

# Quantifying the sensitivity of oscillation experiments to the neutrino mass ordering

Mattias Blennow  
`emb@kth.se`

KTH Theoretical Physics

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## 1 The usual deal

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## 2 The alternative

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## 3 Summary and conclusions

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# Parameter estimation sensitivity

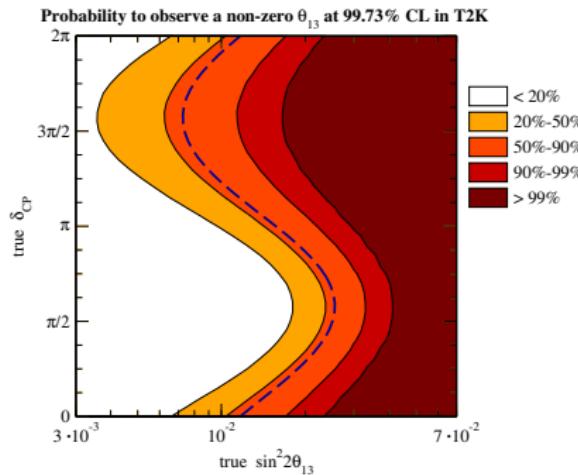
- Define the test statistic “ $\Delta\chi^2$ ”

$$\Delta\chi^2(\theta) = -2 \log \left[ \frac{\mathcal{L}(\theta|d)}{\sup_{\theta'} \mathcal{L}(\theta'|d)} \right]$$

- Assume it is  $\chi^2$  distributed with  $n$  degrees of freedom
- Use the data set without statistical fluctuations (Asimov data)
- Quote result

# The interpretation

- $\Delta\chi^2$  is asymptotically  $\chi^2$   
(Wilks' theorem)
- The Asimov data is representative
- Several requirements, not always fulfilled



Schwetz, Phys.Lett. B648 (2007) 54

# Not a nested hypothesis

- Mass ordering is not nested
- Wilks' theorem not applicable
- Some different choices of test statistic

$$\Delta\chi^2 = \chi_{\text{NO}}^2 - \min \chi^2$$

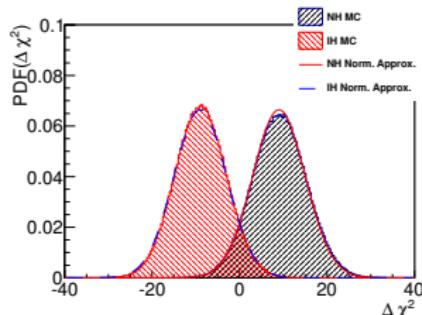
$$T = \chi_{\text{IO}}^2 - \chi_{\text{NO}}^2$$

$$T' = \chi^2(\theta) - \min \chi^2$$

- All have different distributions
- We concentrate on  $T$

$$T \simeq \mathcal{N}(T_0, 2\sqrt{T_0})$$

$T_0$  = value for Asimov data



Qian, et al., Phys.Rev. D86 (2012) 113011

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# Back to basics

- Test hypothesis NO, alternative hypothesis IO → “Can we reject NO if IO is true?”
- Critical region defined by  $T_c$   
*Reject NO if  $T < T_c$*
- Test significance (CL):  $1 - \alpha = 1 - P(T < T_c | \text{NO})$   
 $\alpha$ : *Probability to reject NO if NO is true*
- Test power:  $1 - \beta = P(T < T_c | \text{IO})$   
 $\beta$ : *Probability of failing to reject false NO if IO is true*
- *Both  $\alpha$  and  $\beta$  are relevant*

# What is sensitivity?

## Sensitivity (median)

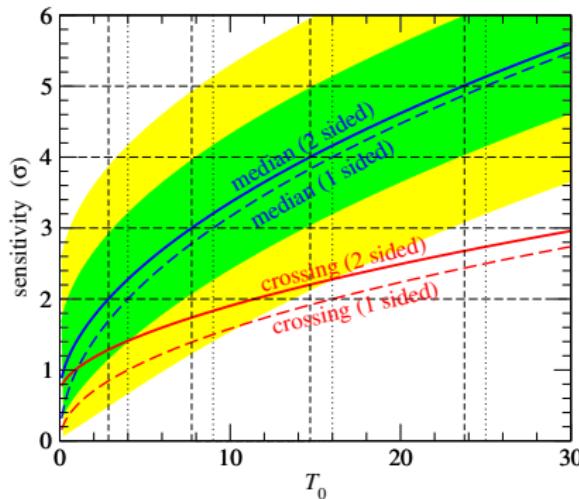
What is the expected rejection of a false ordering? (Given a parameter set)

## Other

- What is the probability of ruling out a false hypothesis at  $x\sigma$ ?
- At what CL is exactly one ordering ruled out?

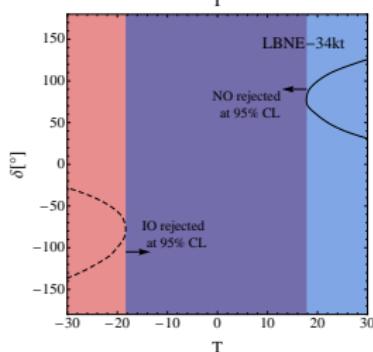
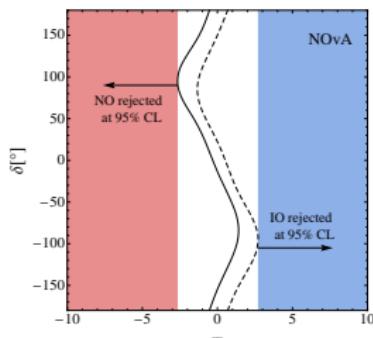
# Simple hypotheses

- Distribution of test statistic independent of parameter values
- Straight forward
- Applicable to reactor experiments



MB, Coloma, Huber, Schwetz, JHEP 03(2014)028

# Composite hypotheses



MB, Coloma, Huber, Schwetz,  
JHEP 03(2014)028

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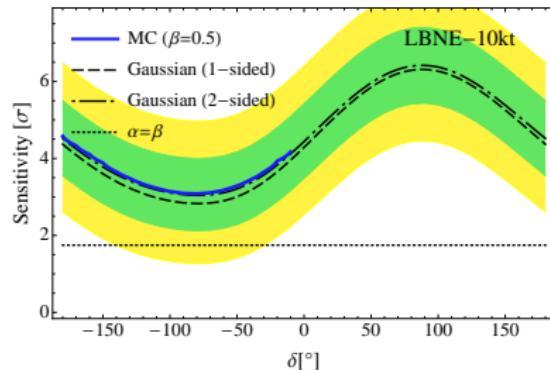
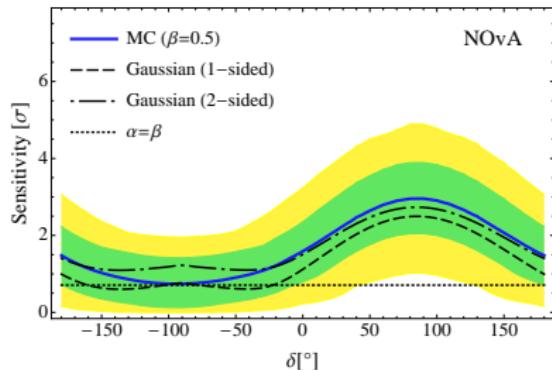
- Distributions have parameter dependence
- To reject = Rejecting all parameter sets ( $\theta_{13}$ ,  $\Delta m_{31}^2$ ,  $\delta$ , etc)
- Distribution of  $T$  is parameter dependent

$$T_c = \min_{\theta} T_c(\theta)$$

- Extensive Monte Carlo simulations
- Approximation (must check validity)

$$T(\theta) = \mathcal{N}(T_0(\theta), 2\sqrt{T_0(\theta)})$$

# Accelerator experiments - results

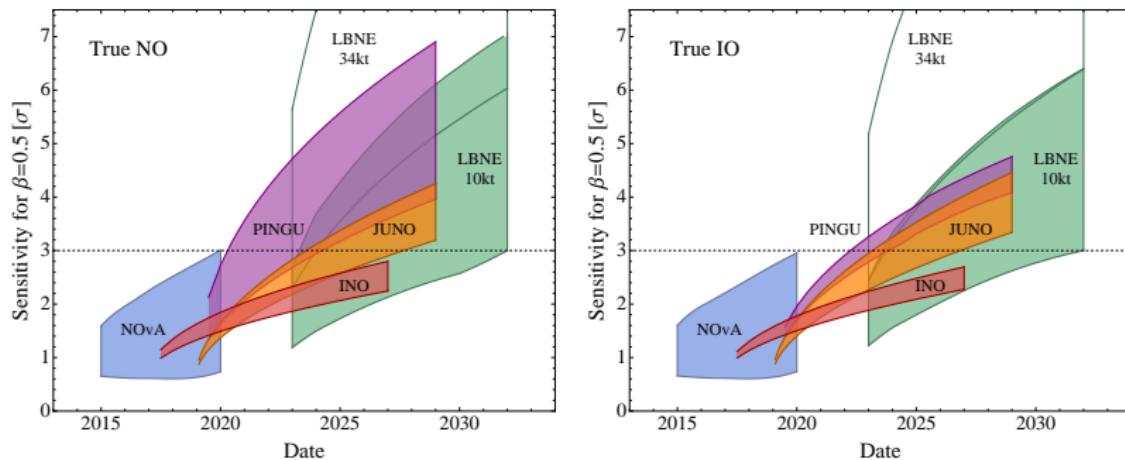


MB, Coloma, Huber, Schwetz, JHEP 03(2014)028

Green: 68 % of outcomes

Yellow: 95 % of outcomes

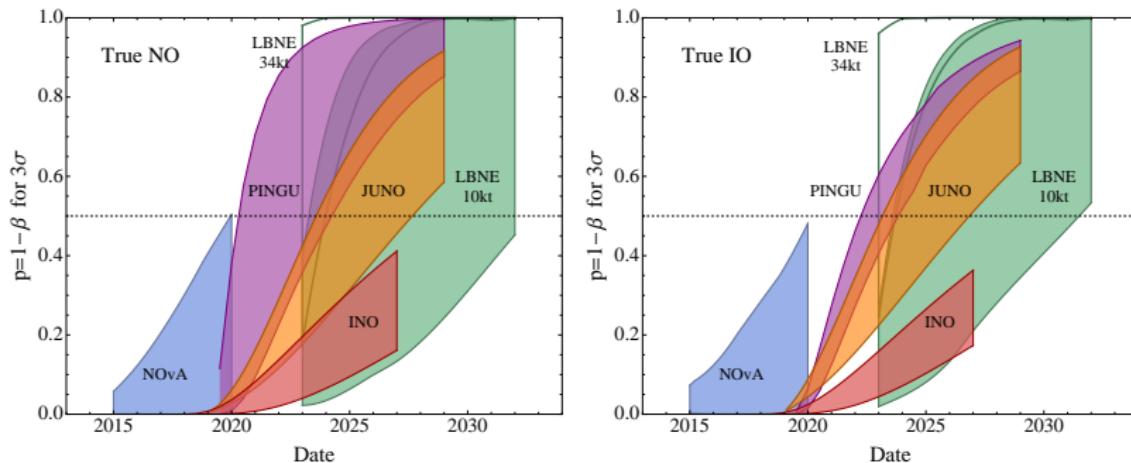
# Comparison of experiments - specific $\beta$



MB, Coloma, Huber, Schwetz, JHEP 03(2014)028

**Note:** Bands have different meanings!

# Comparison of experiments - specific $\alpha$



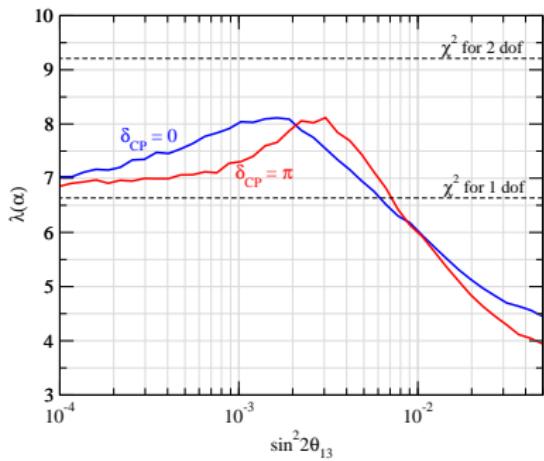
MB, Coloma, Huber, Schwetz, JHEP 03(2014)028

**Note:** Bands have different meanings!

# How to interpret the median sensitivity

- It is *representative* for how well the experiment will do
- 50 % probability of not reaching it
- 50 % probability of *doing better*
- *Not* 50 % probability of “being wrong”
- Not the only relevant quantity, distribution matters (do Brazilian bands!)
- Personal preference: Quote the power  $1 - \beta$  for a target sensitivity

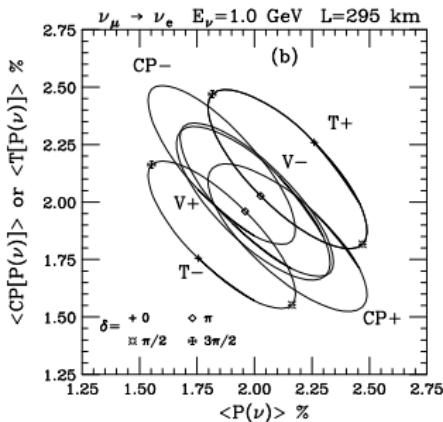
# What about $\delta_{CP}$ ?



Schwetz, Phys.Lett. B648 (2007) 54

- The test statistic may be very different from a  $\chi^2$  distribution
- Checked with full Feldman–Cousins approach
- Several devious details, caveat venditor!

# What about $\delta_{CP}$ ?

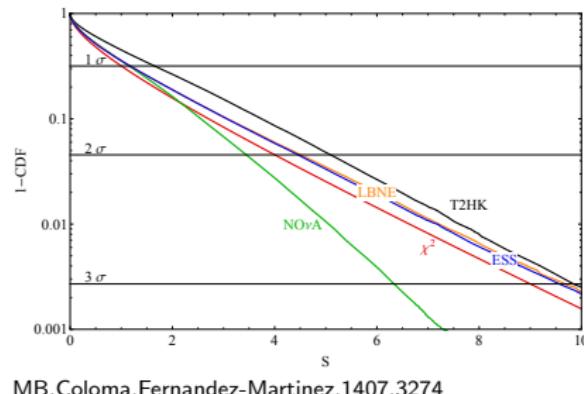


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Minakata, Nunokawa, Parke, PLB537 (2002) 249

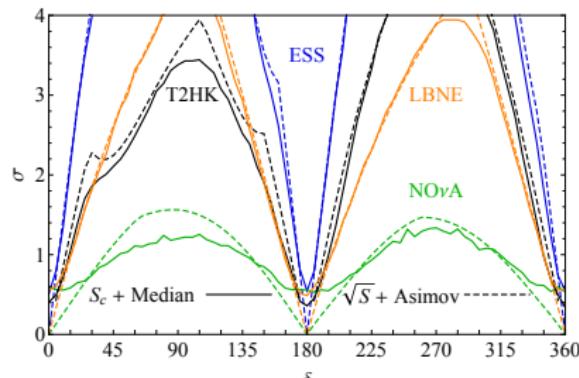
## CP violation sensitivity

- Distribution under CP conserving hypothesis:
    - High precision: Slightly shifted to higher  $\chi^2$
    - Low precision: Shifted to lower  $\chi^2$
  - Median values:
    - High precision: Do not change much
    - Low precision: Shifted to lower  $\chi^2$



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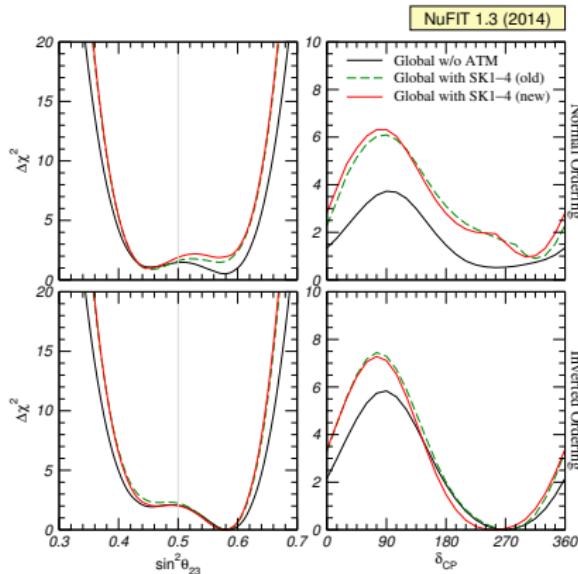
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MB,Coloma,Fernandez-Martinez,1407.3274

# Octant sensitivity

- Degeneracies closer
- Wilk's theorem still violated
- A priori, a dedicated study is needed



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# Summary and conclusions

- Wilks' theorem is not applicable to the neutrino mass ordering, the test statistic is not  $\chi^2$  distributed
- Regardless,  $\sqrt{T_0}$  is still a good approximation of the (median) sensitivity
- Frequentist (as well as Bayesian – not discussed) methods are perfectly applicable
- Critical values will depend on the experiment and underlying assumptions
- May be relevant for  $\delta$ , see also plenary talk by Thomas Schwetz

## 4 Backup frequentist

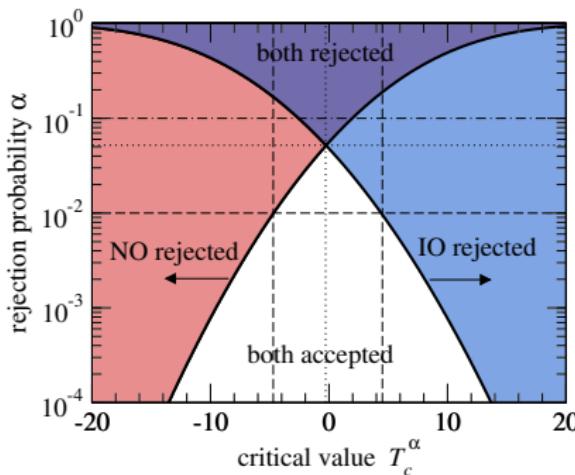
## 5 Backup distributions and results

## 6 Backup Bayesian basics

## 7 Bayesian methods

## Testing both orderings

- Depending on the significance it may be possible to:
    - Reject exactly one hypothesis
    - Reject both hypotheses
    - Not reject any hypothesis
  - It is natural, the true hypotheses *should* be rejected with probability  $\alpha$

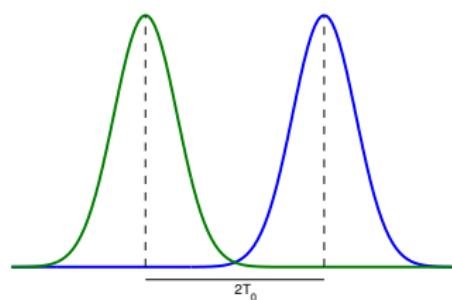


MB, Coloma, Huber, Schwetz, JHEP 03(2014)028

# At the end of the day

For the simple hypotheses:

- Two Gaussians,  
 $H_{\pm} : \mathcal{N}(\pm T_0, 2\sqrt{T_0})$
- For  $H_+$ , *typical* (median) result is  $+T_0$
- $+T_0$  is  $T_0 - (-T_0) = 2T_0$  away from the expected  $H_-$  result
- $2T_0/(2\sqrt{T_0}) = \sqrt{T_0}$



See also: Vitelis, Read, 1311.4076

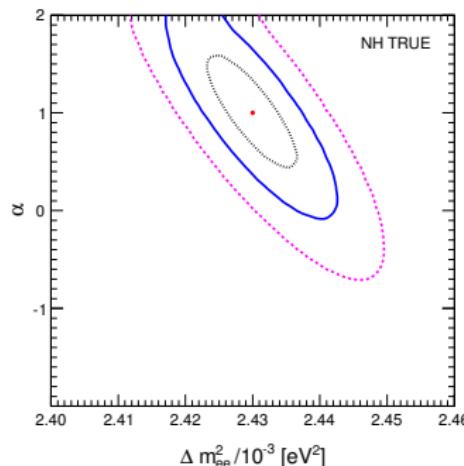
# Interpolating parameters

- Introduce an interpolating continuous parameter  $\alpha$  such that

$$P(\alpha = \pm 1, \theta) = P(\theta, \text{NO/IO})$$

see, e.g., Capozzi, Lisi, Marrone, arXiv:1309.1638

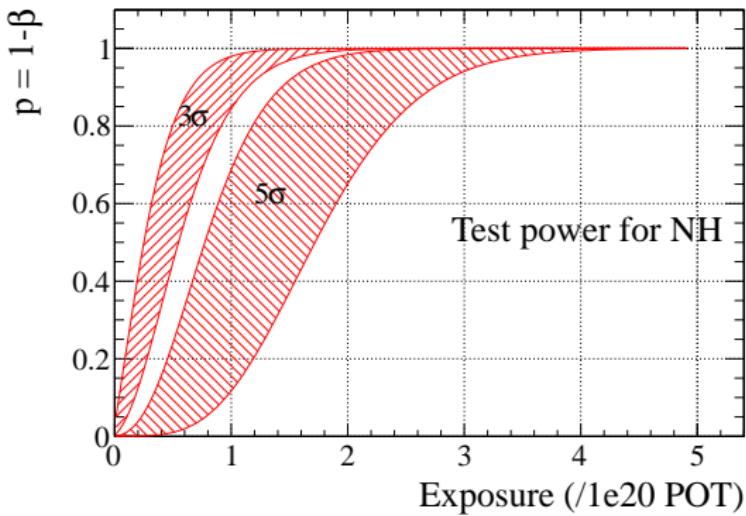
- Wilks' theorem now applies
- Put bounds on  $\alpha$
- If  $\alpha = 1$  ( $-1$ ) is allowed, NO (IO) is allowed



Capozzi, Lisi, Marrone, arXiv:1309.1638

- Personal comment:  $\alpha = 0$  is not special, it being allowed a priori does not affect the rejection of NO/IO

# LBNO predictions



LBNO collaboration, arXiv:1312.6520

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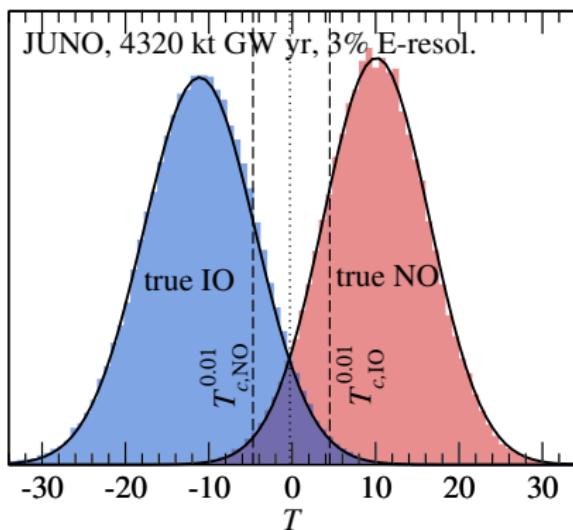
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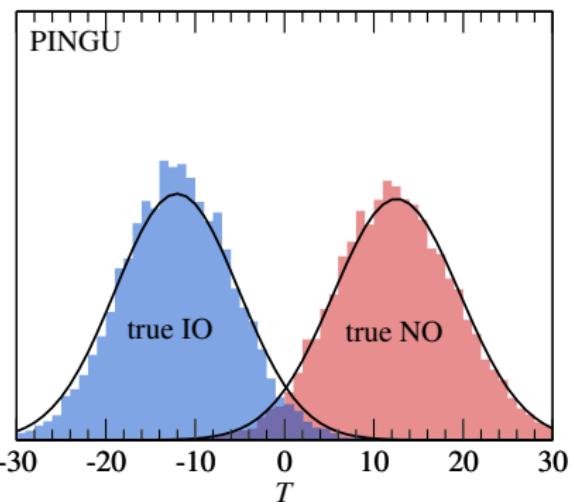
# Reactor experiments

- Distribution of  $T$  essentially parameter independent
- Simple hypotheses
- Gaussian limit well satisfied
- Median sensitivity  $n\sigma$ :

$$n \simeq \sqrt{2} \operatorname{erfc}^{-1} \left[ \frac{1}{2} \operatorname{erfc} \left( \sqrt{\frac{T_0}{2}} \right) \right]$$



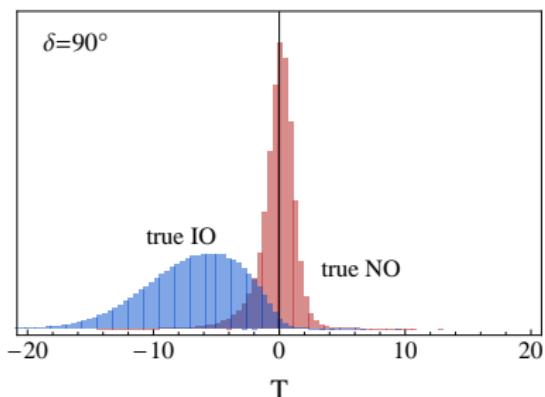
# Atmospheric experiments



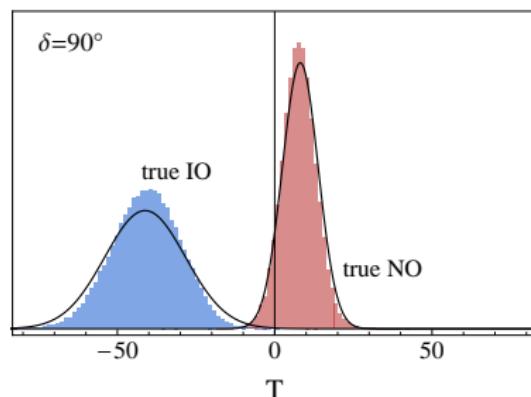
- Distribution of  $T$  mainly depends on  $\theta_{23}$
- Gaussianity is still well satisfied for each  $\theta_{23}$
- Analytic as function of  $T_0(\theta_{23})$
- Boils down to evaluating the Asimov data:  $T_0$

# Accelerator experiments - distributions

- Not always Gaussian!
- Typical for low-sensitivity experiments
- Need to perform Monte Carlo studies for accuracy
- Rejection power *depends on the true parameters*

NO $\nu$ A

LBNE



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# Bayesian basics

- Assign a degree of belief in each hypothesis  $P(H)$
- Update the degree of belief depending on observations
- Bayes' theorem

$$P(A, B) = P(A; B)P(B) = P(B; A)P(A)$$

$$P(A; B) = \frac{P(B; A)P(A)}{P(B)}$$

- Take  $A = \text{hypothesis } H$ ,  $B = \text{data } d$

$$P(H; d) = \frac{\mathcal{L}_H(d)P(H)}{P(d)}$$

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# Bayesian hypothesis testing

- Study the relative degrees of belief in two hypotheses

$$\frac{P(H_1; d)}{P(H_2; d)} = \frac{\mathcal{L}_{H_1}(d)}{\mathcal{L}_{H_2}(d)} \frac{P(H_1)}{P(H_2)}$$

- Strength of the evidence for  $H_1$ :

$$\kappa = 2 \log \left[ \frac{P(H_1; d)}{P(H_2; d)} \right]$$

- Kass-Raftery scale:

Strength of evidence for H	$\kappa$	Posterior odds	Degree of belief
Barely worth mentioning	0 to 2	ca 1 to 3	< 73.11%
Positive	2 to 6	ca 3 to 20	> 73.11%
Strong	6 to 10	ca 20 to 150	> 95.26%
Very strong	> 10	≥ 150	> 99.33%

# What can be said about the future?

- Can compute probability of obtaining evidence at least strength  $\kappa_0$  *for the true ordering*

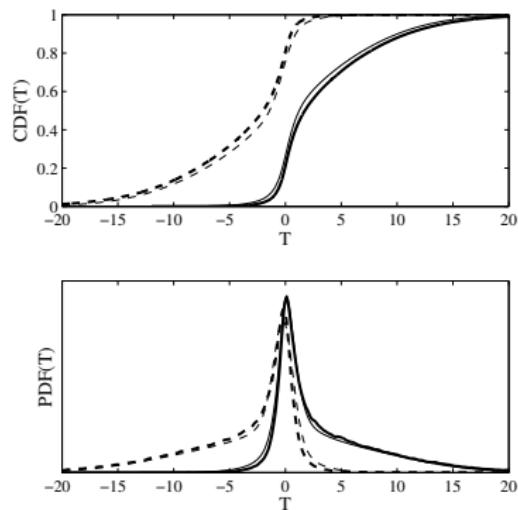
$$P(\kappa_0) = P(\kappa > \kappa_0; H_1)P(H_1) + P(\kappa < -\kappa_0; H_2)P(H_2)$$

- Typical choice  $P(H_1) = P(H_2) = 0.5$
- Takes into account information on oscillation parameters

$$\mathcal{L}_H(d) = \int \mathcal{L}_{H(\theta)}(d)\pi(\theta)d\theta$$

- Compactifies all of the available information to one number
- Easy to simulate through Monte Carlo methods
- Prior dependent

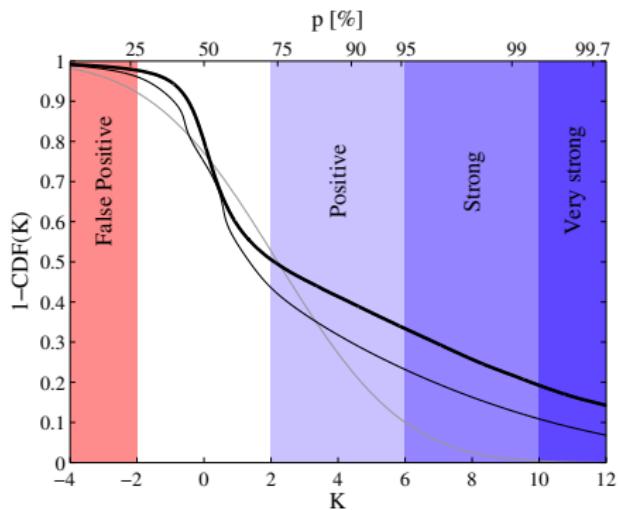
# Example: NO $\nu$ A



- NO $\nu$ A experiment
- GLoBES implementation
- Only  $\delta$  and  $\Delta m_{31}^2$  varying
- Flat and 10 % Gaussian priors, respectively

MB, JHEP 01(2014)139

# Example: NO $\nu$ A, results



- Probability of obtaining evidence with *at least* strength  $K$  for the correct ordering

MB, JHEP 01(2014)139