

Towards a determination of the tau lepton dipole moments

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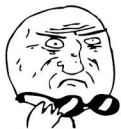
[hep-ph: 1301.5302](#)

[hep-ph: 1310.1081](#)

work in collaboration with:

S. Eidelman, D. Epifanov, L. Mercolli, M. Passera.

Problem:



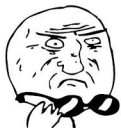
MOTHER OF GOD

Electron

$$a_e = 1\,159\,652\,180.73(28) \cdot 10^{-12}$$

0.24 parts per billion! Hanneke et al, PRL100 (2008) 120801

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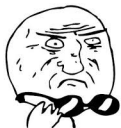
$$a_\mu = 1\,165\,920\,89(63) \cdot 10^{-11}$$

0.54 parts per million! E821-Final Rep: PRD73 (2006) 072003



NOT BAD

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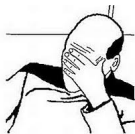
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NOT BAD



Tau

$$-0.052 \leq a_\tau \leq 0.013$$

Not even a test of LO! $\alpha/(2\pi) \approx 0.00116$
DELPHI - EPJC 35 (2004) 159



Outline

$g-2$ and EDM

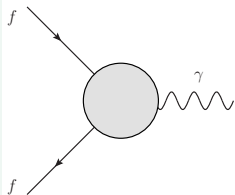
Radiative Leptonic Decays

Feasibility study

Conclusions

Form Factors & Dipole Moments

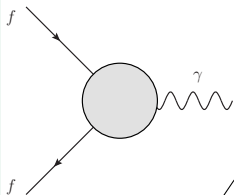
Vertex Function:



$$\Gamma_{ff\gamma}^{\mu}(q^2) = \gamma^{\mu} \left[F_{1V}(q^2) + \gamma_5 F_{1A}(q^2) \right] + \frac{\sigma^{\mu\nu}}{2m_f} q_{\nu} \left[i F_{2V}(q^2) + F_{2A}(q^2) \gamma_5 \right]$$

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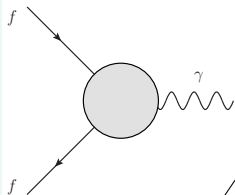


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► Charge: $F_{1V}(q^2 = 0) = Q_f$

Form Factors & Dipole Moments

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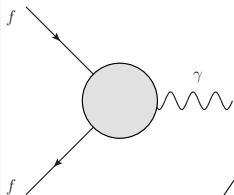


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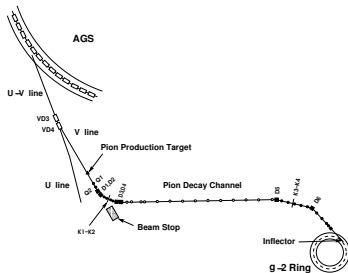
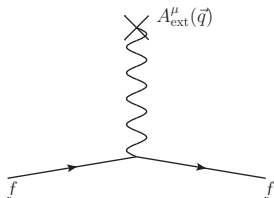


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- ▶ Charge: $F_{1V}(q^2 = 0) = Q_f$
- ▶ $g-2$: $F_{2V}(q^2 = 0) = a_f Q_f$
- ▶ EDM: $F_{2A}(q^2 = 0) = d_f(2m_f)/e$

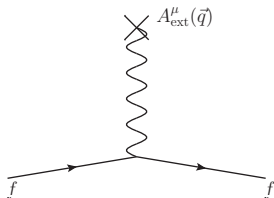
Direct and Indirect Measurement

► Direct measurements

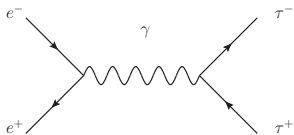


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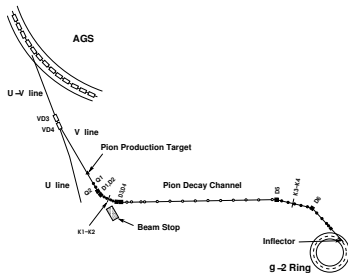
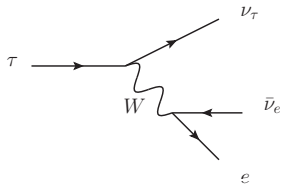
▶ Direct measurements



▶ Indirect measurements



$$q^2 \neq 0$$





Effective Lagrangian Approach

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{\text{SM}} + \left[c_a \frac{eQ_\tau}{4\Lambda} \bar{\tau} \sigma^{\mu\nu} \tau - c_d \frac{i}{2\Lambda} \bar{\tau} \sigma^{\mu\nu} \gamma_5 \tau \right] F_{\mu\nu}$$

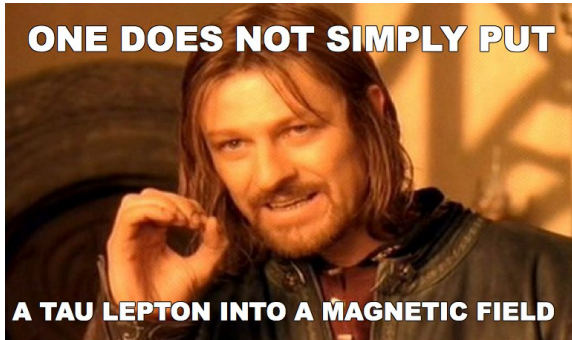
- ▶ The new terms arise after integrating out the heavy degrees of freedom associated to possible NP.

$$a_\tau = a_\tau^{\text{SM}} + c_a \frac{m_\tau}{\Lambda} + \dots$$

$$d_\tau = d_\tau^{\text{SM}} + c_d \frac{1}{\Lambda} + \dots$$

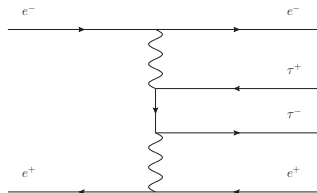
- ▶ The contributions from the two effective operators to a_τ and d_τ are the same for $q^2 = 0$ and $q^2 \neq 0 \ll \Lambda^2$.
- ▶ Only higher dimensional operators would give rise to a difference between these two cases

Radiative Leptonic Decays



Current bounds

- ▶ $e^+ e^- \rightarrow e^+ e^- \tau^+ \tau^-$ at \sqrt{s} between 183 and 208 GeV at LEP2 (the PDG value)
- ▶ The 95% C.L. limit



$$-0.052 \leq a_\tau \leq 0.013 \quad \text{DELPHI - EPJC35 (2004) 159}$$

- ▶ EDM SM estimate $|d_\tau^{\text{SM}}| \leq 10^{-35} e \cdot \text{cm}$!
- ▶ EDM current 95% C.L. limits from $e^+ e^- \rightarrow \tau^+ \tau^-$:

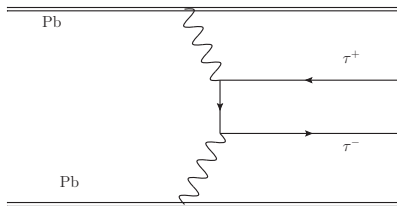
$$\begin{aligned} -2.2 &\leq \text{Re}(d_\tau) \leq 4.5 (10^{-17} e \cdot \text{cm}) \\ -2.5 &\leq \text{Im}(d_\tau) \leq 0.8 (10^{-17} e \cdot \text{cm}) \end{aligned}$$

Belle coll. PLB 551 (2003) 16.

Determination of a_τ : Proposals

► Heavy ions collisions at LHC

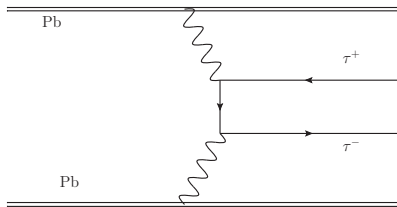
$\text{Pb Pb} \rightarrow \text{Pb Pb} \gamma\gamma \rightarrow \text{Pb Pb} \tau\tau$



F. del Aguila et al, PLB 271 (1991) 256

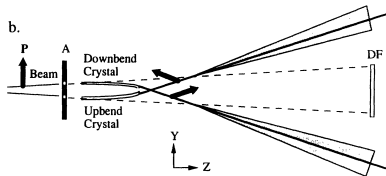
Determination of a_T : Proposals

- ▶ Heavy ions collisions at LHC
 $\text{Pb Pb} \rightarrow \text{Pb Pb } \gamma\gamma \rightarrow \text{Pb Pb } \tau\tau$



F. del Aguila et al, PLB 271 (1991) 256

- ▶ Channeling of polarized τ s in a bent crystal.



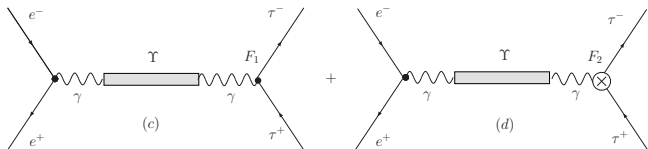
Chen et al. PRL 69 (1992) 3286

Samuel et al. PRL 67 (1991) 668

Determination of a_τ : Proposals

Bernabéu et al. propose the measurement of $F_{2V}(q^2 = M_\Upsilon^2)$ from $e^+e^- \rightarrow \Upsilon \rightarrow \tau^+\tau^-$ production at B factories.

Bernabéu et al. NPB 790 (2008) 160.

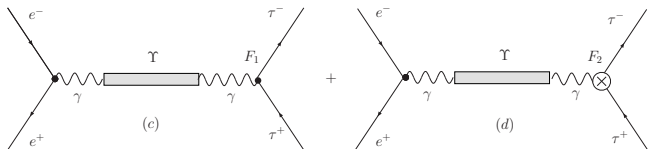


- ▶ Expected sensitivity for Babar+Belle: $4.6 \cdot 10^{-6}$ with $\mathcal{L} = 2 \text{ ab}^{-1}$.

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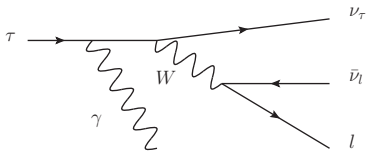
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- ▶ Expected sensitivity for Babar+Belle: $4.6 \cdot 10^{-6}$ with $\mathcal{L} = 2 \text{ ab}^{-1}$.
- ▶ At Belle and Belle II the visible cross section is dominated by non-resonant interaction due to the beam energy spread.

$$\Gamma_{\Upsilon(1S),\Upsilon(2S),\Upsilon(3S)} \sim \mathcal{O}(10 \text{ keV}), \quad \sigma_\epsilon \sim 3 \text{ MeV}$$

Leptonic Radiative Decays of the Tau



$$\tau^\pm \rightarrow \gamma l^\pm \nu_\tau \nu_l$$

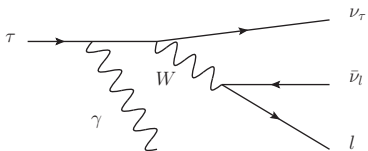
$$\text{with } l = e, \mu$$

- ▶ Suggested by Laursen et al. to search for the a_τ in radiative leptonic τ decays using the phenomenon of radiation zero.
- ▶ The SM tree-level amplitude vanishes in the phase space region:

$$\cos(l, \gamma) = -1, \quad E_l = \frac{m_\tau^2 + m_l^2}{2m_\tau} \quad (\text{in the tau r.f.})$$

M. L. Laursen et al. PRD 29 (1984) 2652

Leptonic Radiative Decays of the Tau



$$\tau^\pm \rightarrow \gamma l^\pm \nu_\tau \nu_l$$

with $l = e, \mu$

- ▶ Extend the strategy to d_τ ,
- ▶ Probe at $\mathcal{O}(10^{-3})$ the parameters

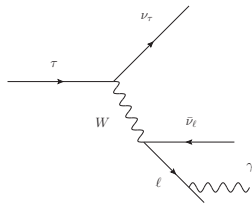
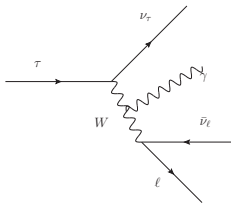
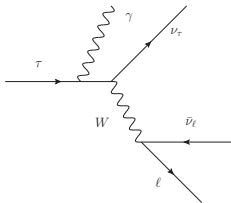
$$\tilde{a}_\tau \equiv c_a \frac{m_\tau}{\Lambda} \quad \text{and} \quad \tilde{d}_\tau \equiv c_d \frac{1}{\Lambda}$$

- ▶ Provide the theoretical framework for such measurement.

QED NLO Corrections to Tau Radiative Decay

- ▶ Polarized differential decay rate $d\Gamma$ up to $\mathcal{O}(10^{-3})$.
- ▶ Dependence on E_l , E_γ , Ω_l and Ω_γ

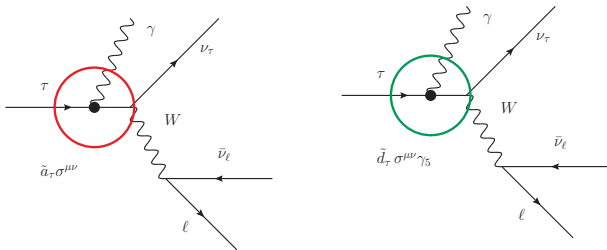
$$d\Gamma = d\Gamma_{\text{LO}} + \left(\frac{m_\tau}{M_W}\right)^2 d\Gamma_W + \tilde{a}_\tau d\Gamma_a + \tilde{d}_\tau d\Gamma_d + \frac{\alpha}{\pi} d\Gamma_{\text{NLO}}$$



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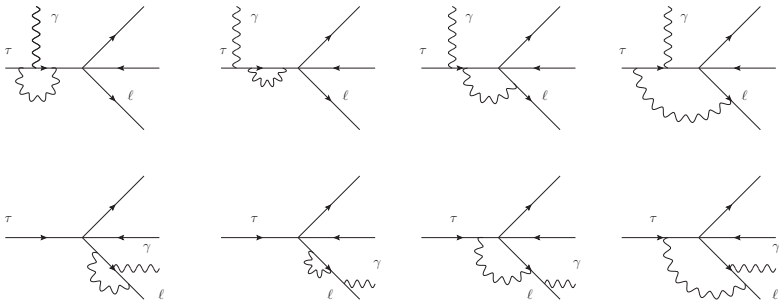
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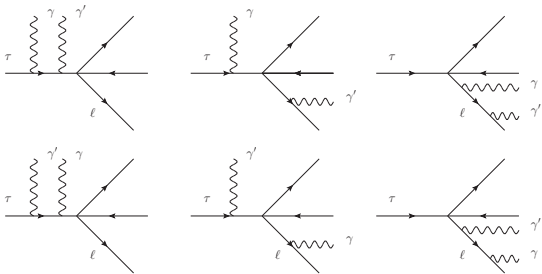
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Branching Ratios: $E_\gamma > 10 \text{ MeV}$

process	LO %	NLO %	MW %
$\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau \gamma$	1.836	-0.183	0.0006
$\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau \gamma$	0.367	-0.009	0.0001

process	B.R. %	exp. B.R. % (PDG)
$\mu^+ \rightarrow e^+ \nu_e \bar{\nu}_\mu \gamma$	1.3	1.4 (4) *
$\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau \gamma$	1.653	$1.75 \pm 0.06 \pm 0.17^\dagger$
$\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau \gamma$	0.358	$0.361 \pm 0.016 \pm 0.035^\dagger$

*Crittenden et al. PR 121 (1964) 1823, † CLEO Coll. PRL 84 (2000) 830

process	B.R. %	exp. B.R. %
$\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau \gamma$	1.653	$1.847 \pm 0.015 \pm 0.052^*$
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*BaBar preliminary results, B. Oberhof

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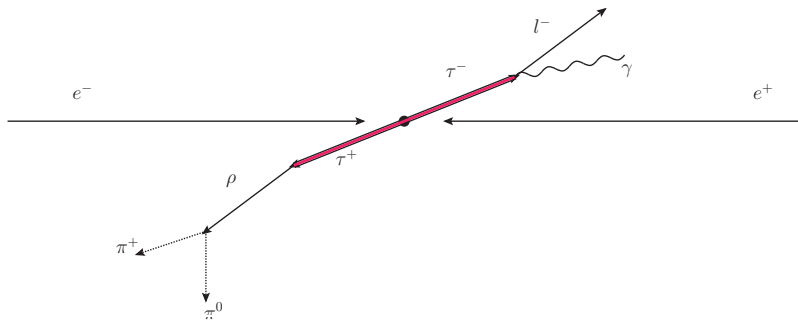
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3.6σ

*BaBar preliminary results, B. Oberhof

Feasibility study





Approach 1: Radiation Zero

- ▶ Study of the radiation zero point.
- ▶ A set of $\tau^+ \tau^-$ events analyzed. ($\tau^\pm \rightarrow l_1^\pm \nu \nu \gamma$, $\tau^\mp \rightarrow l_2^\mp \nu \nu$)
- ▶ Final state: $(l_1^\pm \gamma, l_2^\mp)$ with $l_1, l_2 = e, \mu$ and $l_1 \neq l_2$

$$\cos(l_1, \gamma) < -0.9, \quad 0.1 < \cos(l_2, \gamma), \quad \text{and} \quad E_\gamma > 0.5 \text{ GeV}$$

- ▶ With the whole Belle statistics ($0.9 \times 10^9 \tau$ pairs), the upper \tilde{a}_τ upper bound is



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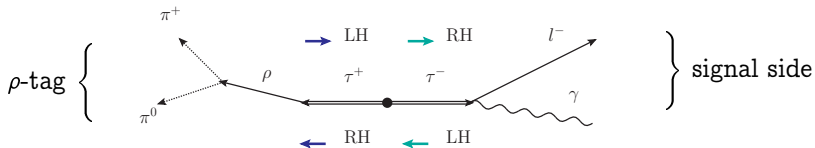
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- ▶ With the whole Belle statistics (0.9×10^9 τ pairs), the upper \tilde{a}_τ upper bound is

$$\text{U.L. } (\tilde{a}_\tau) \simeq 2$$

Approach 2: Unbinned Maximum Likelihood

- ▶ Take advantage of $\tau^\pm \rightarrow \rho^\pm \nu \rightarrow \pi^\pm \pi^0 \nu$ as a spin analyzer.
- ▶ In $\tau^\mp \rightarrow l^\mp \nu \nu \gamma$ we are sensitive to the spin dependent part.
- ▶ 12-dimensional phase space analysis ($l^\mp, \gamma, \pi^\pm, \pi^0$).



- ▶ Developed special generator of the $(l^\mp \nu \nu \gamma, \pi^\pm \pi^0 \nu)$ events
- ▶ Fit the generated event samples corresponding to the amount of data available at Belle and expected at Belle II.

Results

▶ ρ -tag mode, BR= 25.5%

$\tau^\pm \rightarrow \rho^\pm \nu$ only

▶ full-tag mode, BR= 90%

$\tau^\pm \rightarrow \rho^\pm \nu, \tau^\pm \rightarrow \pi^\pm \nu, \tau^\pm \rightarrow \pi^\pm \pi^0 \pi^0 \nu,$
 $\tau^\pm \rightarrow \pi^\pm \pi^+ \pi^- \nu, \tau^\pm \rightarrow e^\pm \nu \nu, \tau^\pm \rightarrow \mu^\pm \nu \nu$

Sensitivity to \tilde{a}_τ and \tilde{d}_τ				
	Re(\tilde{a}_τ)	Im(\tilde{a}_τ)	Re(\tilde{d}_τ)	Im(\tilde{d}_τ)
Belle (ρ -tag)	0.16	0.16	0.15	0.046
Belle-II (ρ -tag)	0.023	0.023	0.021	0.007
Belle (full tag)	0.085	0.085	0.080	0.024
Belle-II (full tag)	0.012	0.012	0.011	0.003
DELPHI [†]	0.017	—	—	—
Belle*	—	—	0.0015	0.0008

[†] DELPHI - EPJC35 (2004) 159

* Belle coll. PLB 551 (2003) 16.

- ▶ Radiative leptonic τ decays can probe τ DipM.
- ▶ We provided with the polarized differential decay rate at NLO in QED plus small W -boson effects.
We found some discrepancies with previous results.
- ▶ The upper limit achievable at Belle via radiation zero phenomenon is only $a_\tau \sim 1$.
- ▶ Analysis in the full phase-space is required.
- ▶ Feasibility study shows that Belle II can ameliorate the current DELPHI result for a_τ .
- ▶ The extraction of τ DipM from $e^+e^- \rightarrow \tau^+\tau^-$ is not excluded.
A careful theoretical reanalysis is needed.
- ▶ A possible dedicated experiment (bent crystal)?
- ▶ What are the prospects at the LHC?



Thanks!



Backup slides



QED NLO Corrections to Tau Radiative Decay

The total differential decay for a polarized τ lepton in the tau r.f. is

$$\frac{d^6\Gamma^{\text{NLO}}}{dx dy d\Omega_l d\Omega_\gamma} = \frac{\alpha G_F^2 m_\tau^5}{(4\pi)^6} \frac{x\beta}{1 + \delta_w(m_\mu, m_e)} \left[G(x, y, c) \right. \\ \left. + x\beta \hat{n} \cdot \hat{p}_l J(x, y, c) + y \hat{n} \cdot \hat{p}_\gamma K(x, y, c) + y x\beta \hat{n} \cdot (\hat{p}_l \times \hat{p}_\gamma) L(x, y, c) \right]$$

where $x = 2E_l/m_\tau$, $y = 2E_\gamma/m_\tau$, $c = \cos\theta_{l\gamma}$. The tau polarization vector $n = (0, \vec{n})$ satisfies $n^2 = -1$ and $n \cdot p_\tau = 0$. The function $G(x, y, c)$, and similarly for J and K , is given by

$$G(x, y, c) = \frac{4}{3yz^2} \left[g_{\text{LO}}(x, y, z) + \frac{\alpha}{\pi} g_{\text{NLO}}(x, y, z; y_{\text{min}}) + \left(\frac{m_\tau}{M_W} \right)^2 g_w(x, y, z) \right]$$

QED NLO Corrections to Tau Radiative Decay

The total differential decay for a polarized τ lepton in the tau r.f. is

$$\frac{d^6\Gamma^{\text{NLO}}}{dx dy d\Omega_l d\Omega_\gamma} = \frac{\alpha G_F^2 m_\tau^5}{(4\pi)^6} \frac{x\beta}{1 + \delta_{\text{W}}(m_\mu, m_e)} \left[G(x, y, c) \right. \\ \left. + x\beta \hat{n} \cdot \hat{p}_l J(x, y, c) + y \hat{n} \cdot \hat{p}_\gamma K(x, y, c) + y x\beta \hat{n} \cdot (\hat{p}_l \times \hat{p}_\gamma) L(x, y, c) \right]$$

where $x = 2E_l/m_\tau$, $y = 2E_\gamma/m_\tau$, $c = \cos\theta_{l\gamma}$. The tau polarization vector $n = (0, \vec{n})$ satisfies $n^2 = -1$ and $n \cdot p_\tau = 0$. The function $G(x, y, c)$, and similarly for J and K , is given by

$$G(x, y, c) = \frac{4}{3yz^2} \left[g_{\text{LO}}(x, y, z) + \frac{\alpha}{\pi} g_{\text{NLO}}(x, y, z; y_{\text{min}}) + \left(\frac{m_\tau}{M_W} \right)^2 g_{\text{W}}(x, y, z) \right]$$

Compared with previous work [A. B. Arbuzov PLB 597 \(2004\) 285](#)



QED NLO Corrections to Tau Radiative Decay

- ▶ Polarized differential decay rate $d\Gamma$ up to $\mathcal{O}(10^{-3})$.
- ▶ Dependence on E_l , E_γ , Ω_l and Ω_γ

$$d\Gamma = d\Gamma_{\text{LO}} + \left(\frac{m_\tau}{M_W}\right)^2 d\Gamma_W + \tilde{a}_\tau d\Gamma_a + \tilde{d}_\tau d\Gamma_d + \frac{\alpha}{\pi} d\Gamma_{\text{NLO}}$$

- ▶ Analytic expression implemented in Fortran

```
subroutine gnlo(resgnlo)
  implicit none
  external NLOfunctions
  double precision resgnlo
```

⋮

Feasibility study: Unbinned maximum likelihood

- ▶ Phase space point $X = (p_l, \Omega_l, p_\gamma, \Omega_\gamma, p_\rho, \Omega_\rho, m_{\pi\pi}^2, \tilde{\Omega}_\pi)$
- ▶ The PDF $\mathcal{P}(\vec{X})$ is constructed from the differential cross section

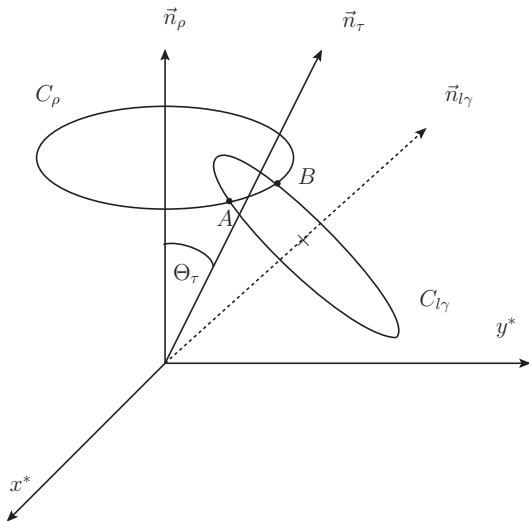
$$\frac{d\sigma}{d\text{PS}}(e^+ e^- \rightarrow \tau^\mp \tau^\pm \rightarrow (l^\mp \nu \nu \gamma, \pi^\pm \pi^0 \nu))$$

- ▶ Unbinned maximum likelihood of the generated events

$$\mathcal{P}(\vec{X} | \tilde{a}_\tau, \tilde{d}_\tau) = \frac{\mathcal{F}_{\tilde{a}_\tau, \tilde{d}_\tau}(\vec{X})}{\int \mathcal{F}_{\tilde{a}_\tau, \tilde{d}_\tau}(\vec{X}) d\vec{X}}$$

where $\mathcal{F}_{\tilde{a}_\tau, \tilde{d}_\tau}(\vec{X})$ is the visible differential cross section.

- ▶ Developed special generator of the $(l^\mp \nu \nu \gamma, \pi^\pm \pi^0 \nu)$ events
- ▶ Fit of generated event samples corresponding to the amount of data available at Belle and expected at Belle II.



- ▶ $\tau \rightarrow \rho\nu$:
 $(p_\tau - p_\rho)^2 = 0$
- ▶ $\tau \rightarrow \gamma l\nu\bar{\nu}$:
 $M_{\nu\bar{\nu}} > 0$