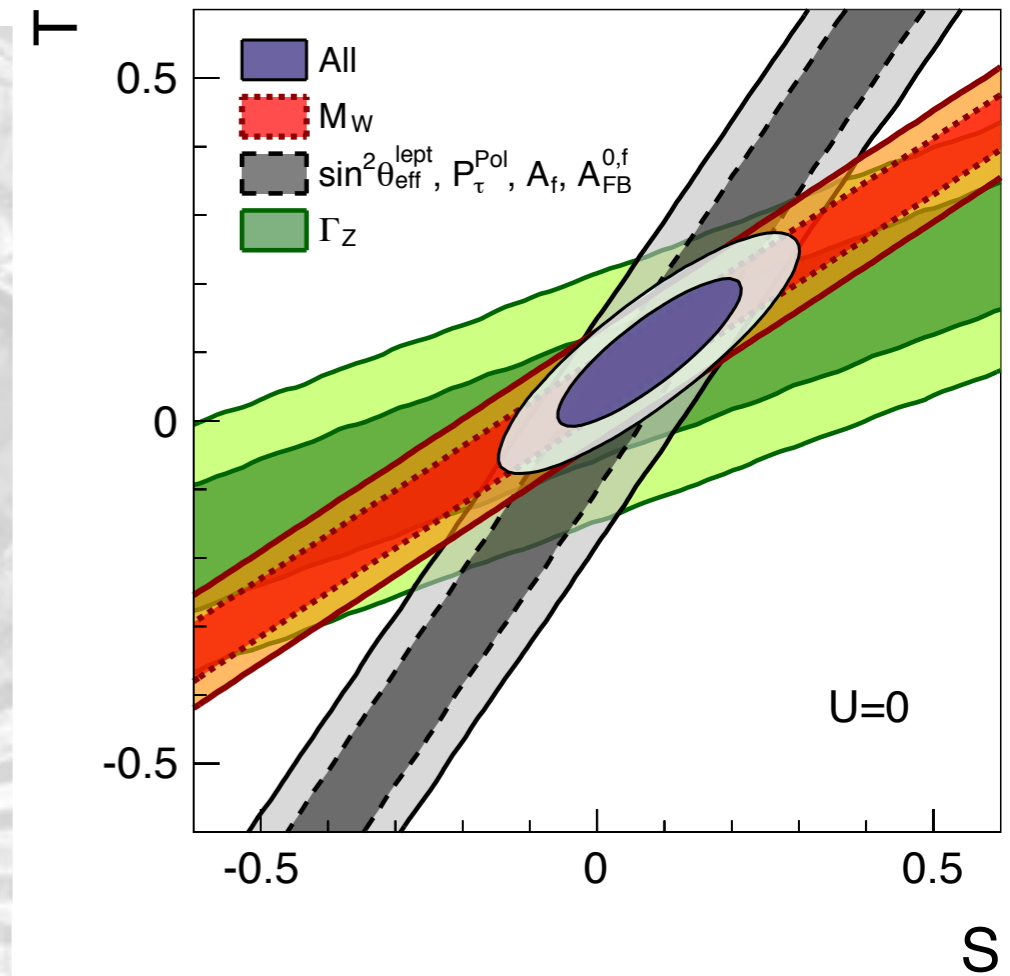
Maurizio Pierini
Caltech

SO FAR WE FOUND A RESONANCE



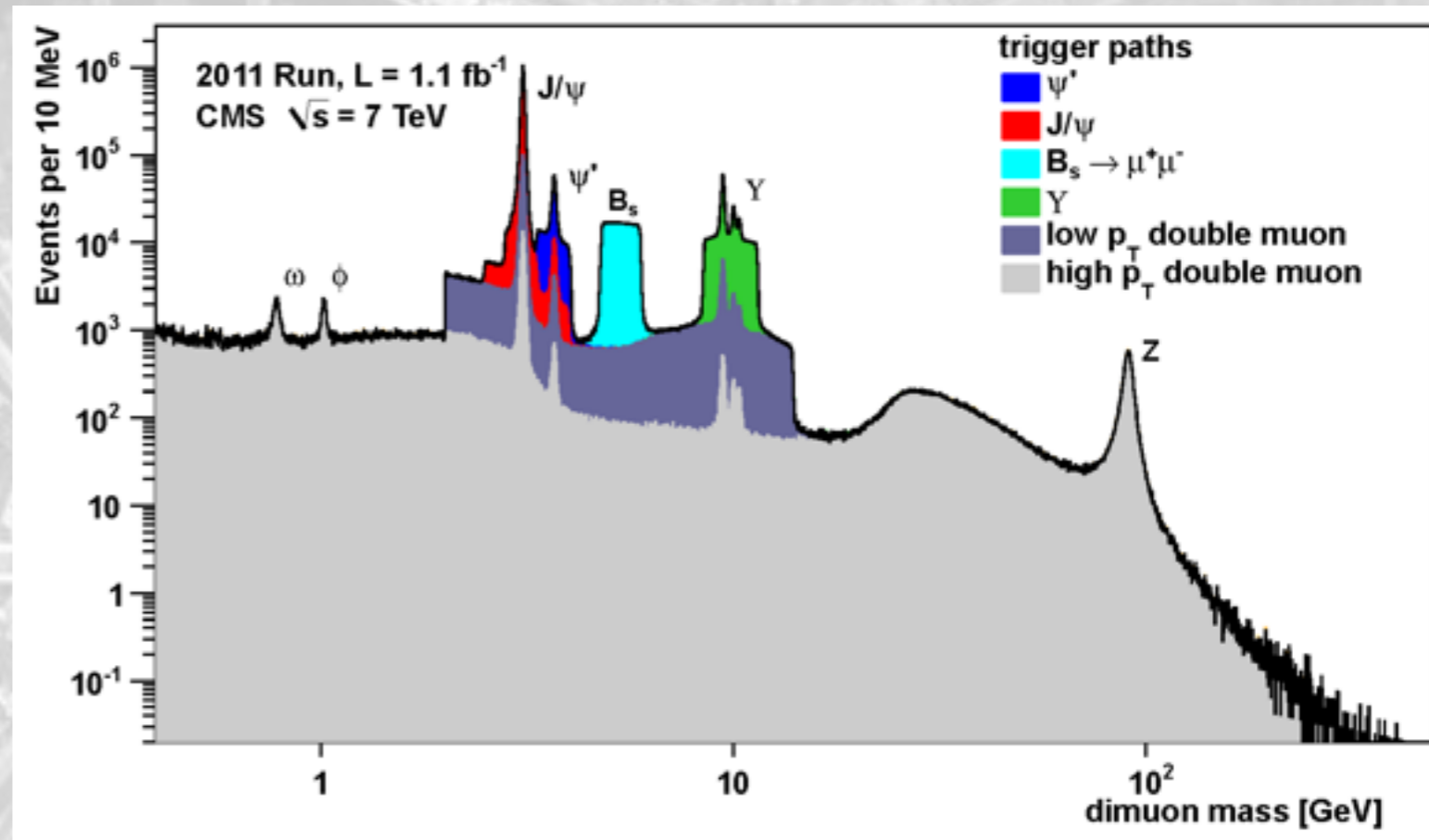
COULD THERE BE MORE?

- Sure... Why not?
- But we know from pre-LHC results that they should/could to be heavier than $O(1 \text{ TeV})$
 - Not seen in previous experiments (e.g, Tevatron)
 - Strong bounds from EW precision physics
- ... and/or have to come with SM-like features
 - no sign of NP in B/D/K physics pushed flavor-violating NP to PeV scale
 - even assuming (Next-to-)Minimal Flavor Violation, the bounds do not generically relax below the TeV



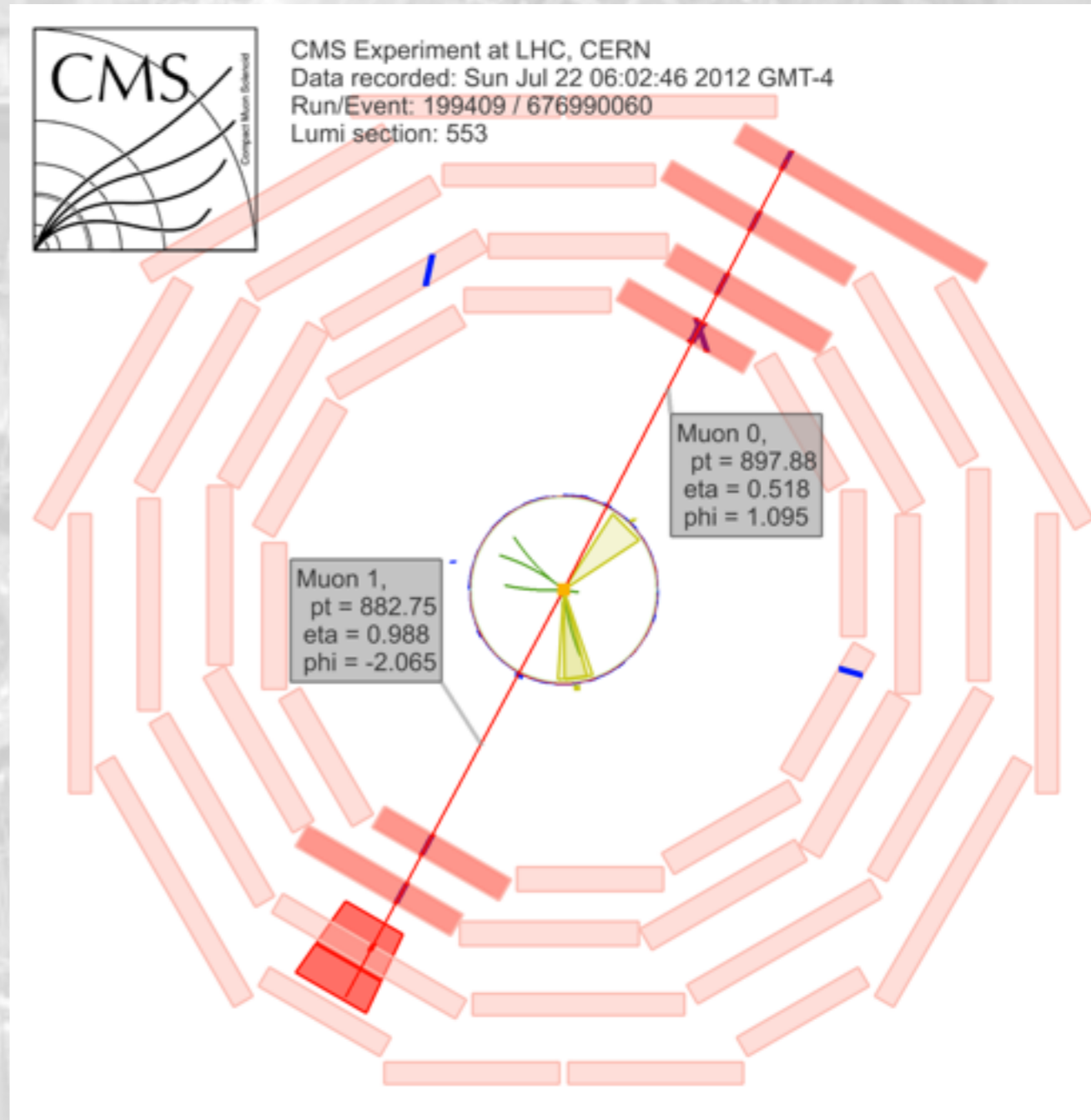
THE EXPERIMENTAL CHALLENGE

- For masses above 1 TeV, new resonances would decay to high- p_T particles
- Experimentally, the high p_T regime is different than what expiated for Higgs physics
- Electrons mainly measured by calorimeter (not tracker)
- Muons have large p_T uncertainties
- Jets not necessarily come from QCD (quarks and gluons)
- Dedicated reconstruction techniques in place, optimized to keep good performances at large- p_T
- **The peculiarity of ATLAS in CMS is the excellent performances in a large range of p_T (~ 1 GeV to ~ 1 TeV)**



HIGH-PT MUONS

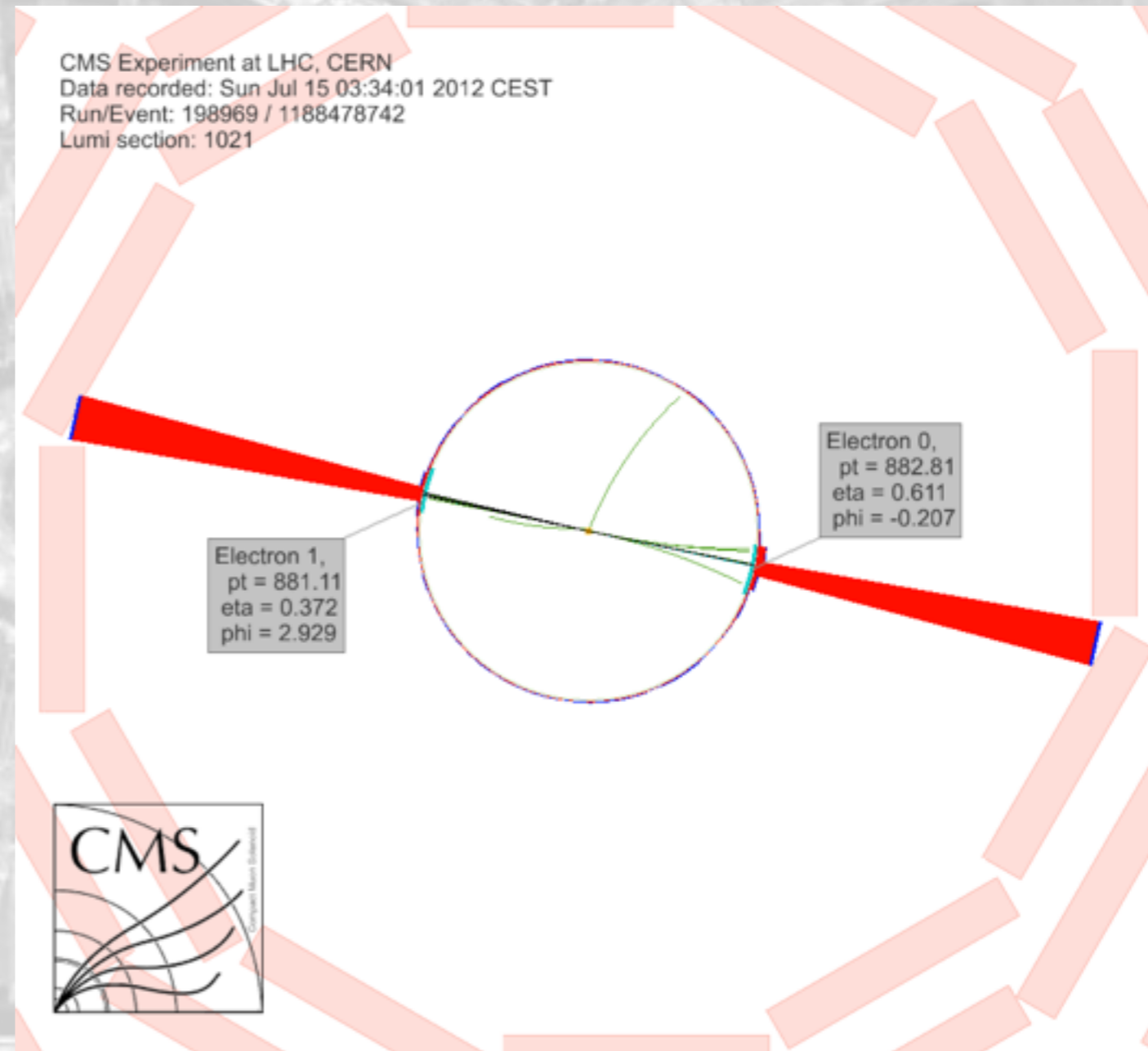
- Muon momenta are measured through the bending in the magnetic fields
- The bending is reduced at large muon momenta
- For high- p_T muons, the precision deteriorates
- Unlike for W/Z /top/ H , muon final states are not the golden channel for very-high- p_T physics
- Despite the resolution, high- p_T muons are an excellent discovery tool



$$\sigma_{p_T}/p_T \sim p_T/(qBL)$$

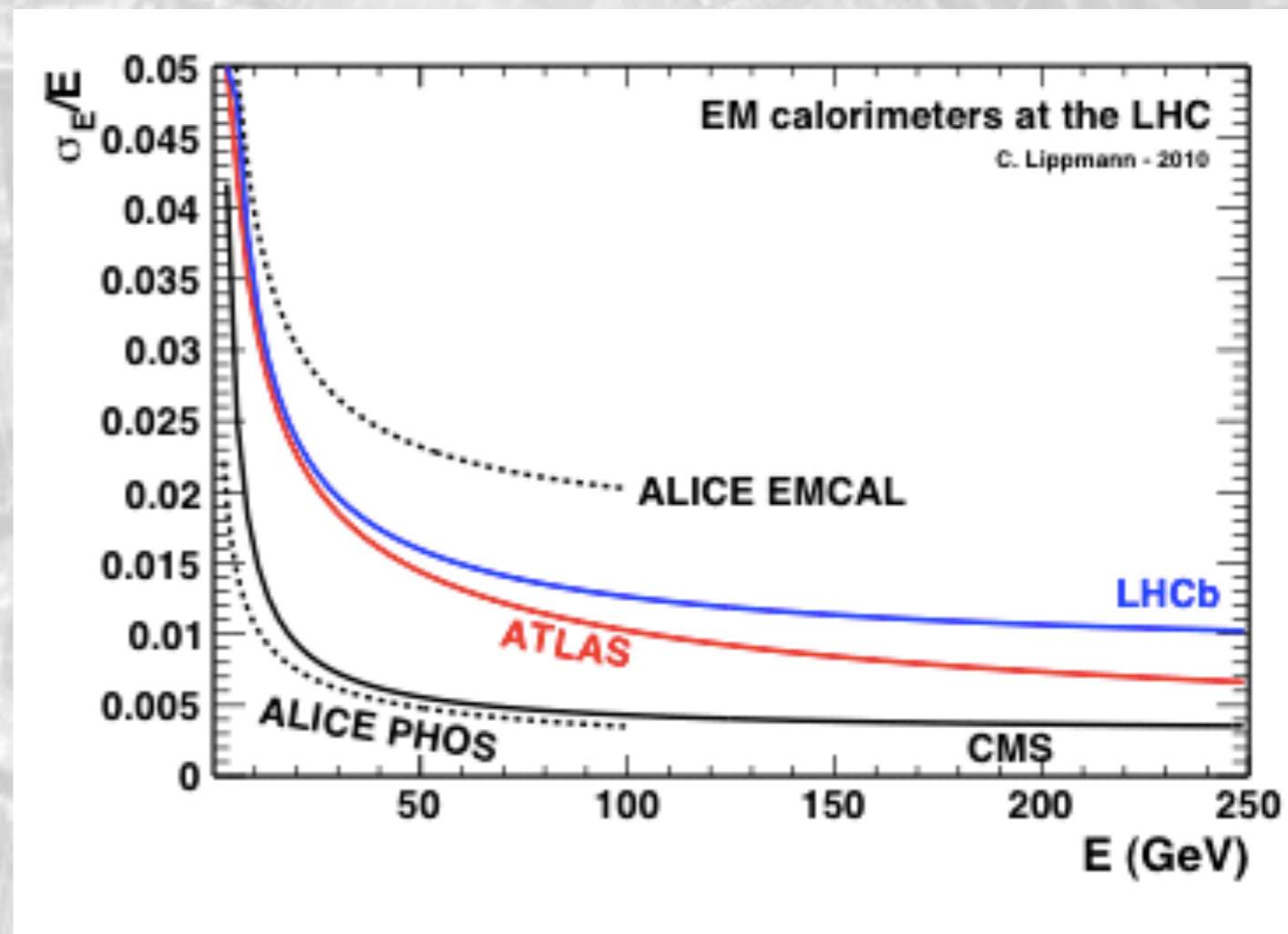
HIGH-PT ELECTRONS

- Electron momenta are measured in the tracker and in the calorimeter
- The resolution of the calorimeter improves with energy, giving a better S vs B discrimination above 1 TeV
- Electrons (and photons) are excellent tools to search for heavy resonances and measure their masses



HIGH-PT ELECTRONS

- Electron momenta are measured in the tracker and in the calorimeter
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- Electrons (and photons) are excellent tools to search for heavy resonances and measure their masses



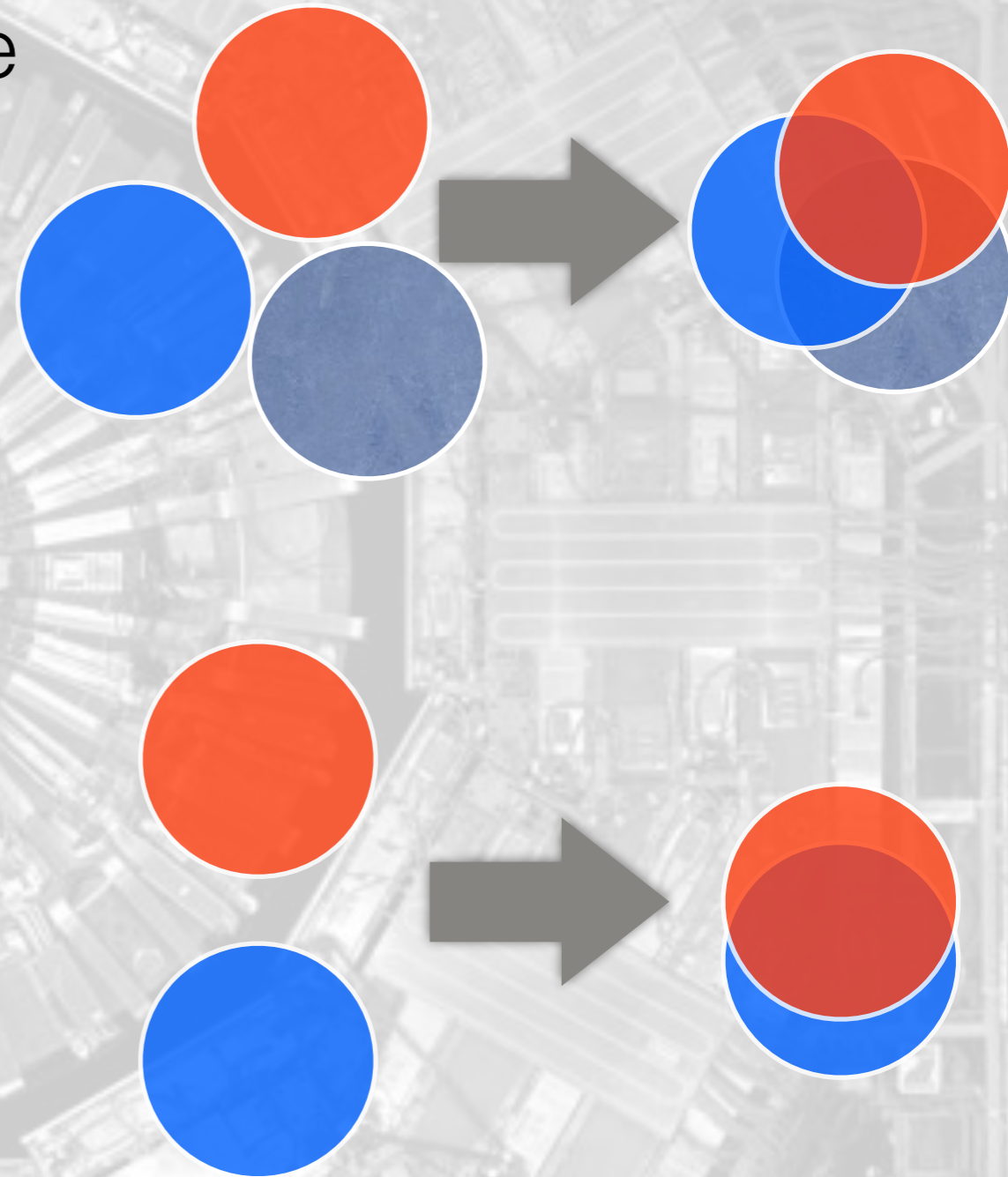
$$\frac{\sigma(E)}{E} = \frac{a}{\sqrt{E}} \oplus \frac{b}{E} \oplus c$$

$a \sim 0.027 \text{ GeV}^{1/2}$
 $b < 200 \text{ MeV}$
 $c \sim 0.005$



JETS BEYOND QUARKS & GLUONS

- High- p_T top/W/H/Z can be boosted enough for the final-state jets to merge into a single jet
- These jets are special: particles cluster in multiple “poles” in (η, ϕ)
→ jet substructure
- Jet reconstruction to be pushed to the next level
- QCD multijet background becomes an issue for more than just hadronic searches



EXAMPLES OF RESONANCE SEARCHES

- Search for XVV resonances ($V=W/Z$)
 - Fully-hadronic final states and substructure
 - But also final states with leptons
- Ditop, dilepton and dijet resonance searches

DISCLAIMER:

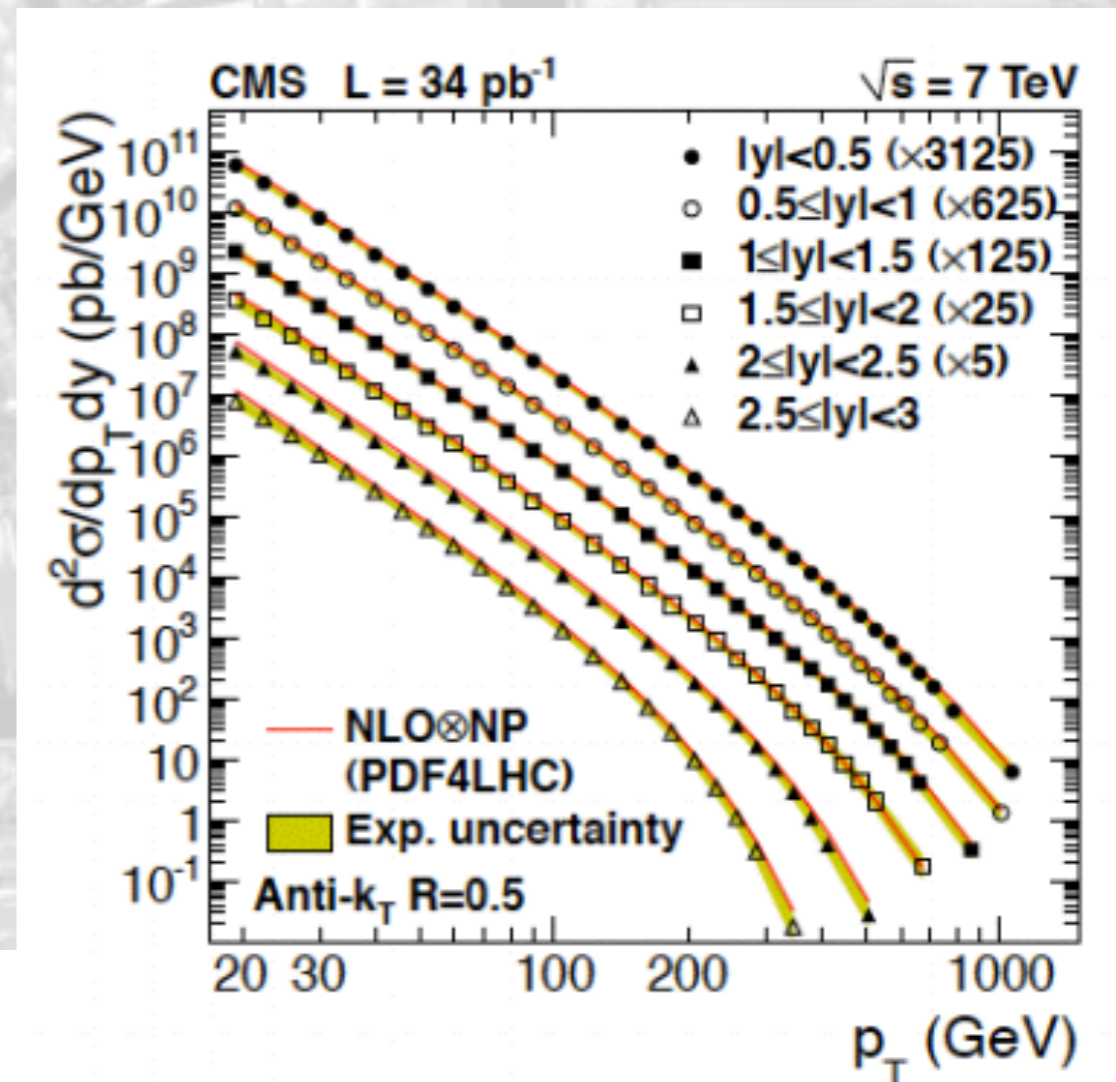
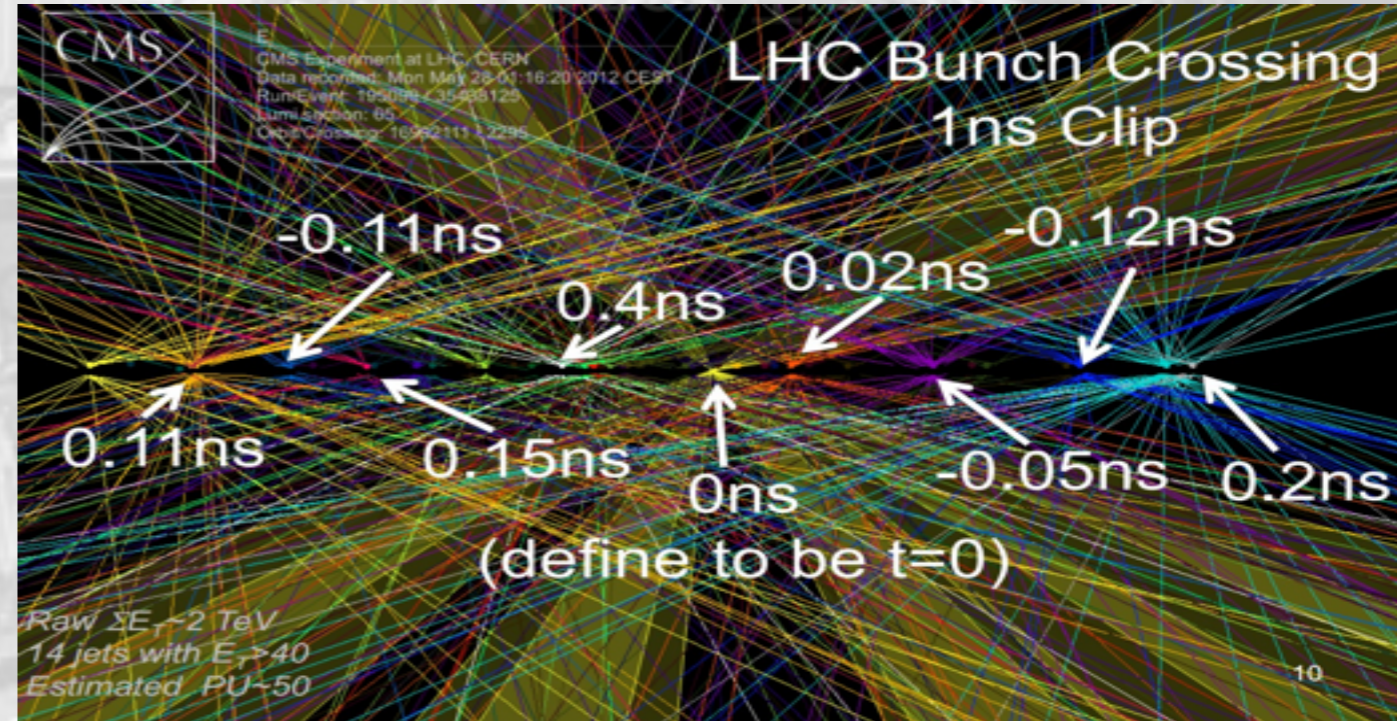
- I mainly show CMS results
- I also show some ATLAS result for comparison (which you saw this morning already)
- When I don't show ATLAS results, I am not implying that they don't exist. I was just lazy...



**Jets, Jet
substructure, and
 $X \rightarrow \nu\nu$ searches**

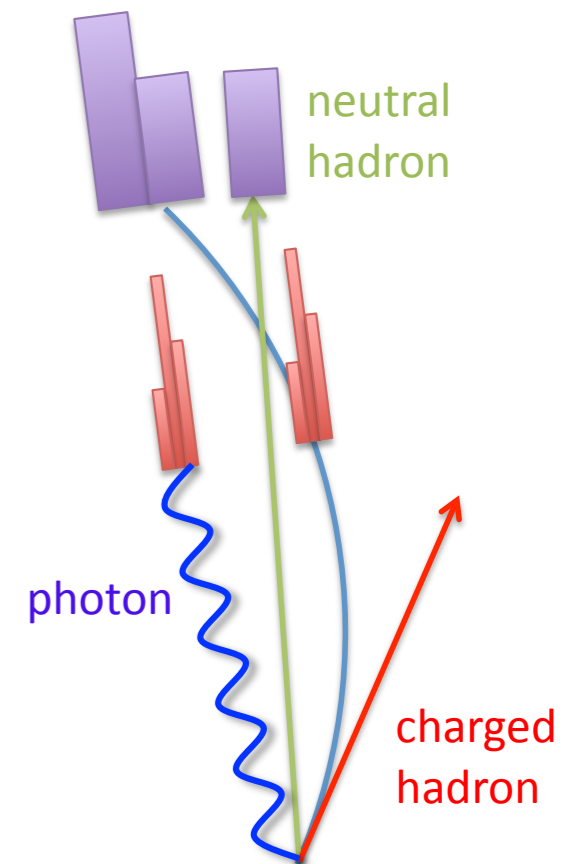
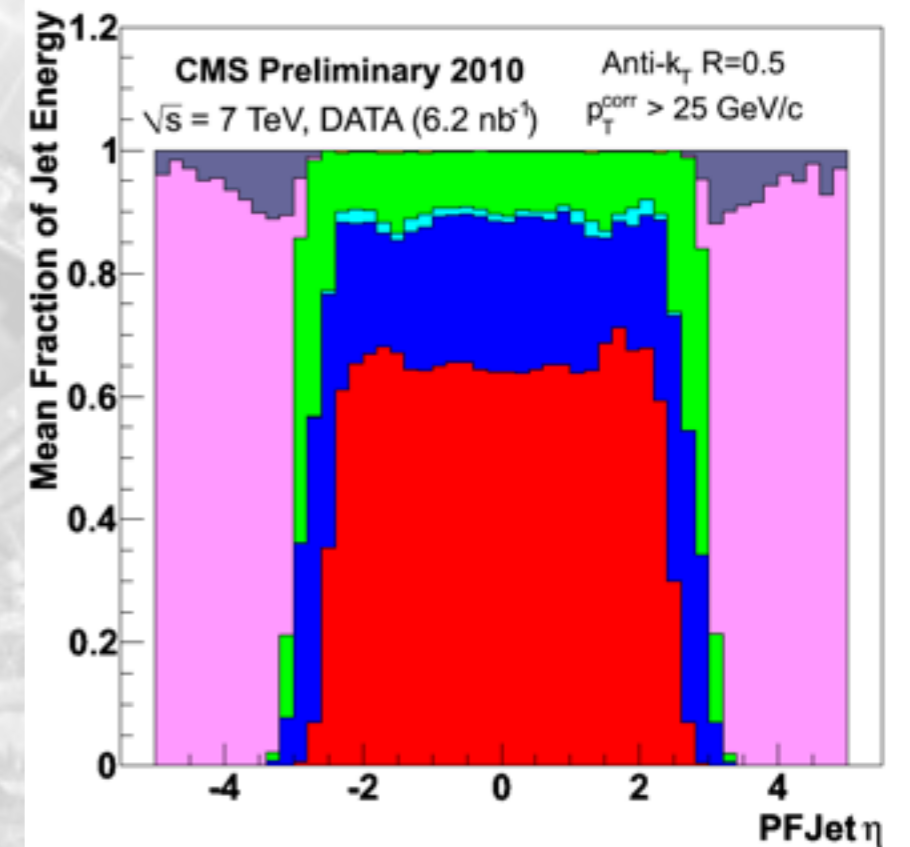
PRECISION PHYSICS WITH JETS

- QCD physics at the LHC is precision physics
 - Multijet NLO calculations in event generators
 - Solid jet definition
 - Handling of pileup
- During LHC run I, intense program of measurements
 - “Validated” the tools describing jet physics (e.g., event generators and detector simulation)
 - Set the basis to search for New Physics with jets



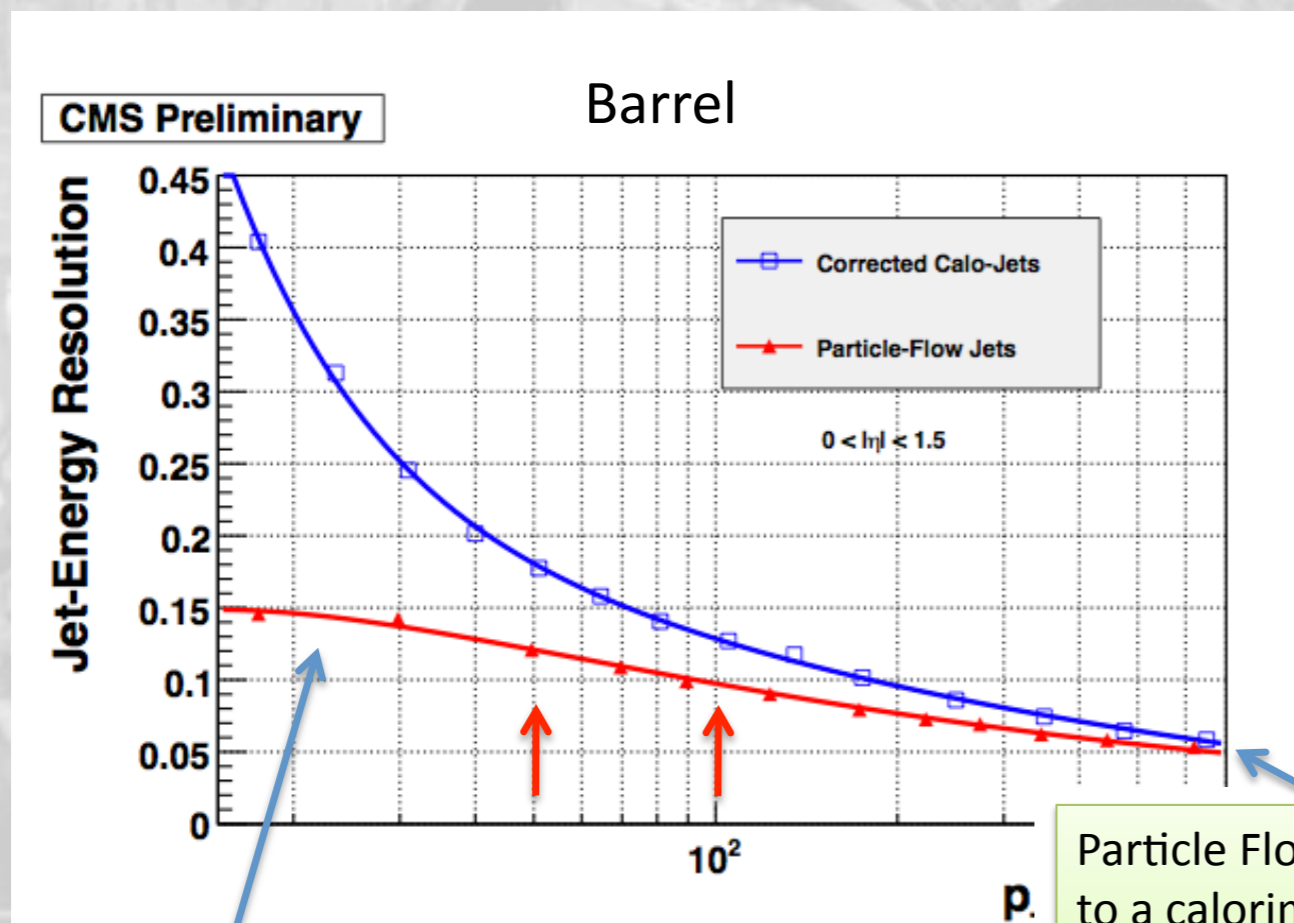
JETS OF WHAT?

- What if your HCAL is not that great? Physics at rescue
 - 60% of the jet is made of **charged hadrons**
 - a largely fraction of the neutral hadrons is pions, which decay to **photon** pairs before your HCAL
- One can then
 - reconstruct individual hadrons, matching tracks to their energy deposit (the good p_T resolution compensated for the bad energy resolution of the HCAL)
 - use the ECAL to reconstruct the energy deposit there
 - Live with the poor resolution on **neutral hadrons** (e.g., neutrons, KL, etc)



PARTICLE FLOW

- Combine the information from all detectors to reconstruct single particles
- Provides lists of particles (e, μ, γ , charged and neutral hadrons)
- Improves HCAL resolution with tracker



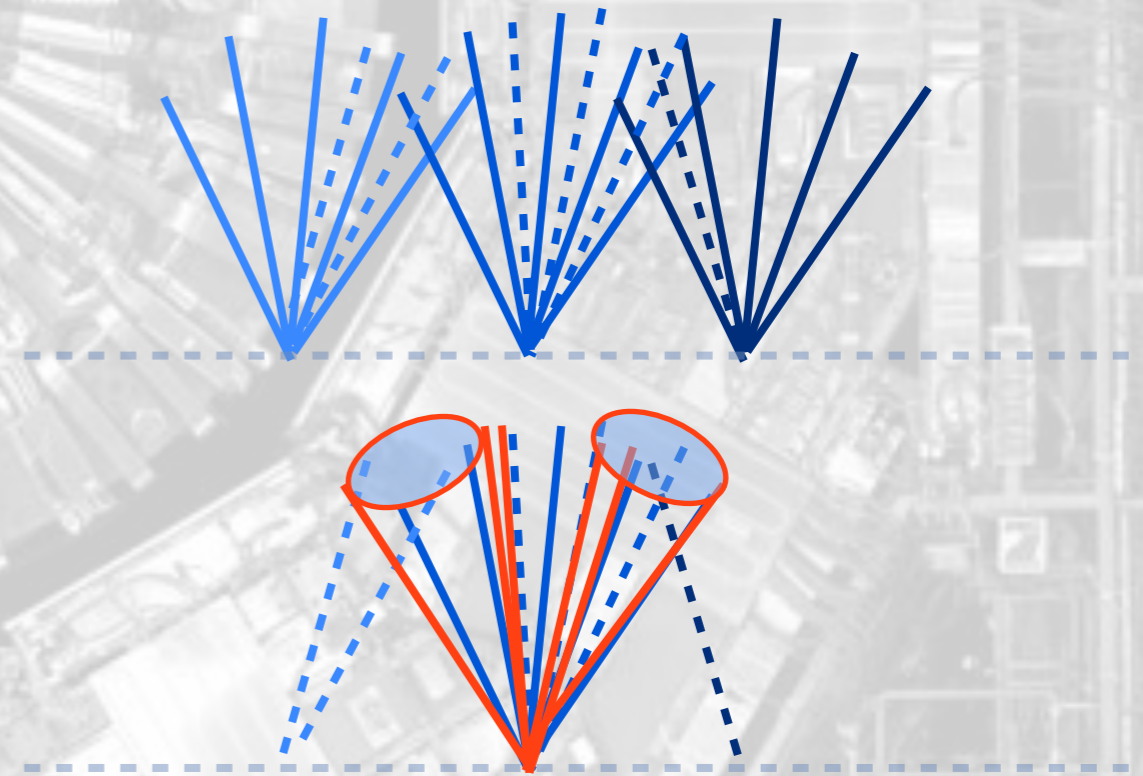
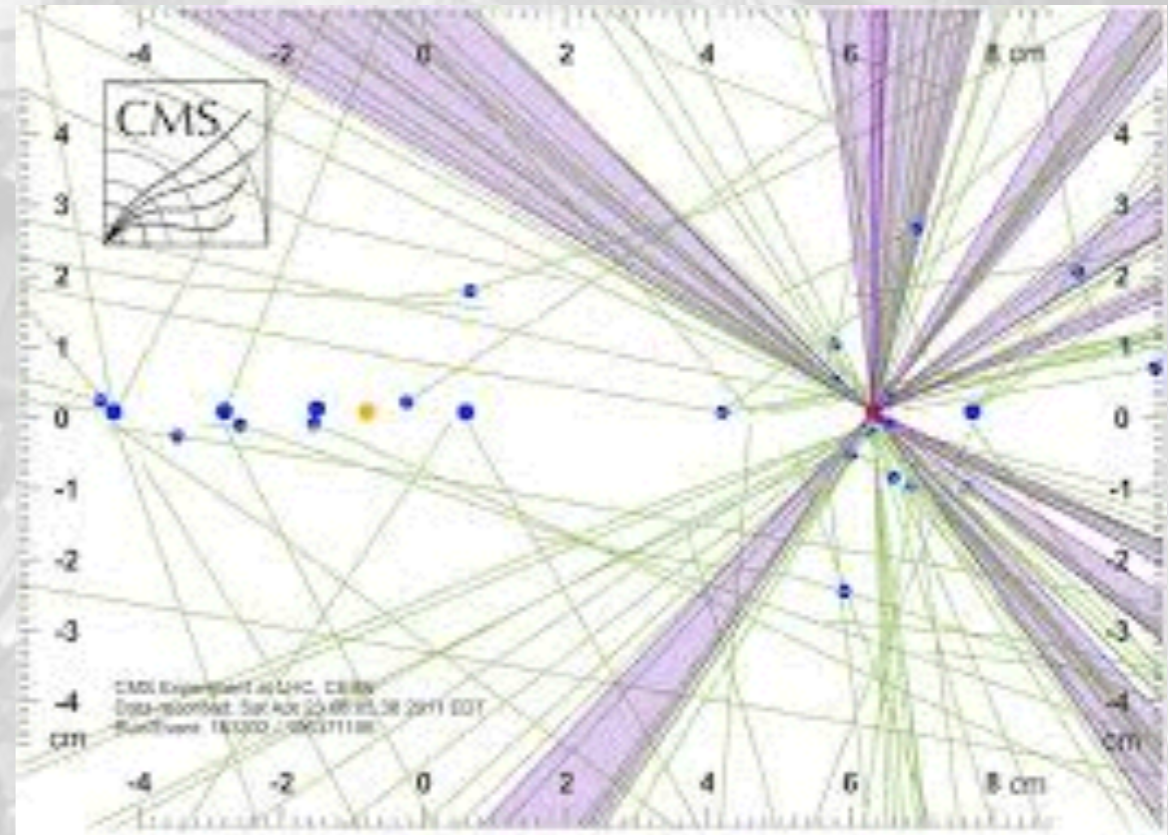
Factor 2 improvement at low p_T

Particle Flow converges to a calorimetric measurement at high p_T

- Replace the HCAL granularity with tracker granularity (important for jet substructure)
- The final result is a list of particles, similar to a generator-level study

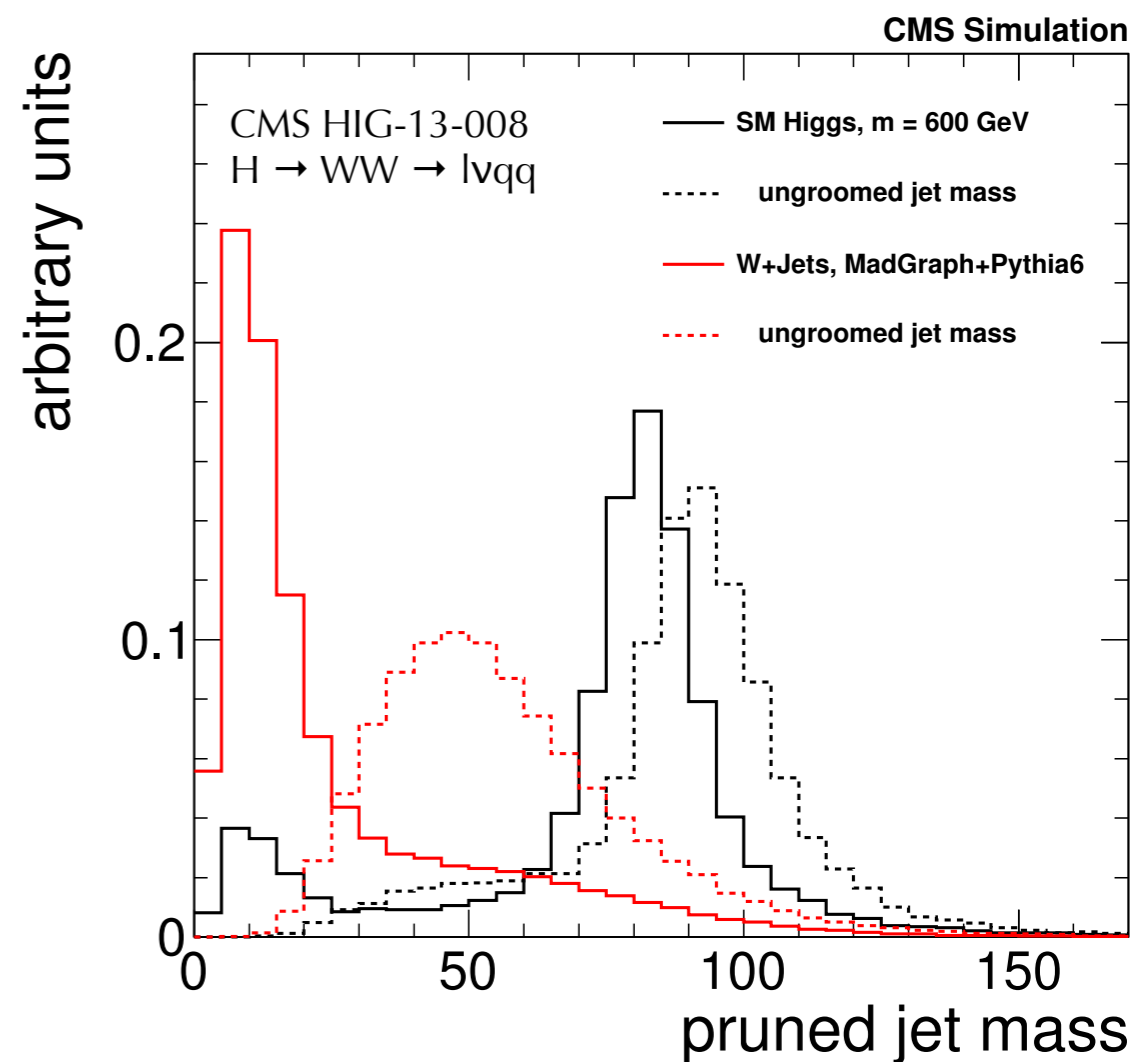
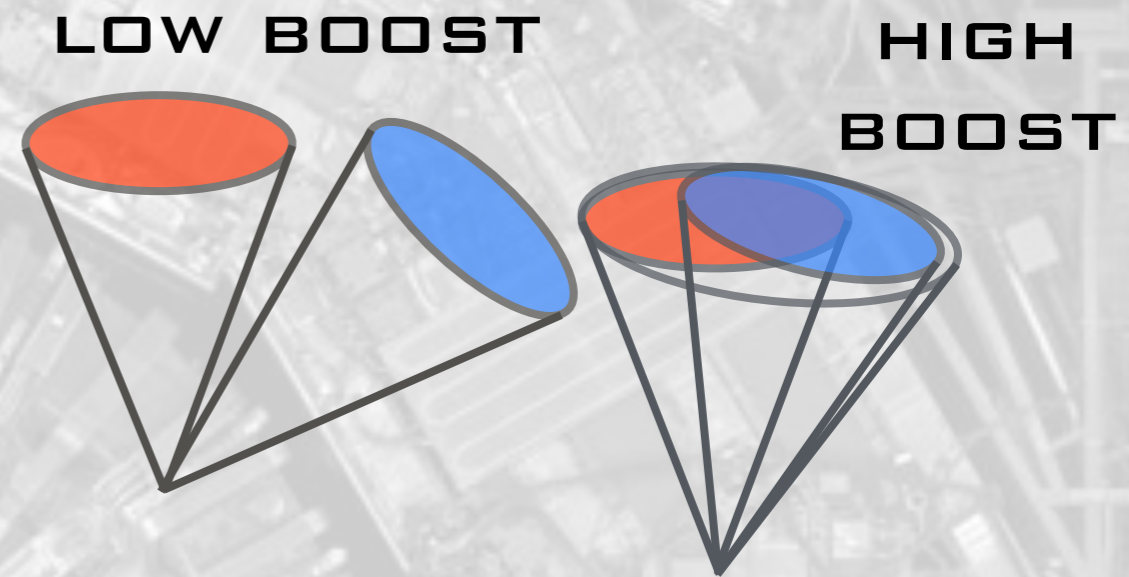
THE PILEUP PROBLEM

- Pileup is particularly problematic for jets
- particles from different vertices overlap in the calorimeter area
- the jets are then to be found on top of a diffuse noise from additional collisions
- These collisions are typically soft (a high- p_T collision is a rare event)
- Several methods put in place to limit the impact of pileup on the jet reconstruction (charged hadron subtraction, jet vertex fraction, jet area subtraction, ...)



A SPECIAL KIND OF JET

- W/Z/H/top can decay to fully hadronic final states
- For large enough p_T , the decay products might merge into a single jet
- These jets are special, as their properties (e.g., the mass) are not those of a QCD jet

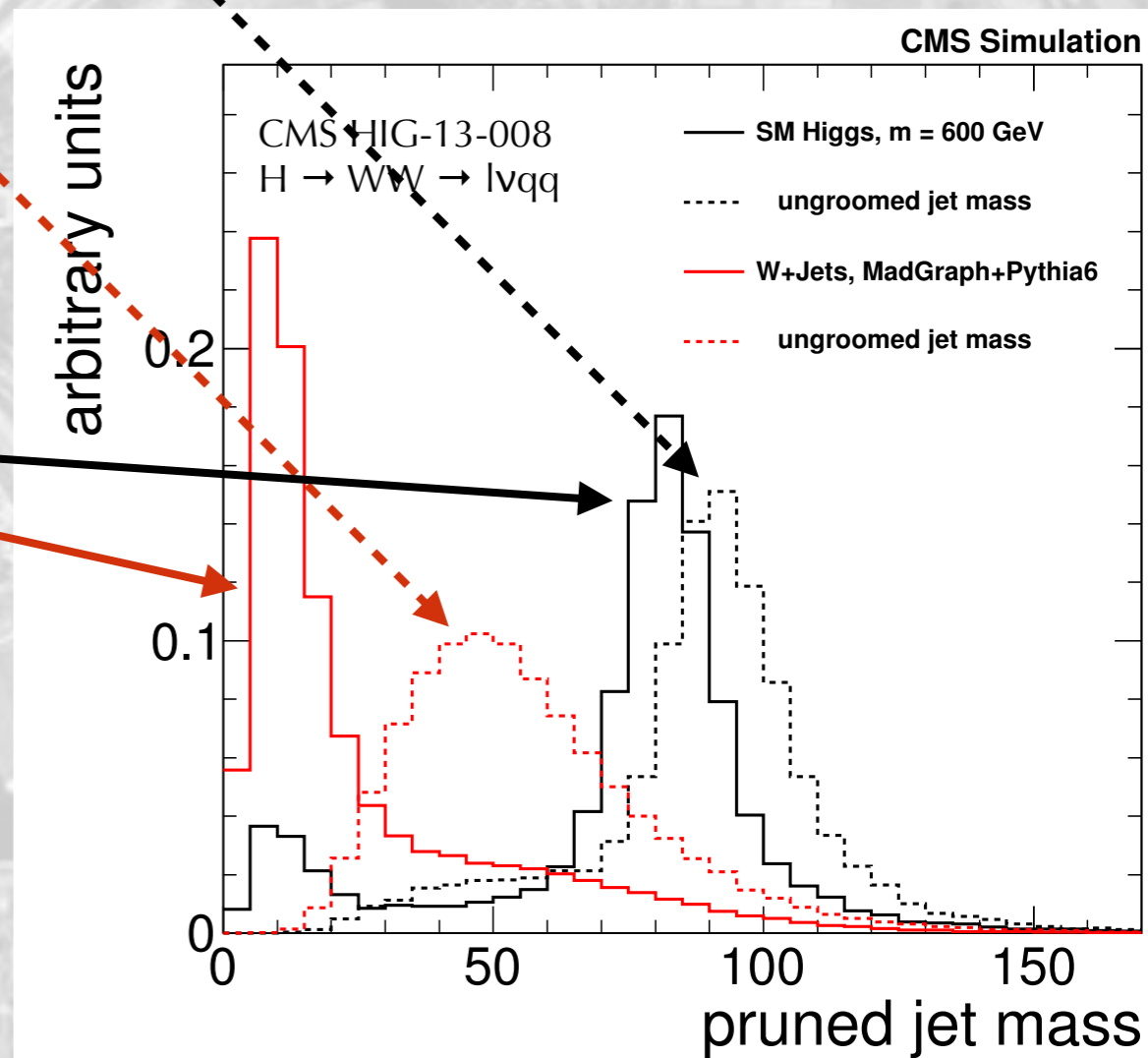
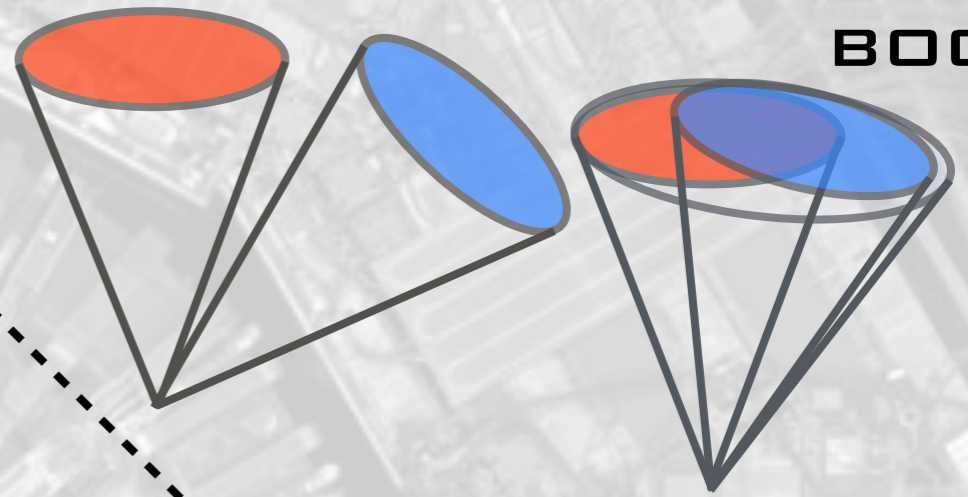


JET GROOMING

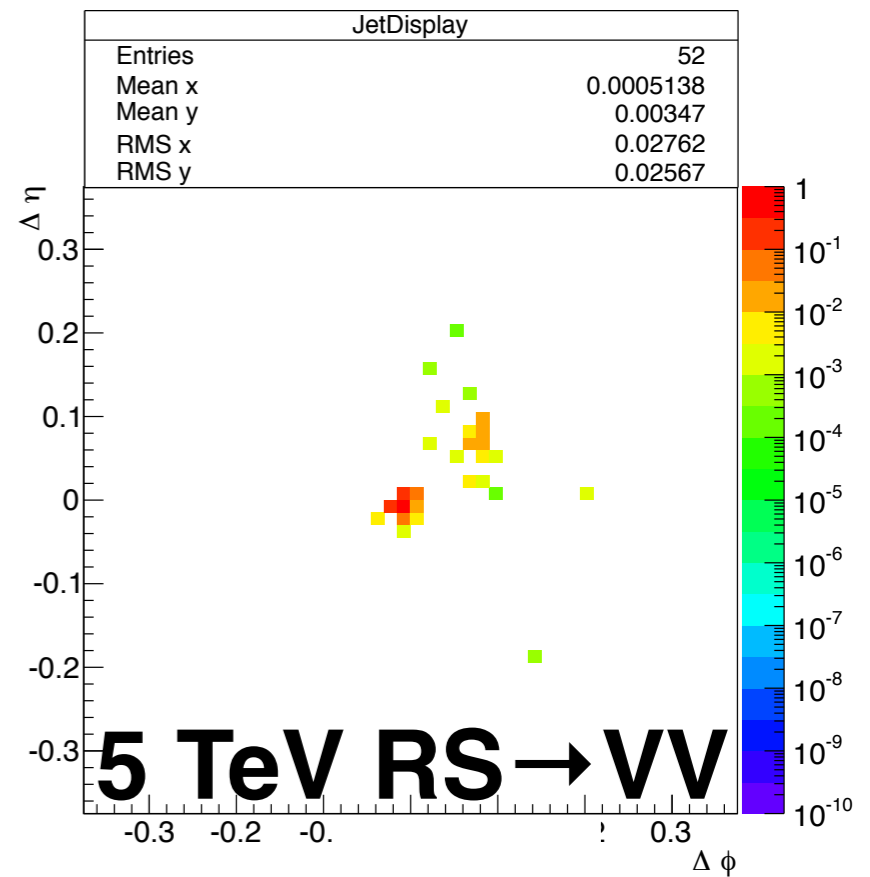
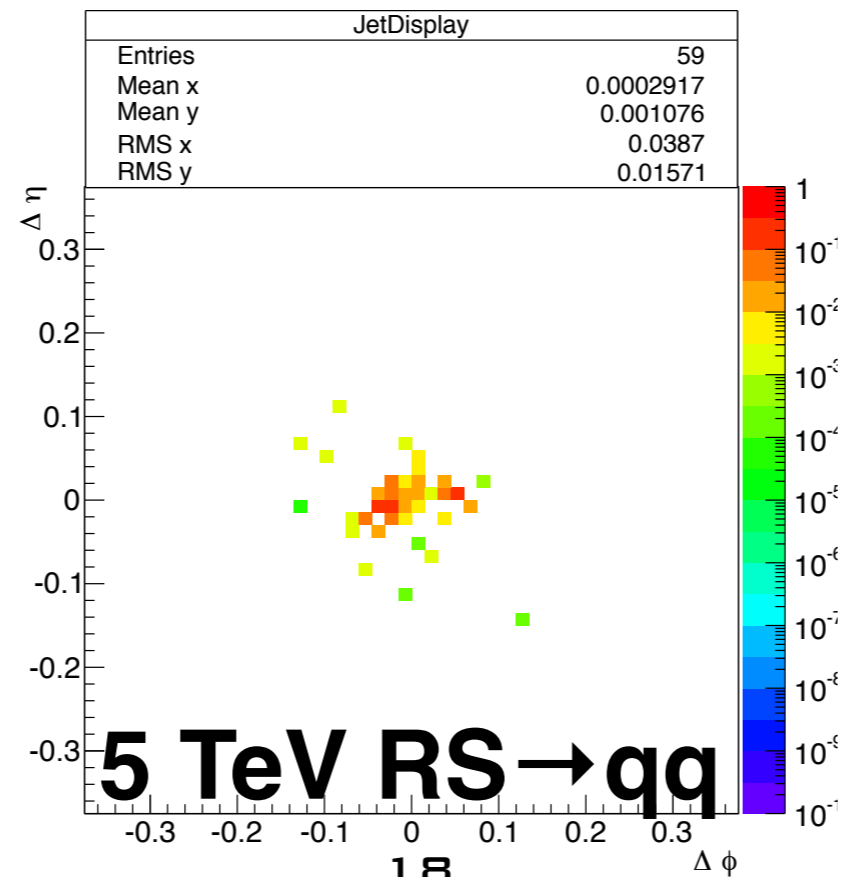
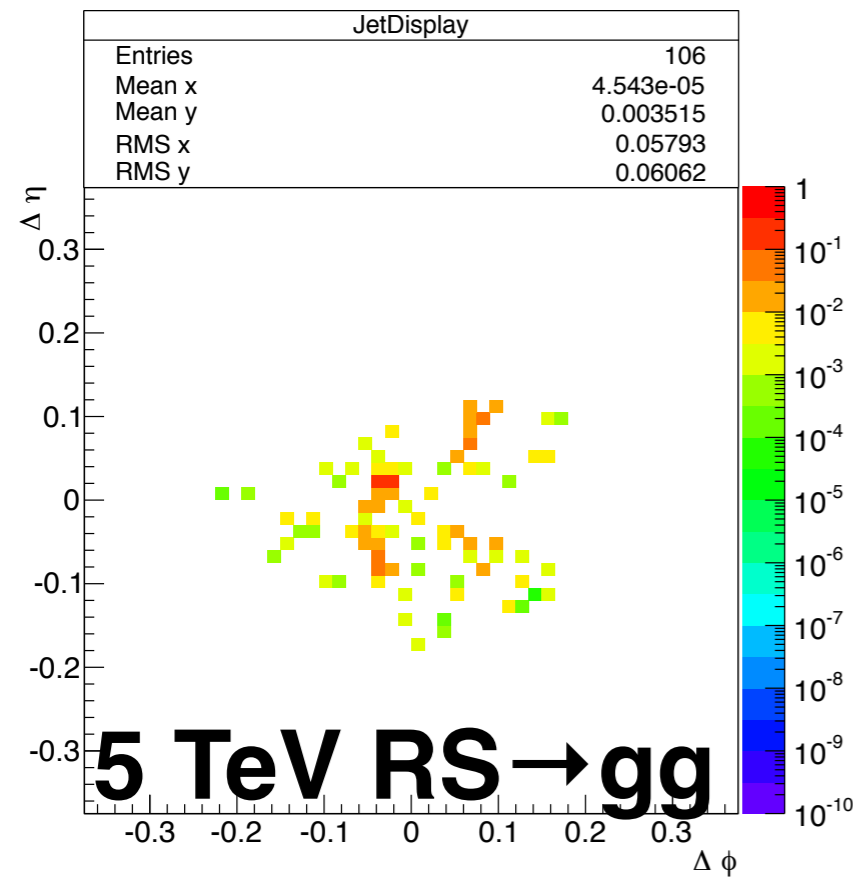
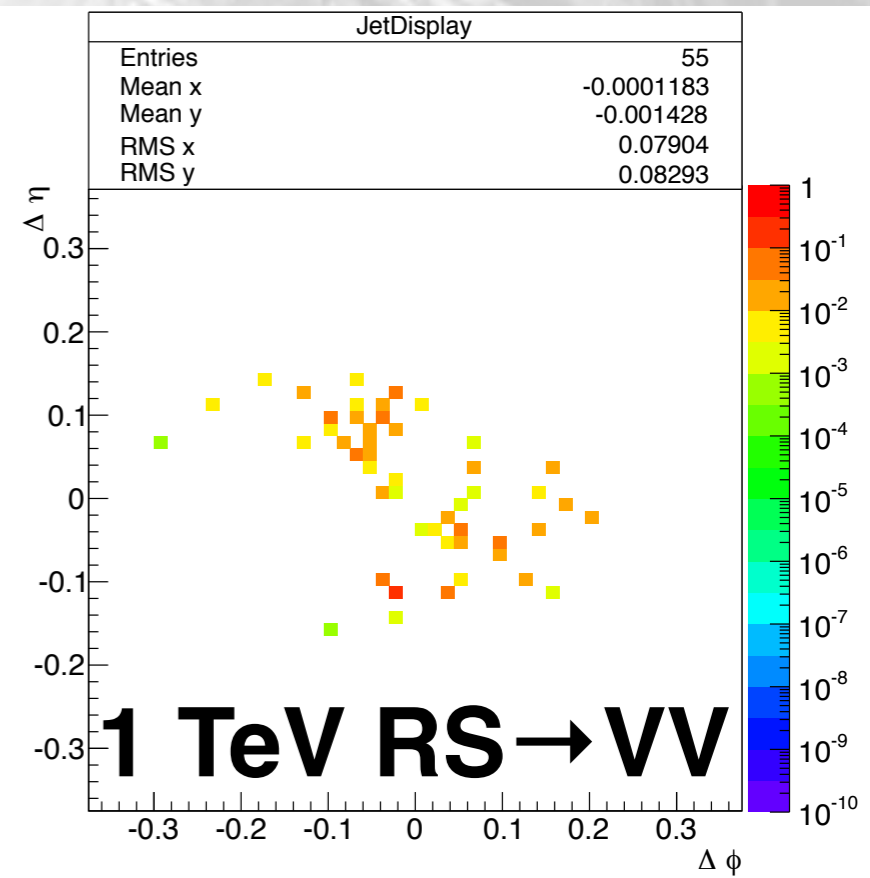
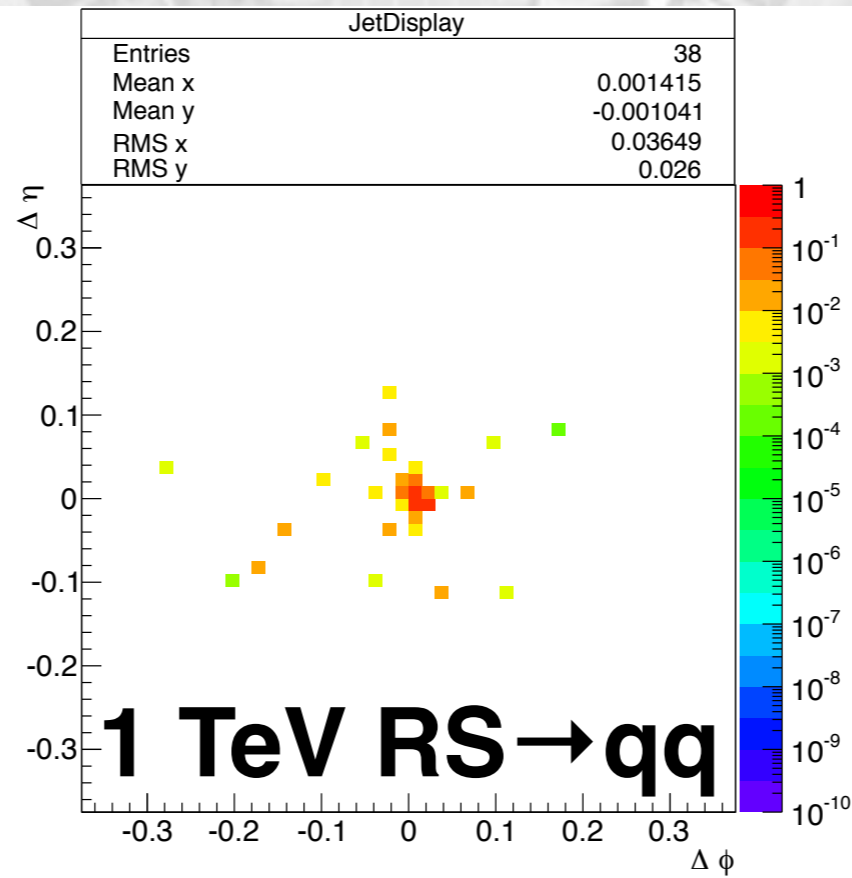
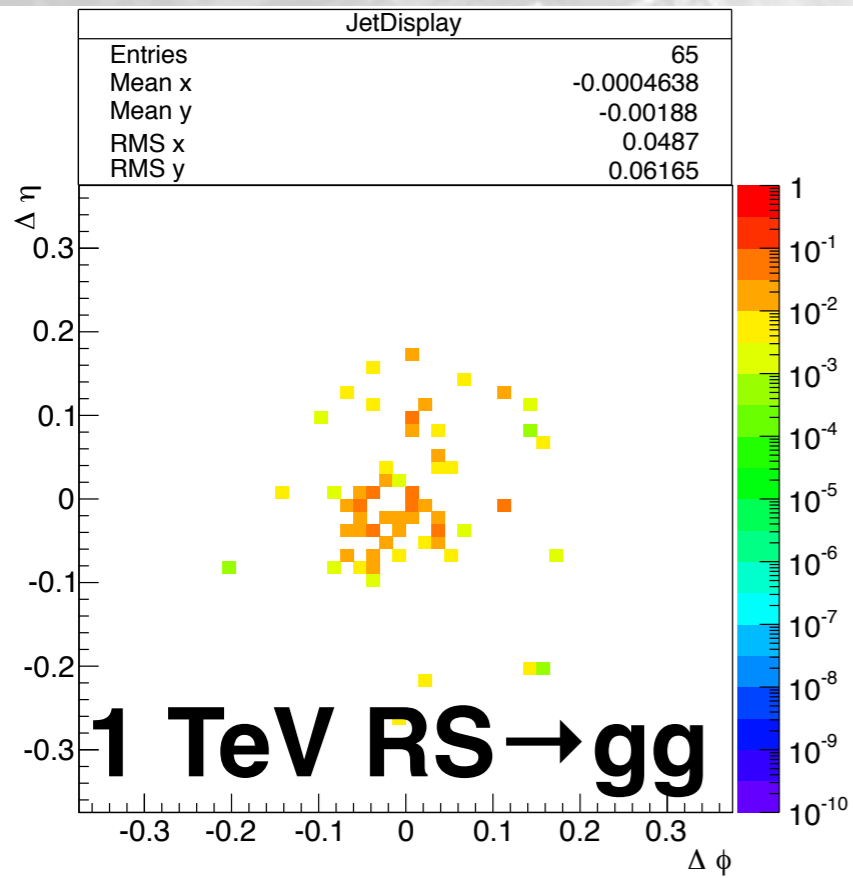
- Care is needed in defining the “right” mass
 - the plain mass for a jet from a colourless object (W/Z/H) peaks in the right spot
 - for a QCD jet, the mass depends on the jet p_T (i.e. it is NOT the right mass)
- The “right” mass comes from grooming (removal of soft and collinear radiation + pileup)
- Several techniques proposed and ad data opted for data by ATLAS & CMS
- More than the mass: the jet constituents have multipole distribution (due to the number of partons starting the showering)

LOW BOOST

HIGH BOOST

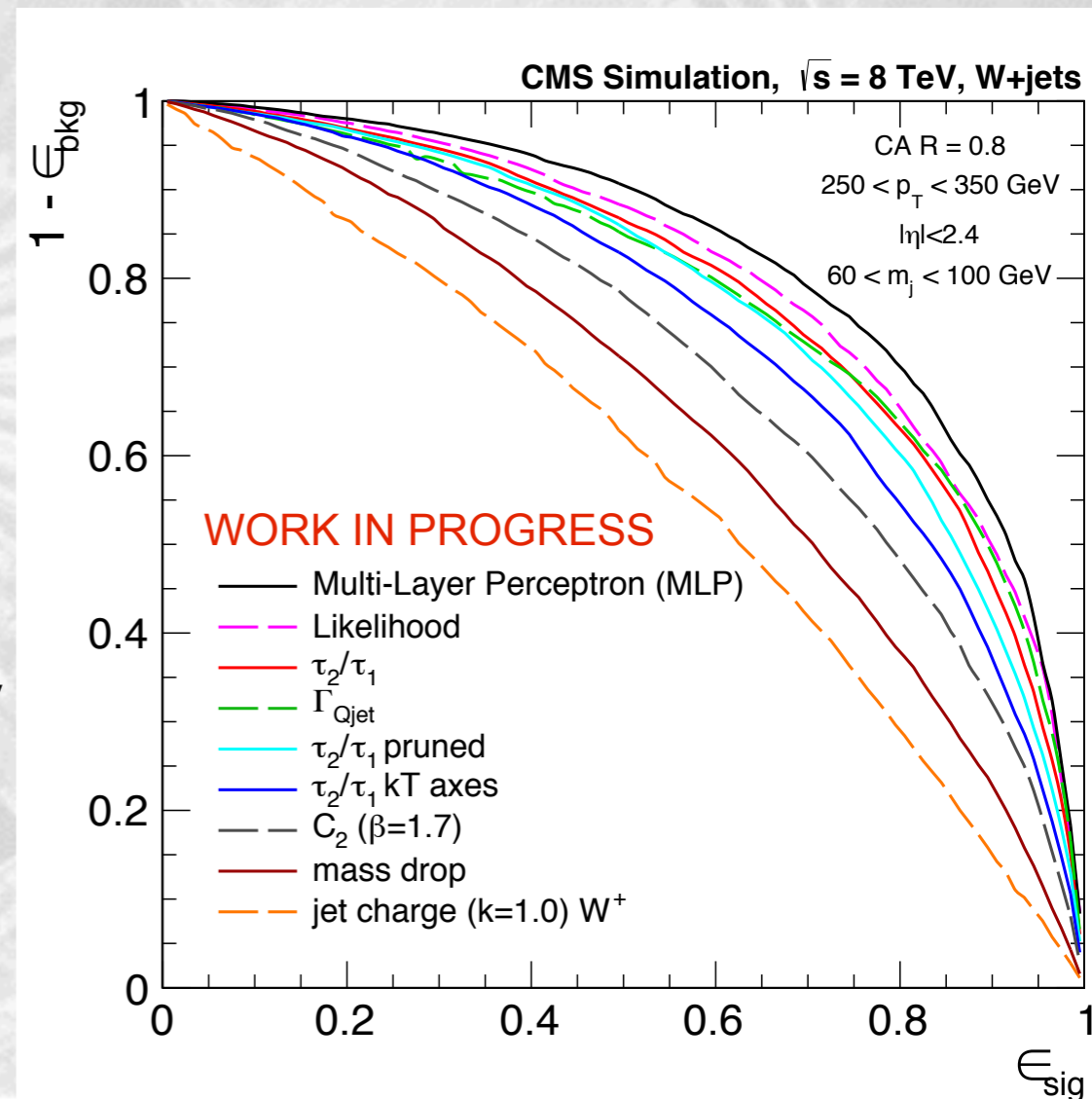


MORE THAN QCD JETS



SUBSTRUCTURE VARIABLES

- One wants to separate a multi-prong from a single-prong angular distribution of the decay products around the jet axis
- Several variables “on the market” to exploit this difference (see for instance CMS and ATLAS papers)
- A few lessons learned
 - the jet mass cut does much of the job
 - the higher the p_T , the smaller the QCD background. Better use substructure as an event classifier, not to loose efficiency
- In CMS, Particle Flow is important here (despite the high- p_T regime). Angular resolution matters, and tracker is better for that

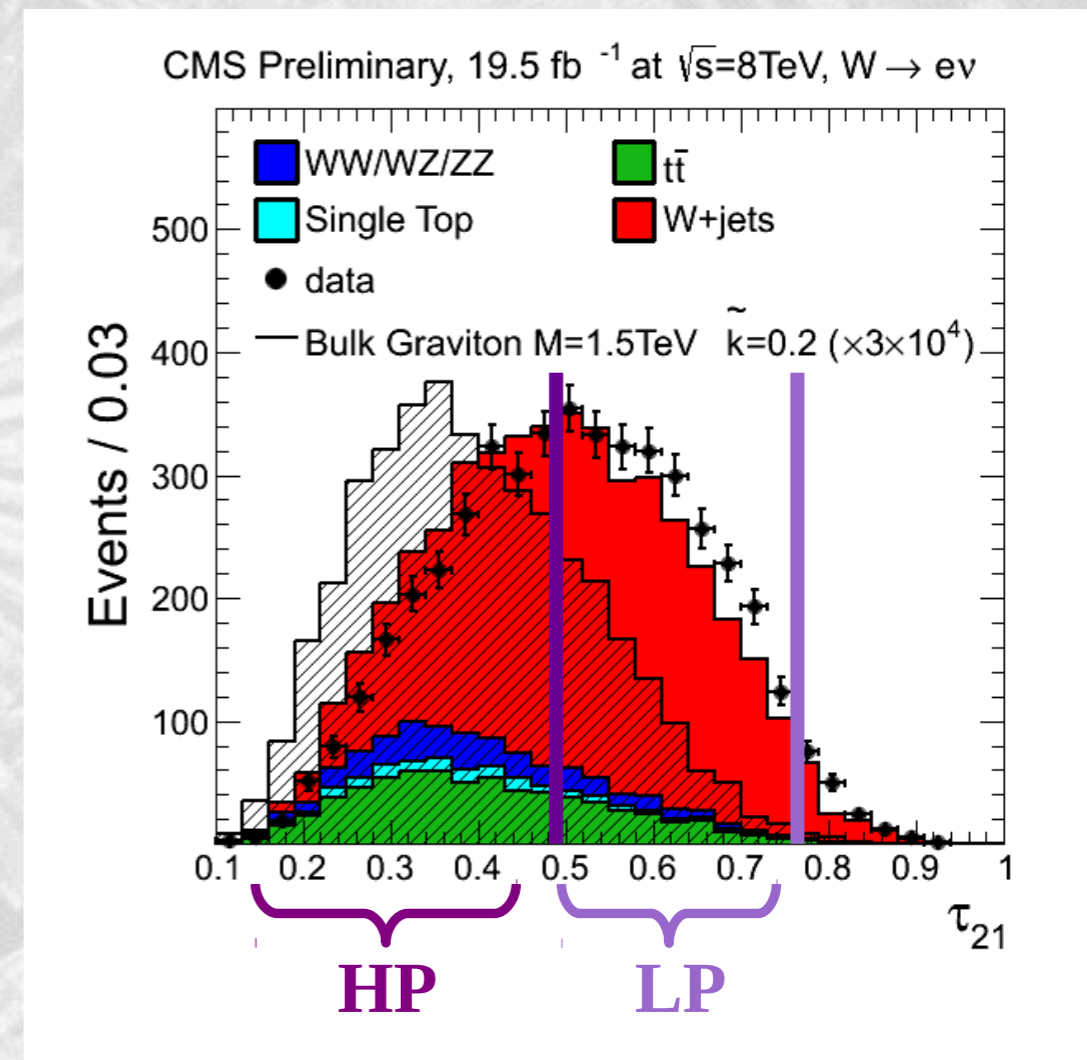


[CMS-PAS-JME-13-006](#)

JET N-SUBJETTINESS

- N-subjettiness is smaller if the constituents of a jet can be arranged in N subjets
- One can use the variable to test different hypotheses (e.g. V vs top vs QCD jet)
- In real life, ratios are particularly useful to categorize events (High purity vs Low Purity) after a mass cut on the jet is applied
- Correlation with jet mass and PU effects tend to reduce the discrimination power
- Polarization also matters (e.g. separation more effective for V_L than V_T)

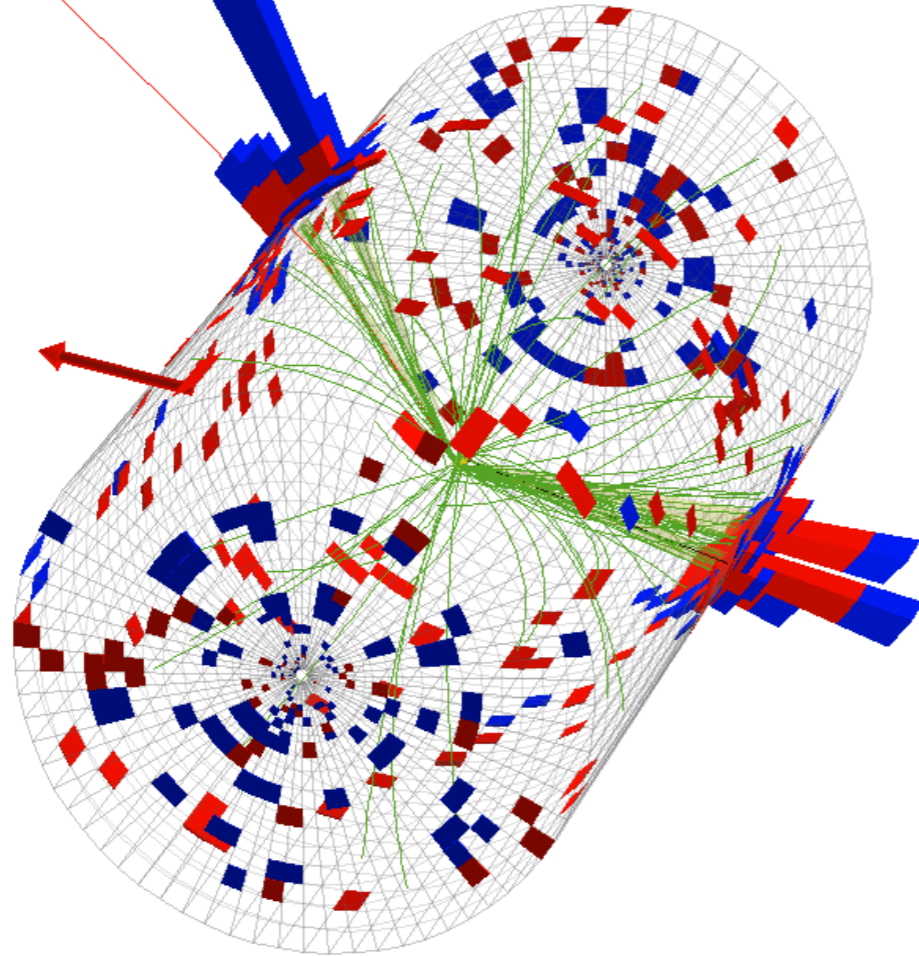
$$\tau_N = \frac{1}{d_0} \sum_k p_{T,k} \min \{ \Delta R_{1,k}, \Delta R_{2,k}, \dots, \Delta R_{N,k} \}$$



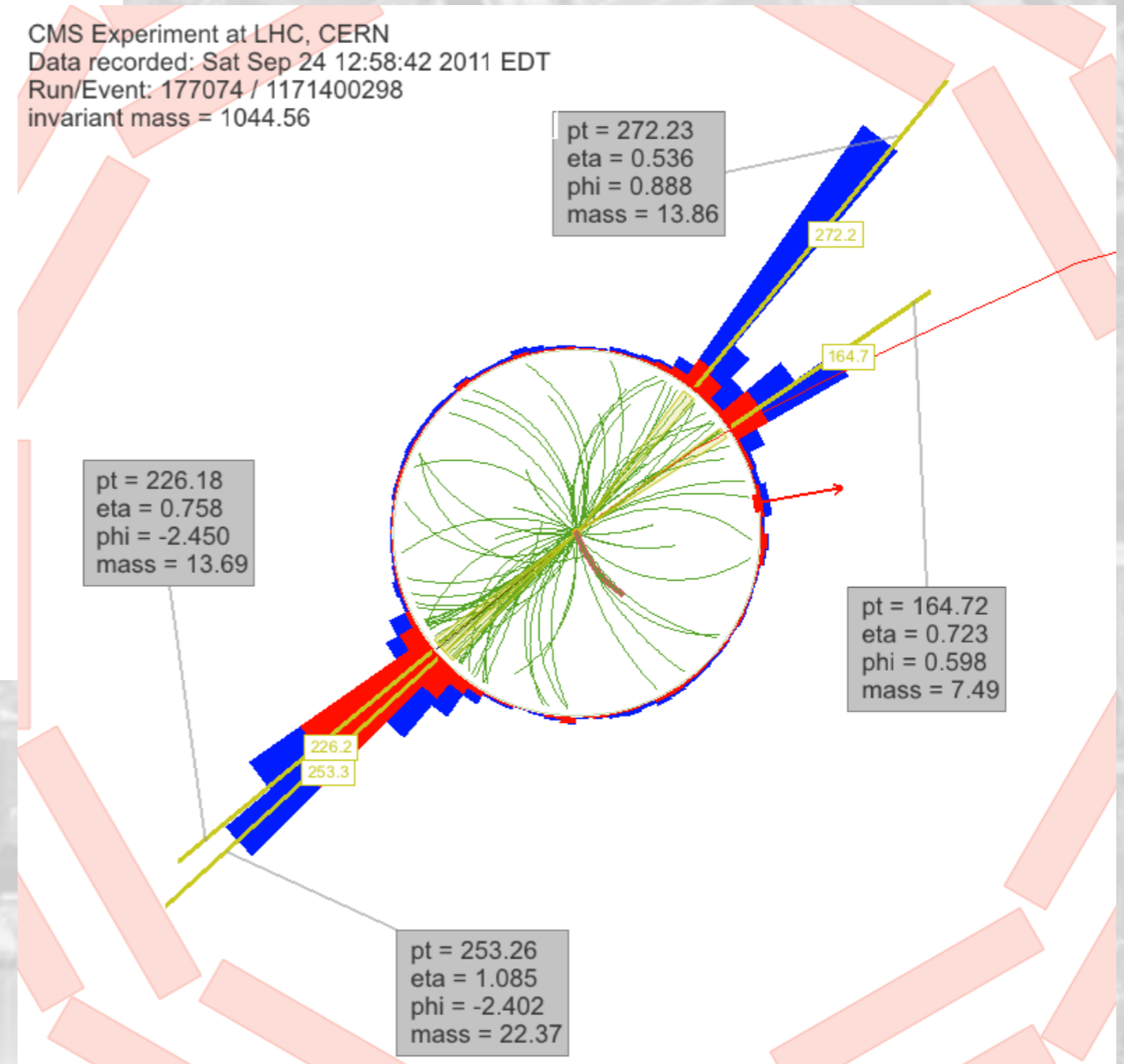
<http://arxiv.org/abs/1011.2268>

A DOUBLE-TAG EVENT

CMS Experiment at LHC, CERN
Data recorded: Sat Sep 24 12:58:42 2011 EDT
Run/Event: 177074 / 1171400298
invariant mass = 1044.56

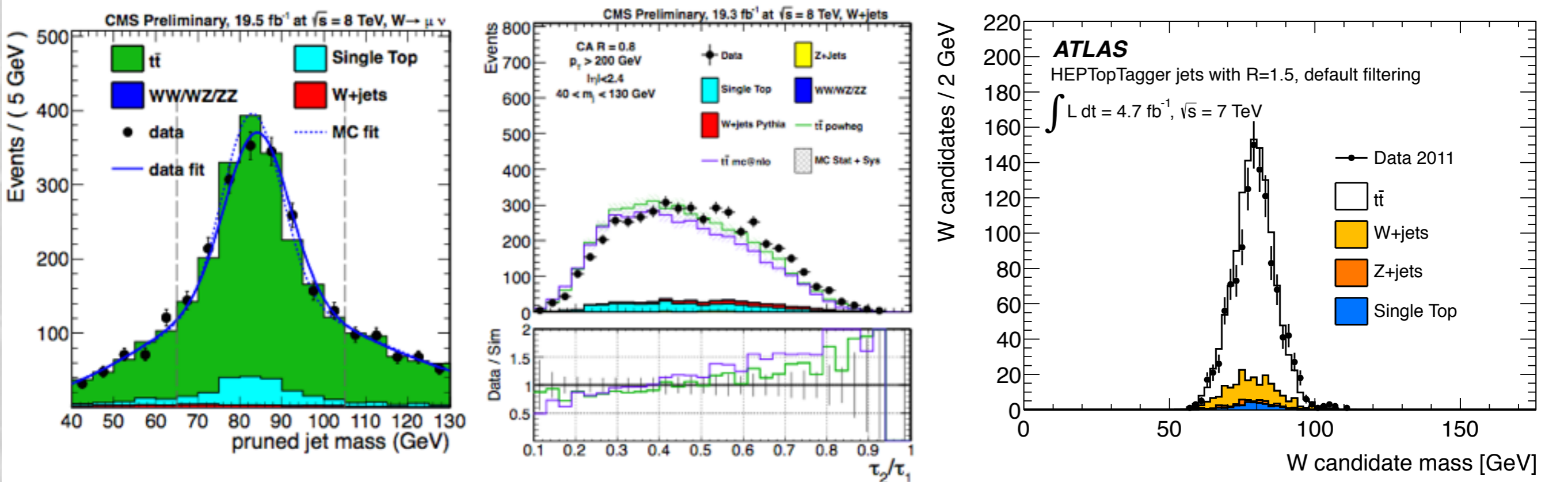


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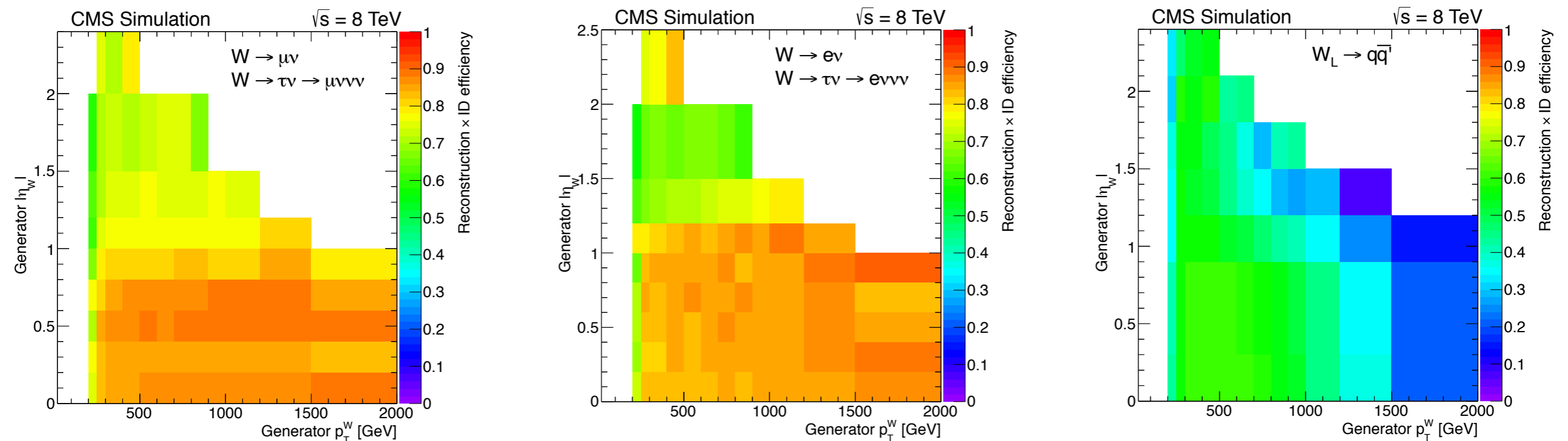
VALIDATION WITH DATA

- boosted $t\bar{t}$ reconstructed as one b-Jet+1lepton recoiling against one bjet and one jet (the W candidate)
- Peak in the jet mass: we are seeing boosted Ws
- Study substructure variables data vs MC
- MC get substructure quite right ($\sim 5\%$ systematic on predicted efficiency)

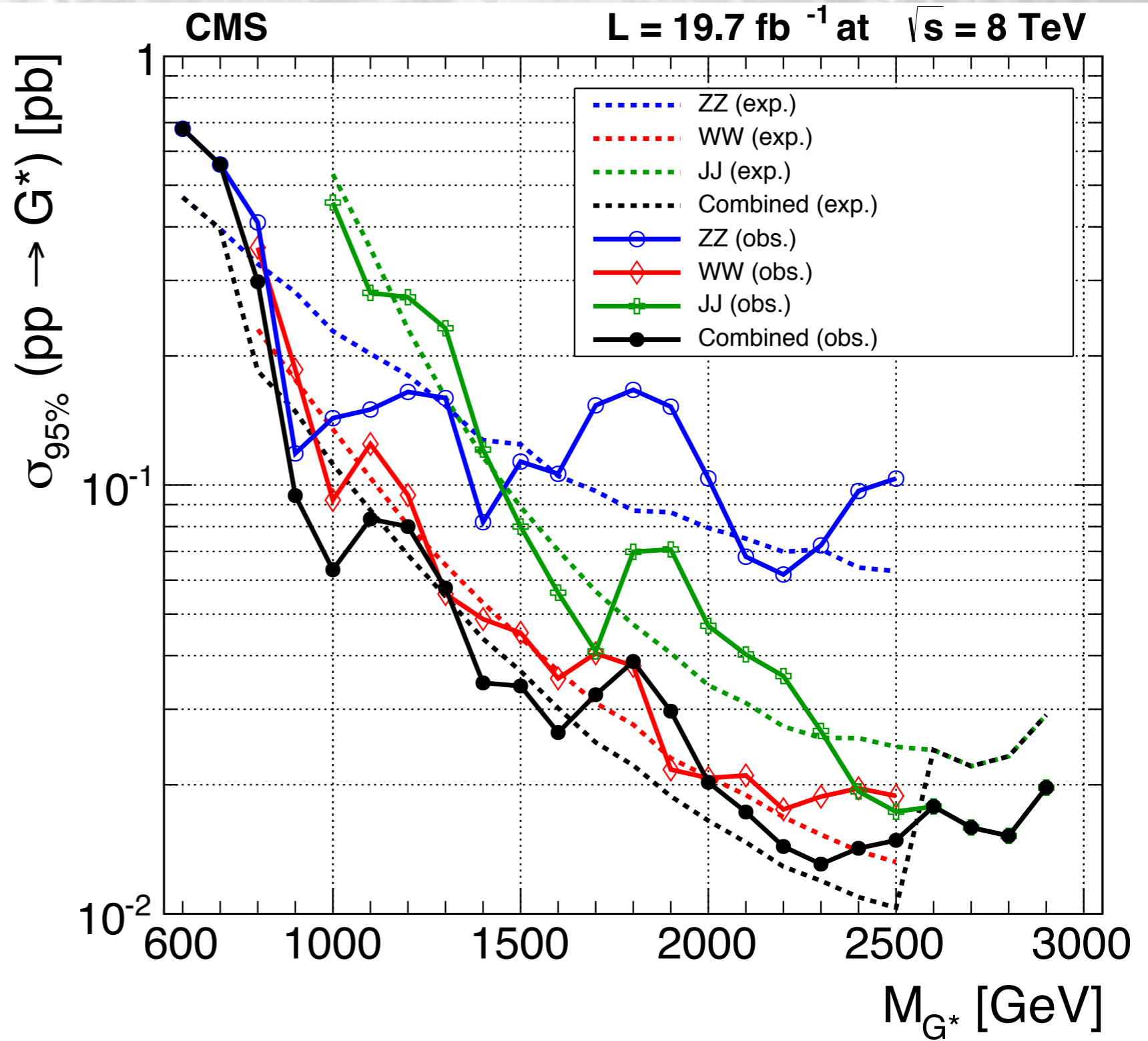
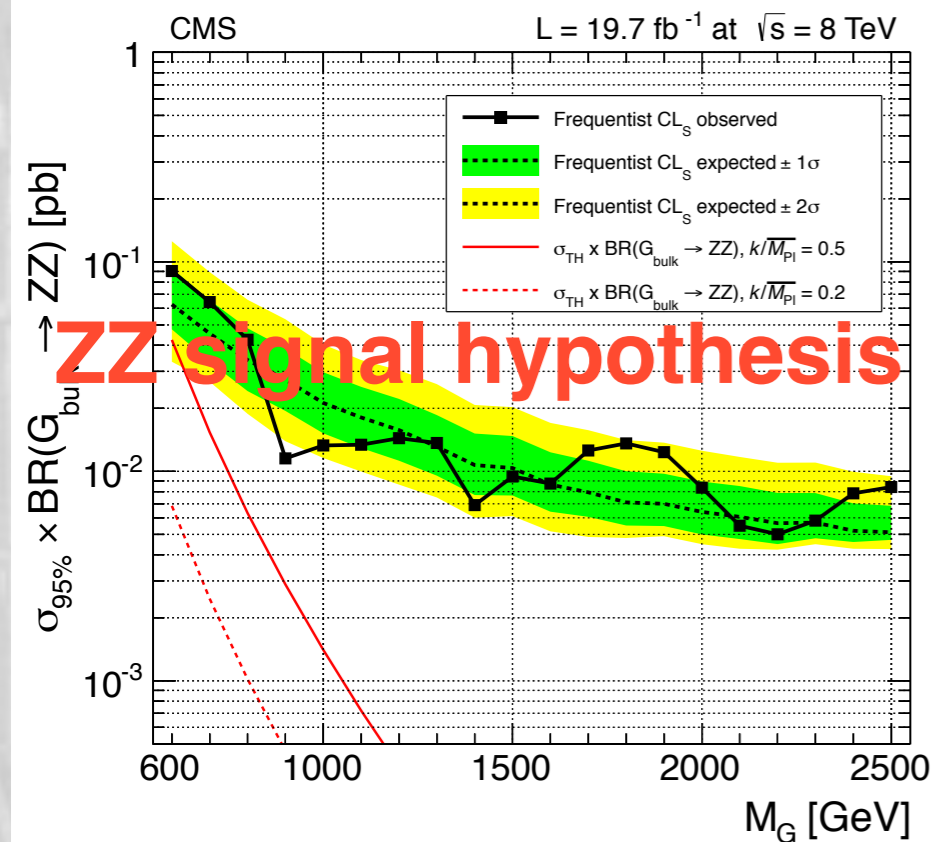
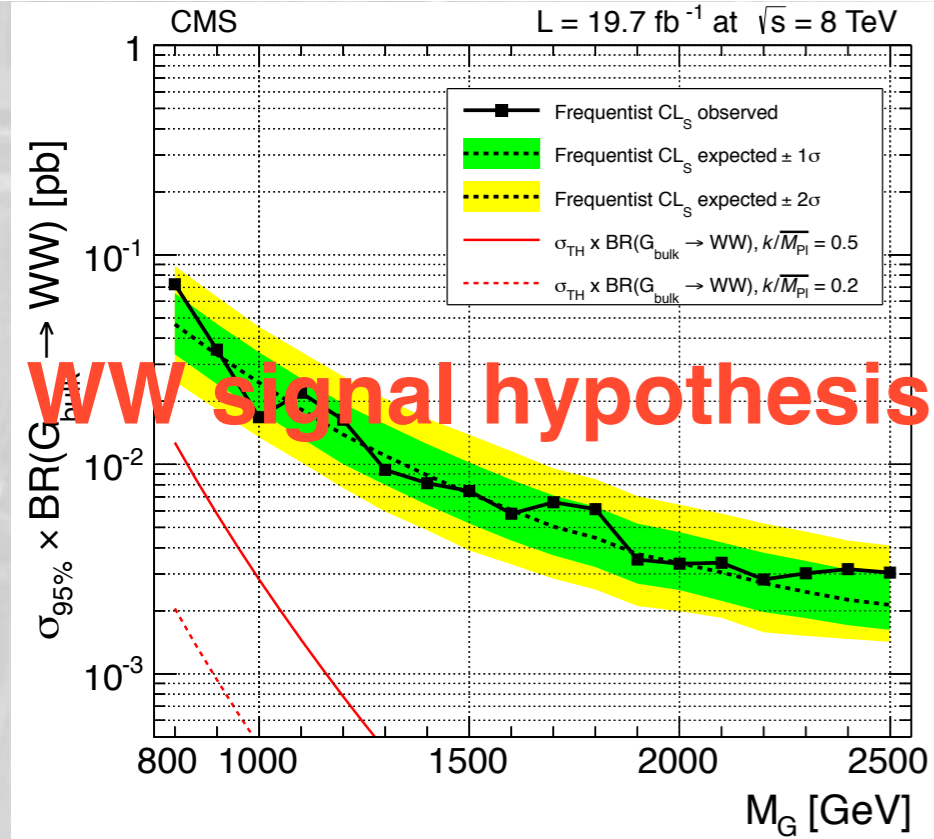


THEORY “OUTREACH”

- We spent some time to provide the information to reinterpret these results in specific models
- You have to generate resonance production w/o detector effects (run Madgraph/Pythia/...) and decay to VV
- We provide the efficiency vs the boson p_T and η for leptonic and hadronic final states (including the cross-feed background for leptons from τ decays). Efficiency quoted for V_L and reduced by 15% for V_T

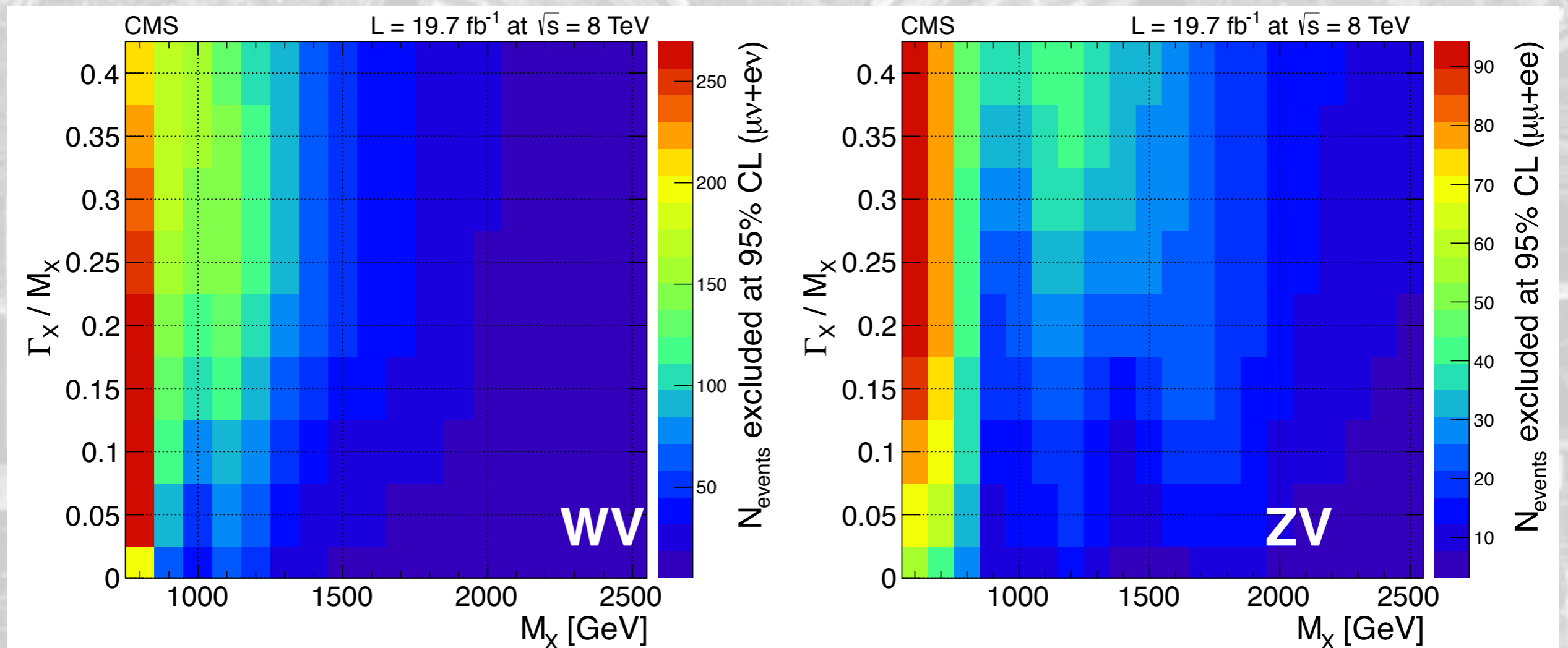


RESULTS WITH 8TeV DATA



RESULTS WITH 8TeV DATA

- Limits derived for generic masses and relative widths
- Provided (in paper Appendix) the details on how to derive these results with the provided “outreach” information (look for “Simplified limits”)

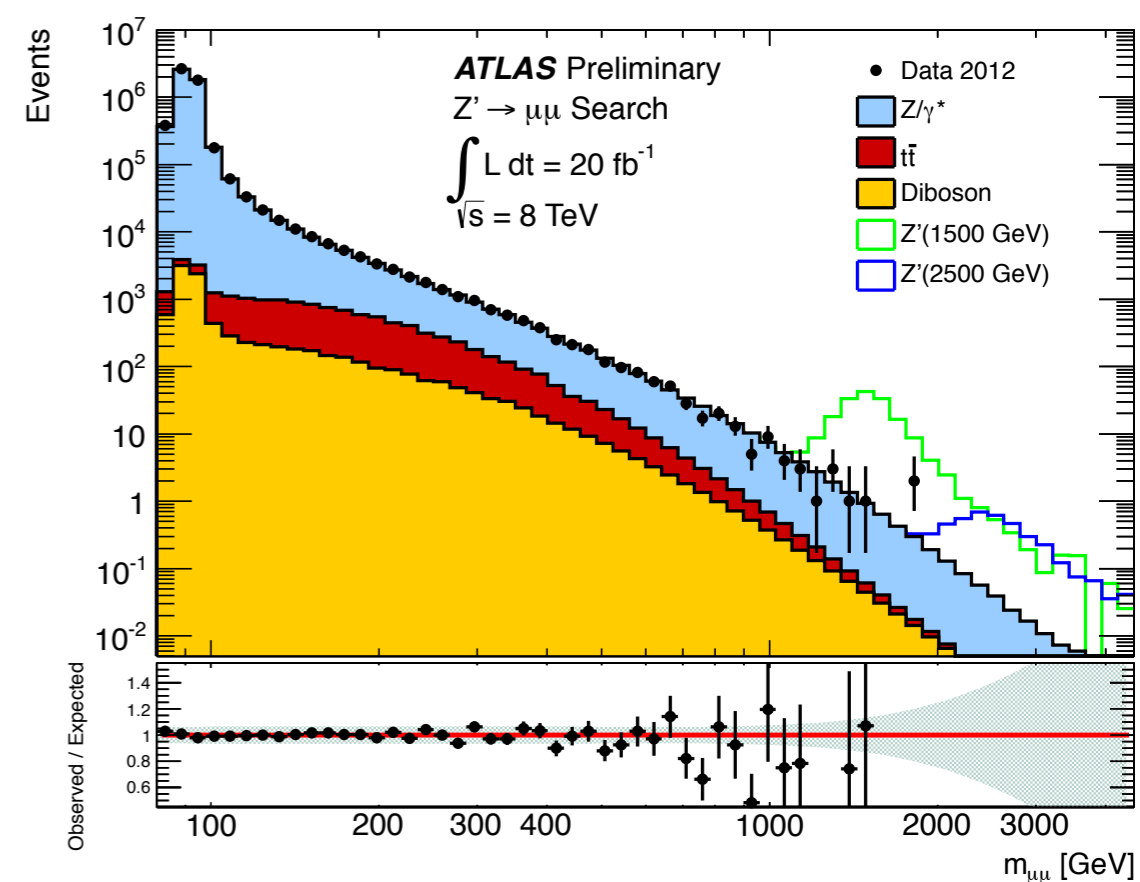
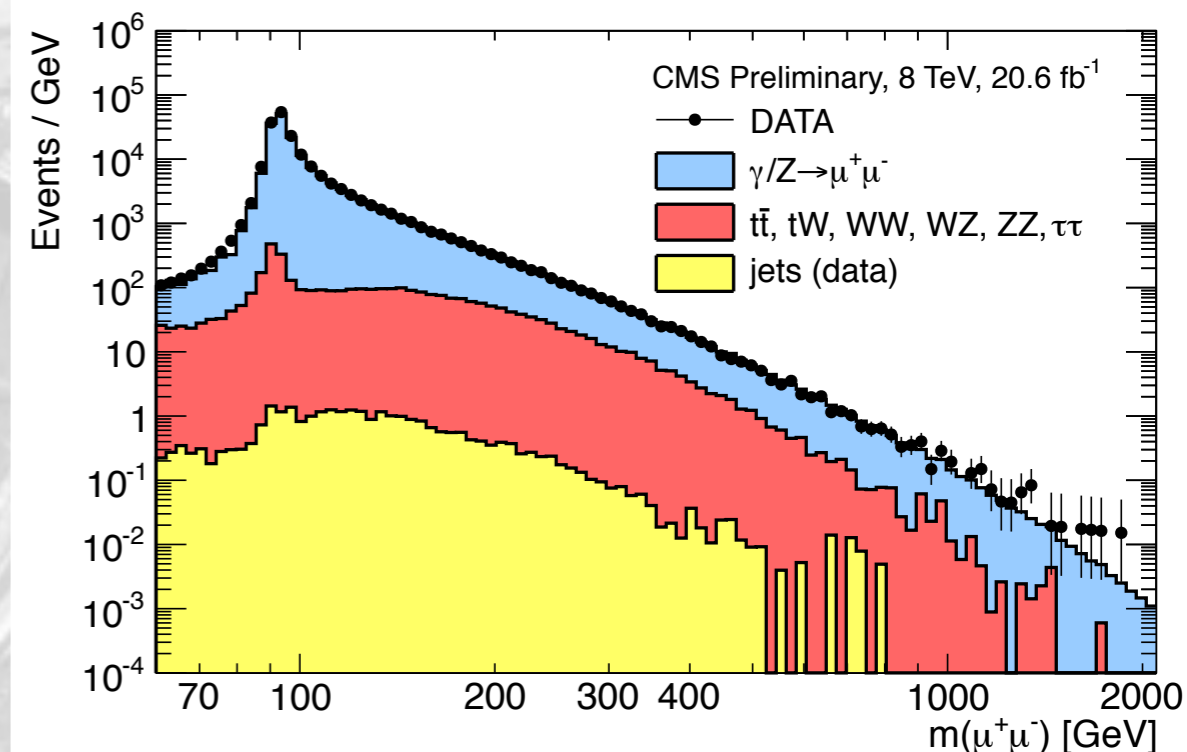




$X \rightarrow ll/jj$ Searches

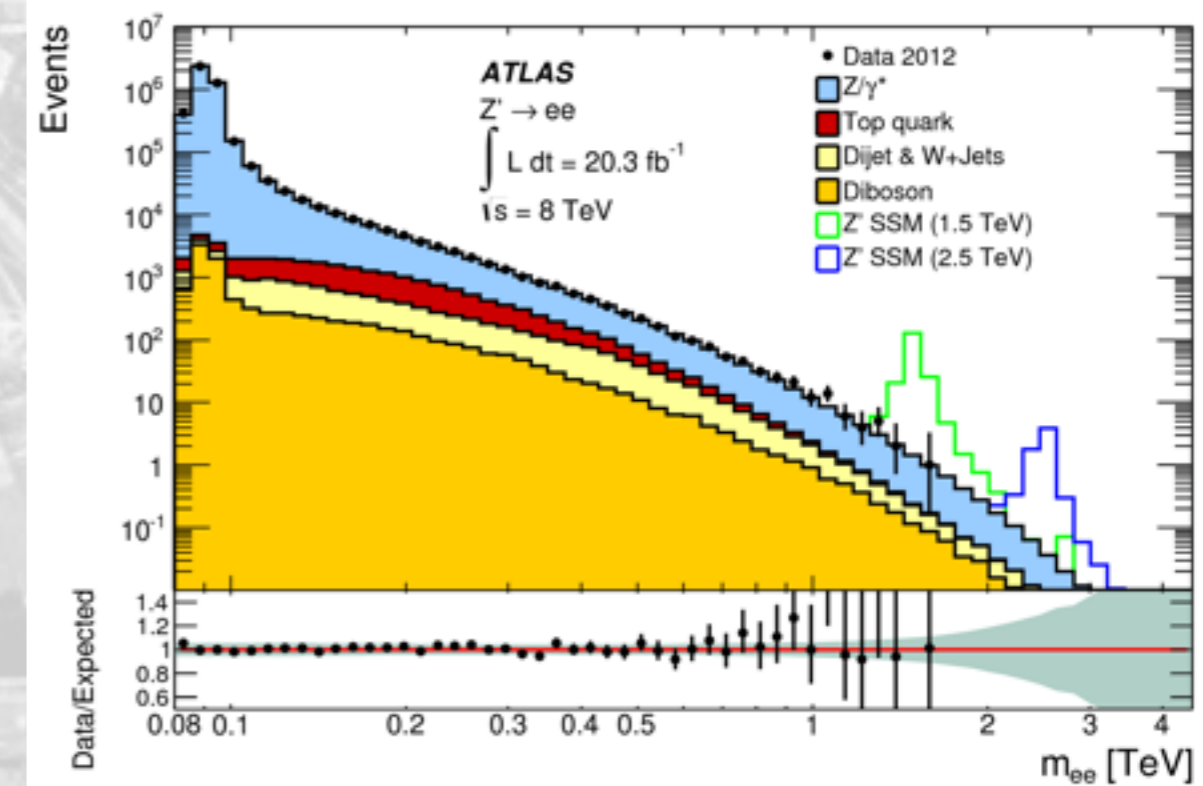
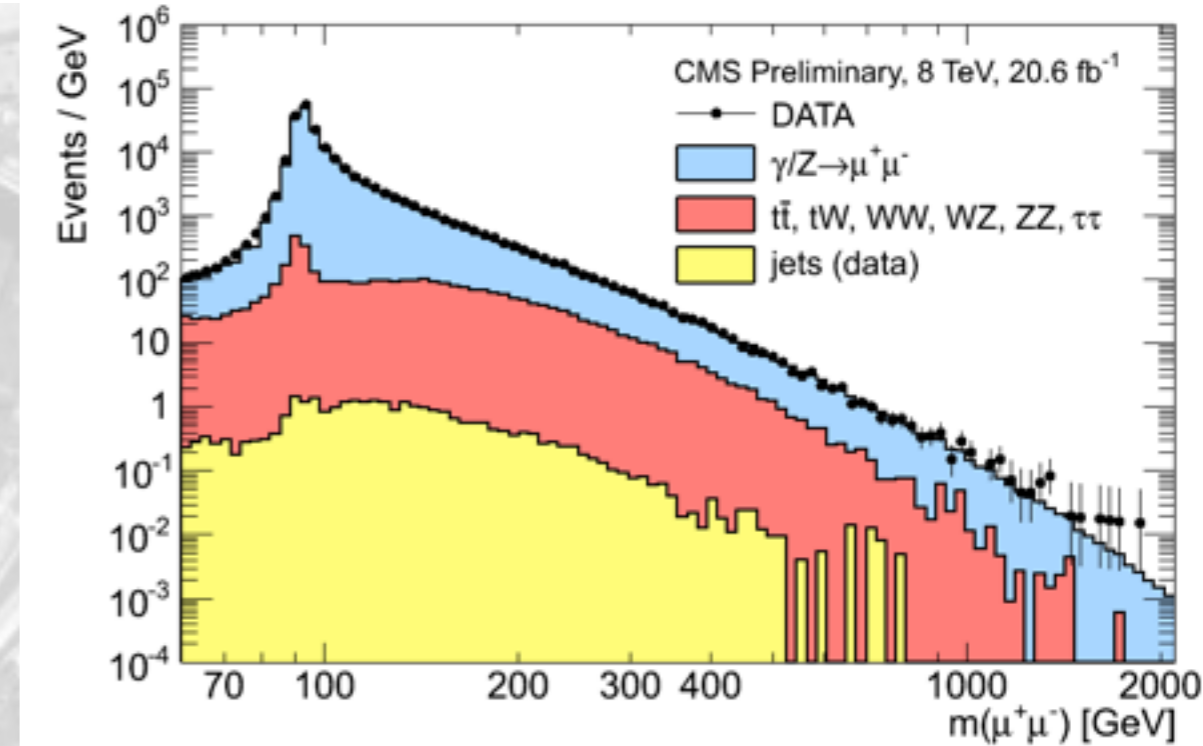
HIGH-PT MUONS

- Muon momenta are measured through the bending in the magnetic fields
- The bending is reduced at large muon momenta
- For high-pT muons, the precision deteriorates
- Unlike the case of measurements with W/Z/top/H, muon final states are not the golden channel for this physics
- Despite the resolution, high-pT muons are an excellent discovery tool

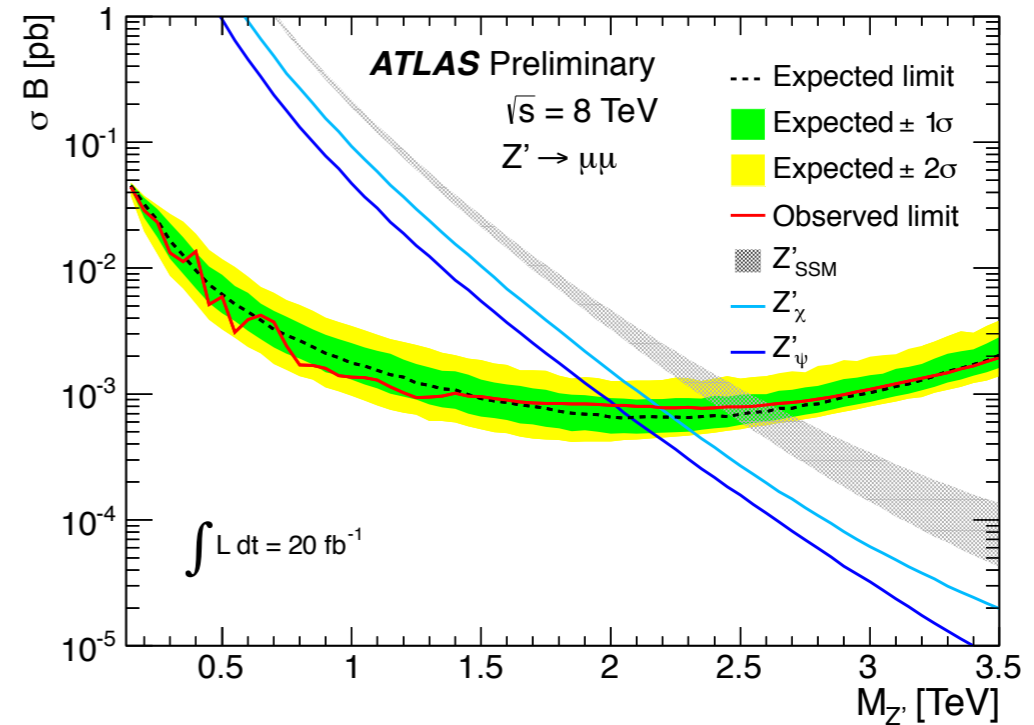
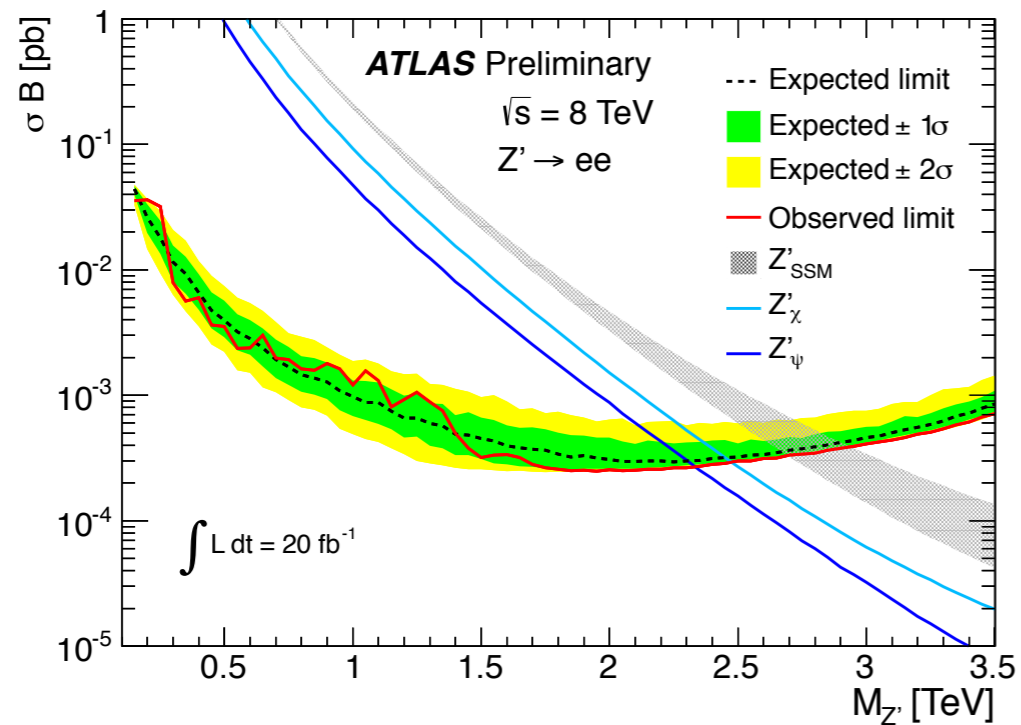
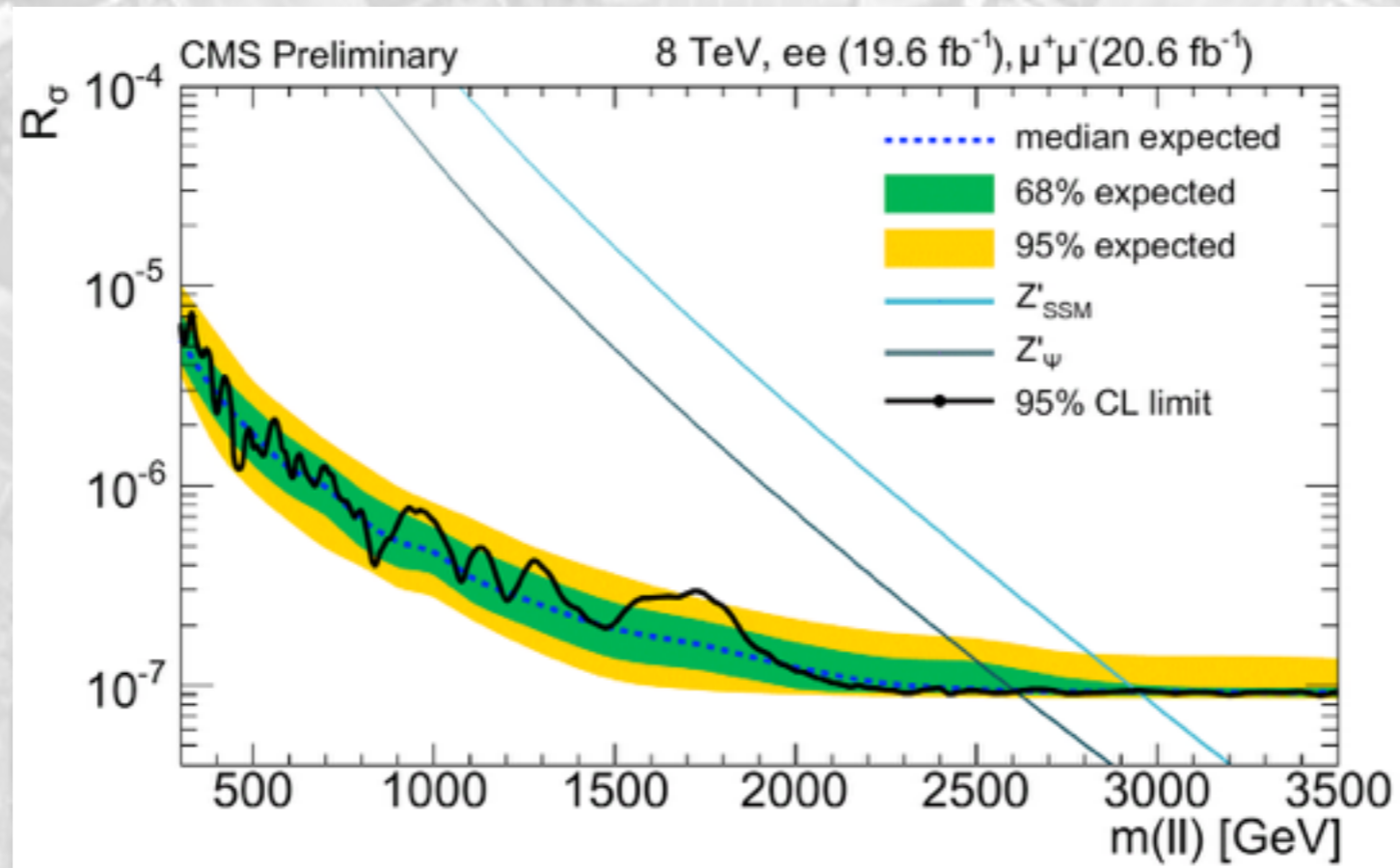


HIGH-PT ELECTRONS

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- The resolution of the calorimeter improves with energy, giving a better S vs B discrimination above 1 TeV
- Electrons (and photons) are excellent tools to search for dark resonances and measure their masses

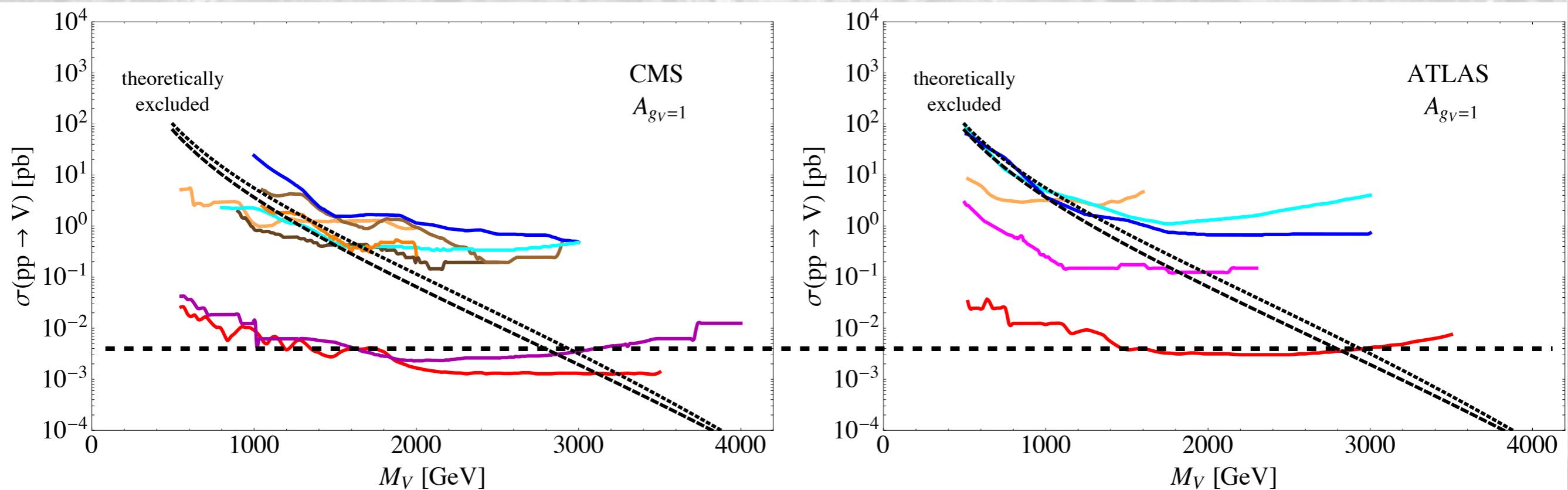


DILEPTON SEARCH



DILEPTON SEARCH

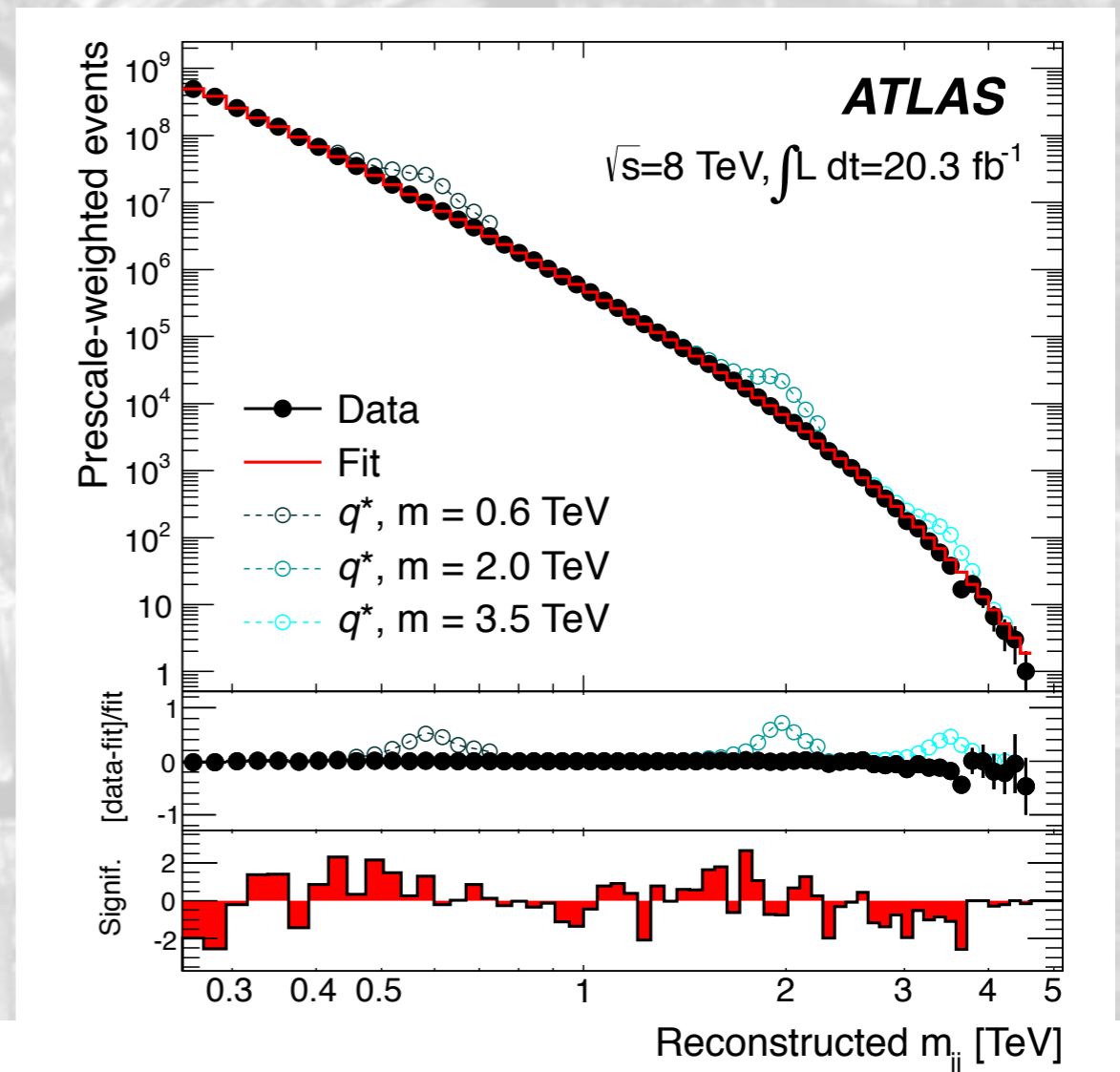
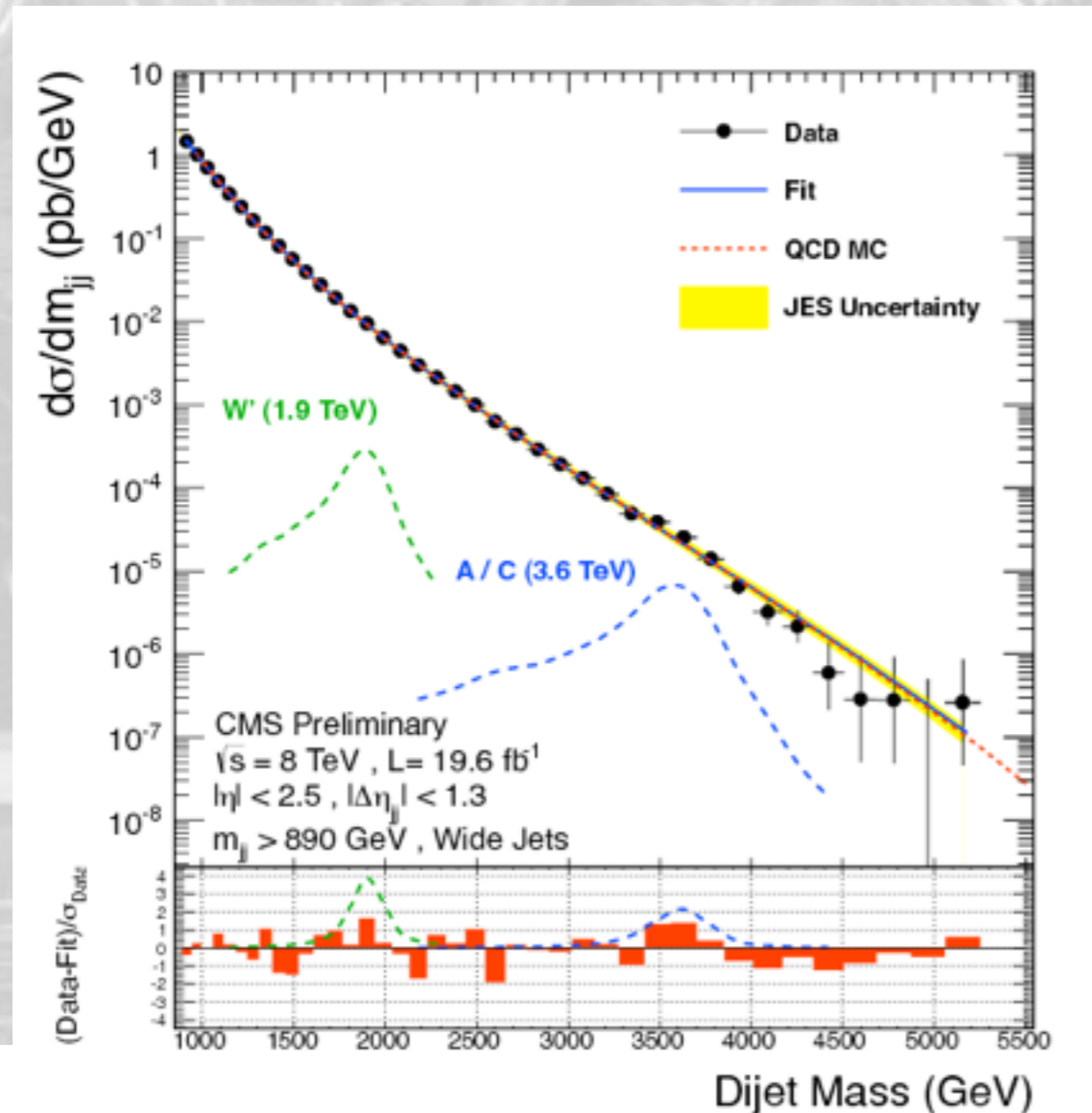
- Despite what it might look like, no contradiction here
 - Same limit at ~ 1800 GeV
 - CMS expected a better limit, but sees a $\sim 2\sigma$ excess
 - ATLAS did not expect to exclude below the observed limit



<http://arxiv.org/abs/arXiv:1402.4431>

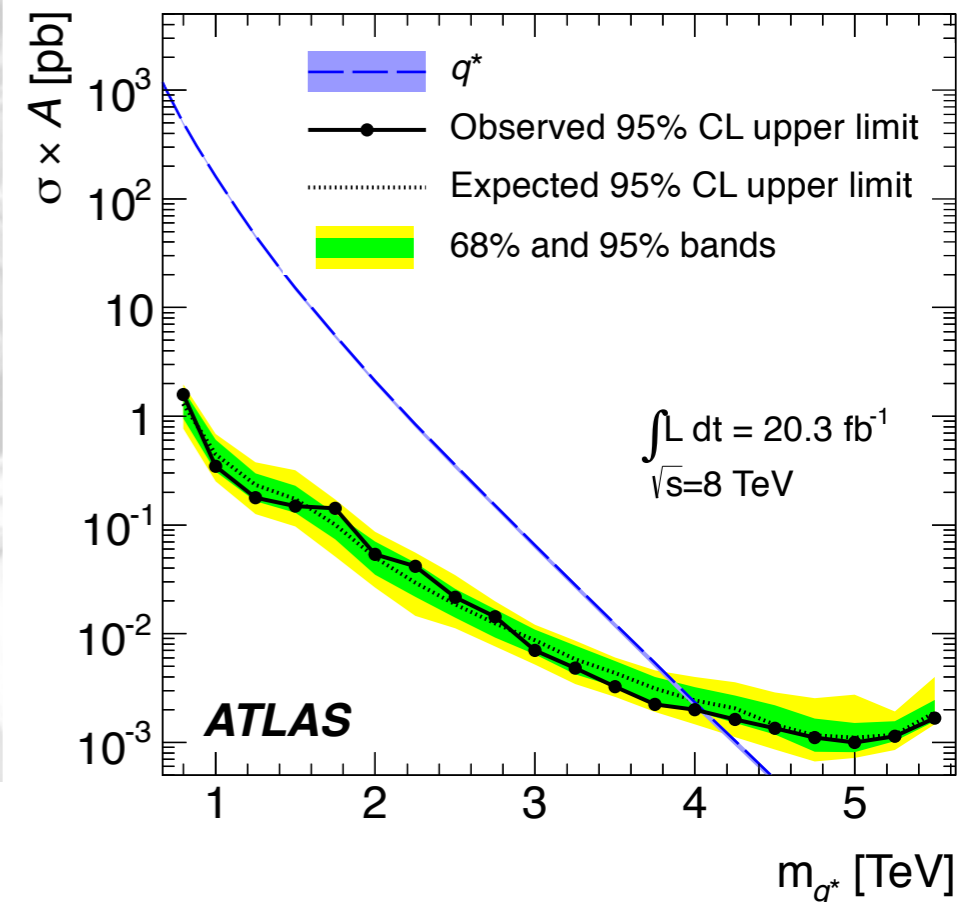
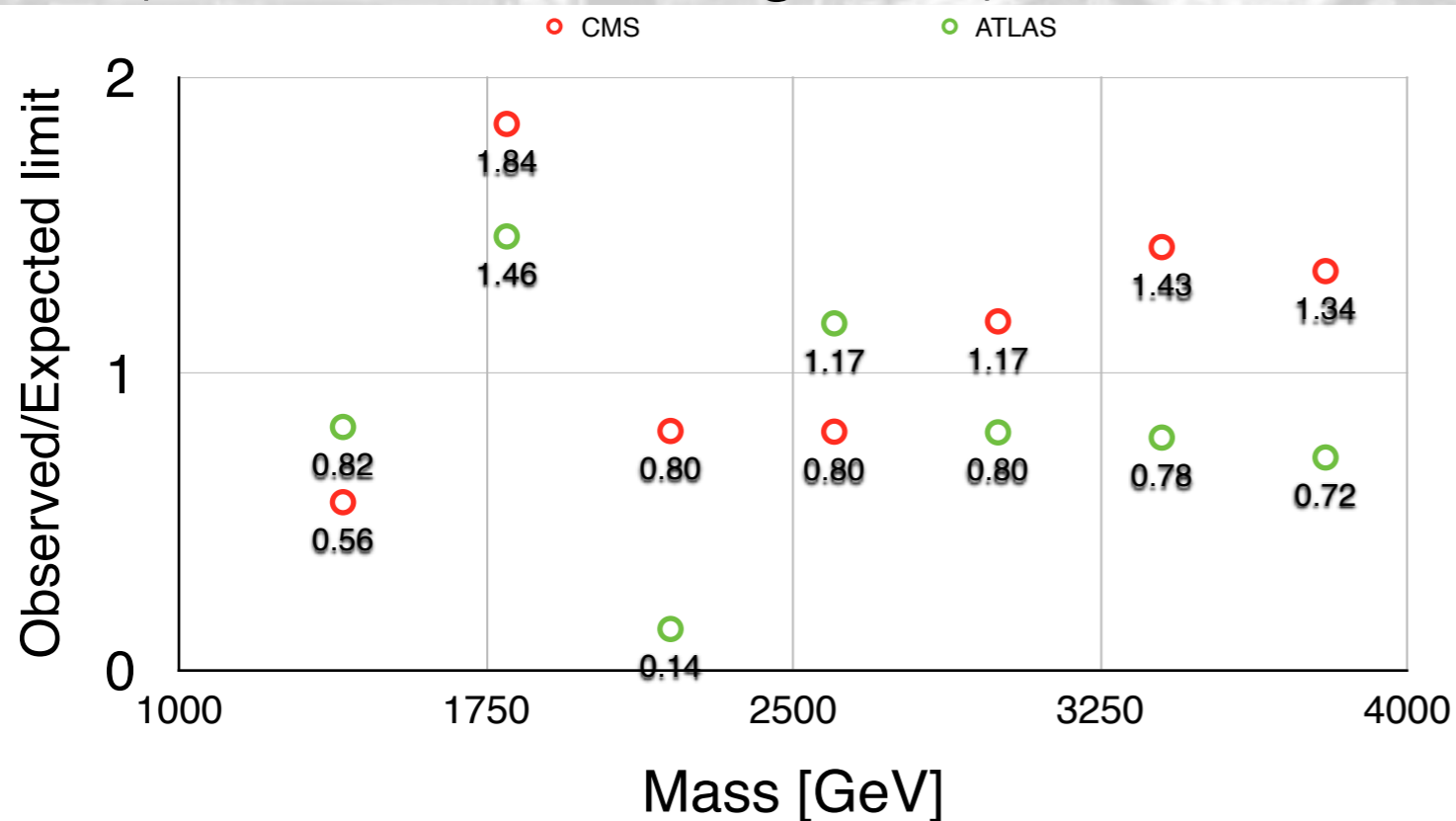
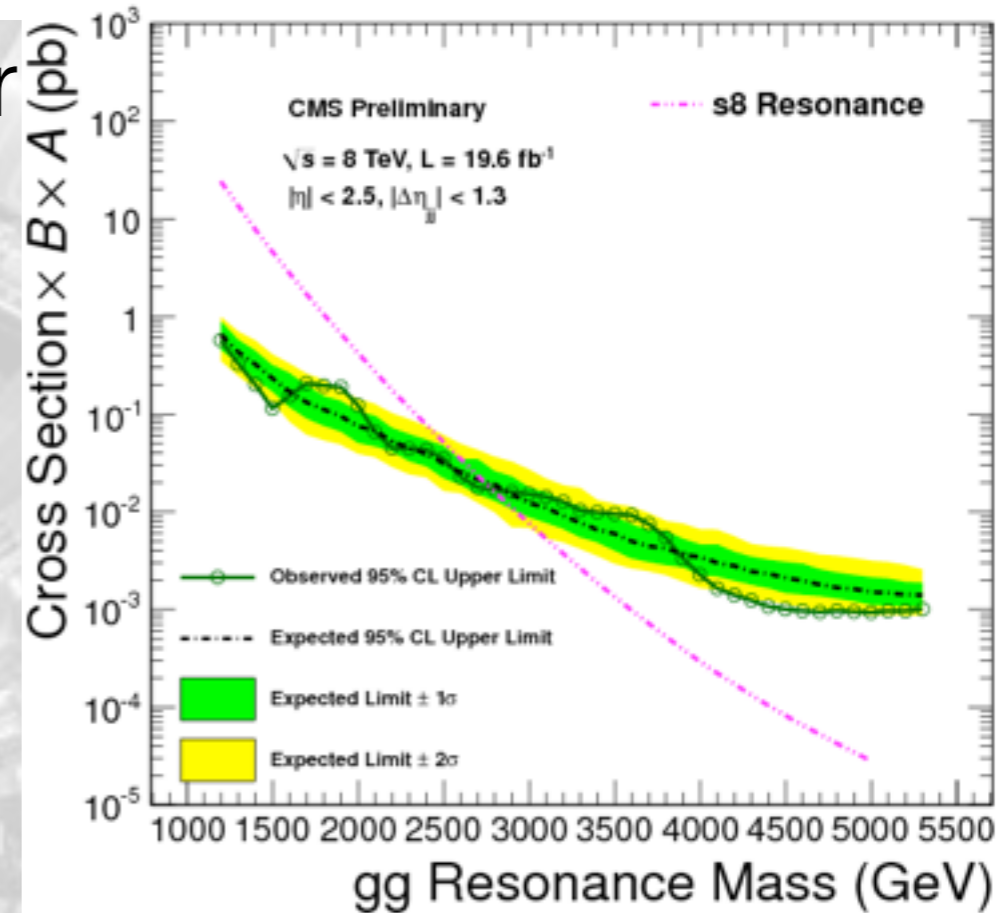
DIJET SEARCH

- The fastest search at a hadron collider
 - Parameterize the QCD dijet mass spectrum
 - Look for a bump on the falling spectrum
 - Main challenge from trigger
 - High-pT jets in the final state



DIJET SEARCH

- The fastest search at a hadron collider
 - No excess observe
 - Result interpreted as 95% exclusion on narrow resonances
 - Result depends on the final state (qq/qq/gg)
 - Also interpretation for broad resonance (available or coming soon)



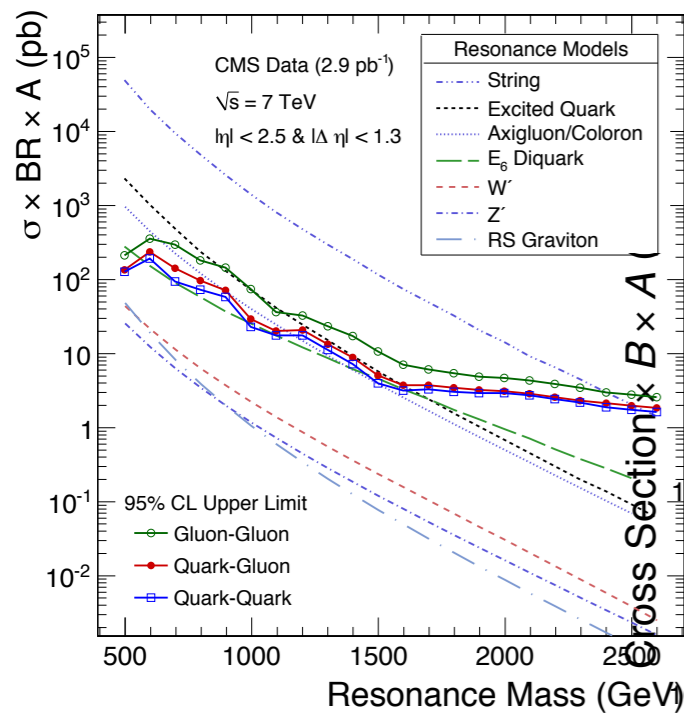


Not Only Heavy Resonances

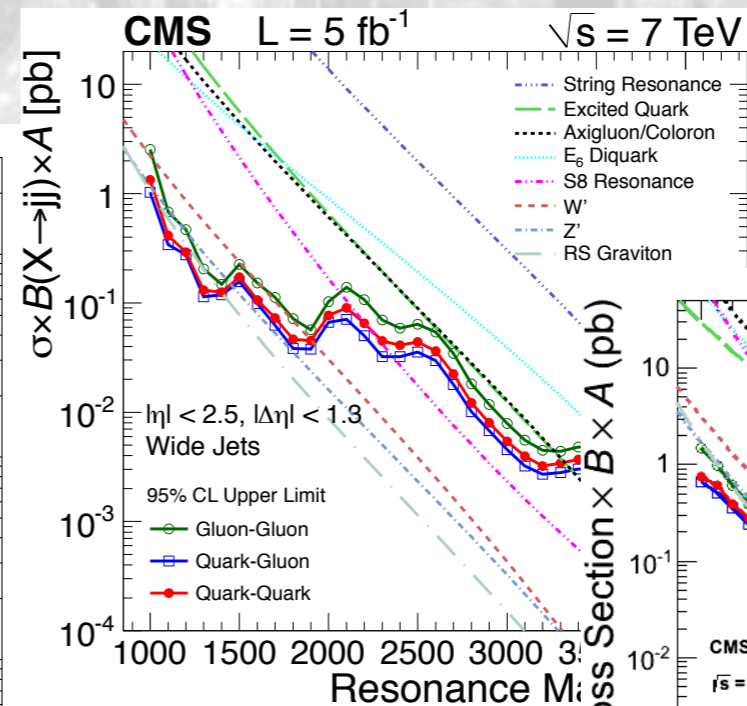
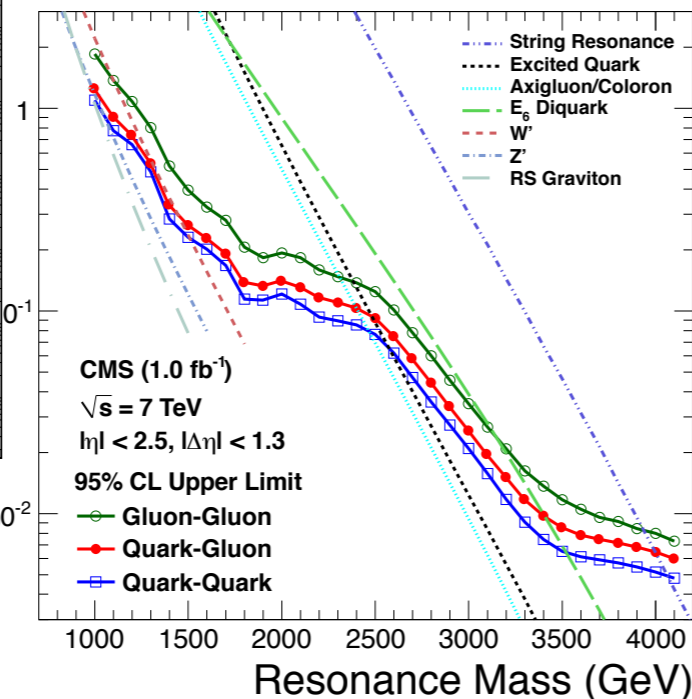
BLIND SPOTS

- Hadronic searches are limited by trigger & disk space
 - we cannot keep all the events with jets (as almost any event has jets)
 - Loose triggers are used at low-peak luminosity
 - When luminosity increases, the triggers become tighter
 - Combinatoric makes the peak reconstruction more complicated

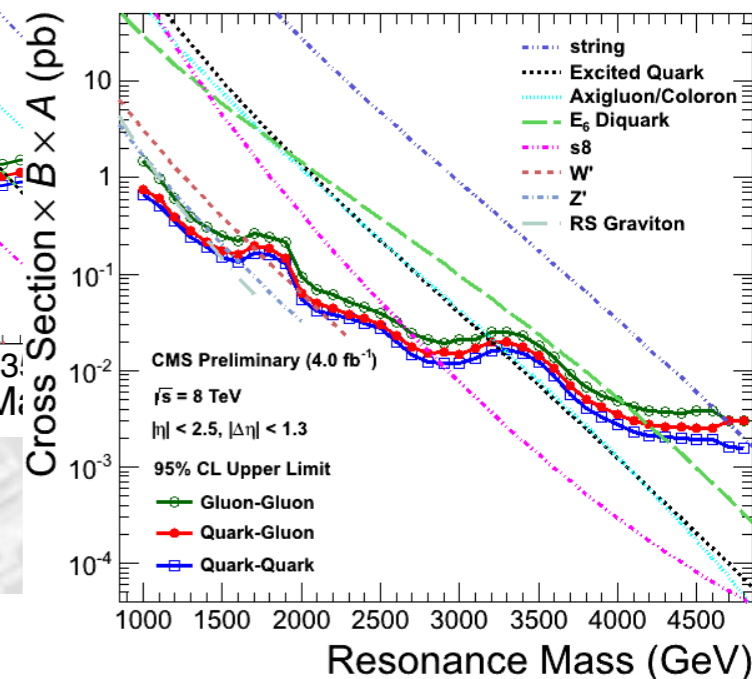
3 pb^{-1} 7 TeV



1 fb^{-1} 7 TeV



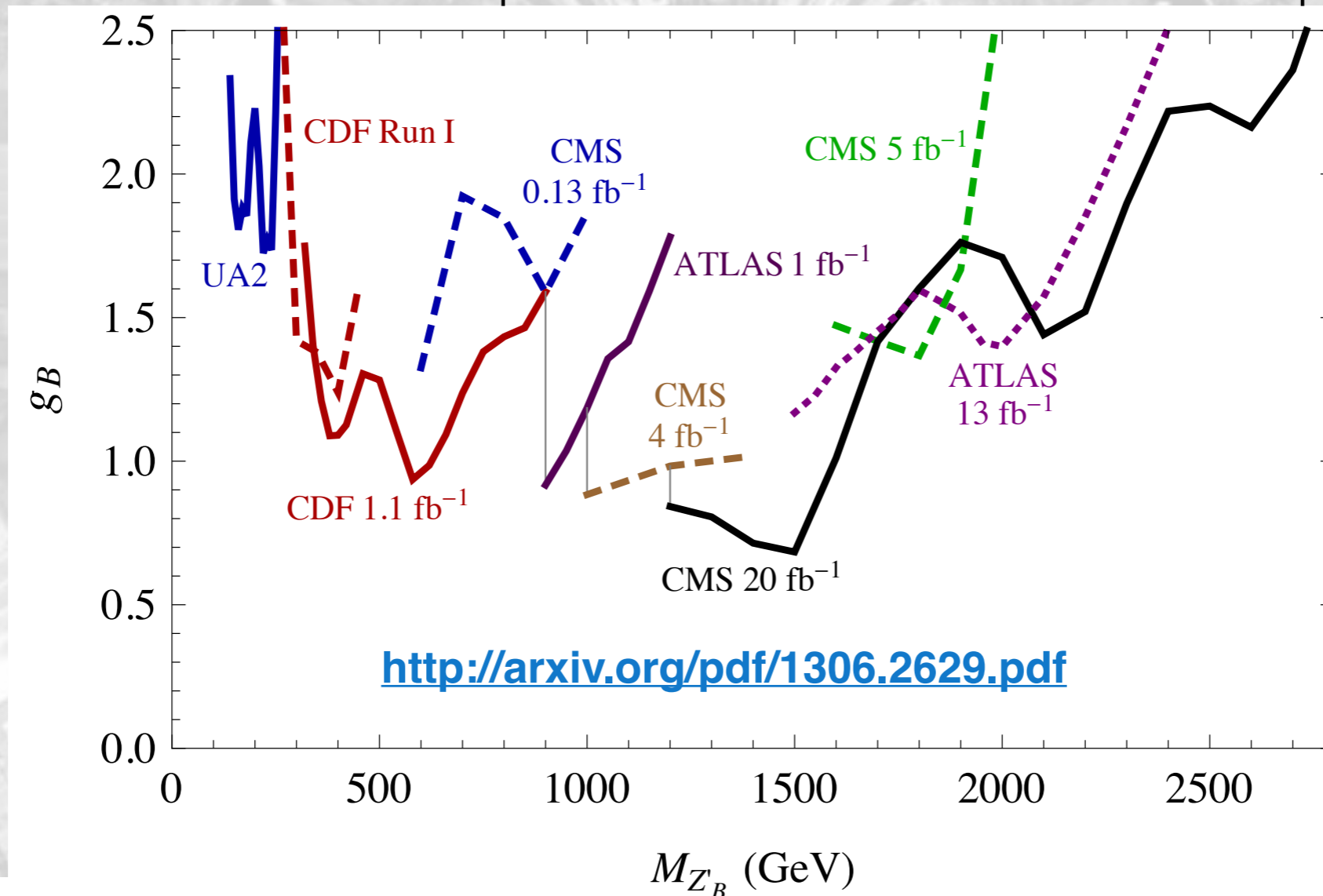
3 fb^{-1} 8 TeV



5 fb^{-1} 7 TeV

BLIND SPOTS

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BLIND SPOTS

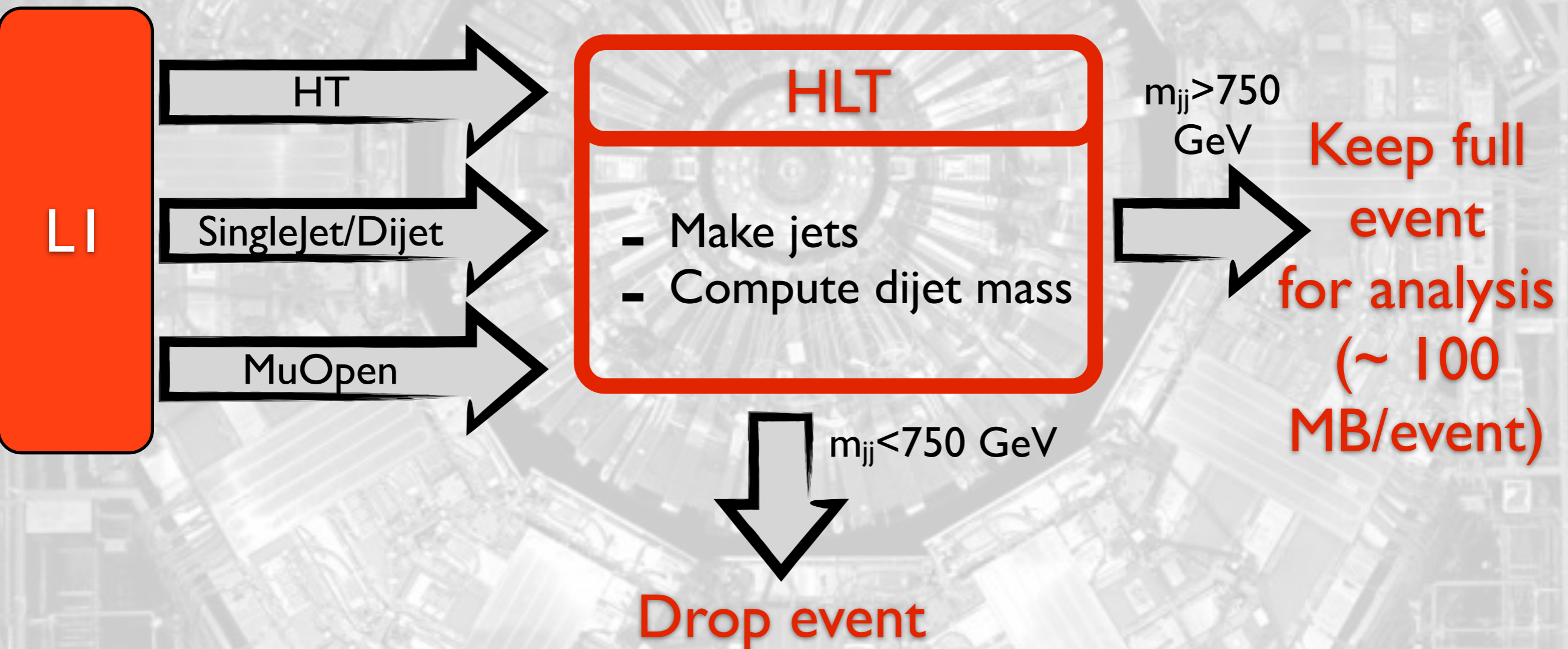
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 - Combinatoric makes the peak reconstruction more complicated

| E [| L [| MINMASS [| σ MINMASS |
|---------------|---------------|---------------------|--|
| 7 | 0.003 | 400 | ~ 200 |
| 7 | 1 | 900 | ~ 1 |
| 7 | 5 | 1000 | ~ 1 |
| 8 | 3 | 1000 | ~ 0.6 |

Unexplored territory left behind with increasing luminosity

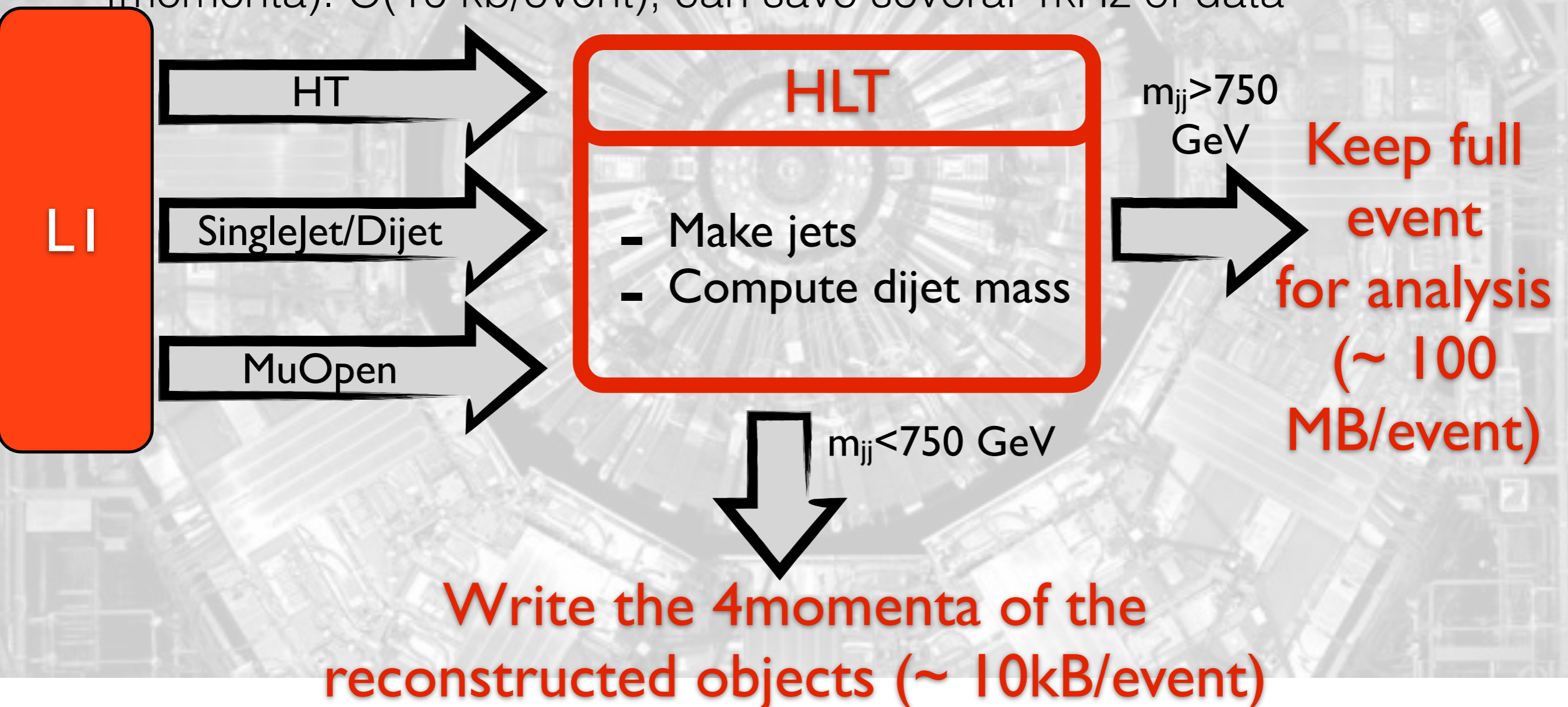
STANDARD DATA TAKING

- The moment we drop the event, we know everything about it (jet momenta, MET, maybe lepton momenta) with a coarser precision than offline (faster reconstruction → less precise reconstruction)
- When we access the event for analysis, sometimes we need a small part of this information (e.g., dijet search only needs dijet 4momenta)

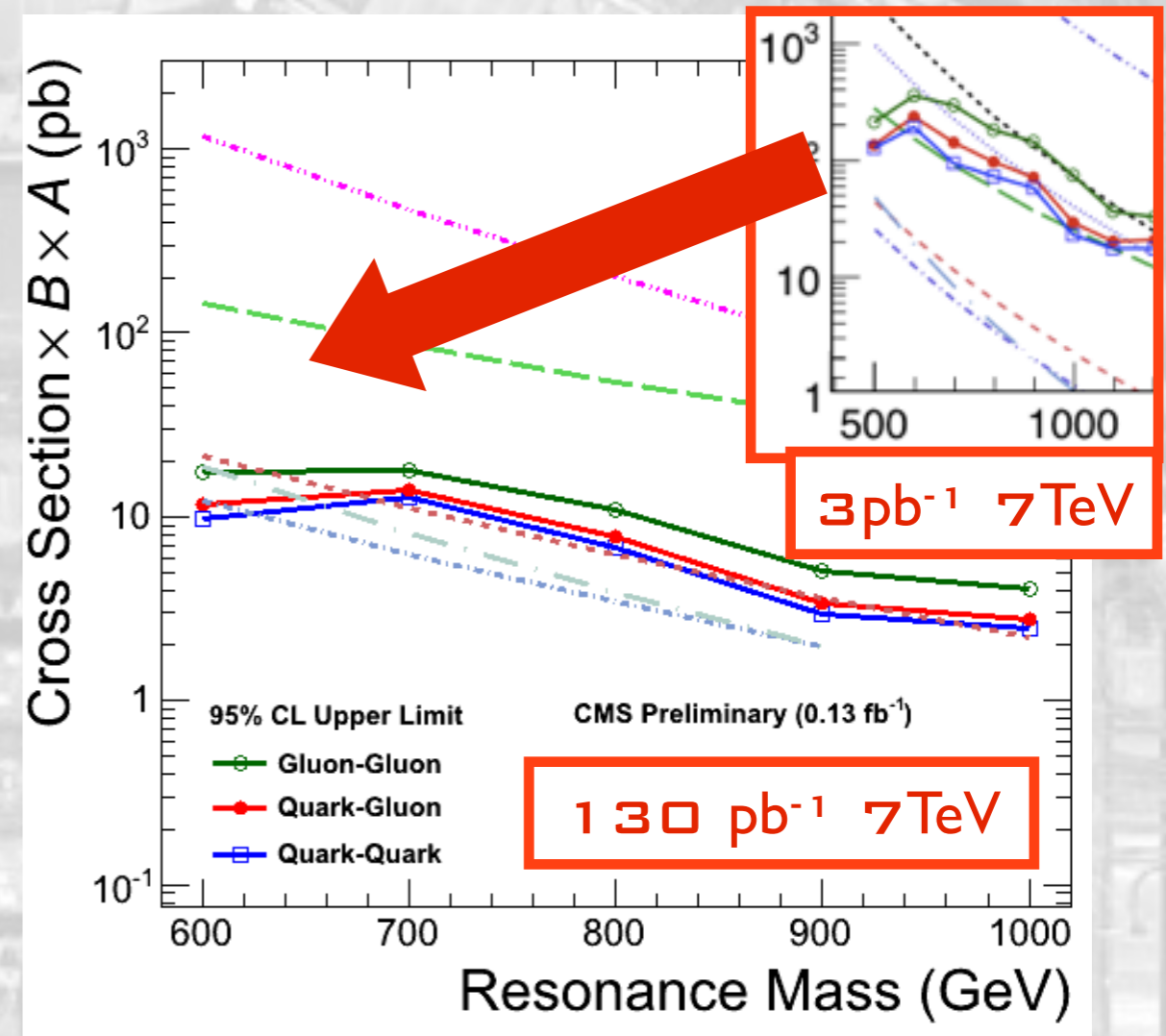
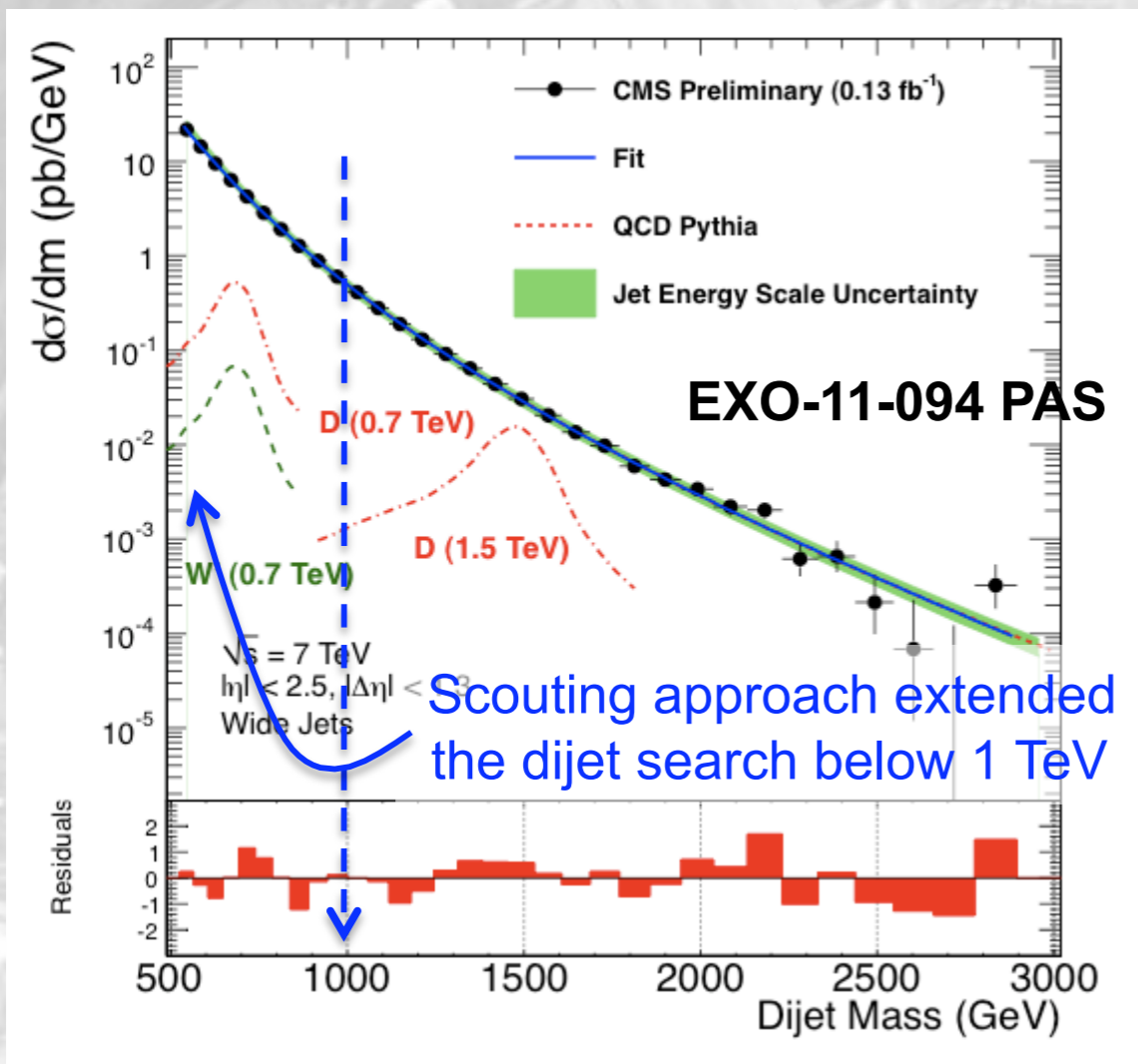


DATA SCOUTING

- We don't save events because we cannot write them on disk
- We write on disk more event information that we (sometimes) need
- Keep ALL the events, with a small customised data format (only jet 4momenta). $O(10 \text{ kb/event})$, can save several 1kHz of data



THE DIJET DATA SCOUTING



- 16 hour Run at the end of 2011 run (7TeV)
- Collected ~4 times the statistics we had in 2010 (35 pb⁻¹) with equivalent trigger
- Improved the limit published in 2010 by one order of magnitude
- 18 fb⁻¹ results@8TeV to be released soon

CONCLUSIONS

- The search for heavy resonances has been a productive front during Run I
- Ingredients are in place for Run II
 - high-pT object reconstruction
 - Jet substructure
- Despite the large pT regime, pileup was an issue (e.g. for substructure) and this will be worse. We are getting ready for that
- Interesting to look at with first few fb^{-1} (a few 2σ excesses around $m_X \sim 1800 \text{ GeV}$)
- Keeping open eyes on low-mass regime too (e.g. improving data scouting)



Backup Slides

TOP TAGGER: TYPE1 TOP

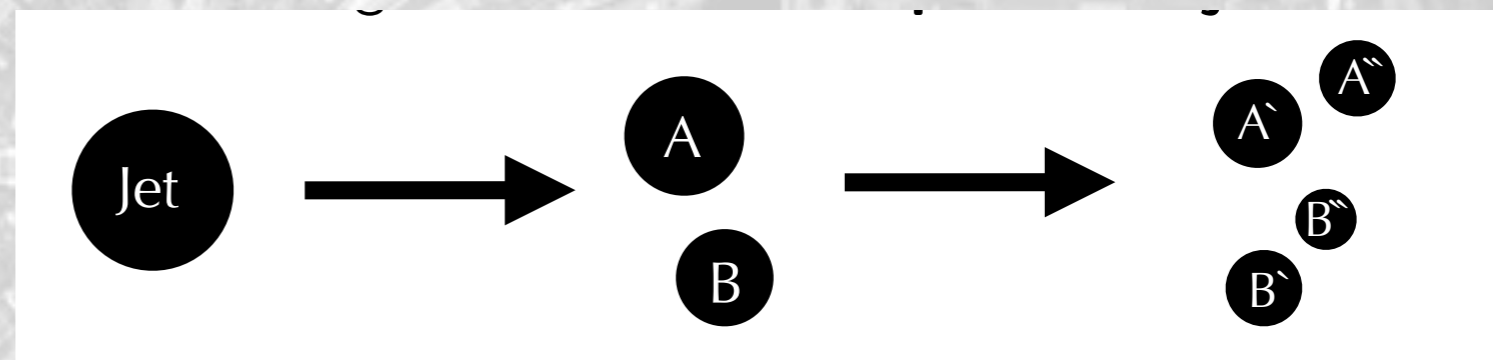
CMS-PAS-B2G-12-005



Reconstruct fully hadronic top quarks merging into one jet

- Start from CA jet with $R=0.8$ and $p_T > 400$ GeV

- Decluster in two stages to identify up to 4 jets



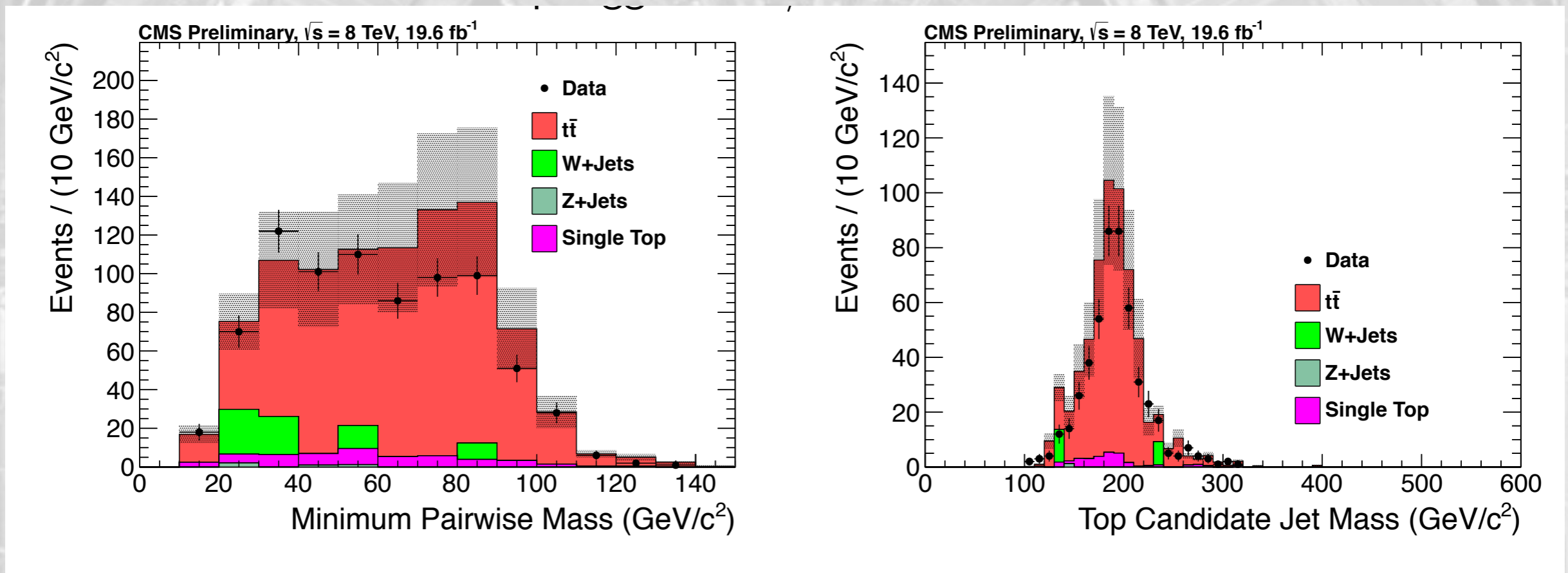
- Remove jets failing requirements and iterate

$$\left\{ \begin{array}{l} \text{subjet } p_T > 5\% \text{ jet } p_T \\ \text{Adjacency: } \Delta R(C1, C2) > 0.4 - 0.0004 p_T(C) \end{array} \right.$$

- Apply requirements on jet mass (compatible to m_{top}) and minimum subjet pair mass (compatible to W)

TOP TAGGER: TYPE1 TOP

- Use semileptonic $t\bar{t}$ events: 1lep+1b (tag) and 1 CA jet (probe)
- Compare the tagging performances data vs MC



M_{min} (W candidate in top jet)

top mass (for $M_{\text{min}} > 50 \text{ GeV}$)

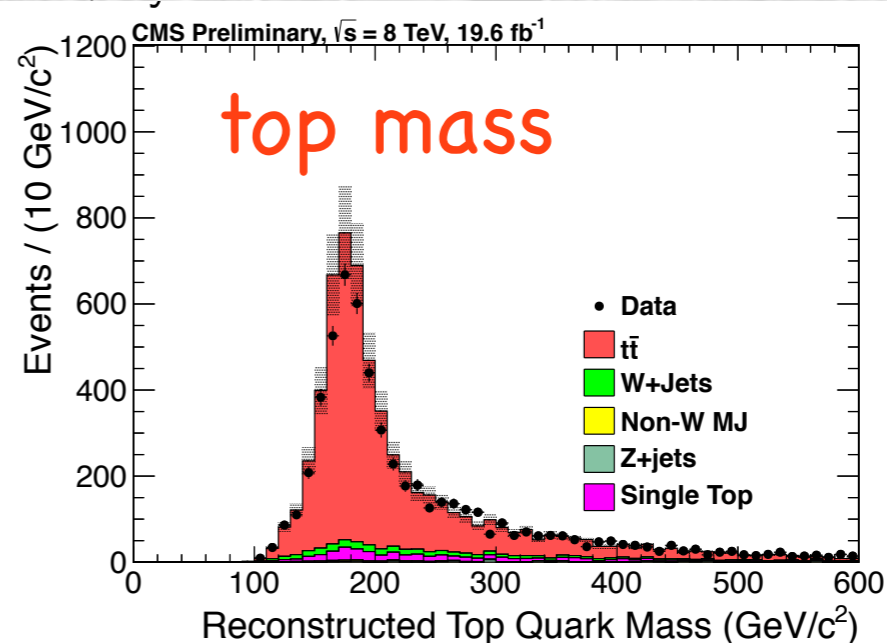
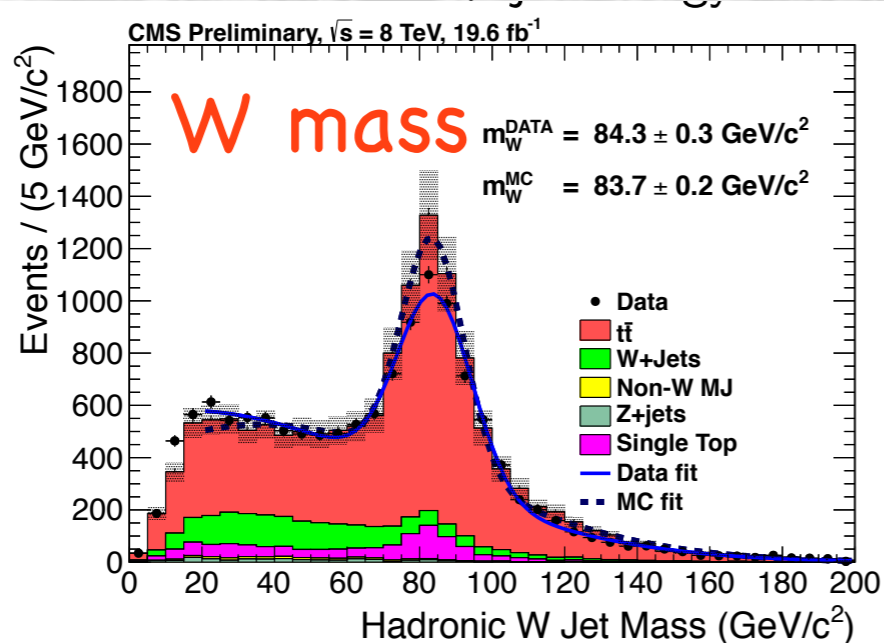
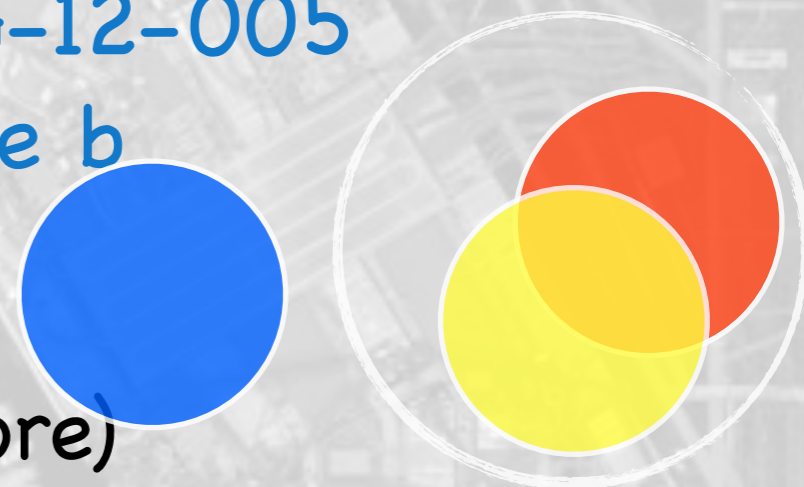
- Data / MC = 0.94 ± 0.03 (Gray area shows the normalization uncertainty)

TOP TAGGER: TYPE I TOP

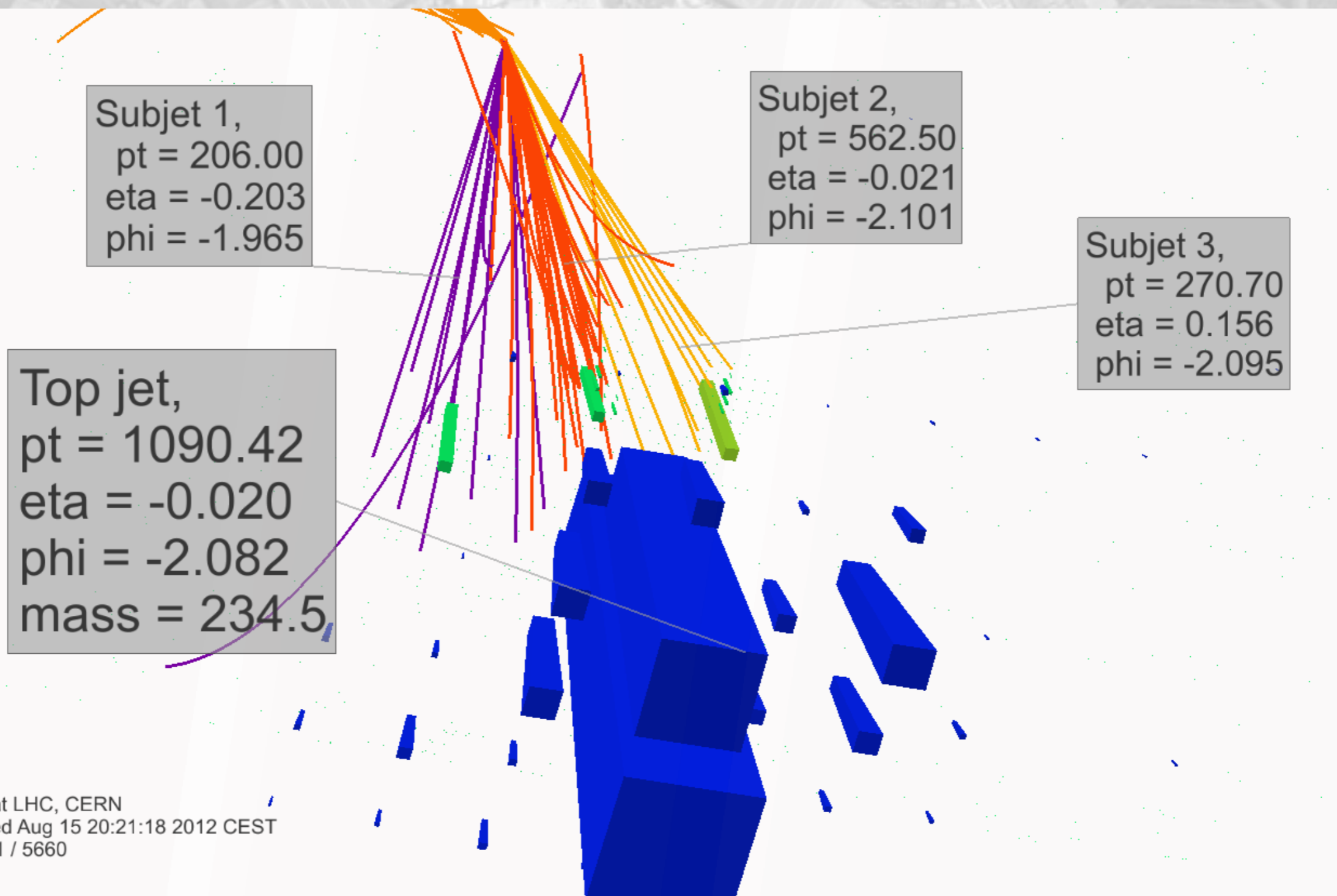
CMS-PAS-B2G-12-005

Reconstruct fully hadronic top quarks as one b jet and one merged-W jet

- Using V tagger for the W (as described before)
- Reconstruct the top adding W 4-momentum to a close jet (no b-tag applied)
- Validate the technique with semileptonic $t\bar{t}$

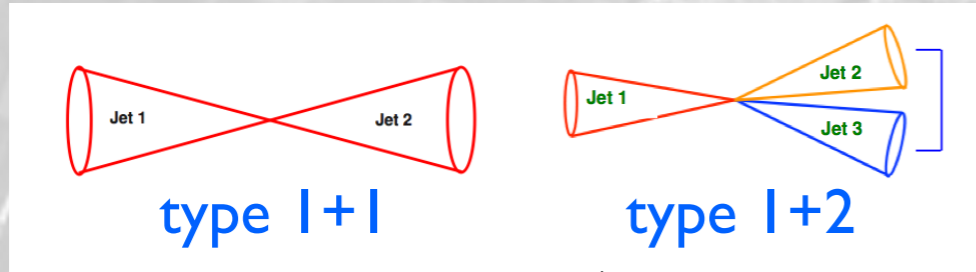


WHAT WE TOP-TAG



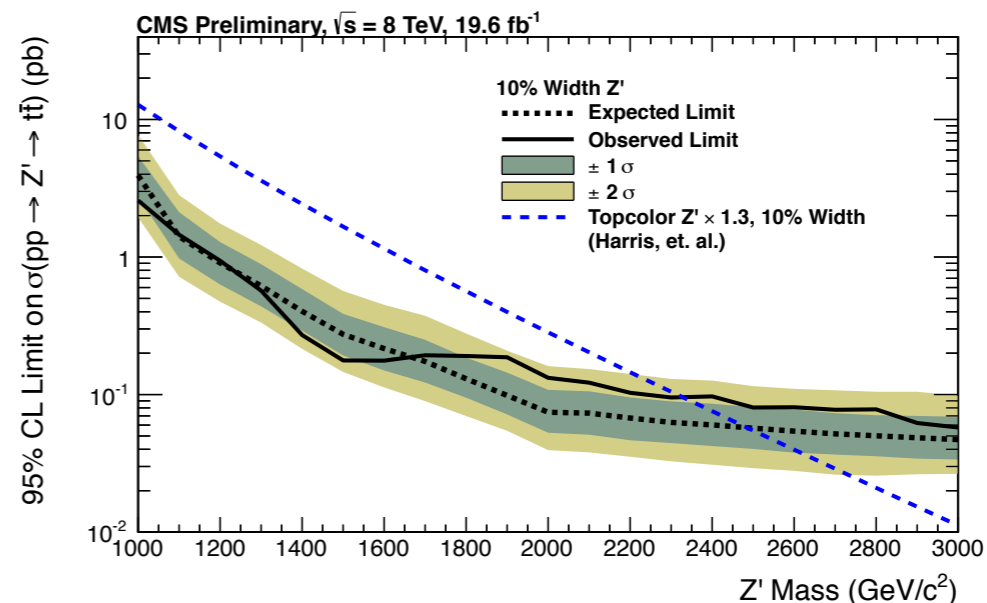
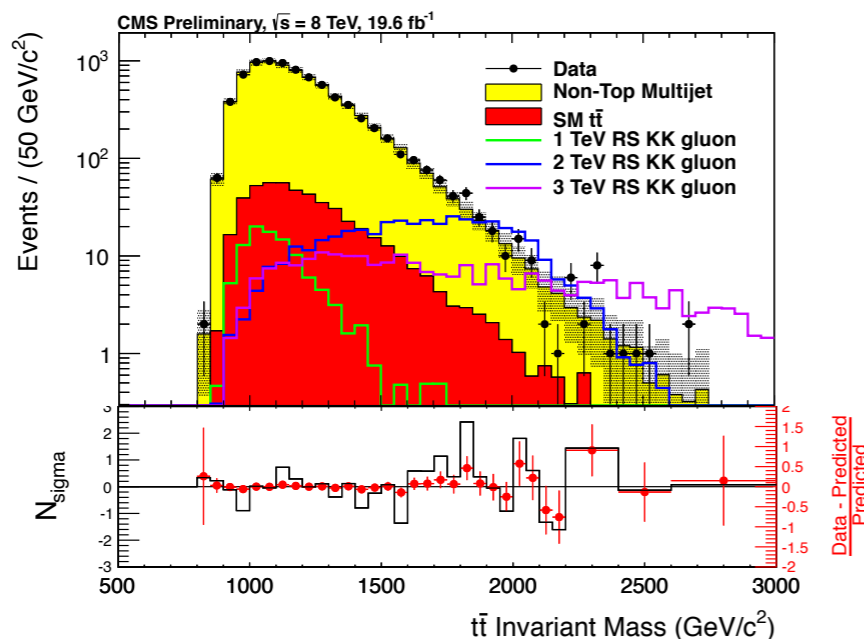
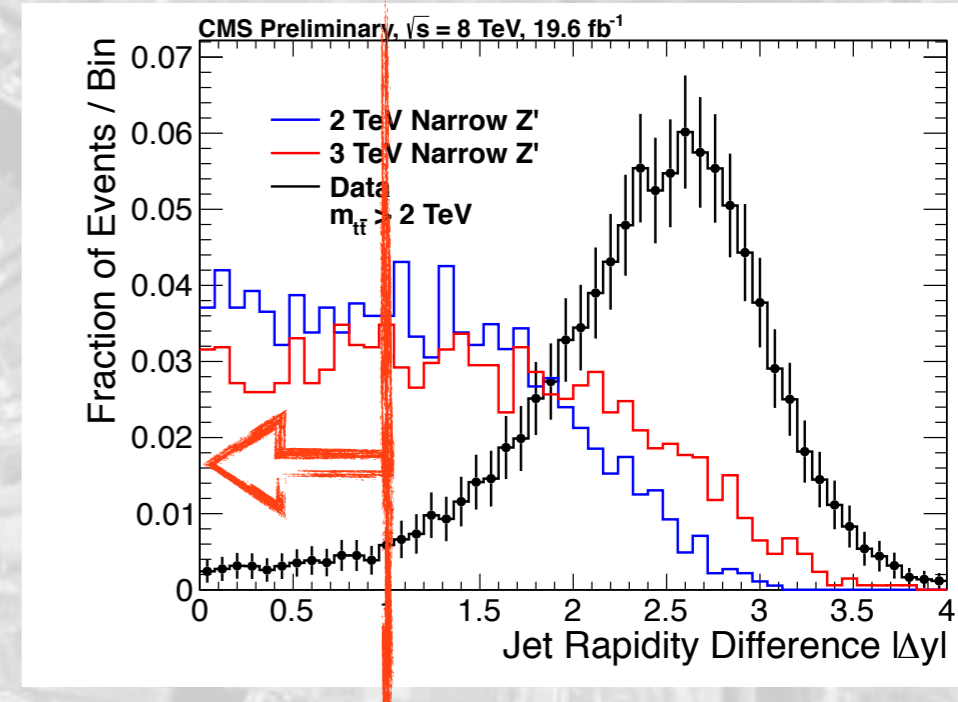
FULLY HADRONIC $Z' \rightarrow t\bar{t}$

Consider two event topologies: type-I + type I or type-I + type-II



Select events by jet-rapidity difference

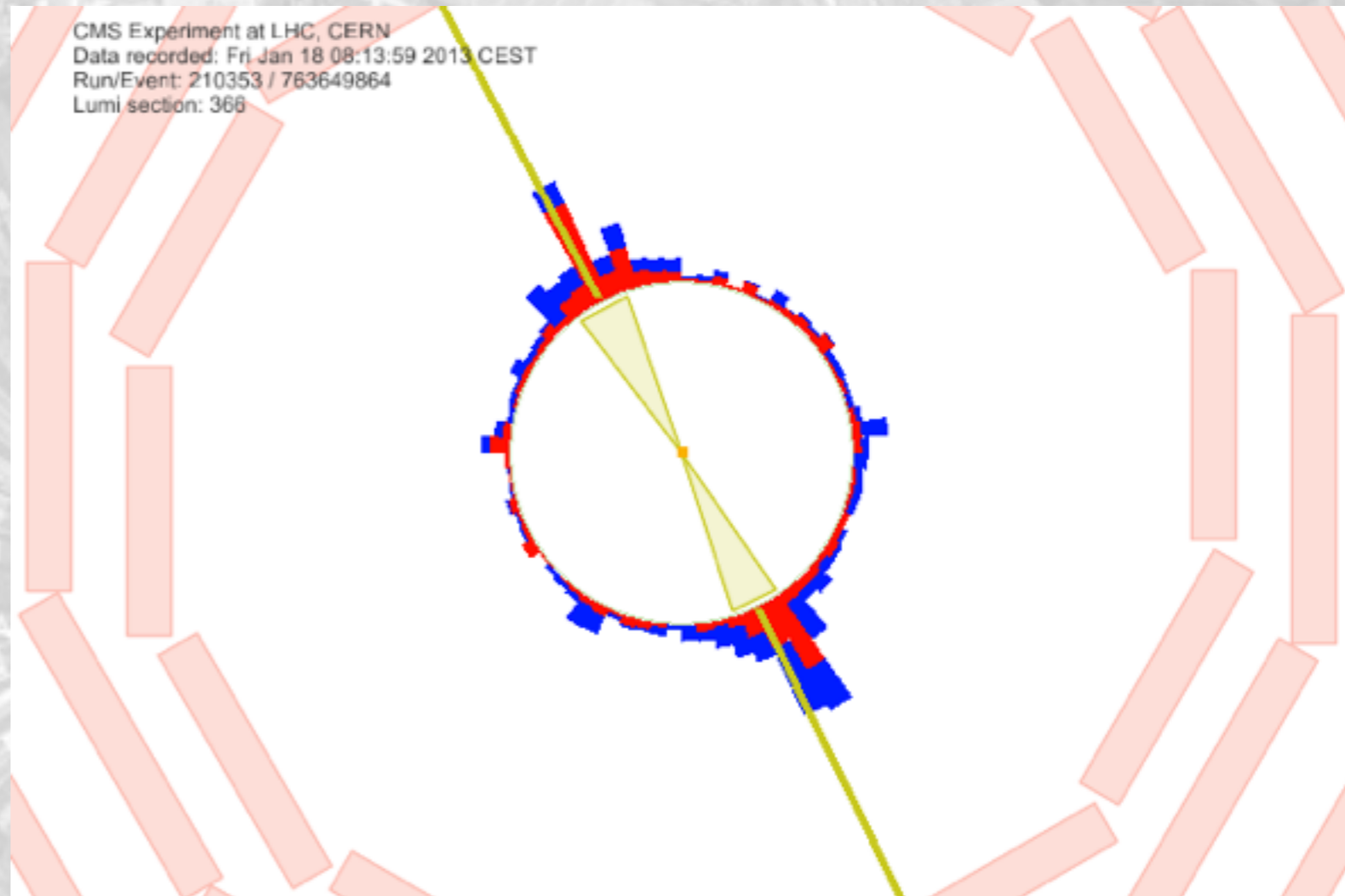
Study the ditop mass spectrum, looking for a bump on a falling spectrum



WHAT WE MEAN BY JET

- The concept of a jet is intuitive

- A quark or gluon is produced
- It showers other quarks and gluons
- The quarks hadronize to hadrons
- Because of momentum conservation, the hadrons are close to the original parton



- We then imagine a jet as a cone around the parton

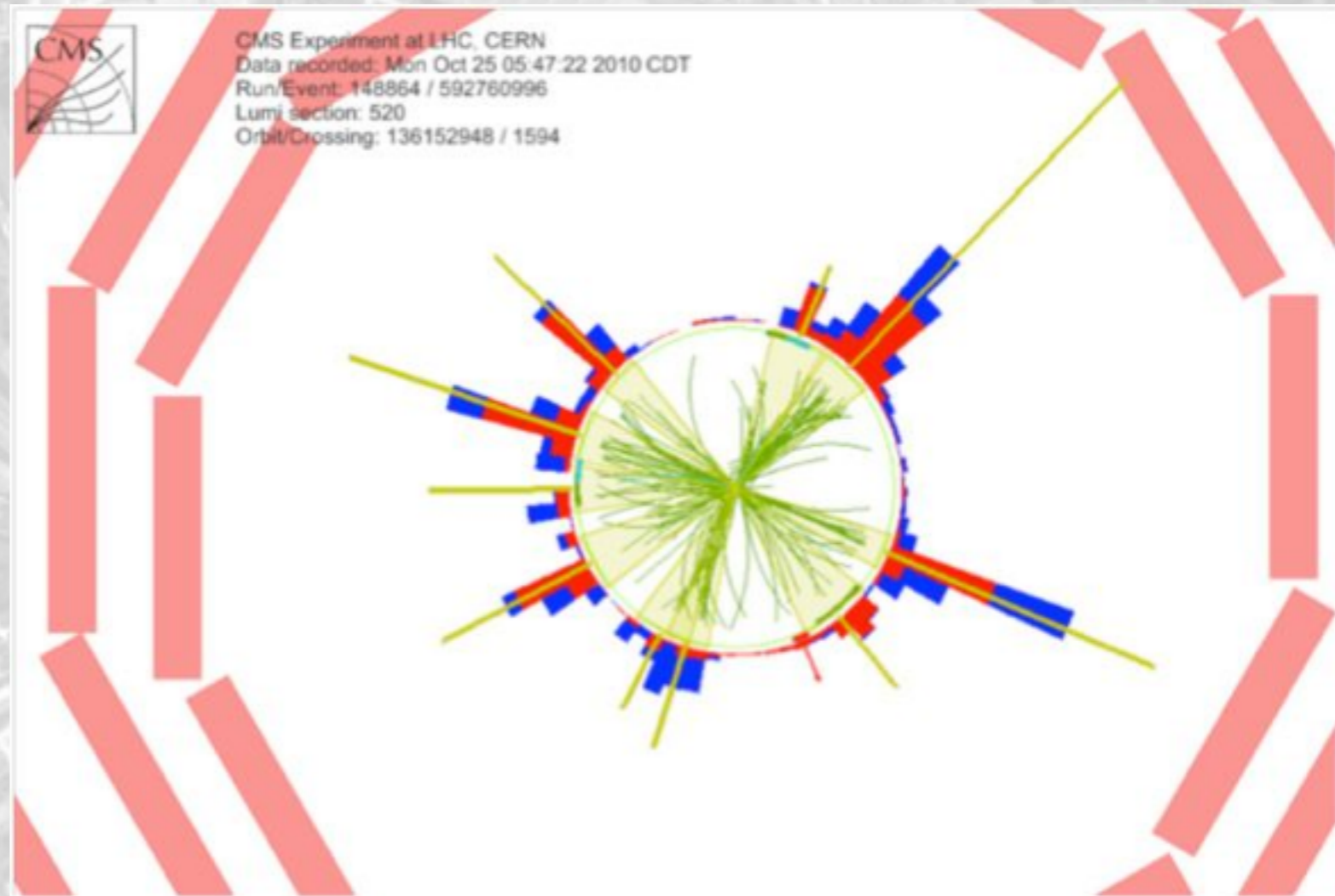
- But QCD is not geometry

- **How many jets do you see?**

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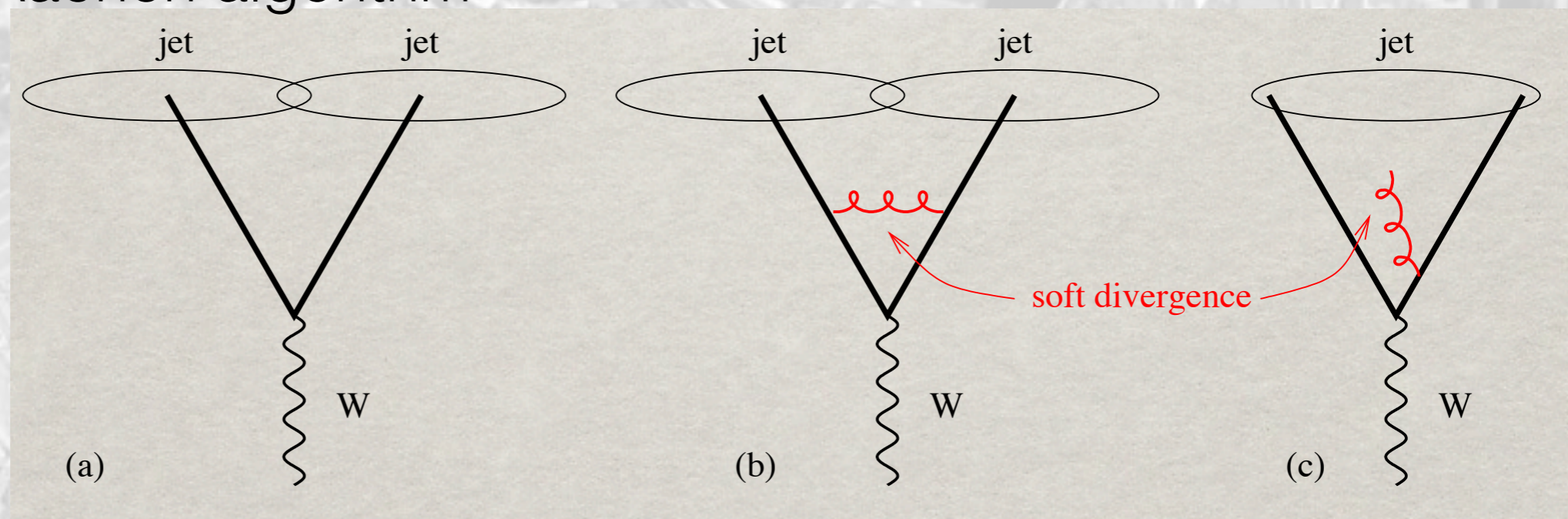
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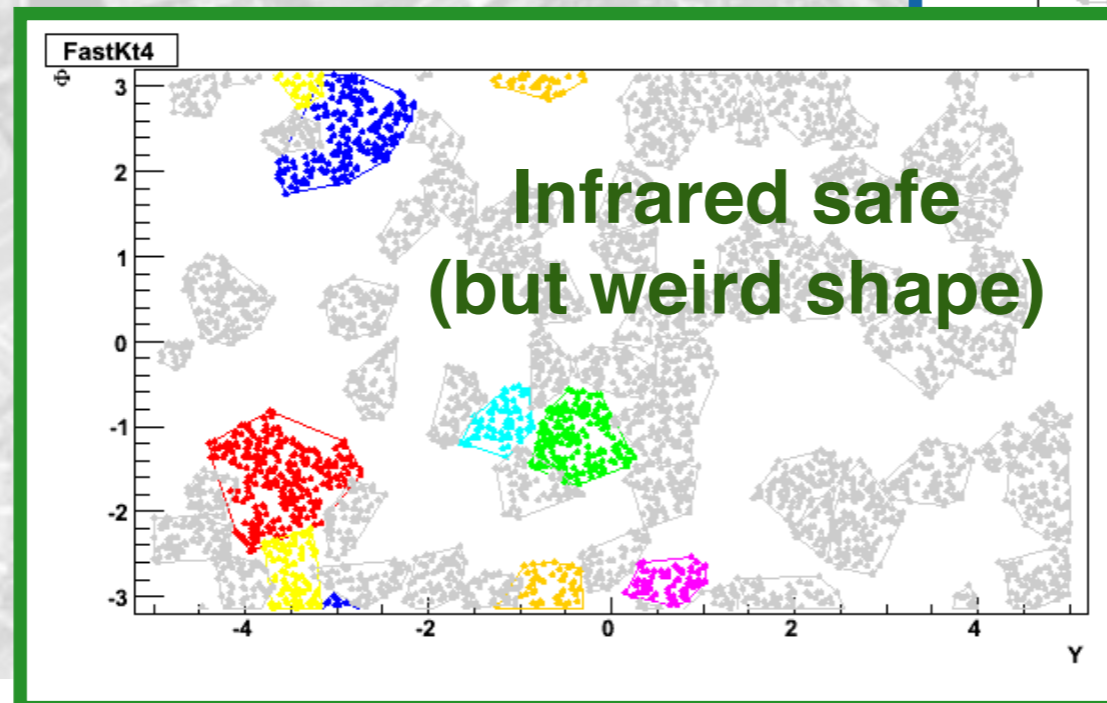
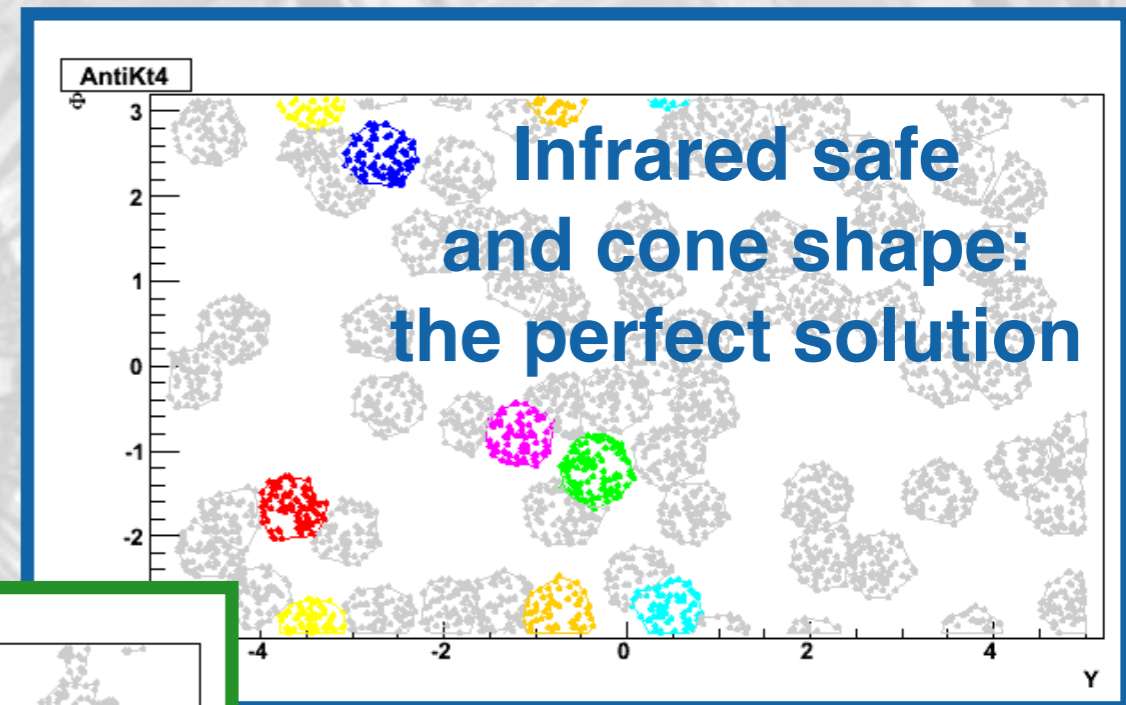
MORE THAN A CONE

- For some time, cone algorithms were used to reconstruct jets
- This created problems to compute QCD processes (e.g., differential cross section vs jet p_T). The number of reconstructed jets was fragile vs. soft divergences
- Several solutions proposed in literature
 - kT algorithm
 - Cambridge-Aachen algorithm
 - Siscone
 - Anti-kT



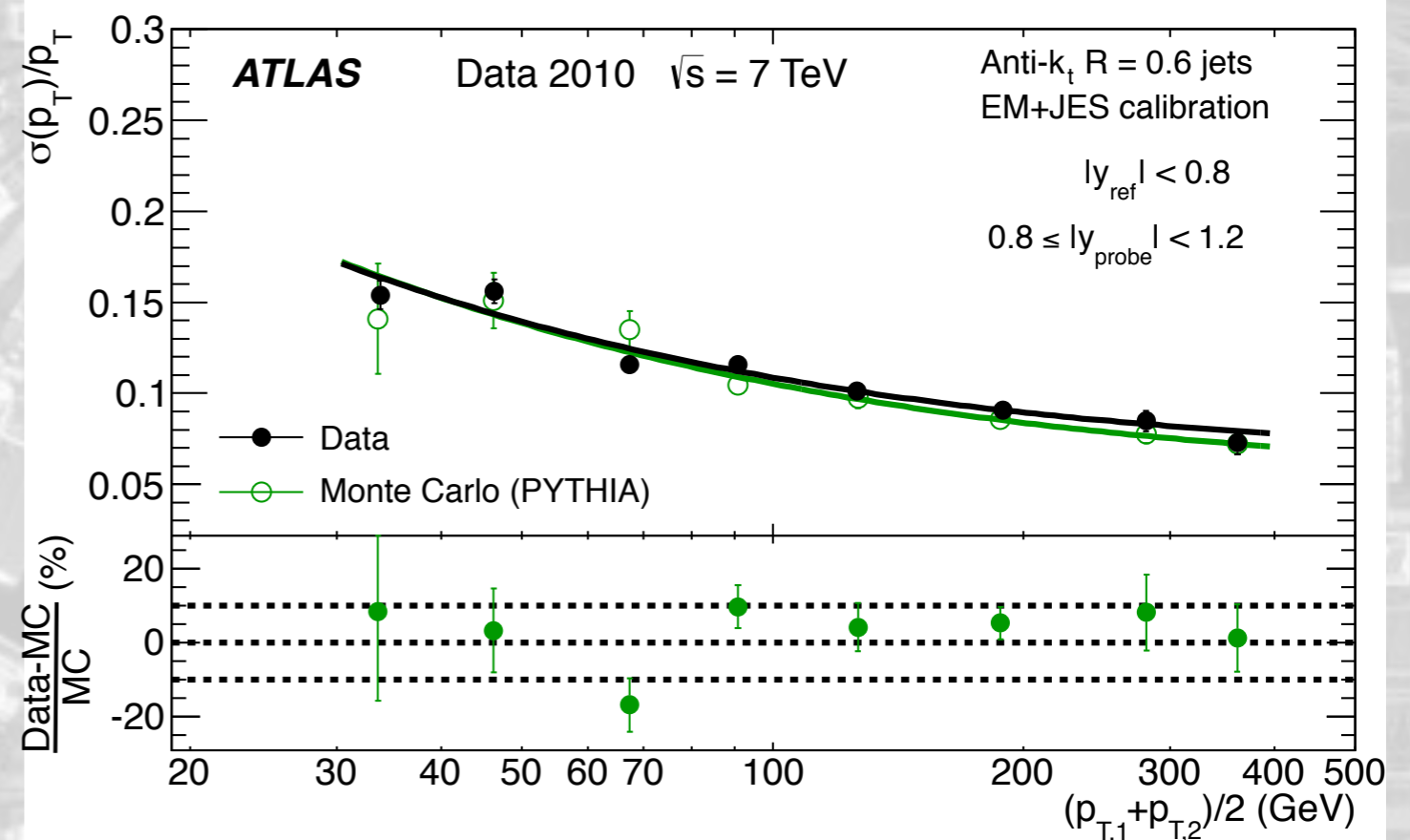
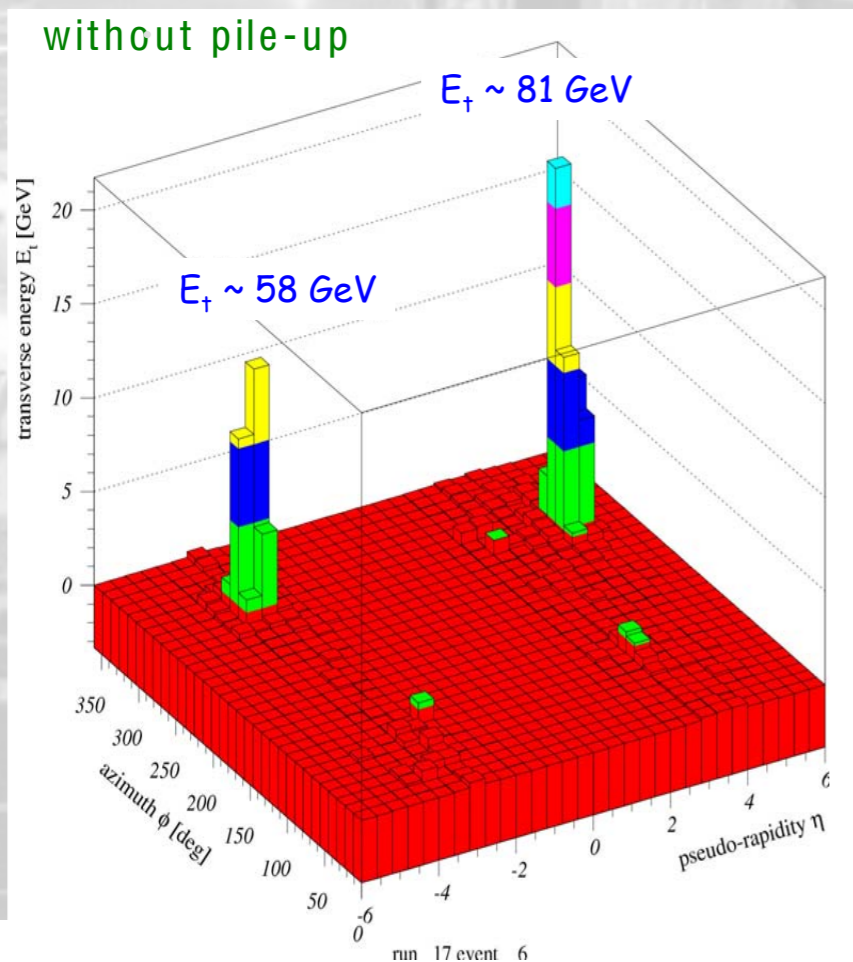
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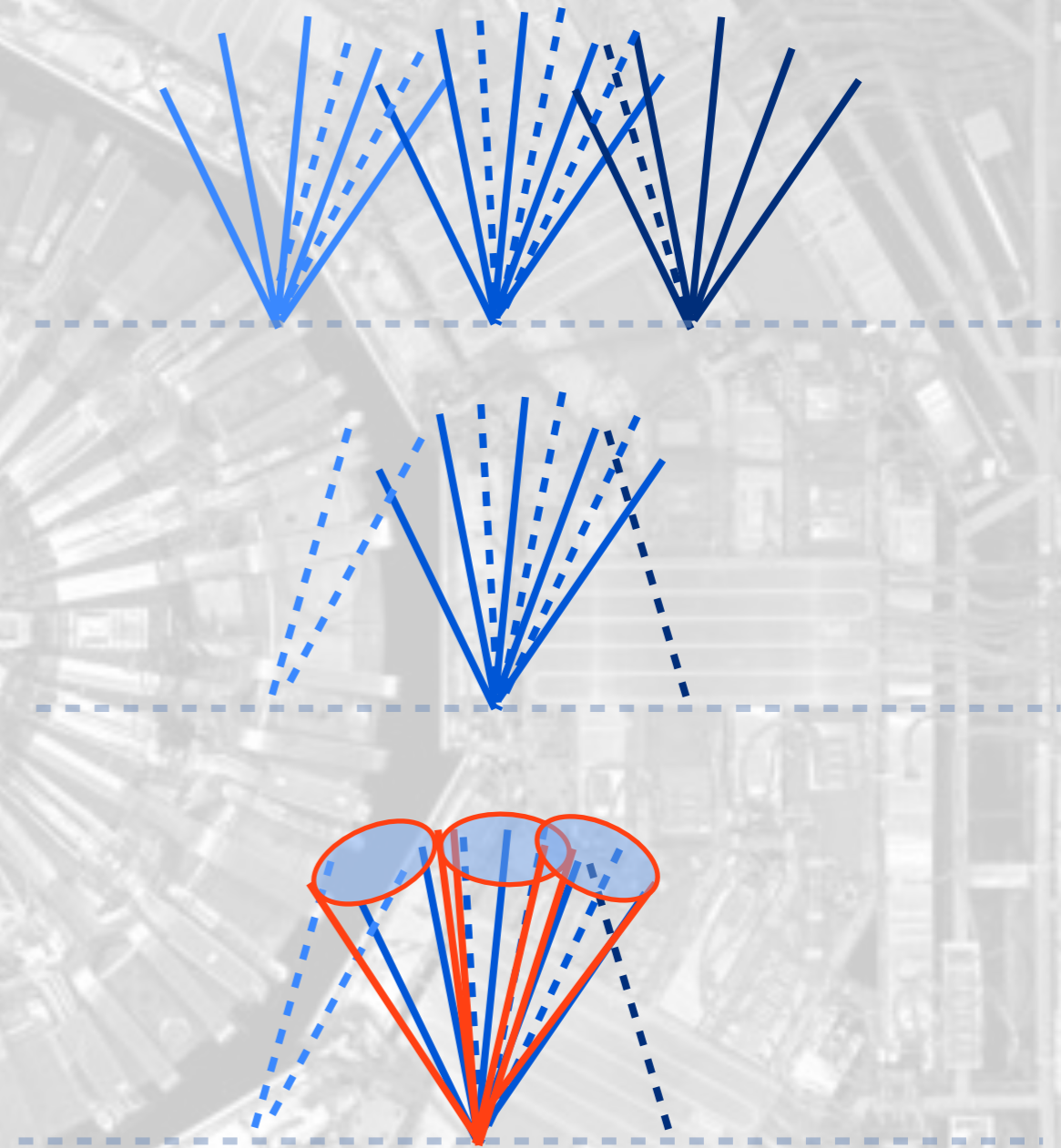
JETS OF WHAT?

- Historically, jets are made putting together energy deposits in the calorimeter
 - Deposits in ECAL and HCAL are put together into clusters
 - clusters are made into jets using a jet algorithm
- Good technique, as long as
 - the calorimeter has the dept to contain the shower
 - the energy resolution is good for both ECAL and HCAL (as in ATLAS)



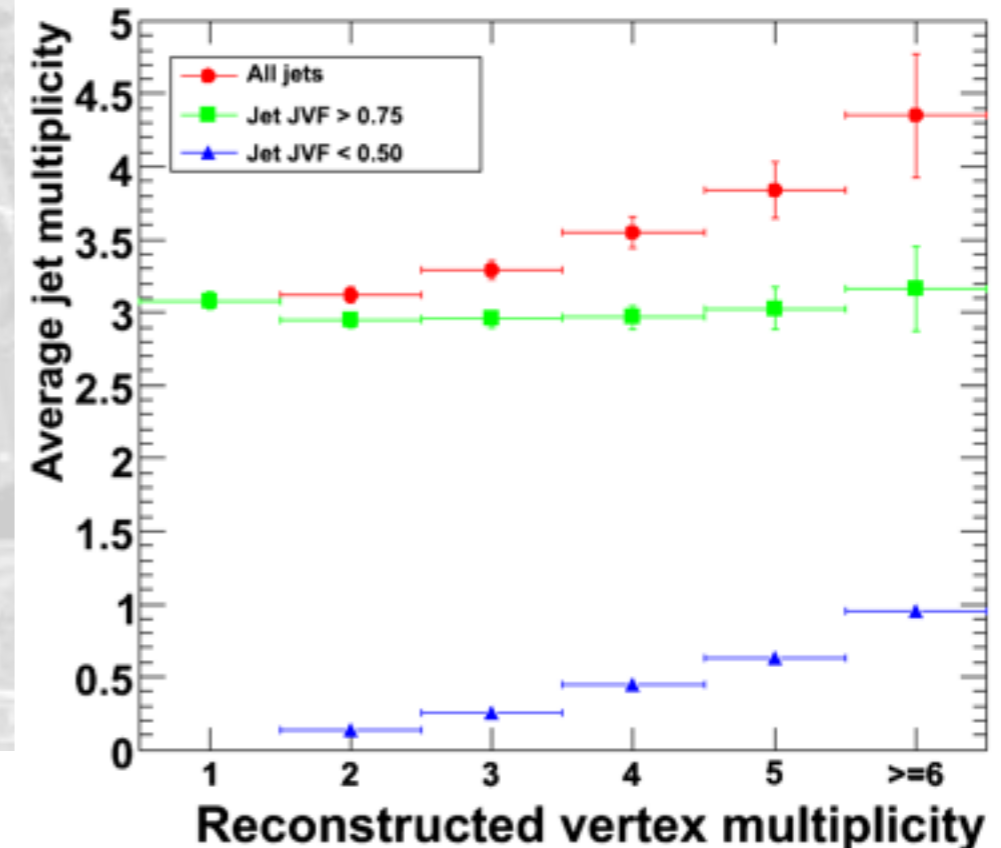
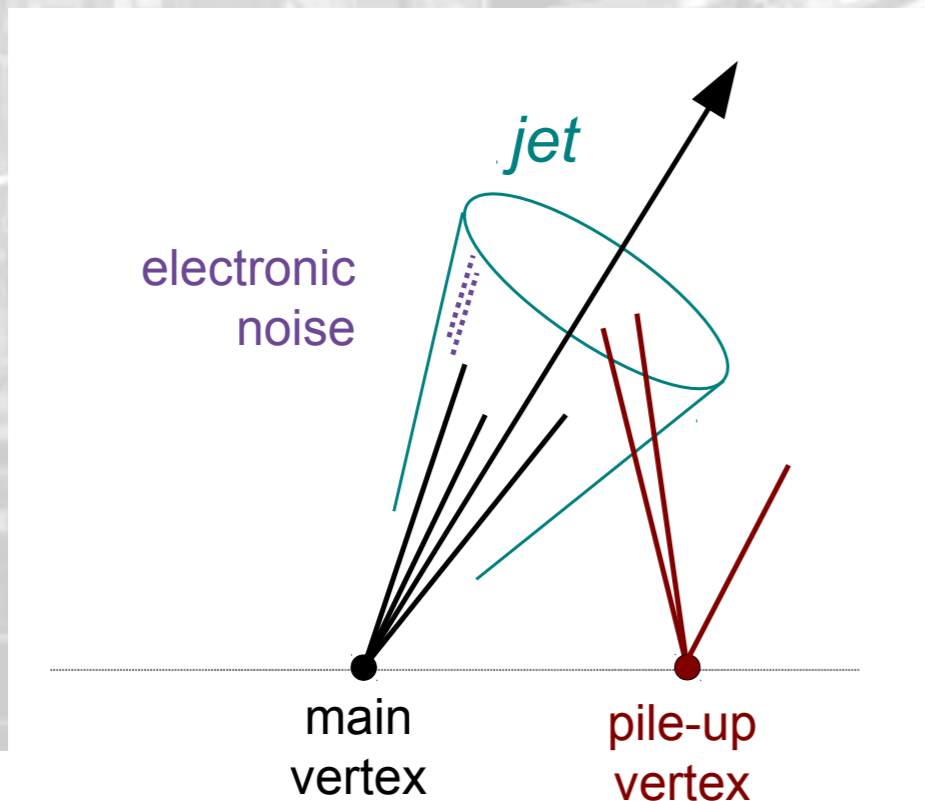
CHARGE HADRON SUBTRACTION

- We know where charged particles come from
- we can reconstruct vertices from tracks
- for each track, we know the vertex it comes from
- For an event, we know (usually) which is the “interesting” vertex (highest associated energy) which triggered the event
- It is then easier to remove the charged hadrons from the wrong vertices, before clustering the jets



CHARGE HADRON SUBTRACTION

- ATLAS does not use PFJets.
- Tracks vertexing is used to perform CHS a posteriori
 - Identify the tracks inside a jet “cone”
 - Compute the fraction of track momentum coming from the primary vertex
 - If this number is small, the jet is most likely due to pileup and it is rejected



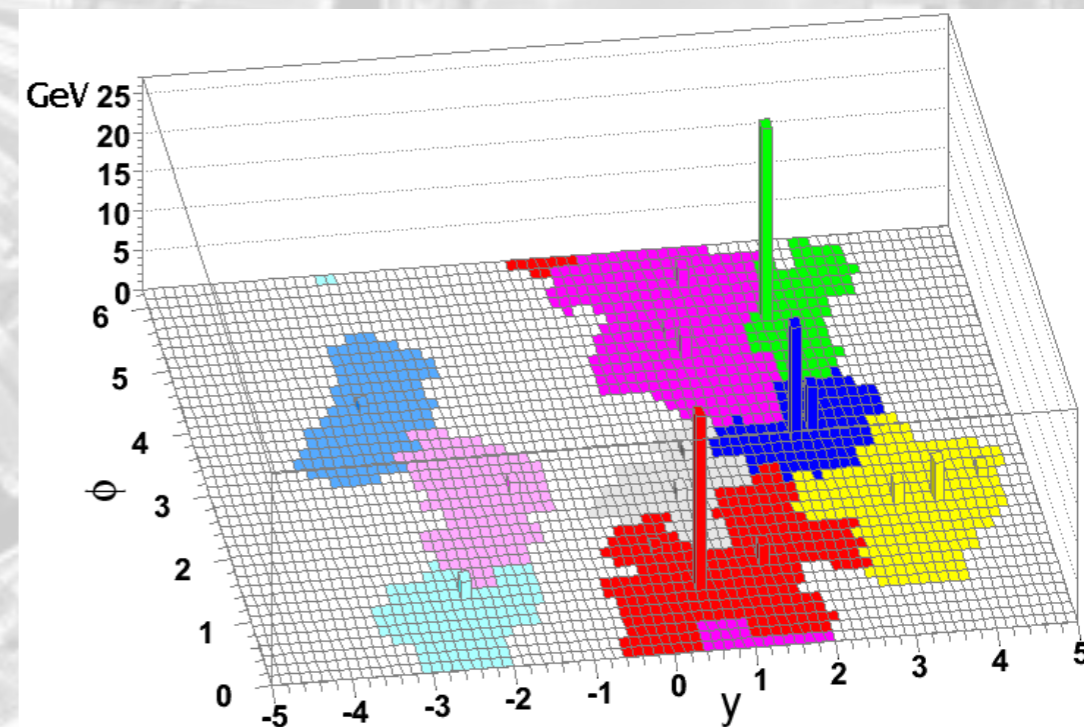
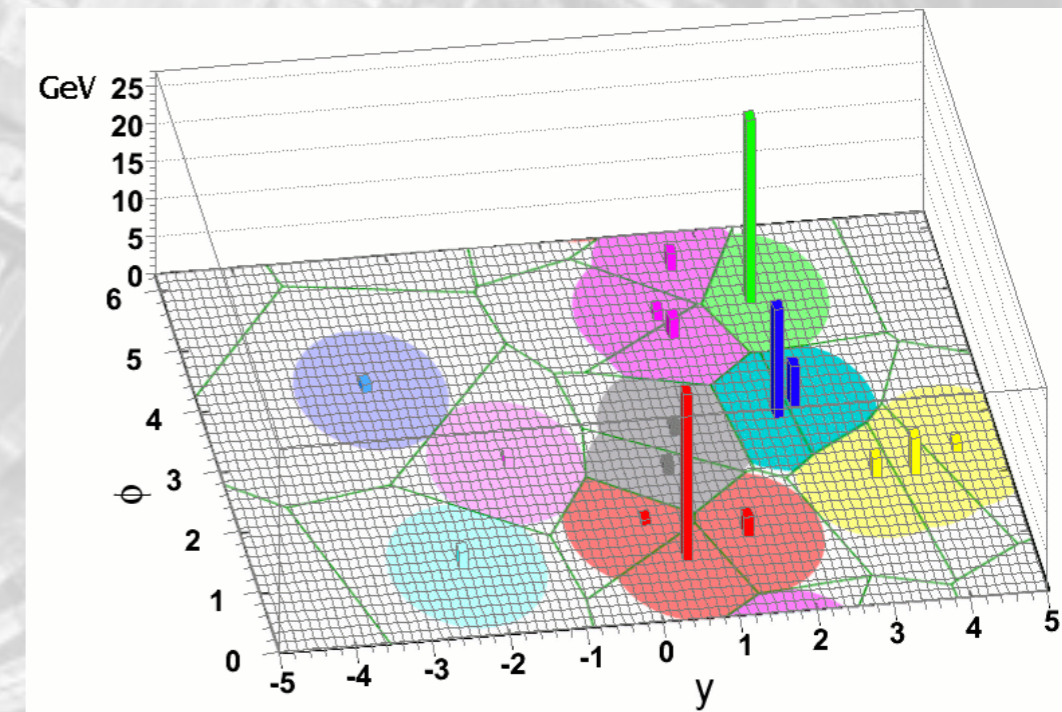
JET AREA SUBTRACTION

- When an exact subtraction of PU contamination is not possible, one can correct this in average
- measure the average energy deposit per unit of area ρ
- measure the jet area A
- determine the jet p_T by subtraction

$$p_{tj}^{(\text{sub})} = p_{tj} - A_j \rho$$

- The concept of a jet area is subject to definition

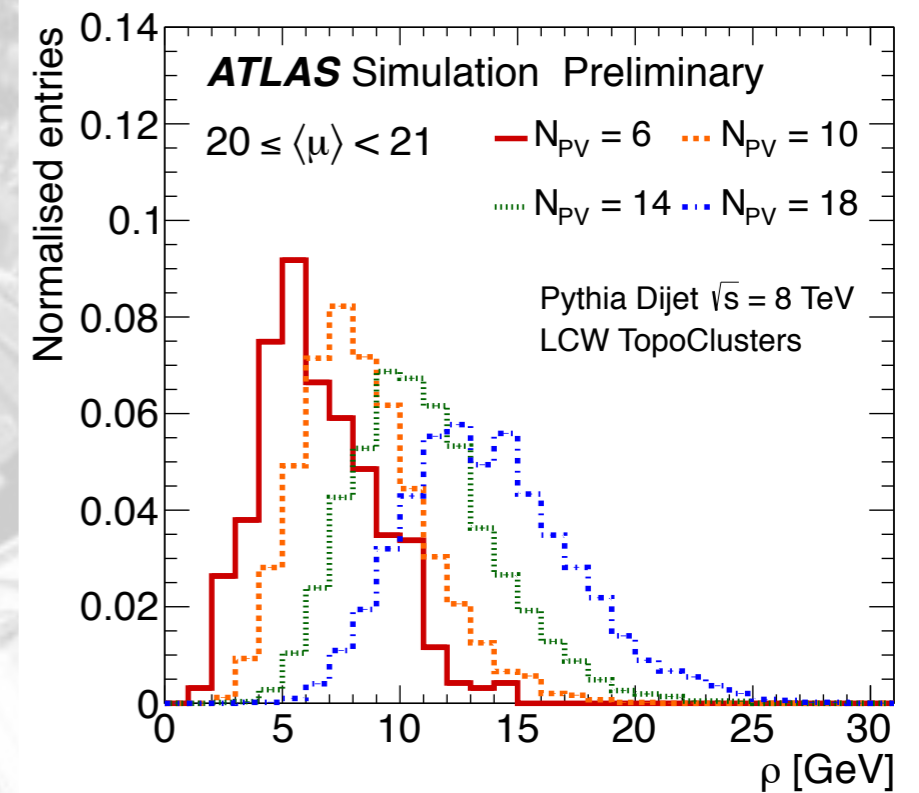
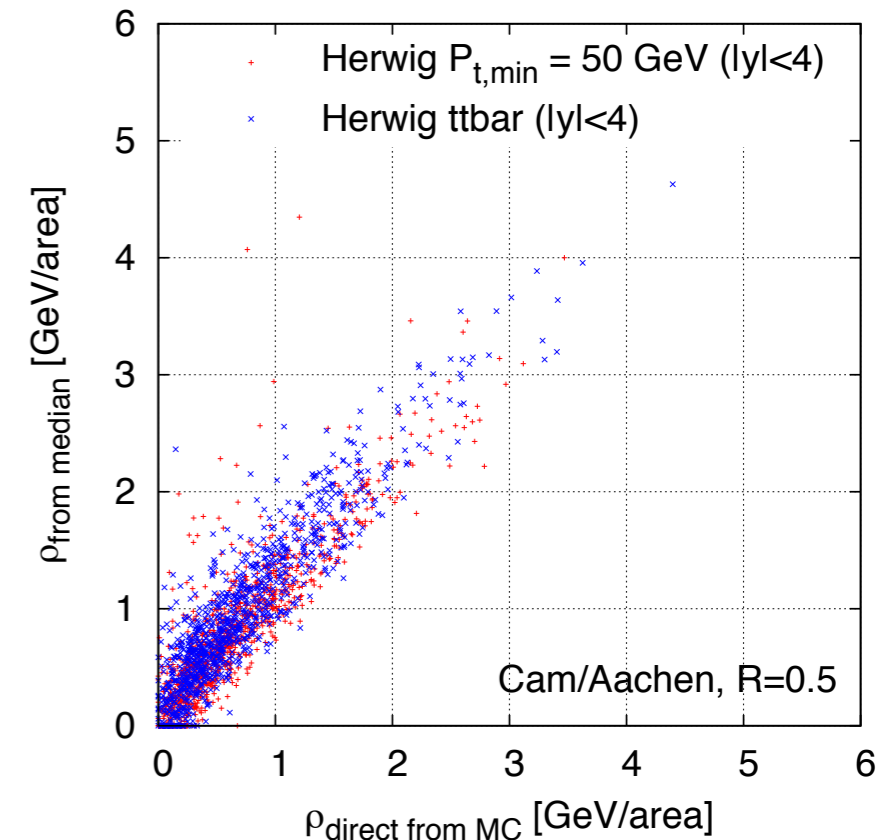
arxiv.org/abs/0802.1188



JET AREA SUBTRACTION

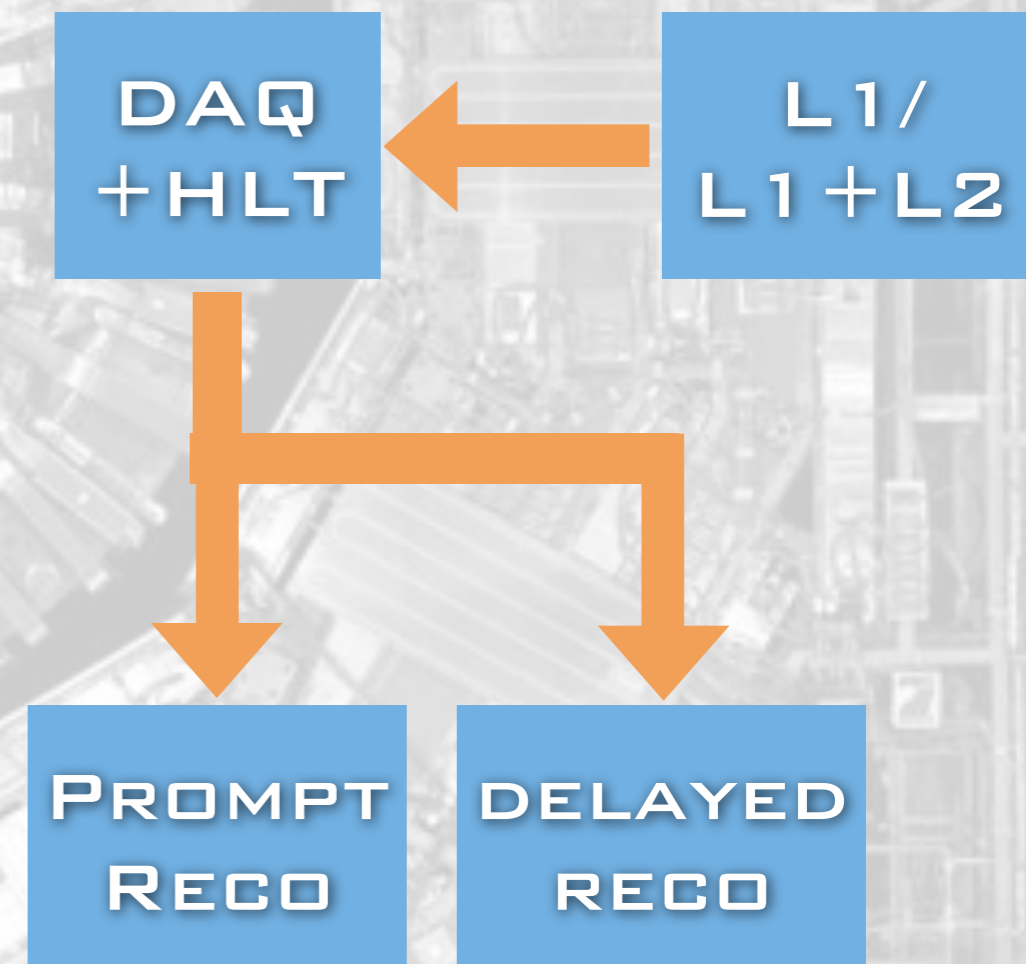
- Once the area is determined, one needs to estimate the energy density
- For this, one should consider that the list of jets clustered in an event is made of
 - a few high-pT jets
 - a long list of soft jets, coming from clustering the PU particles
- Under the assumption of a uniform diffuse PU contribution
 - each soft jet has energy ρA
 - ρ can then be estimated as

$$\rho = \text{median} \left[\left\{ \frac{p_{tj}}{A_j} \right\} \right]$$



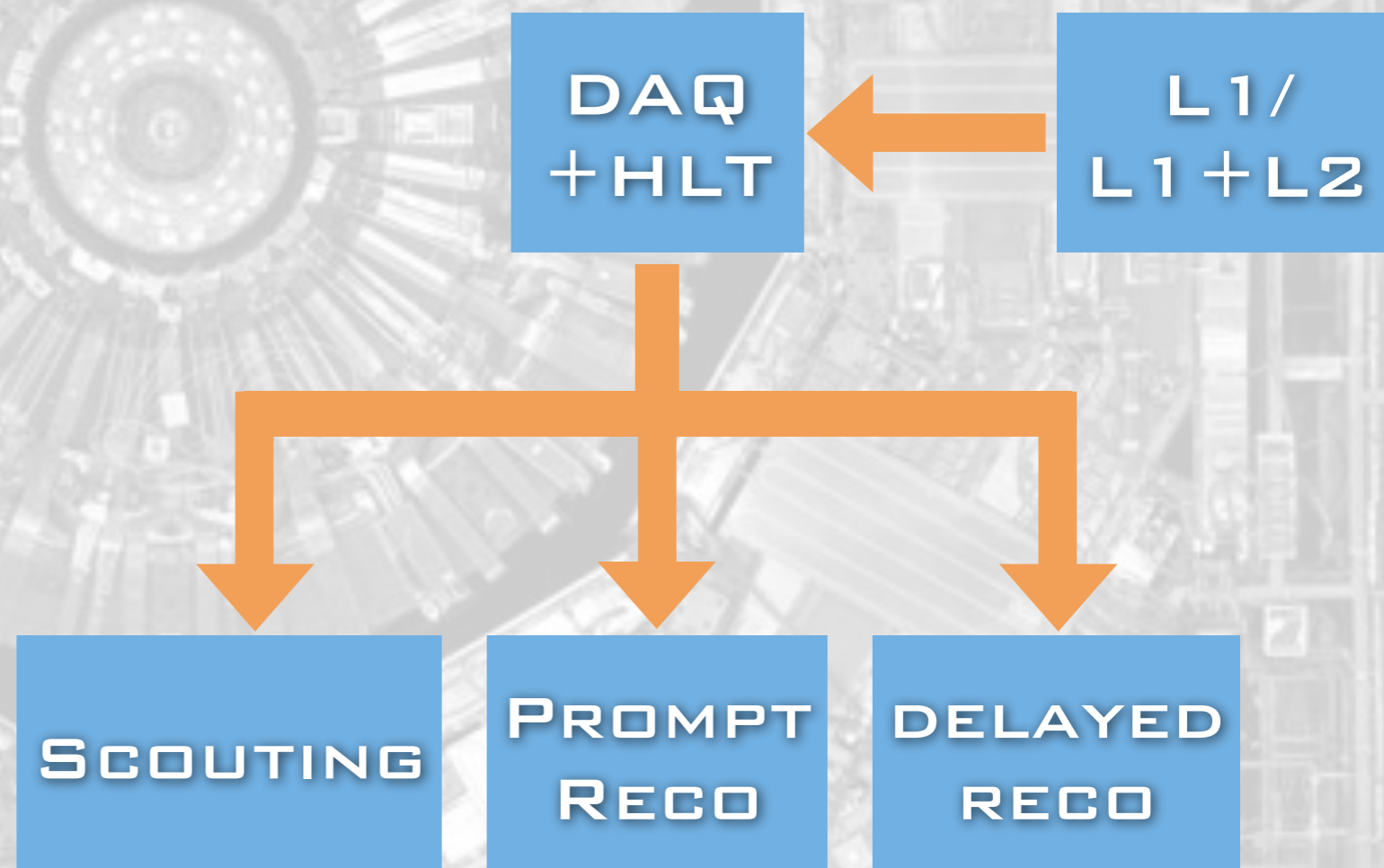
RATE VS DISK I/O

- At the beginning of the run, we considered 300 events/sec the maximum we could get out of the trigger
- It was suddenly realised that the trigger can handle much more
- The real problem was after
 - enough disk to write the data to?
Set the limit to 1 kHz, considering the available disk, and the event size ($O(300 \text{ Mb/event})$ after reconstruction)
 - enough CPU to process the data and reconstruct? OK for 2012 (full 2013/2014 shutdown)
- Data parking was introduced



BEYOND DATA PARKING

- We don't save events because we cannot write them on disk
- We write on disk more event information that we (sometimes) need
- Keep ALL the events, with a small customised data format (only jet 4momenta). $O(10 \text{ kb/event})$, can save several 1kHz of data
- Event reco not as good as normal
 - not all the analyses cannot be done this way
 - when resolution is not great (e.g., with jets) this is not a big issue
- Can do things we could not do otherwise (i.e., better than nothing)

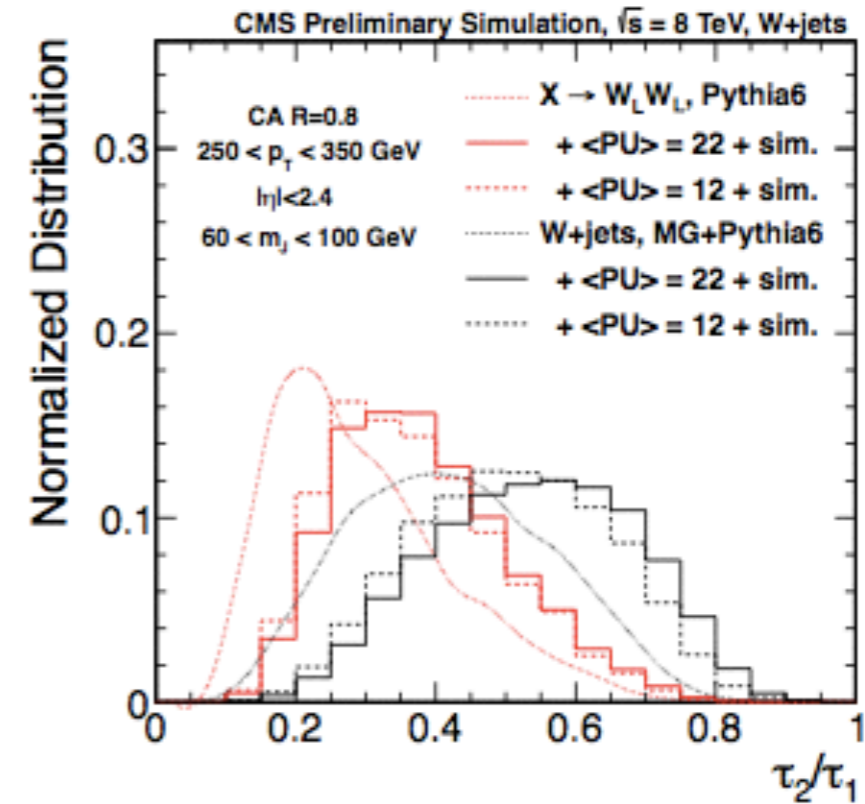
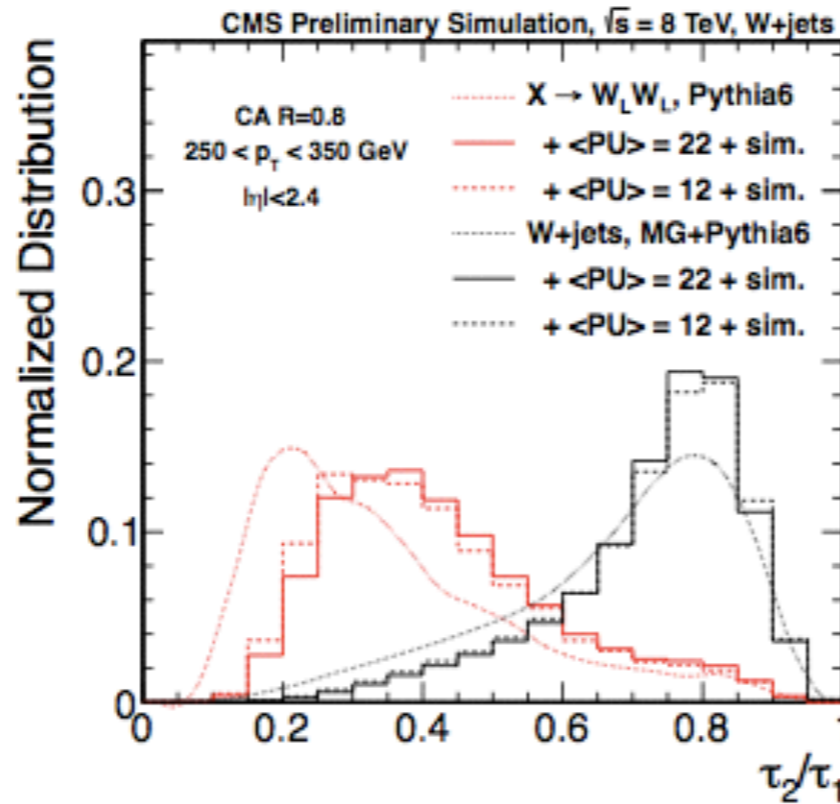


MORE ON SUBSTRUCTURE

Before Mass cut

After Mass cut

Not Pruned



Pruned

