



UCL

Status of the AlCap Experiment

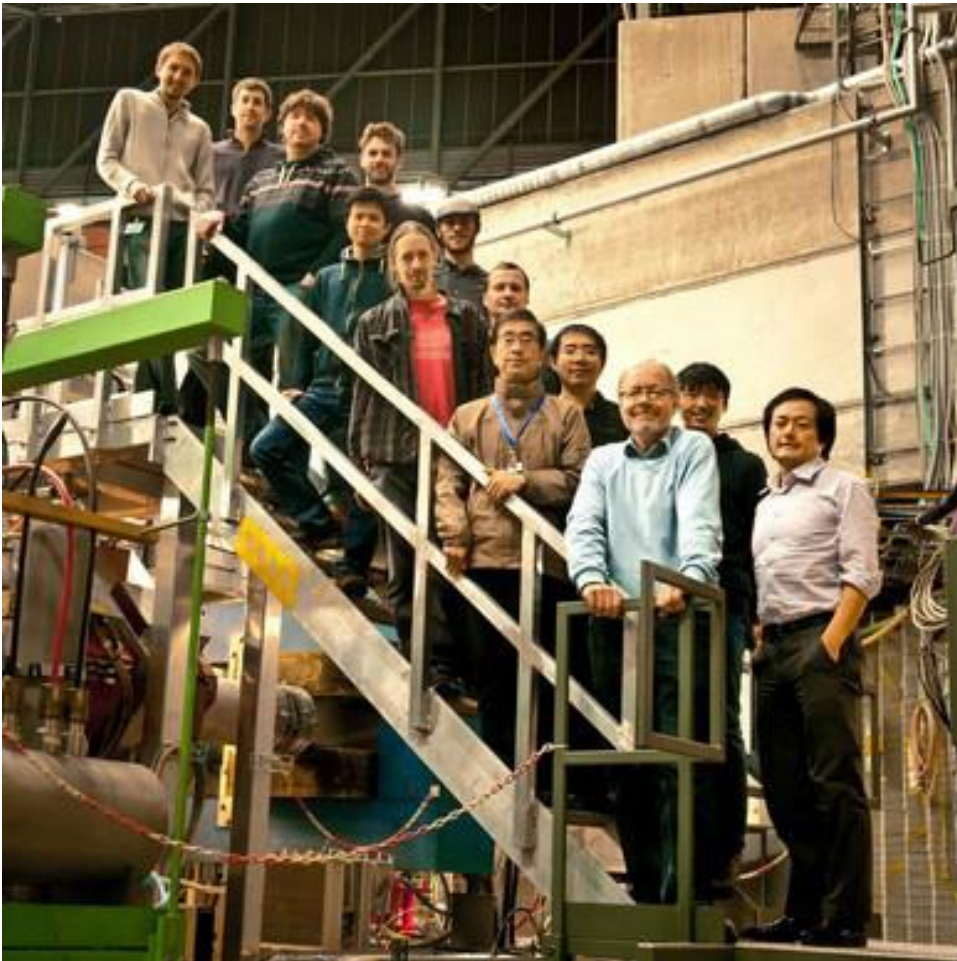
**Phillip Litchfield,
UCL**

1. **Who** is AlCap?
M110
2. **Why** do AlCap?
M11A
3. **What** is AlCap?
M11SF
4. **Where** is AlCap?
M11G1G
5. **How** does AlCap work?
E10A
6. **What next** for AlCap?
M11SF U11G1G

The AlCap Collaboration

from COMET
Osaka University,
IHEP China,
Imperial College London,
University College London,

from Mu2e
Argonne NL,
Boston University,
Brookhaven NL,
University of Houston,
University of Washington,



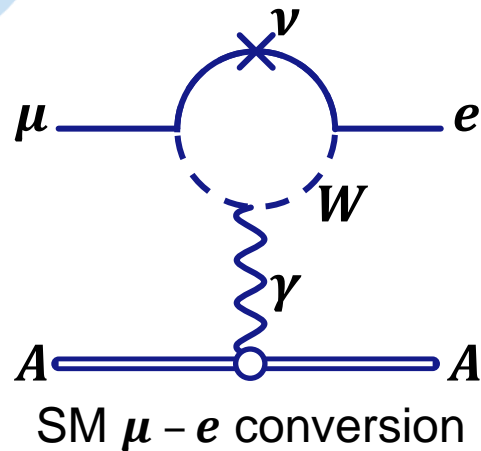
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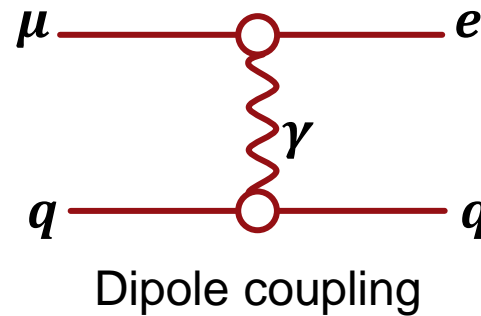
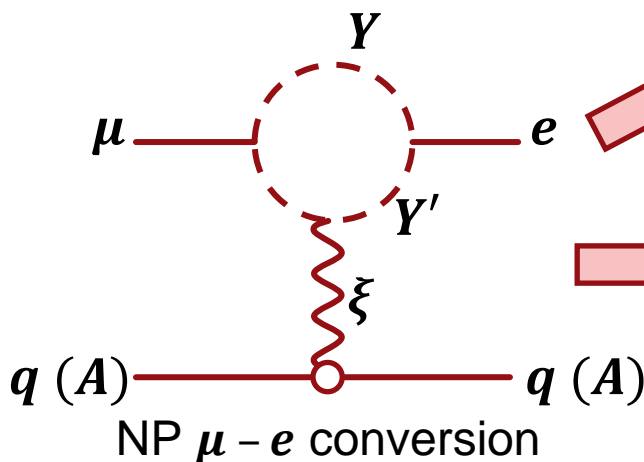


μ to e conversion

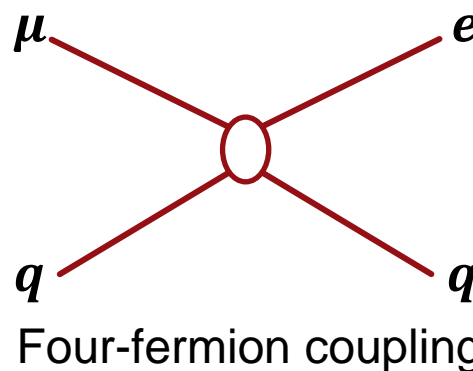


In the **SM** $\mu A \rightarrow e A$ is heavily suppressed because of the mass disparity between the W and neutrino.

In **new physics** scenarios this does not usually apply, and other diagrams typically give CLFV much higher than the SM.

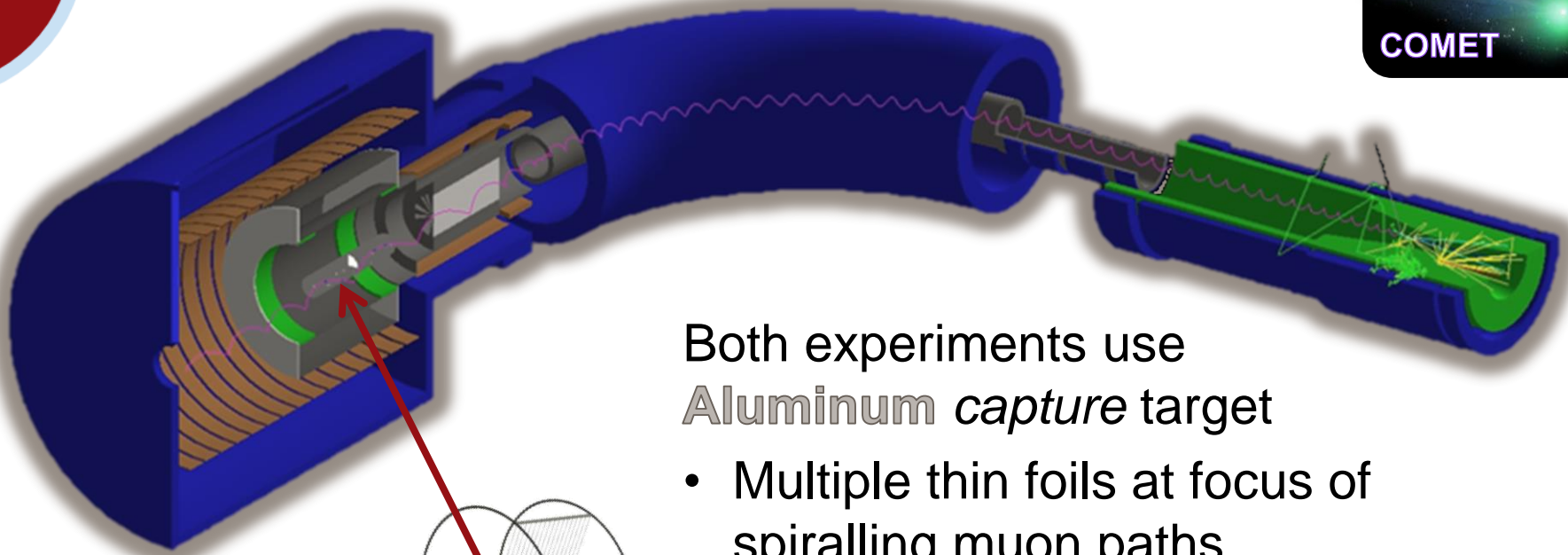


$$\mathcal{L}_d \sim \frac{m_\mu}{\Lambda^2} \bar{\mu} \sigma_{\mu\nu} e \cdot F^{\mu\nu}$$



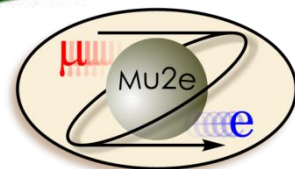
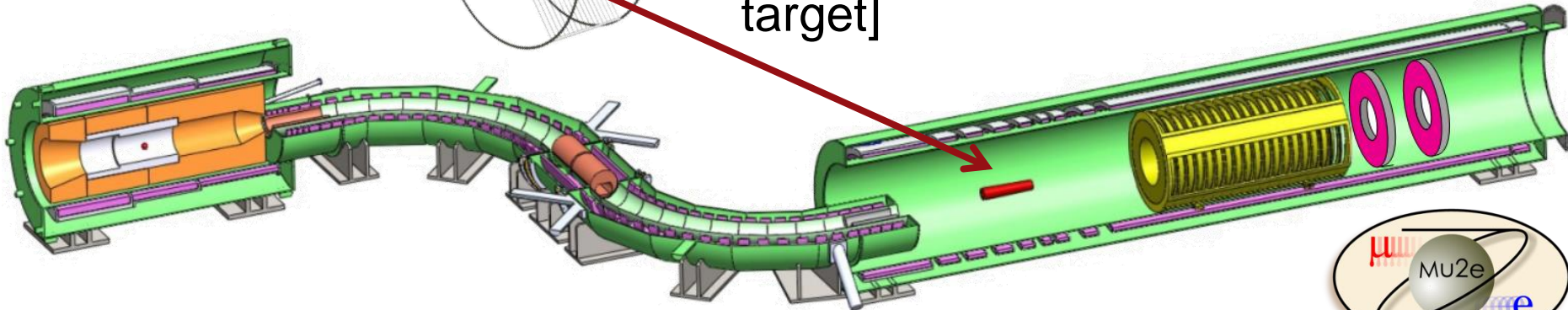
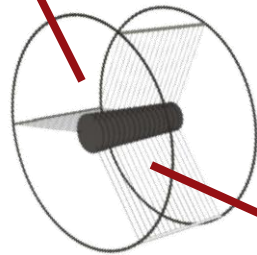
$$\mathcal{L}_4 \sim \frac{1}{\Lambda^2} \bar{\mu} \gamma_\mu e \cdot \bar{q} \gamma_\mu q$$

COMET Phase-I and Mu2e

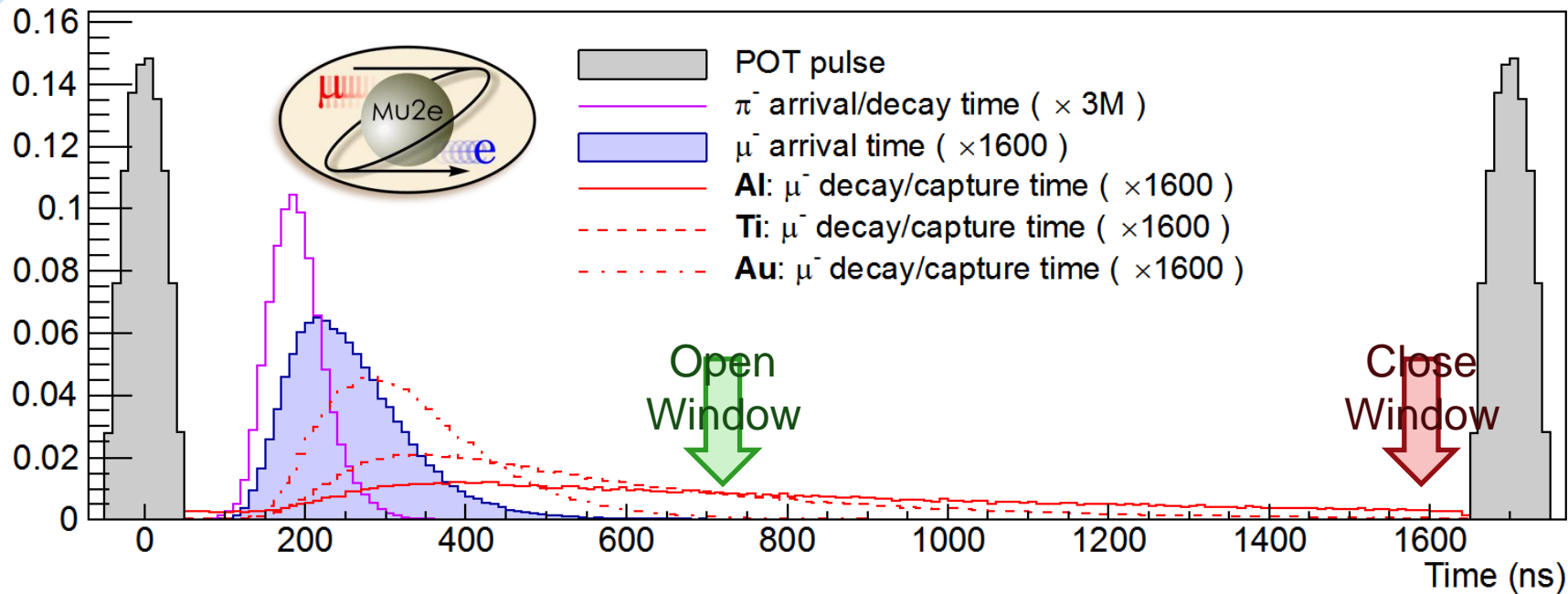


Both experiments use *Aluminum capture target*

- Multiple thin foils at focus of spiralling muon paths
- [Not the high-power production target]



Why aluminium?



High Z materials are good, as it increases capture cross section

- Previous experiment (SINDRUM-II) used gold.

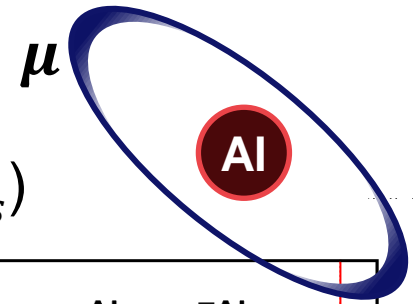
But for prompt BG rejection the new experiments will use pulsed beam and a late ($\geq 700\text{ns}$) window.

Therefore need a lighter (i.e. low Z) target material for longer muon lifetimes. \rightarrow aluminium

Muon capture on Aluminium

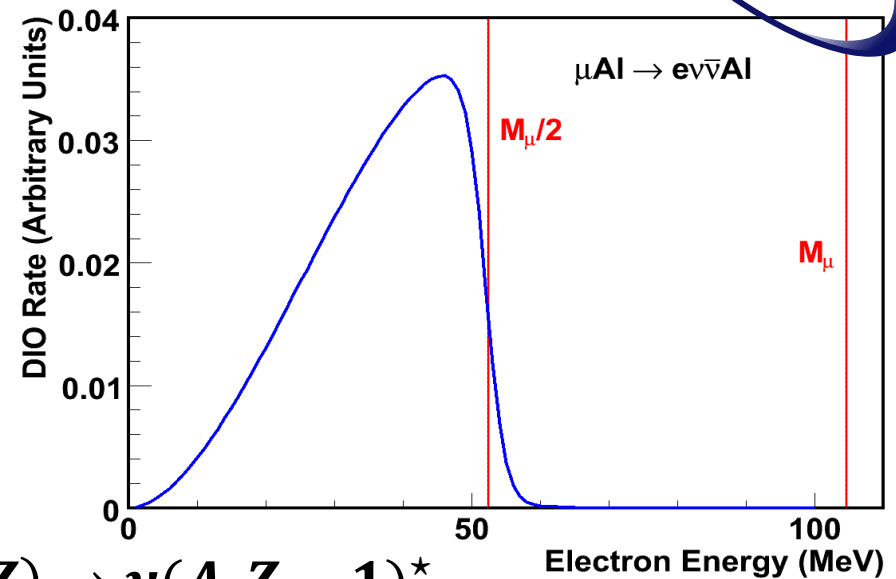
Muons allowed stop in the Aluminium target

- $\mu A \rightarrow eA$ conversion from 1s orbital gives a **mono-energetic electron** at 105MeV ($\approx m_\mu - B_{1s}^\mu$)



‘Normal’ decays are backgrounds:

- Decay in Orbit [DIO]: $\mu \rightarrow e\nu\bar{\nu}$
For a free muon, cuts off at $\frac{1}{2}m_\mu$, but bound state has a small tail up to m_μ



- **Nuclear muon capture:** $\mu(A, Z) \rightarrow \nu(A, Z - 1)^*$
 - The resulting nucleus is in an excited state.
 - **De-excitation can emit γ, p, n, \dots** (i.e. pretty much anything)
 - Uh-oh...

Charged particle emission after muon capture.

- Protons are a major component of the single-hit rates in the tracking chambers for both Mu2e and COMET Phase-I.
- Measure both the total rate and the energy spectrum

Gamma and X-ray emission after muon capture.

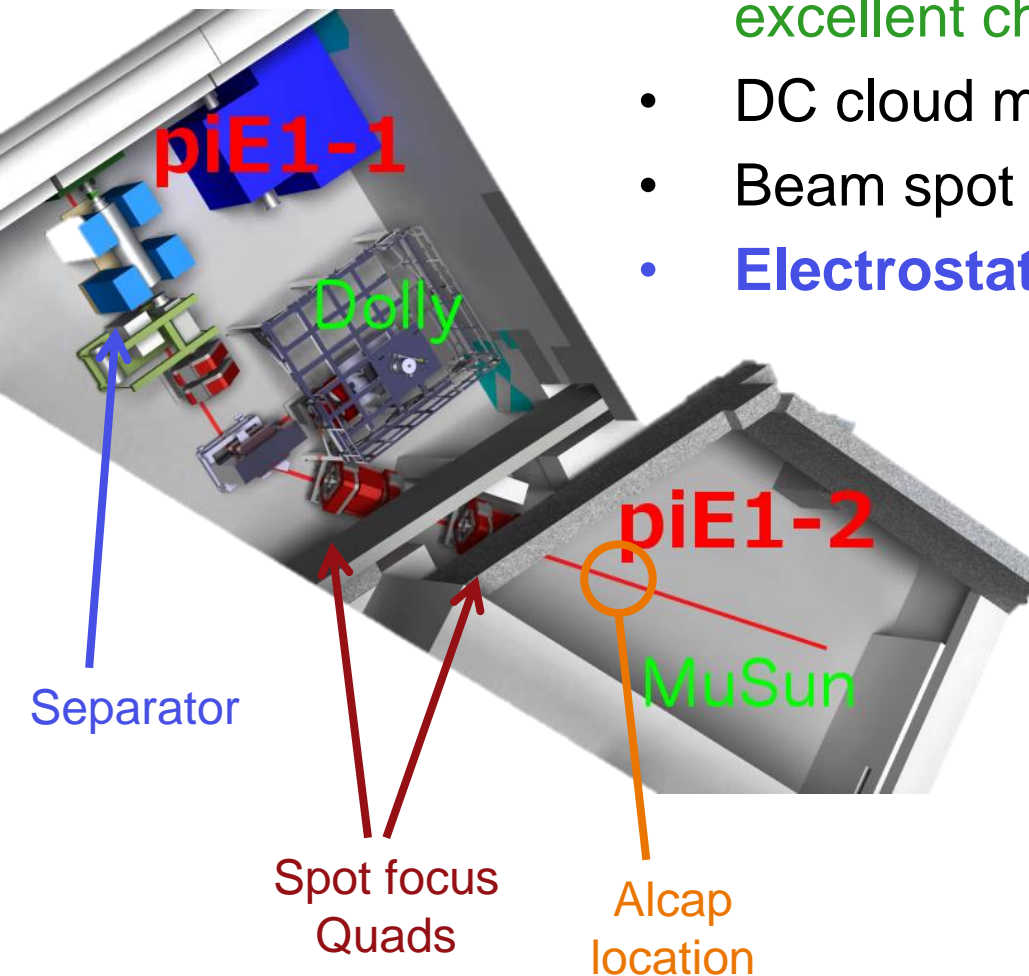
- Measure X-rays from the muonic atomic cascade, in order to provide the muon-capture normalization.
 - Normalise the charged particle rate measurement
 - Verify method for Mu2e and COMET experiments.

Neutron emission after muon capture.

- Important for determining backgrounds in the Mu2e/COMET detectors and evaluating the radiation damage to electronic components. Also may affect layout of CR veto counters.

The 2013 run used the $\pi E1$ beamline at PSI:

- New beamline, developed by MuSun, with **excellent characteristics** at low energy.
- DC cloud muon beam
- Beam spot diameter **< 2cm**
- **Electrostatic separator** reduces e^- contamination **< 10%**



AlCap configuration:

- Momentum \sim **28MeV/c**
- Momentum bite \sim **1%**
- Mainly in **-ive** polarity.
(μ capture expt!)

AlCap muon beam

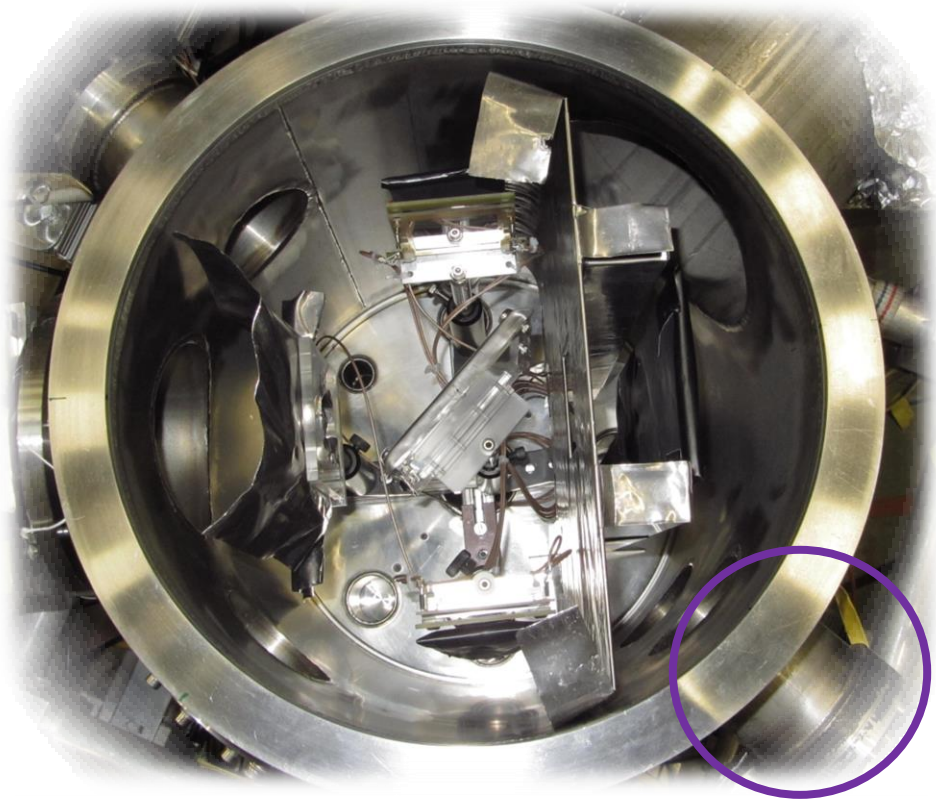
Most of the experiment is housed in a (~ 30cm diameter) vacuum chamber.

- Pump attached to **side port** to reduce pressure to below 10^{-4} mbar to prevent sparking and reduce energy loss in flight.

The **beam** momentum is tuned near 28 MeV/c so that muons stop in the target placed at the centre of the chamber.



Mylar beam window is of standard PSI design



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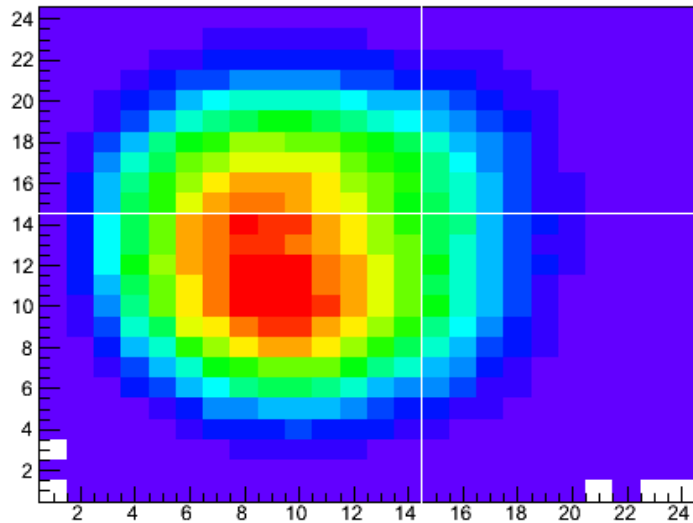


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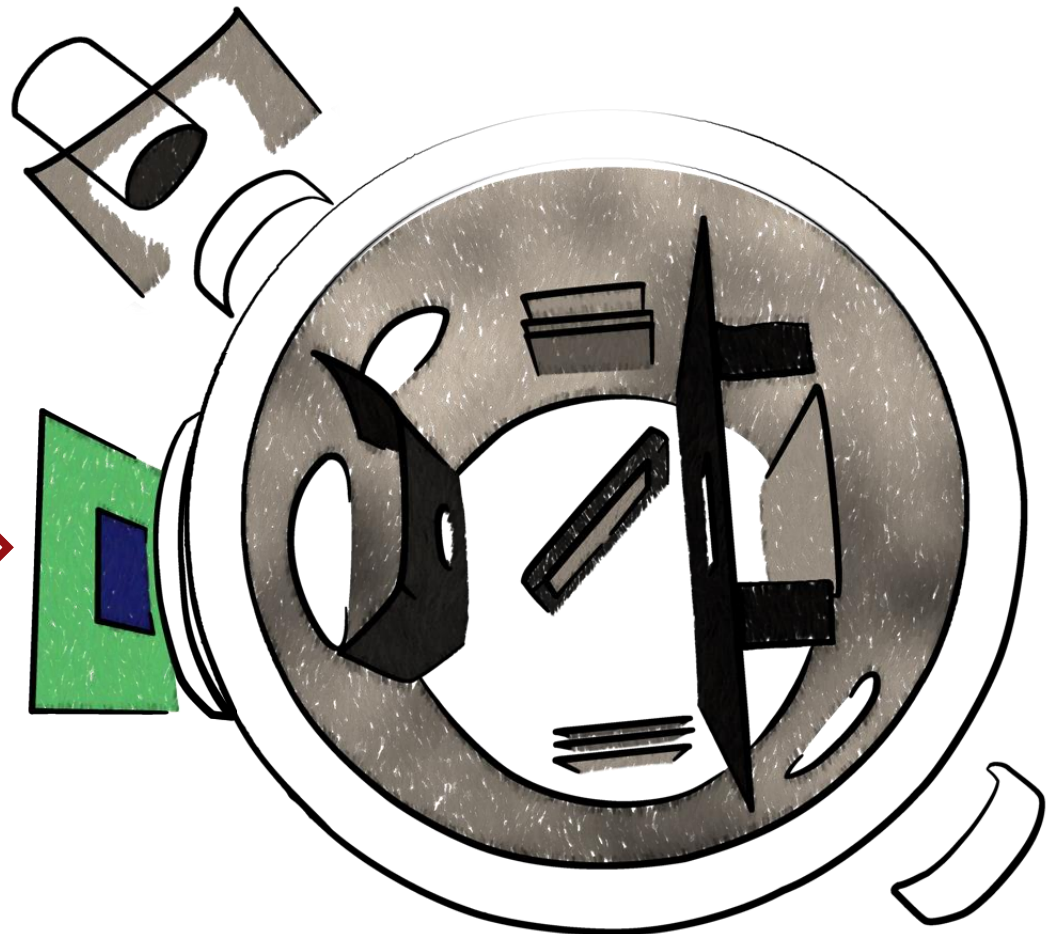
Beam monitoring

Upstream of the beam window, a **scintillator** paddle is used to tag charged particles in the beam. This is also used in the offline analysis.



μ beam

▲ Just after the scintillator, there is a **wire chamber** to measure the beam profile in x & y

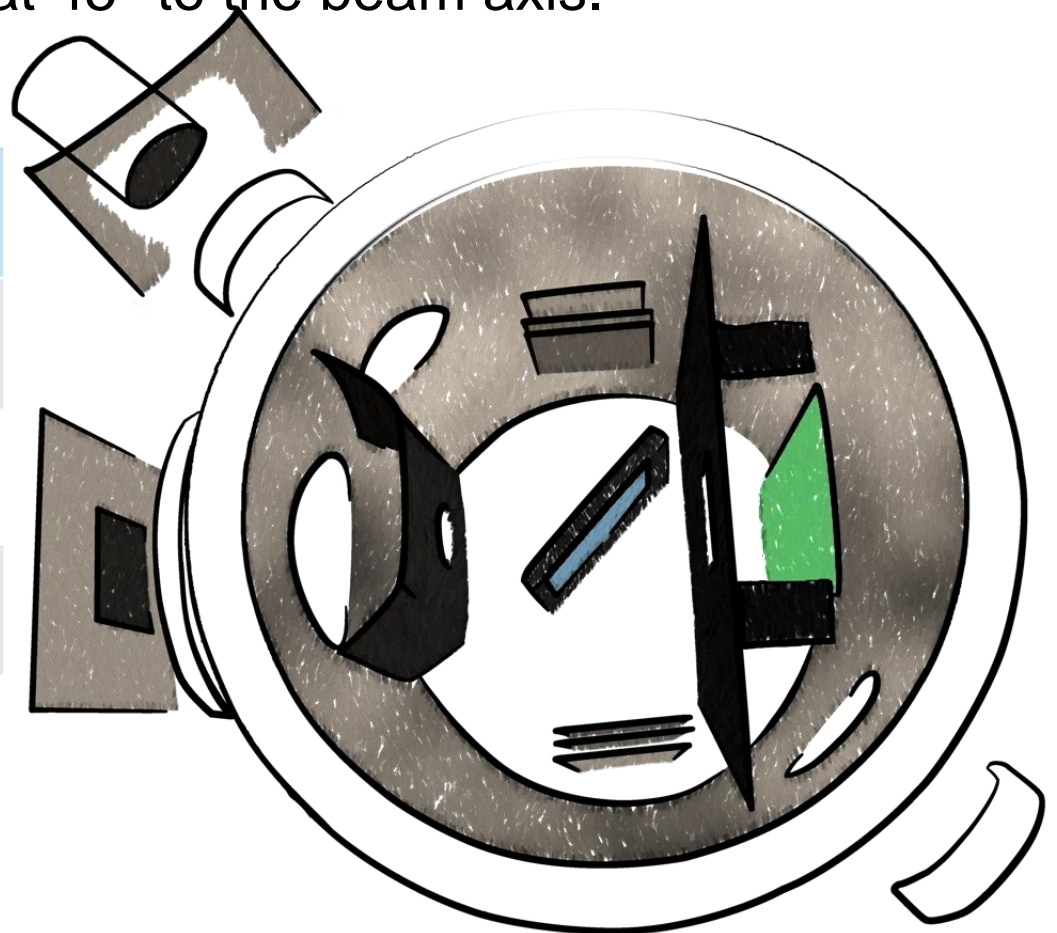


Target and veto

Behind the target is a **veto scintillator** for **electrons**, which are more penetrating at this energy.

The **target** sits at the centre, at 45° to the beam axis.
Several targets were used:

| Material | Thickness | Notes |
|----------|--------------------|--------------------|
| Si | 1500 μm | 'Right arm' |
| Si | 65 μm | Mostly passive (*) |
| Al | 100 μm | |
| Al | 50 μm | |

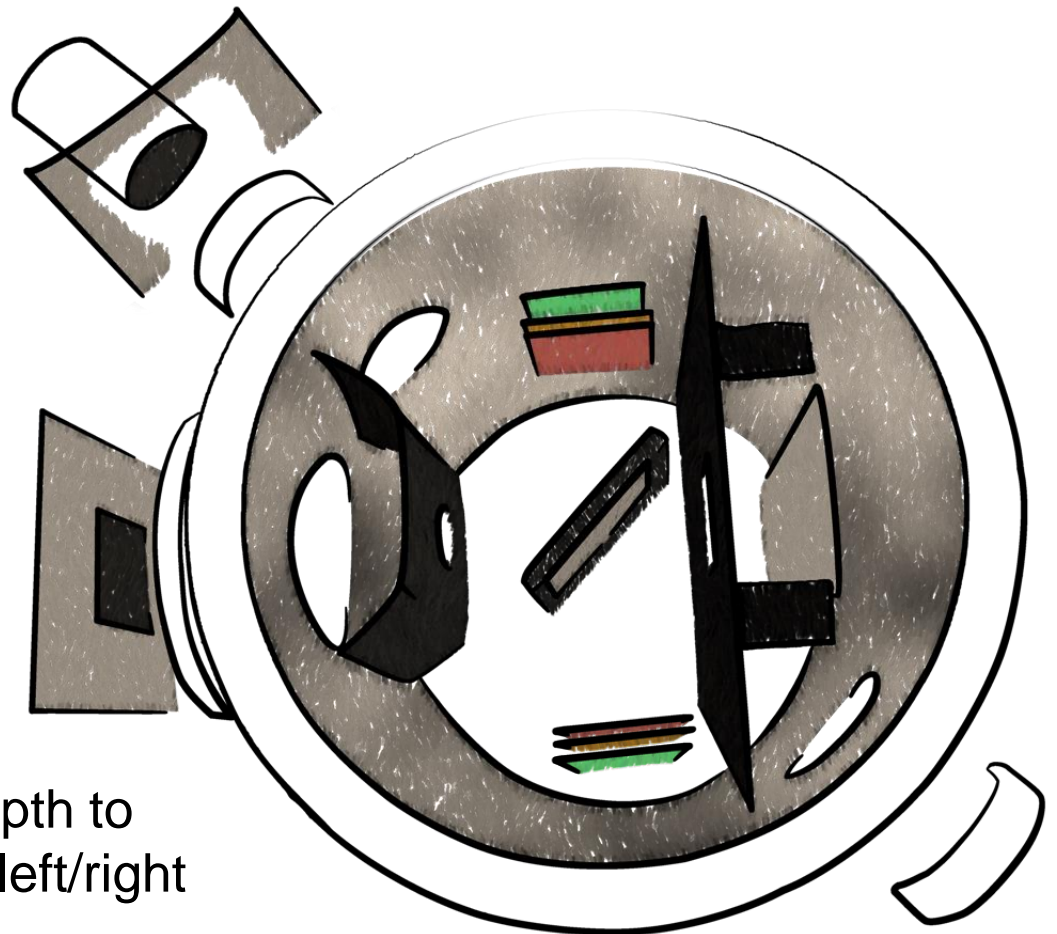


At $\pm 90^\circ$ to the beam axis are two detector arms, consisting of:

- **Thin 'transmission' silicon**, $65\mu\text{m}$, with 2×2 segmentation to measure δE of emitted particles
- **Thick silicon** 1.5mm, to measure overall energy
- **Scintillator** paddles to tag escaping particles

The arms are:

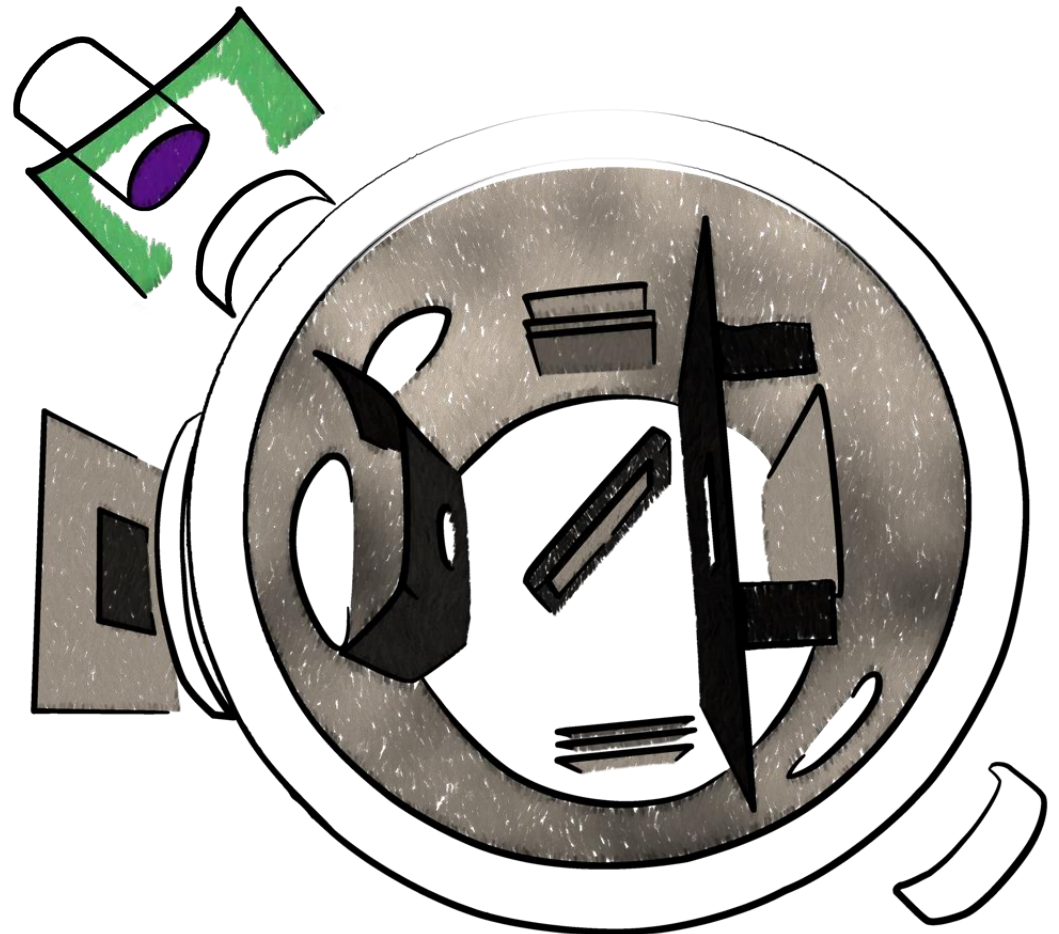
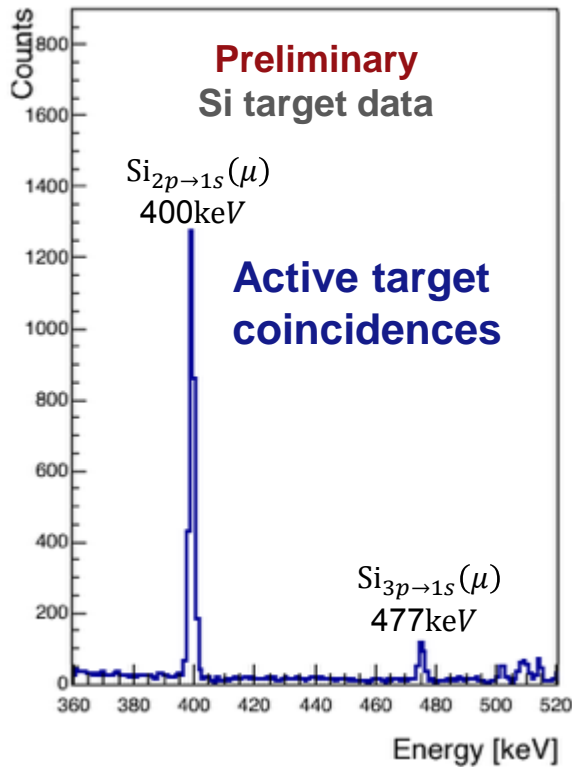
- Symmetric to the beam
 - Equalises BG
- At 45° to opposite faces of target
 - Allows beam penetration depth to be estimated by comparing left/right



Germanium detector

At one port is a **HP-Ge detector**, and associated **scintillator** paddle

- Used to measure the capture rate by looking for the muonic $2p \rightarrow 1s$ at 347keV (Al) or 400keV (Si)



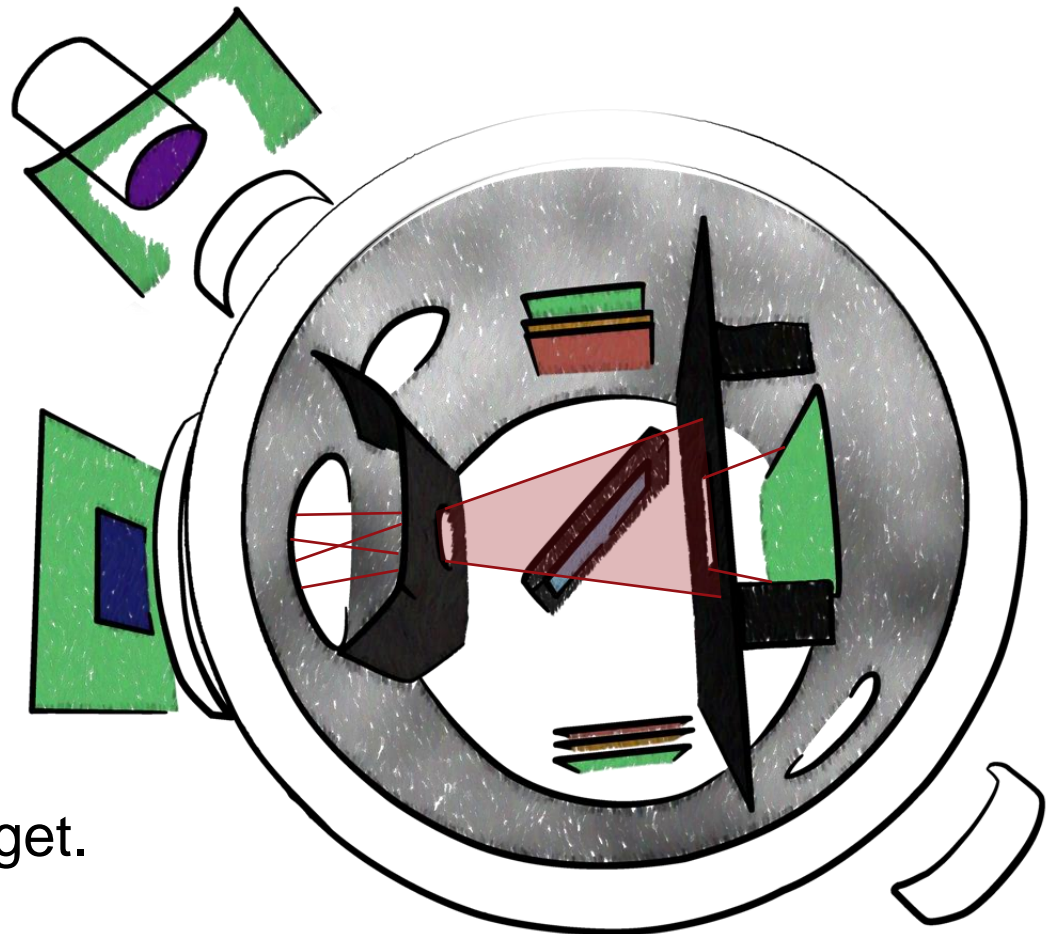
Shielding

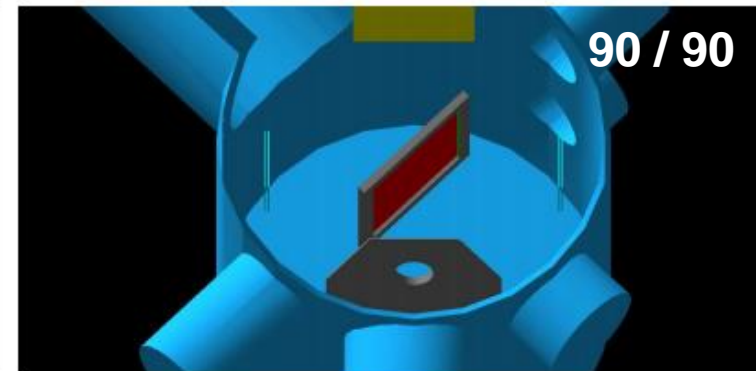
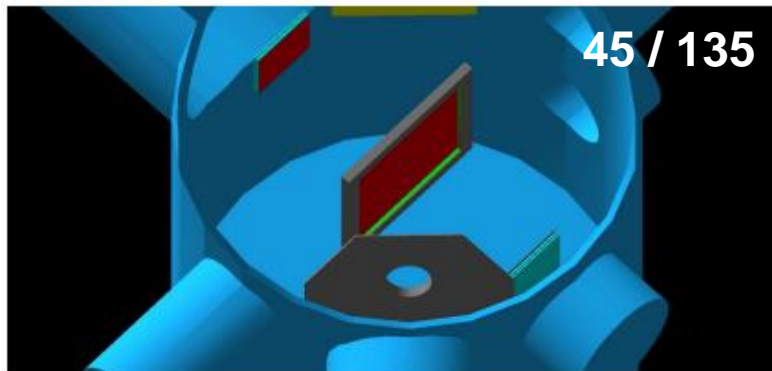
Shielding and geometry received careful consideration, based on experience from 2009 run:

- We added lead ($\tau \sim 82\text{ns}$) shielding **upstream** and **downstream** of the target, and on the **target mount**



- Goal is that the only **beam-irradiated** areas visible to detectors are made of lead, or are the target.

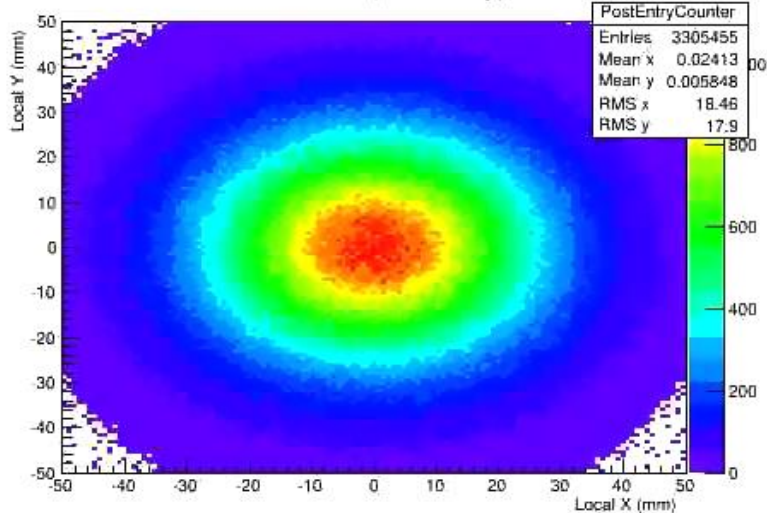




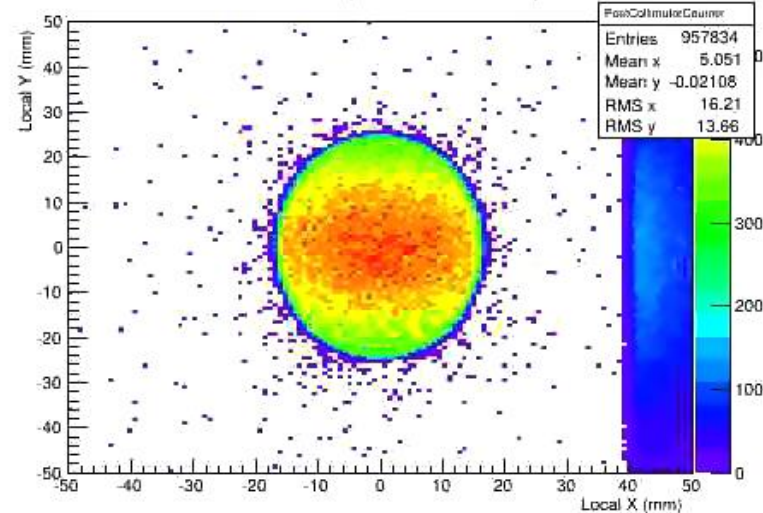
▲ Comparison of detector orientations

Simulation of upstream shield/collimator ▼

Hit Position (PostEntry)

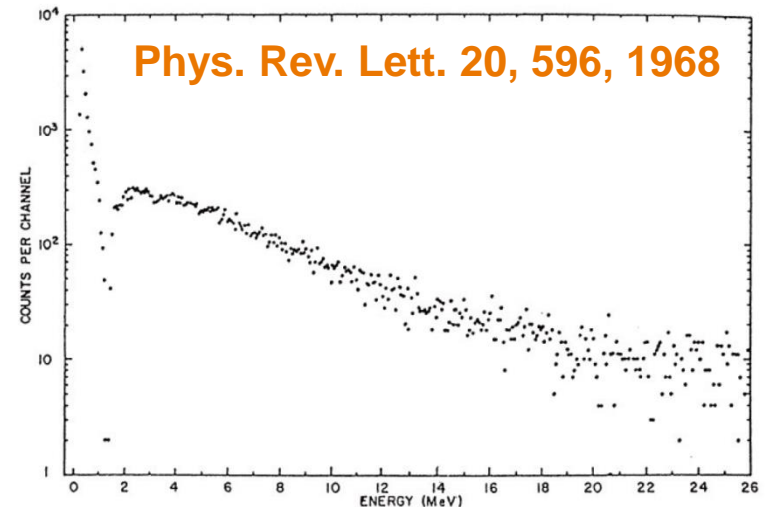
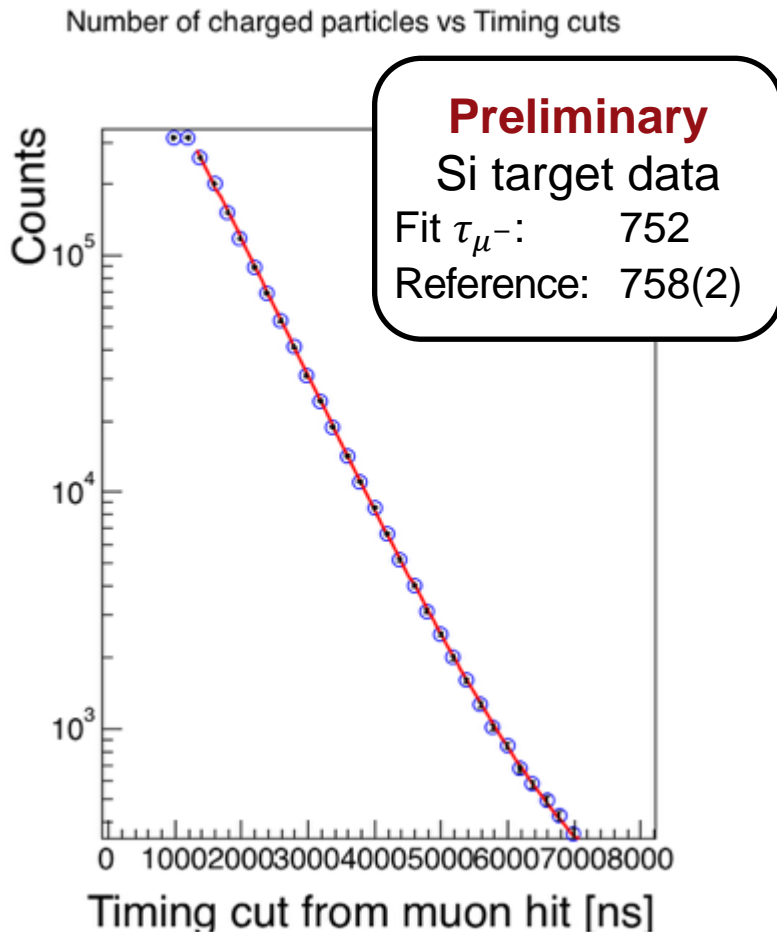


Hit Position (PostCollimator)



Analysis is currently focussing on data taken with the active targets.

- Goal is to reproduce known spectra then use what we learn from these to 'bootstrap' analyses of the passive targets.



- More active Si data with thin target would have been preferable...
Using thick Si where possible

Initial Analysis II

Silicon packages measure the particle energy twice.

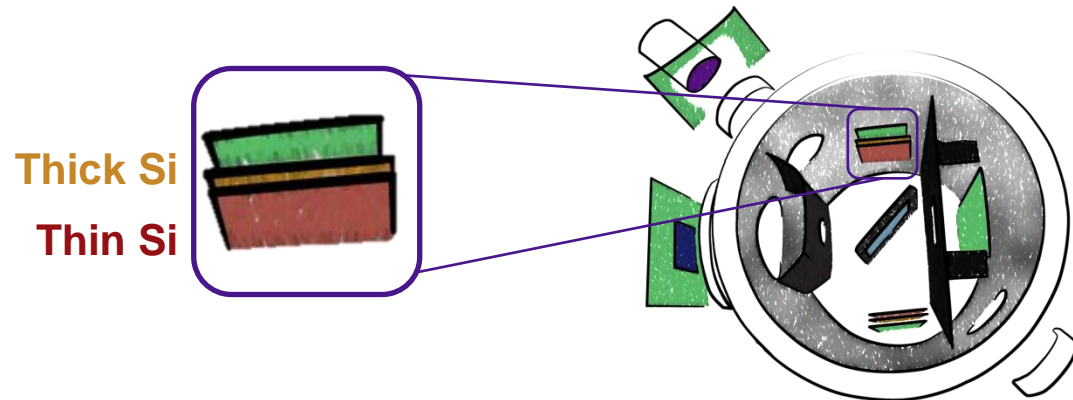
Since first Si is thin, energy deposit is:

$$\delta E \simeq \frac{dE}{dx} x_{\text{thin}}$$

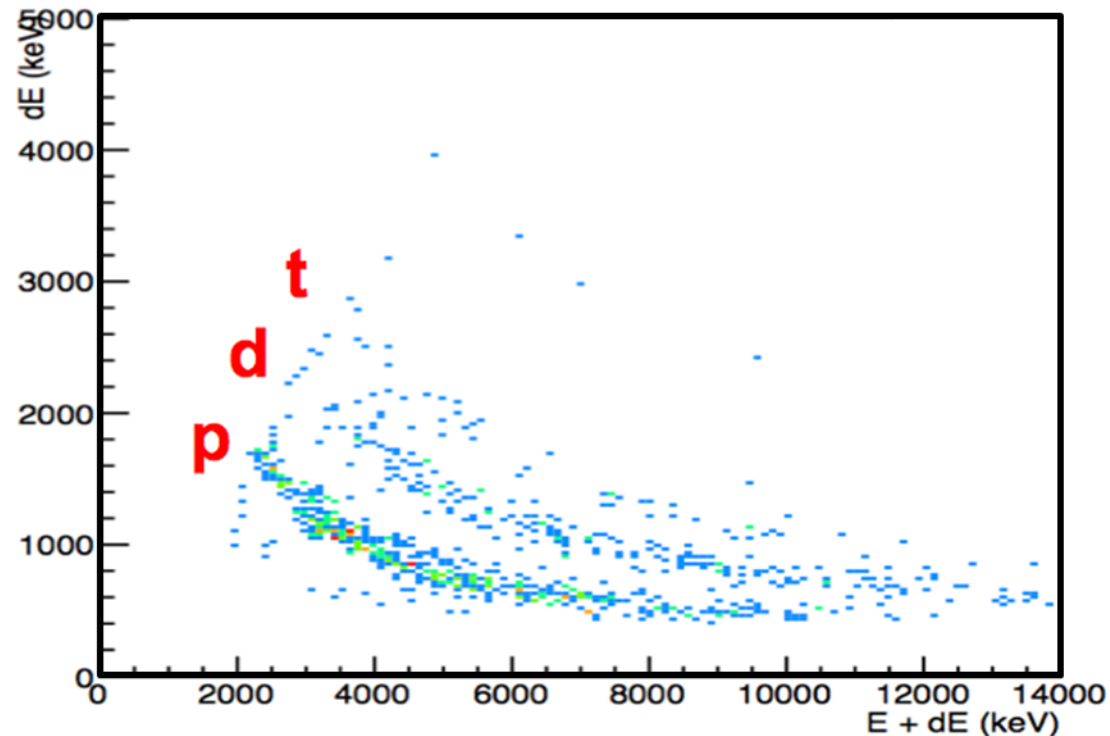
And since second Si is thick (and no scintillator veto) it measures the remaining energy E . So:

$$\delta E(E + \delta E) \propto \frac{dE}{dx} (E + \delta E)$$

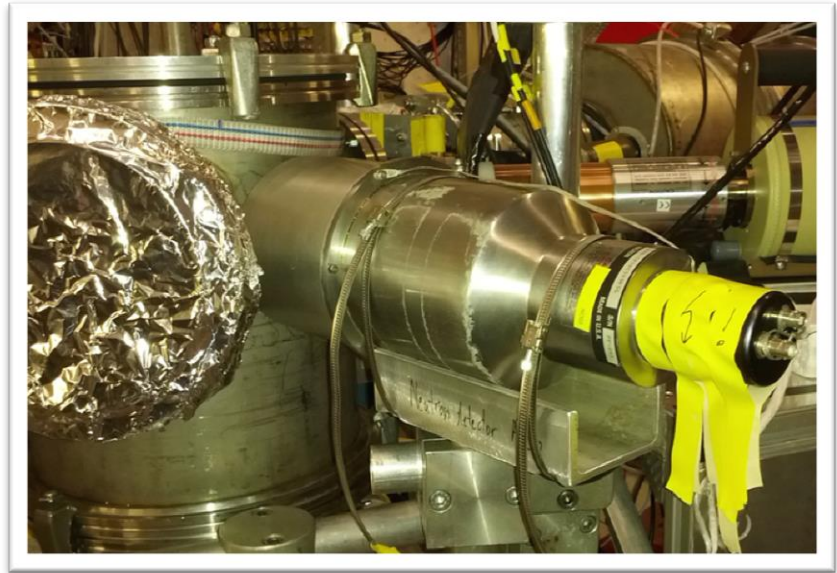
which can be used for charged-particle PID ►



dEdx, Al100, left, 1 - 6μs from μSc hit

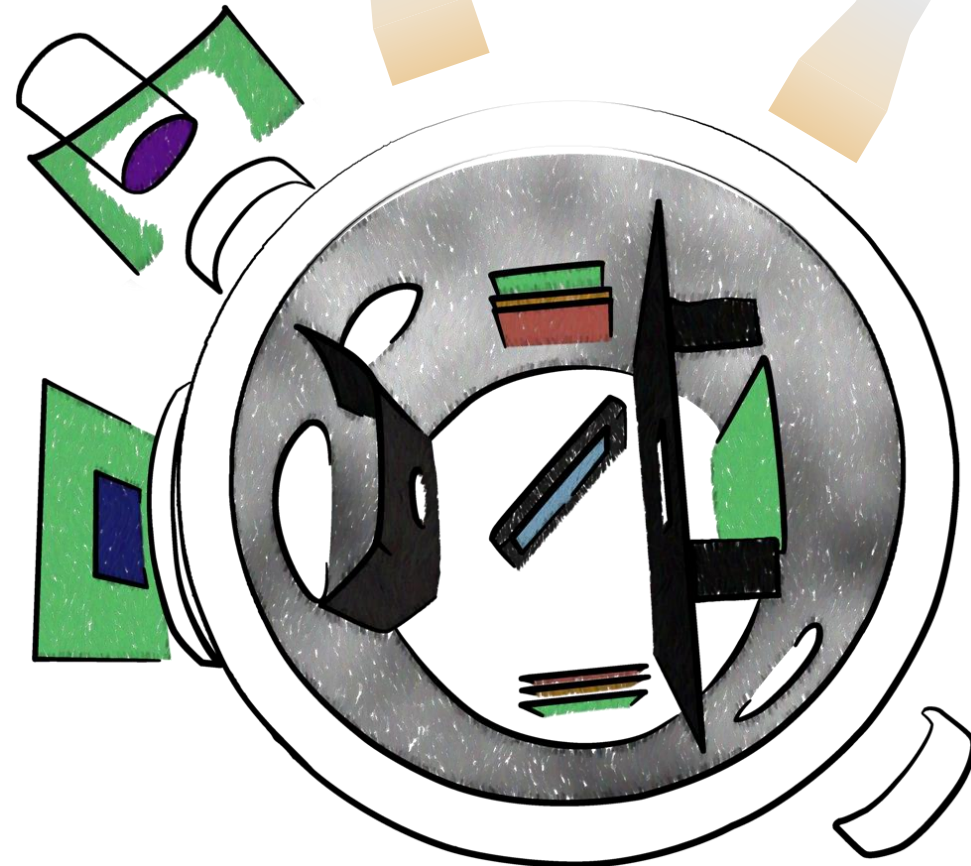


Some preliminary work testing and running with neutron detectors was completed



‘Proof of principle’ analysis is being done on limited data...

...but full data and analysis will come from a later run



Another run?

Run seems to have been successful, and analysis is progressing:

- Problems mostly limited to electronics noise
- Most of the 'Golden' data sets have passed low level quality checks

Should be able to meet goals of **gamma ray normalisation**, and some **charged particle rates/spectra**. But...

- More active Si data would make it possible to do **detailed cross checks**.
- Would like to take more data for serious **neutron studies**.
- Would also like to obtain data with **titanium target** for possible use later in Mu2e and COMET programs

Currently considering making a request for PSI beam in 2015

- It is desirable to take another run before too many students & postdocs move on.

End



