

Leptons & QCD

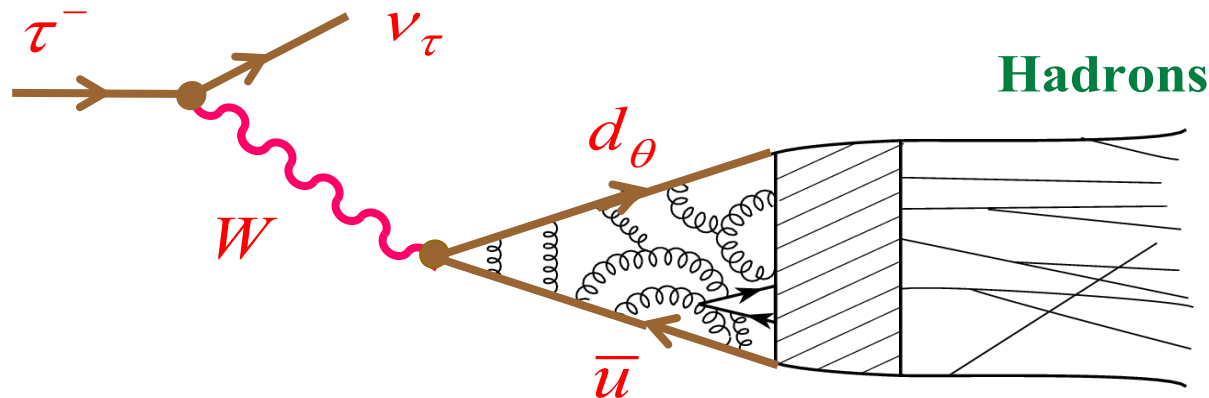
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IFIC, Valencia

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Aachen, Germany, 15-19 September 2014



HADRONIC TAU DECAY



$$d_\theta = V_{ud} d + V_{us} s$$

Only lepton massive enough to decay into hadrons

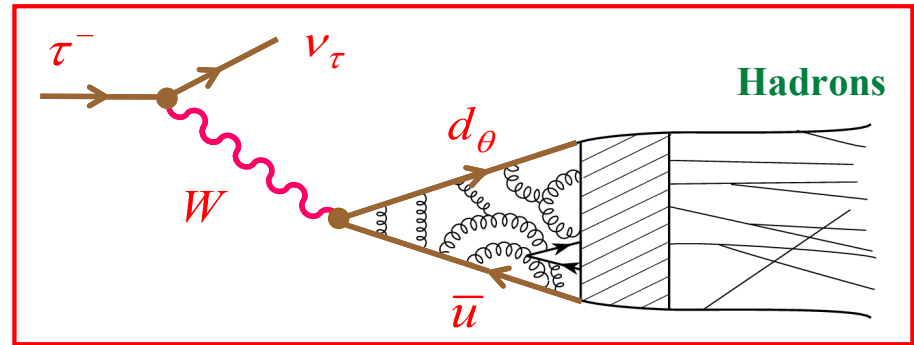
$$R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau + \text{Hadrons})}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e)} \approx N_C \quad ; \quad R_\tau = \frac{1 - B_e - B_\mu}{B_e} = 3.636 \pm 0.011$$

$$R_\tau = \frac{1}{B_e^{\text{univ}}} - 1.97256 = 3.6397 \pm 0.0076 \quad ; \quad R_\tau = \frac{\text{Br}(\tau^- \rightarrow \nu_\tau + \text{Hadrons})}{B_e^{\text{univ}}} = 3.6326 \pm 0.0084$$

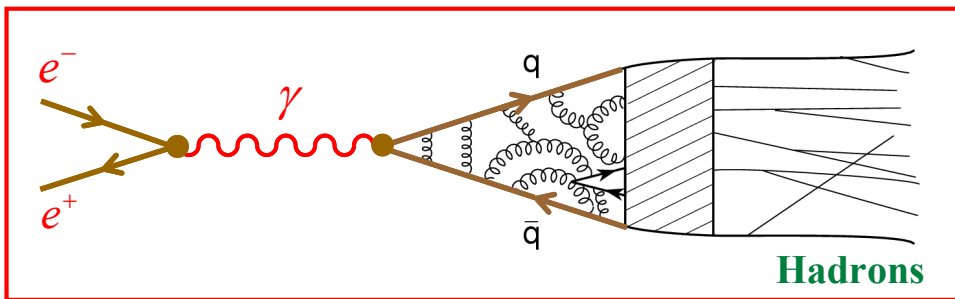
Only Lepton Massive Enough to Decay into Hadrons

$\tau^- \rightarrow \nu_\tau H^-$ probes the hadronic V-A current

$$\langle H^- | \bar{d}_\theta \gamma^\mu (1 - \gamma_5) u | 0 \rangle$$

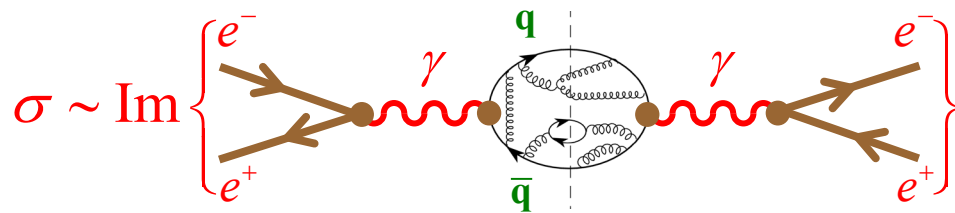


$e^+e^- \rightarrow H^0$ probes the hadronic electromagnetic current



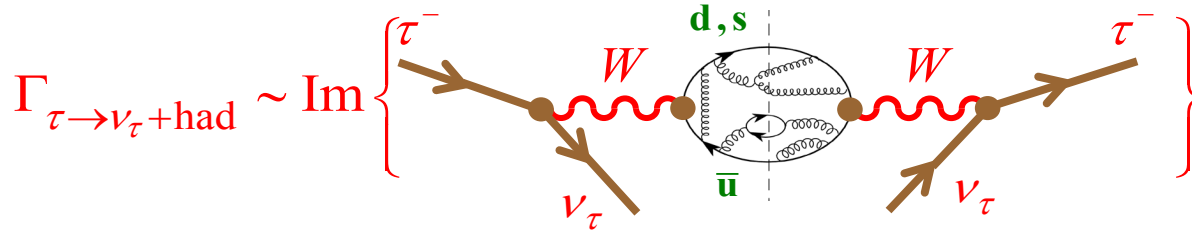
$$\langle H^0 | \sum_q Q_q \bar{q} \gamma^\mu q | 0 \rangle$$

Isospin:
$$\frac{\Gamma(\tau^- \rightarrow \nu_\tau V^-)}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e)} = \frac{3 \cos^2 \theta_C}{2\pi\alpha^2} S_{EW} \int_0^1 dx (1-x)^2 (1+2x) x \sigma_{e^+e^- \rightarrow V^0}(x m_\tau^2)$$



$$\frac{\sigma(e^+e^- \rightarrow \text{had})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 12\pi \text{Im} \Pi_{\text{em}}(s)$$

$$\Pi_{\text{em}}^{\mu\nu}(q) \equiv i \int d^4x e^{iqx} \langle 0 | T [J_{\text{em}}^\mu(x) J_{\text{em}}^\nu(0)] | 0 \rangle = (-g^{\mu\nu} q^2 + q^\mu q^\nu) \Pi_{\text{em}}(q^2)$$



$$R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau + \text{had})}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e)} = 12\pi \int_0^{m_\tau^2} \frac{ds}{m_\tau^2} \left(1 - \frac{s}{m_\tau^2}\right)^2 \left[\left(1 + 2\frac{s}{m_\tau^2}\right) \text{Im} \Pi^{(1)}(s) + \text{Im} \Pi^{(0)}(s) \right]$$

$$\Pi^{(J)}(s) \equiv |V_{ud}|^2 [\Pi_{ud,V}^{(J)}(s) + \Pi_{ud,A}^{(J)}(s)] + |V_{us}|^2 [\Pi_{us,V}^{(J)}(s) + \Pi_{us,A}^{(J)}(s)]$$

$$\Pi_{ij,J}^{\mu\nu}(q) \equiv i \int d^4x e^{iqx} \langle 0 | T [J_{ij}^\mu(x) J_{ij}^\nu(0)^\dagger] | 0 \rangle = (-g^{\mu\nu} q^2 + q^\mu q^\nu) \Pi_{ij,J}^{(1)}(q^2) + q^\mu q^\nu \Pi_{ij,J}^{(0)}(q^2)$$

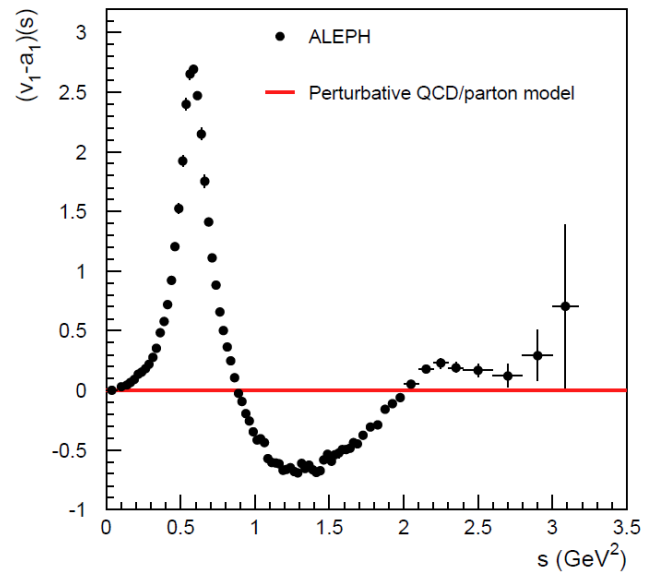
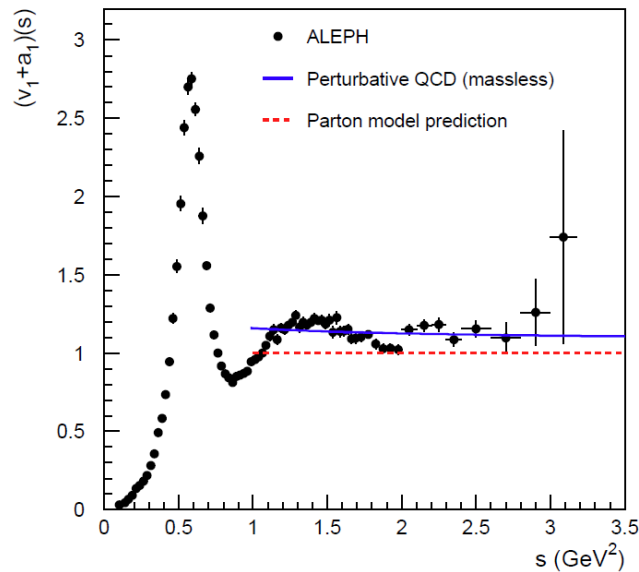
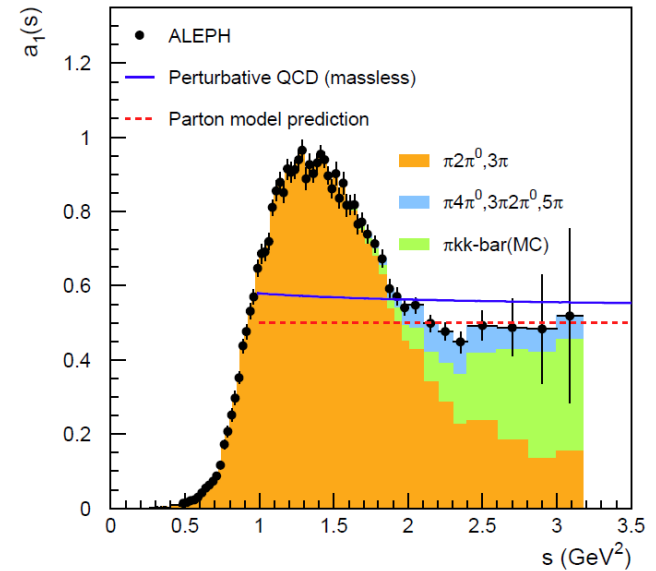
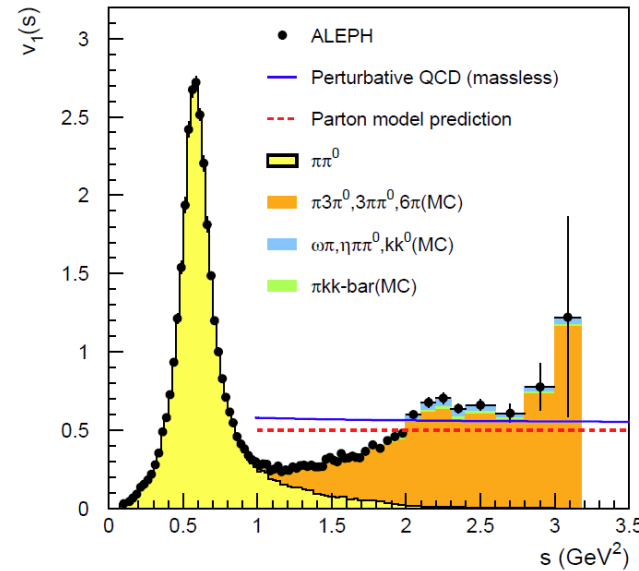
SPECTRAL FUNCTIONS

Davier et al, 1312.1501

$$v_1(s) = 2\pi \text{Im} \Pi_{ud,V}^{(0+)}(s)$$

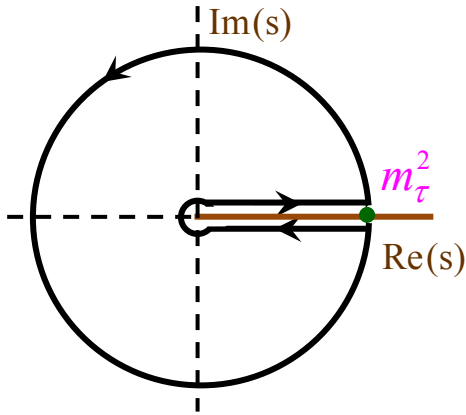
$$a_1(s) = 2\pi \text{Im} \Pi_{ud,A}^{(0+)}(s)$$

**BF data
needed**



$$R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau + \text{had})}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e)} = 12\pi \int_0^1 dx (1-x)^2 \left[(1+2x) \text{Im} \Pi^{(1)}(x m_\tau^2) + \text{Im} \Pi^{(0)}(x m_\tau^2) \right]$$

$$x \equiv s/m_\tau^2$$



$$R_\tau = 6\pi i \oint_{|x|=1} dx (1-x)^2 \left[(1+2x) \Pi^{(0+1)}(x m_\tau^2) - 2x \Pi^{(0)}(x m_\tau^2) \right]$$

$$\Pi^{(J)}(s) = \sum_{D=2n} \frac{C_D^{(J)}(s, \mu) \langle O_D(\mu) \rangle}{(-s)^{D/2}}$$

OPE

$$R_\tau = N_C S_{EW} (1 + \delta_P + \delta_{NP}) = R_{\tau,V} + R_{\tau,A} + R_{\tau,S}$$

$$S_{EW} = 1.0201 (3)$$

Marciano-Sirlin, Braaten-Li, Erler

$$\delta_{NP} = -0.0064 \pm 0.0013$$

Fitted from data (Davier et al)

$$\delta_P = a_\tau + 5.20 a_\tau^2 + 26 a_\tau^3 + 127 a_\tau^4 + \dots \approx 20\%$$

Baikov-Chetyrkin-Kühn

$$a_\tau \equiv \alpha_s(m_\tau) / \pi$$

Perturbative ($m_q=0$)

$$-s \frac{d}{ds} \Pi^{(0+1)}(s) = \frac{1}{4\pi^2} \sum_{n=0} K_n \left(\frac{\alpha_s(-s)}{\pi} \right)^n$$

$$K_0 = K_1 = 1, \quad K_2 = 1.63982, \quad K_3 = 6.37101, \quad K_4 = 49.07570$$

Baikov-Chetyrkin-Kühn '08

$$\longrightarrow \delta_P = \sum_{n=1} K_n A^{(n)}(\alpha_s) = a_\tau + 5.20 a_\tau^2 + 26 a_\tau^3 + 127 a_\tau^4 + \dots$$

Le Diberder-Pich '92

$$A^{(n)}(\alpha_s) \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} (1 - 2x + 2x^3 - x^4) \left(\frac{\alpha_s(-s)}{\pi} \right)^n = a_\tau^n + \dots \quad ; \quad a_\tau \equiv \alpha_s(m_\tau) / \pi$$

Power Corrections

$$\Pi_{\text{OPE}}^{(0+1)}(s) \approx \frac{1}{4\pi^2} \sum_{n \geq 2} \frac{C_{2n} \langle O_{2n} \rangle}{(-s)^n}$$

Braaten-Narison-Pich '92

$$C_4 \langle O_4 \rangle \approx \frac{2\pi}{3} \langle 0 | \alpha_s G^{\mu\nu} G_{\mu\nu} | 0 \rangle$$

$$\delta_{\text{NP}} \approx \frac{-1}{2\pi i} \oint_{|x|=1} dx (1 - 3x^2 + 2x^3) \sum_{n \geq 2} \frac{C_{2n} \langle O_{2n} \rangle}{(-x m_\tau^2)^n} = -3 \frac{C_6 \langle O_6 \rangle}{m_\tau^6} - 2 \frac{C_8 \langle O_8 \rangle}{m_\tau^8}$$

Suppressed by m_τ^6 [additional chiral suppression in $C_6 \langle O_6 \rangle^{V+A}$]

Perturbative Uncertainty on $\alpha_s(m_\tau)$

$$-s \frac{d}{ds} \Pi^{(0+1)}(s) = \frac{1}{4\pi^2} \sum_{n=0} K_n a(-s)^n$$

$$\delta_P = \underbrace{\sum_{n=1} K_n A^{(n)}(\alpha_s)}_{\text{CIPT}} = \underbrace{\sum_{n=0} r_n a_\tau^n}_{\text{FOPT}} \quad r_n = K_n + g_n$$

$$A^{(n)}(\alpha_s) \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} (1-2x+2x^3-x^4) a_\tau (-x m_\tau^2)^n = a_\tau^n + \dots \quad ; \quad a_\tau \equiv \alpha_s(m_\tau) / \pi$$

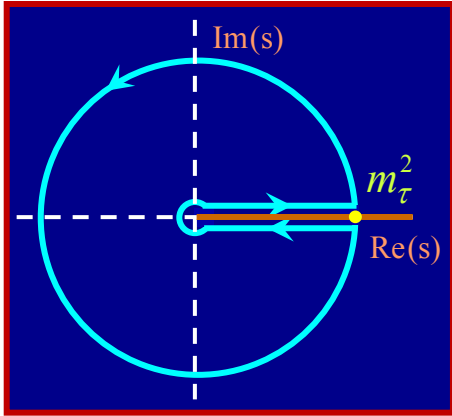
n	1	2	3	4	5
K_n	1	1.6398	6.3710	49.0757	
g_n	0	3.5625	19.9949	78.0029	307.78
r_n	1	5.2023	26.3659	127.079	

The dominant corrections come from the contour integration

Le Diberder- Pich 1992

Large running of a_s along the circle $s = m_\tau^2 e^{i\varphi}$, $\varphi \in [0, 2\pi]$

$$A^{(n)}(a_\tau) \equiv \frac{1}{2\pi i} \oint_{|x|=1} \frac{dx}{x} (1-2x+2x^3-x^4) a_\tau (-x m_\tau^2)^n = a_\tau^n + \dots \quad ; \quad a_\tau \equiv \alpha_s(m_\tau)/\pi$$



$$A^{(1)}(a_\tau) = a_\tau - \frac{19}{24} \beta_1 a_\tau^2 + \left[\beta_1^2 \left(\frac{265}{288} - \frac{\pi^2}{12} \right) - \frac{19}{24} \beta_2 \right] a_\tau^3 + \dots$$

$$a(-s) \simeq \frac{a_\tau}{1 - \frac{\beta_1}{2} a_\tau \log(-s/m_\tau^2)} = \frac{a_\tau}{1 - i \frac{\beta_1}{2} a_\tau \phi} = a_\tau \sum_n \left(i \frac{\beta_1}{2} a_\tau \phi \right)^n \quad ; \quad \phi \in [0, 2\pi]$$

FOPT expansion only convergent if $\alpha_\tau < 0.14$ (0.11) [at 1 (3) loops]

Experimentally $\alpha_\tau \approx 0.11$ \longrightarrow **FOPT should not be used**
(divergent series)

FOPT suffers a large renormalization-scale dependence (Le Diberder- Pich , Menke)

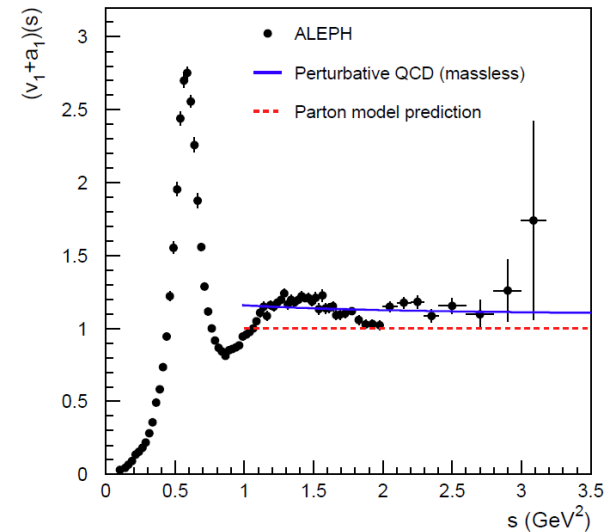
The difference between FOPT and CIPT grows at higher orders

Spectral Function Distribution

■ Moments:

$$R_{\tau}^{kl}(s_0) \equiv \int_0^{s_0} ds \left(1 - \frac{s}{s_0}\right)^k \left(\frac{s}{m_{\tau}^2}\right)^l \frac{dR_{\tau}}{ds}$$

Sensitivity to power corrections (k,l)



The non-perturbative contribution to R_{τ} can be obtained from the invariant-mass distribution of the final hadrons

$$\delta_{\text{NP}} = -0.0064 \pm 0.0013$$

Davier et al. (ALEPH data)

■ Fitting the Spectral Function itself: Duality Violations

$$\text{Im } \Pi(s) = \kappa e^{-\gamma s} \sin[\alpha + \beta s]$$

$$\delta_{\text{NP}} = -0.003 \pm 0.012$$

Boito et al. (OPAL data)

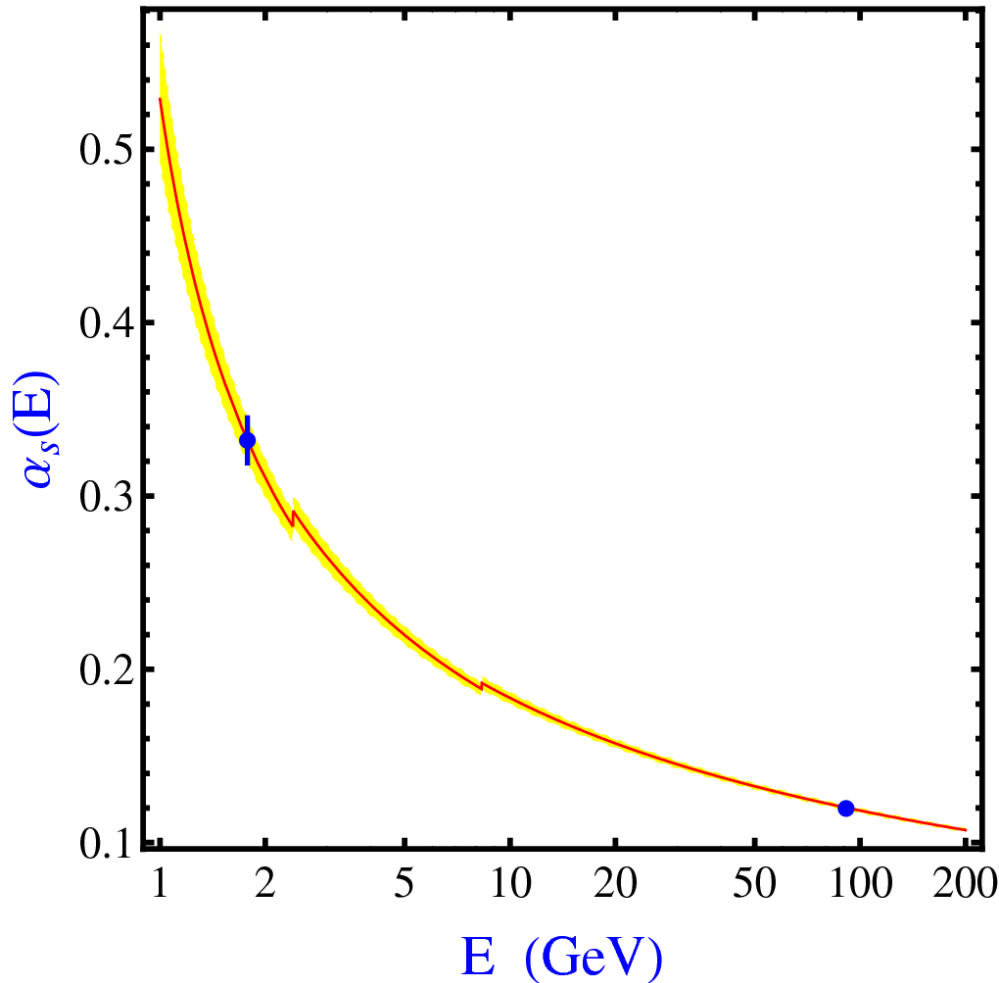
Recent $\alpha_s(m_\tau)$ Analyses

Reference	Method	δ_{NP}	δ_P	$\alpha_s(m_\tau)$	$\alpha_s(m_Z)$
Baikov et al	CIPT, FOPT		0.1998 (43)	0.332 (16)	0.1202 (19)
Davier et al'14	CIPT, FOPT	- 0.0064 (13)	-	0.332 (12)	0.1199 (15)
Beneke-Jamin	BSR + FOPT	- 0.007 (3)	0.2042 (50)	0.316 (06)	0.1180 (08)
Maltman-Yavin	PWM + CIPT	+ 0.012 (18)	-	0.321 (13)	0.1187 (16)
Menke	CIPT, FOPT		0.2042 (50)	0.342 (11)	0.1213 (12)
Narison	CIPT, FOPT		-	0.324 (08)	0.1192 (10)
Caprini-Fischer	BSR + CIPT		0.2037 (54)	0.322 (16)	-
Abbas et al	IFOPT		0.2037 (54)	0.338 (10)	
Cvetič et al	β_{exp} + CIPT		0.2040 (40)	0.341 (08)	0.1211 (10)
Boito et al	CIPT, DV	- 0.002 (12)	-	0.347 (25)	0.1216 (27)
	FOPT, DV	- 0.004 (12)		0.325 (18)	0.1191 (22)
Pich'14	CIPT	- 0.0064 (13)	0.2014 (31)	0.342 (13)	0.1213 (14)
	FOPT			0.320 (14)	0.1187 (17)
Pich'14	CIPT, FOPT	- 0.0064 (13)	0.2014 (31)	0.332 (13)	0.1202 (15)

CIPT: Contour-improved perturbation theory
 FOPT: Fixed-order perturbation theory
 BSR: Borel summation of renormalon series
 IFOPT: Improved FOPT

β_{exp} : Expansion in derivatives of α_s (β function)
 PWM: Pinched-weight moments
 CIPTm: Modified CIPT (conformal mapping)
 DV: Duality violation (OPAL only)

Present Status



$$\alpha_s(m_\tau^2) = 0.332 \pm 0.013$$



$$\alpha_s(M_Z^2) = 0.1202 \pm 0.0015$$

$$\alpha_s(M_Z^2)_{Z_{\text{width}}} = 0.1197 \pm 0.0028$$

**The most precise test of
Asymptotic Freedom**

$$\alpha_s^\tau(M_Z^2) - \alpha_s^Z(M_Z^2) = 0.0005 \pm 0.0015_\tau \pm 0.0028_Z$$

V_{us} Determination

$$|V_{us}|^2 = \frac{R_{\tau,S}^{00}}{\frac{R_{\tau,V+A}^{00}}{|V_{ud}|^2} - \delta R_{\tau,\text{th}}^{00}}$$

$$\delta R_{\tau}^{kl} \approx 24 \frac{m_S^2(m_{\tau}^2)}{m_{\tau}^2} \Delta_{kl}(\alpha_S)$$

$$\delta R_{\tau,\text{th}}^{00} \equiv \underbrace{0.1544 (37)}_{J=0} + \underbrace{0.086 (32)}_{m_S(2 \text{ GeV}) = 94 (6) \text{ MeV}} = 0.240 (32)$$

$$R_{\tau,S}^{00} = 0.1614 (28)$$

$$R_{\tau,V+A}^{00} = 3.4712 (79)$$

$$|V_{ud}| = 0.97425 (22)$$



$$|V_{us}| = 0.2173 \pm 0.0020_{\text{exp}} \pm 0.0010_{\text{th}}$$

$$R_{\tau,S}^{00} = 0.1665 (34)$$



$$|V_{us}| = \mathbf{0.2207 \pm 0.0023_{\text{exp}} \pm 0.0011_{\text{th}}}$$

$$\mathbf{K_{13}}: |V_{us}| = 0.2239 \pm 0.0009 \quad [f_+(0) = 0.9661 \pm 0.0032]$$

The τ could give the most precise V_{us} determination

Cabibbo-suppressed τ Decays

HFAG

X_S^-	$\text{Br}(\tau^- \rightarrow \nu_\tau X_S^-)$ (%)	X_S^-	$\text{Br}(\tau^- \rightarrow \nu_\tau X_S^-)$ (%)
K^-	(0.6955 ± 0.0096)	$K^- \eta$	(0.0153 ± 0.0008)
$K^- \pi^0$	(0.4322 ± 0.0149)	$K^- \pi^0 \eta$	(0.0048 ± 0.0012)
$K^- 2\pi^0$ (ex. K^0)	(0.0630 ± 0.0222)	$\pi^- \bar{K}^0 \eta$	(0.0094 ± 0.0015)
$K^- 3\pi^0$ (ex. K^0, η)	(0.0419 ± 0.0218)	$K^- \omega$	(0.0410 ± 0.0092)
$\pi^- \bar{K}^0$	(0.8206 ± 0.0182)	$K^- \phi$ ($\phi \rightarrow K\bar{K}$)	(0.0037 ± 0.0014)
$\pi^- \bar{K}^0 \pi^0$	(0.3649 ± 0.0108)	$K^- \pi^- \pi^+$ (ex. K^0, ω)	(0.2923 ± 0.0068)
$\pi^- \bar{K}^0 2\pi^0$	(0.0269 ± 0.0230)	$K^- \pi^- \pi^+ \pi^0$ (ex. K^0, ω, η)	(0.0411 ± 0.0143)
$\bar{K}^0 h^- h^+ h^-$	(0.0222 ± 0.0202)		
$\text{Br}(\tau^- \rightarrow \nu_\tau X_S^-) = (2.875 \pm 0.050)\%$			\rightarrow $(2.967 \pm 0.060)\%$

Unaccounted modes:

$$1 - \sum_i \text{Br}(\tau^- \rightarrow \nu_\tau^- X_i) = (0.0704 \pm 0.1060)\%$$

Antonelli-Cirigliano-Lusiani-Passemar:

$$\Gamma(K \rightarrow \mu \nu_\mu)$$

$$K_{13} + \tau \rightarrow \nu_\tau (K\pi)^- \text{ spectra}$$



$$\text{Br}(\tau^- \rightarrow \nu_\tau K^-) = (0.713 \pm 0.003)\%$$

$$\text{Br}(\tau^- \rightarrow \nu_\tau K^- \pi^0) = (0.471 \pm 0.018)\%$$

$$\text{Br}(\tau^- \rightarrow \nu_\tau \bar{K}^0 \pi^-) = (0.857 \pm 0.030)\%$$

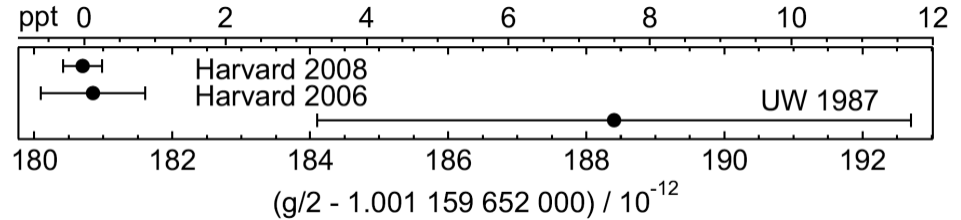


$$R_{\tau,S}^{00} = 0.1665 (34)$$

Electron Anomalous Magnetic Moment

Hanneke-Fogwell-Gabrielse '08

$$a_e = 0.001\,159\,652\,180\,73\ (28)$$



$$a_e^{\text{QED}} = \sum_{n=1} \left(\frac{\alpha}{\pi} \right)^n a_e^{(2n)}$$

$$a_e^{(2n)} = A_1^{(2n)} + A_2^{(2n)}(m_e/m_\mu) + A_2^{(2n)}(m_e/m_\tau) + A_3^{(2n)}(m_e/m_\mu, m_e/m_\tau)$$

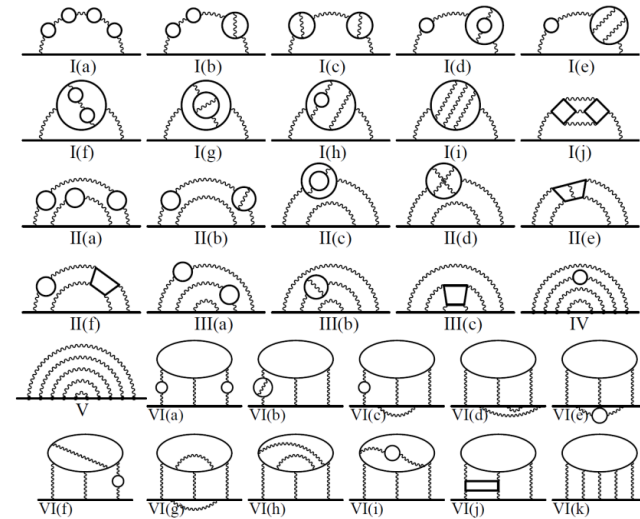
Aoyama-Hayakawa-Kinoshita-Nio

$$A_1^{(8)} = -1.9106\ (20) \quad , \quad A_2^{(8)}(m_e/m_\mu) = 9.222\ (66) \cdot 10^{-4}$$

$$A_2^{(8)}(m_e/m_\tau) = 8.24\ (12) \cdot 10^{-6} \quad , \quad A_3^{(8)}(m_e/m_\mu, m_e/m_\tau) = 7.465\ (18) \cdot 10^{-7}$$

$$A_1^{(10)} = 9.16\ (58) \quad , \quad A_2^{(10)}(m_e/m_\mu) = -3.82\ (39) \cdot 10^{-3}$$

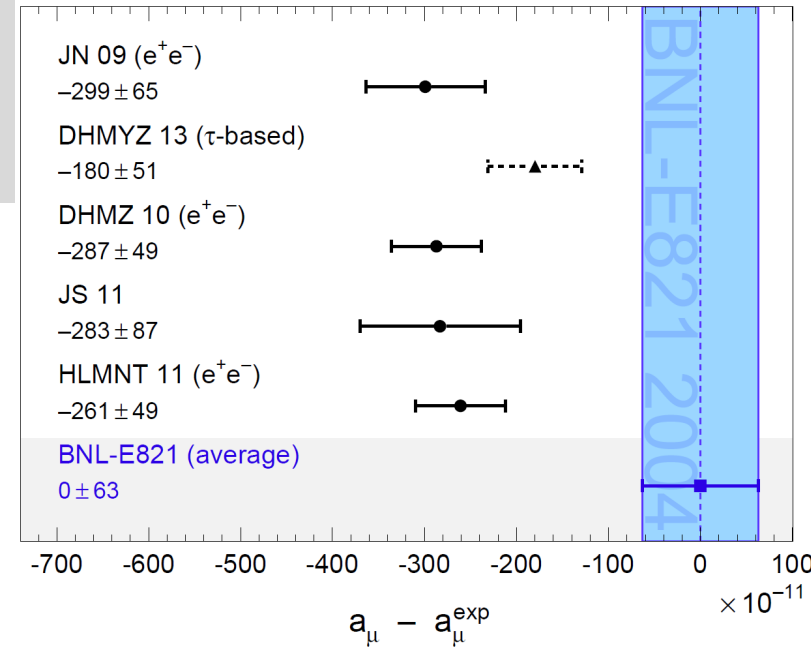
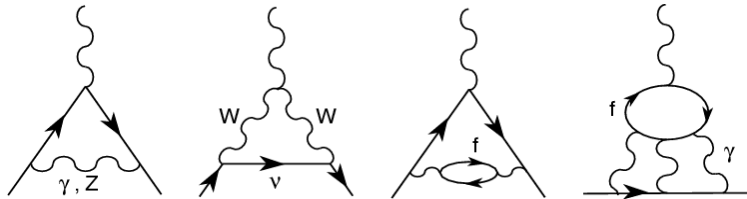
$$a_e^{\text{QCD}} = 1.685\ (33) \times 10^{-12} \quad , \quad a_e^{\text{Weak}} = 0.0297\ (5) \times 10^{-12}$$



$$\alpha_{\text{Rb2010}}^{-1} = 137.035\,999\,049\ (90) \quad \longrightarrow \quad a_e^{\text{th}} = 0.001\,159\,652\,181\,78\ (6)_{8\text{th}}\ (4)_{10\text{th}}\ (3)_{\text{had}}\ (77)_{\alpha}$$

$$a_e \quad \longrightarrow \quad \alpha_{a_e}^{-1} = 137.035\,999\,173\,6\ (68)_{8\text{th}}\ (46)_{10\text{th}}\ (26)_{\text{had}}\ (331)_{\text{exp}} \quad [0.25\ \text{ppb}]$$

μ Anomalous Magnetic Moment



$$a_\mu^{\text{exp}} = (11\,659\,208.9 \pm 6.3) \cdot 10^{-10}$$

BNL-E821

$$\begin{aligned}
 10^{10} \cdot a_\mu^{\text{th}} &= 11\,658\,471.895 \pm 0.008 && \text{QED} && \text{Aoyama-Hayakawa-Kinoshita-Nio} \\
 &+ 15.4 \pm 0.1 && \text{EW} && \text{Gnendiger et al, Czarnecki et al, Knecht et al} \\
 &+ 697.4 \pm 5.3 && \text{hvp} && (703.0 \pm 4.4)_\tau, (692.3 \pm 4.2)_{e^+e^-} \quad \text{Davier et al,} \\
 & && && \text{Hagiwara et al, Jegerlehner-Nyffeler} \\
 &- 8.6 \pm 0.1 && \text{hvp NLO+NNLO} && \text{Kurz et al, Hagiwara et al, Krause} \\
 &+ 10.5 \pm 2.6 && \text{light-by-light} && \text{de Rafael-Prades-Vainshtein, Knecht et al,} \\
 & && && \text{Melnikov-Vainshtein, Nyffeler, Bijnens et al,} \\
 & && && \text{Hayakawa et al, Goecke et al, Roig et al,} \\
 & && && \text{Masjuan-Vanderhaeghen} \\
 \\
 &= 11\,659\,186.6 \pm 5.9 && (11\,659\,192.2 \pm 5.1)_\tau, (11\,659\,181.5 \pm 4.9)_{e^+e^-} \\
 a_\mu^{\text{exp}} - a_\mu^{\text{th}} &= 2.7 \sigma && 1.9 \sigma && 3.4 \sigma
 \end{aligned}$$

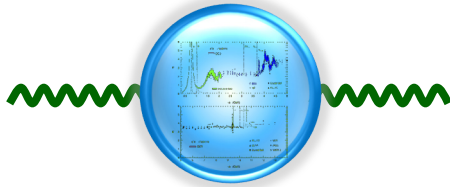
Future Challenge

$$\Delta a_\mu = 1.6 \cdot 10^{-10}$$

(0.14 ppm)

FNAL, J-PARC

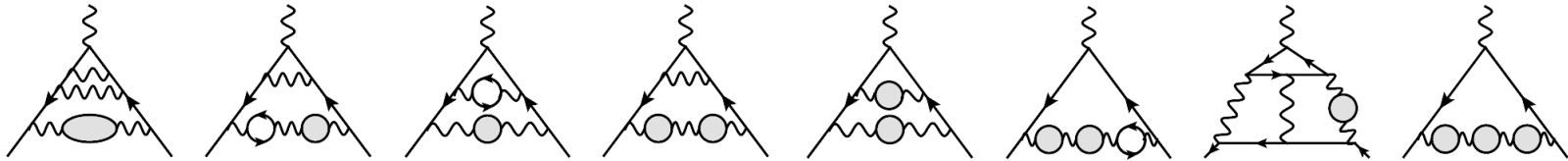
□ Hadronic Vacuum Polarization



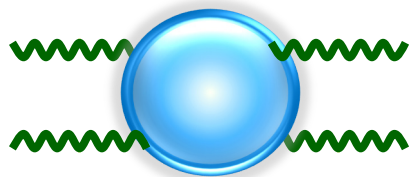
- Improved data
- Radiative return
- Isospin breaking
- Improved theoretical tools
- Lattice simulations

□ NNLO HVP

Kurz et al



□ Light-by-Light



- Lattice Blum et al
- Dispersive approach Colangelo et al, Pauk-Vanderhaeghen
- Analytical methods de Rafael-Prades-Vainshtein, Knecht et al,
Melnikov-Vainshtein, Nyffeler, Bijmans et al, Hayakawa et al, Goecke et al, Roig et al, Masjuan-Vanderhaeghen

τ Anomalous Magnetic Moment

Difficult to measure!

$$a_{\tau}^{\text{exp}} = (-0.018 \pm 0.017)$$

DELPHI

$$-0.007 < a_{\tau}^{\text{New Phys}} < 0.005$$

González-Springer, Santamaria, Vidal '00 (LEP/SLD data)

Eidelman, Passera

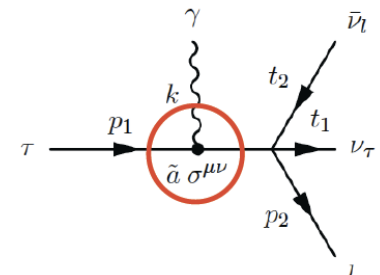
$$\begin{aligned}
 10^8 \cdot a_{\tau}^{\text{th}} &= 117\,324 \pm 2 && \text{QED} \\
 &+ 47.4 \pm 0.5 && \text{EW} \\
 &+ 337.5 \pm 3.7 && \text{hvp} \\
 &+ 7.6 \pm 0.2 && \text{hvp NLO} \\
 &+ 5 \pm 3 && \text{light-by-light} \\
 &= \mathbf{117\,721 \pm 5}
 \end{aligned}$$

Enhanced sensitivity to new physics: $(m_{\tau}/m_{\mu})^2 = 283$

	Electron	Muon	Tau
$a^{\text{EW}}/a^{\text{HAD}}$	1/56	1/45	1/7
$a^{\text{EW}}/\delta a^{\text{HAD}}$	1.6	3	10

Essentially unknown

May be accessible at BFs through radiative leptonic decays (Fael et al)



SUMMARY



- Very precise determination of α_s from τ decays

$$\alpha_s(m_\tau^2) = 0.332 \pm 0.013 \quad \longrightarrow \quad \alpha_s(M_Z^2) = 0.1202 \pm 0.0015$$

- The τ could give the most precise V_{us} determination

$$|V_{us}| = 0.2207 \pm 0.0023_{\text{exp}} \pm 0.0011_{\text{th}} \quad (\text{present } \tau \text{ data} + \text{K data})$$

$$\mathbf{K}_{13}: \quad |V_{us}| = 0.2239 \pm 0.0009 \quad [f_+(0) = 0.9661 \pm 0.0032]$$

- Many QCD tests with τ data: **Chiral Dynamics**

- Improved predictions for a_μ^{QCD} needed

Theoretical and experimental challenge

Backup Slides

Renormalons

$$D(s) \equiv -s \frac{d}{ds} \Pi^{(0+1)}(s) = \frac{1}{4\pi^2} \sum_{n=0} K_n a(-s)^n$$

Asymptotic series

Borel Summation:

$$B(t) \equiv \sum_{n=0} K_{n+1} \frac{t^n}{n!} \quad \longrightarrow \quad D(s) = \frac{1}{4\pi^2} \left\{ 1 + \int_0^\infty dt e^{-t/a(-s)} B(t) \right\}$$

However, $B(t)$ has pole singularities at

- $u \equiv -\beta_1 t/2 = +n \quad (n \geq 2)$

Infrared Renormalons

- $u \equiv -\beta_1 t/2 = -n \quad (n \geq 1)$

Ultraviolet Renormalons

IR - n Renormalon





Ambiguity:

$$\delta D(s) \sim \left(\frac{\Lambda^2}{-s} \right)^n$$

Renormalon Hypothesis: Asymptotics already reached

Modelling a better behaved FOPT

(Beneke – Jamin)

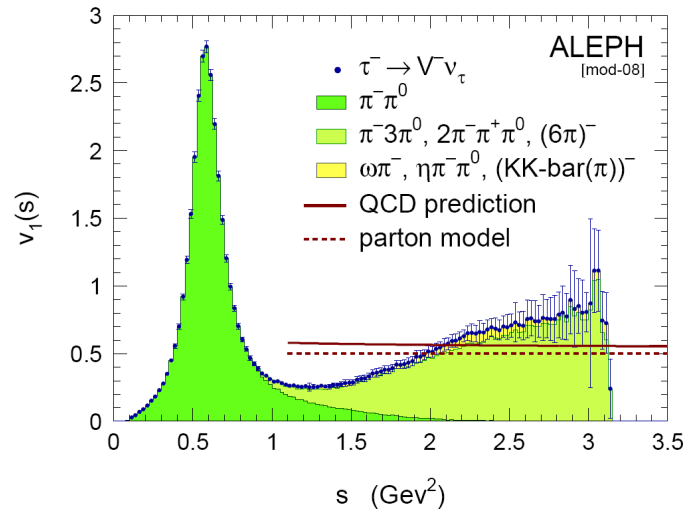
- Large higher-order K_n corrections could cancel the g_n ones
Happens in the “large- β_0 ” approximation (UV renormalon chain)
- $D = 4$ corrections very suppressed in R_τ
 $n = 2$ IR renormalons can do the job ($K_n \approx -g_n$)
- No sign of renormalon behaviour in known coefficients
 $n = -1, 2, 3$ renormalons + linear polynomial
5 unknown constants fitted to K_n ($2 \leq n \leq 5$). $K_5 = 283$ assumed
- **Borel summation:** large renormalon contributions. Smaller α

Nice model of higher orders. But too many different possibilities ...

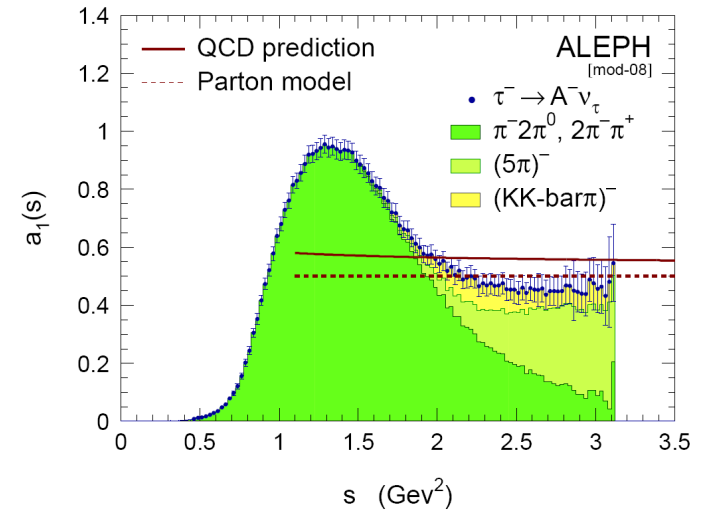
(Descotes-Genon – Malaescu)

SPECTRAL FUNCTIONS

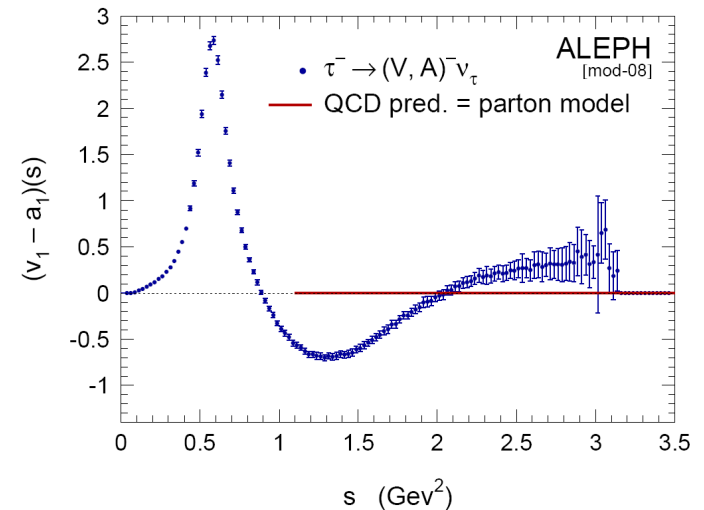
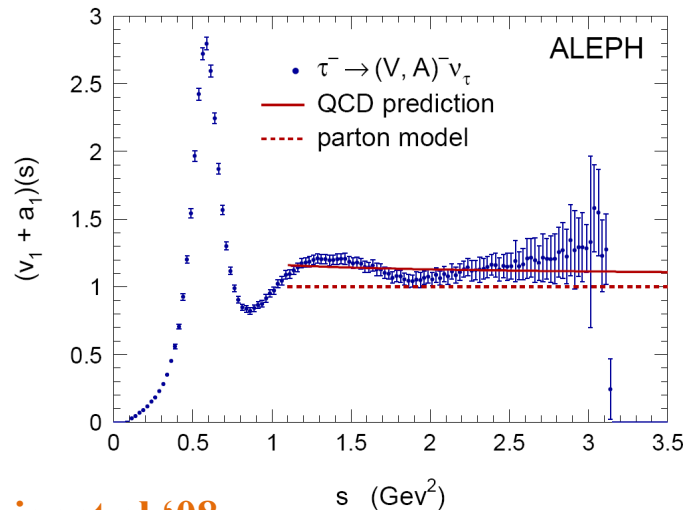
$$v_1(s) = 2\pi \text{Im} \Pi_{ud,V}^{(0+1)}(s)$$



$$a_1(s) = 2\pi \text{Im} \Pi_{ud,A}^{(0+1)}(s)$$



BF data needed



Davier et al '08