

# VERTEX 2014

23<sup>rd</sup> International Workshop on Vertex Detectors

Mácha Lake | Doksy, Czech Republic

September 15 - 19, 2014



# Simulation of Irradiated Si Detectors

**Ranjeet Dalal, Kirti Ranjan, Ashutosh Bhardwaj, Geetika Jain,  
Kavita Lalwani**

**Department of Physics and Astrophysics,  
University of Delhi, INDIA**

# Contents

- Simulation approaches up to now and some interesting observations
- Simulation framework
- Bulk damage model and its simple implications
- Surface damage implementations into simulations
  - Surface Oxide Charge density ( $Q_F$ )
  - Interface Trap density ( $N_{it}$ )
- Simulations of Interstrip resistance ( $R_{int}$ ) for irradiated Si strip sensors
- Electric field simulations for irradiated Si strip sensors
- Summary

# Simulation approaches for irradiated sensors : Up to now !

- Either considered Surface damage only or considered Bulk damage only
- Only Oxide charge density ( $Q_F$ ) variation was taken into account for Surface damage;

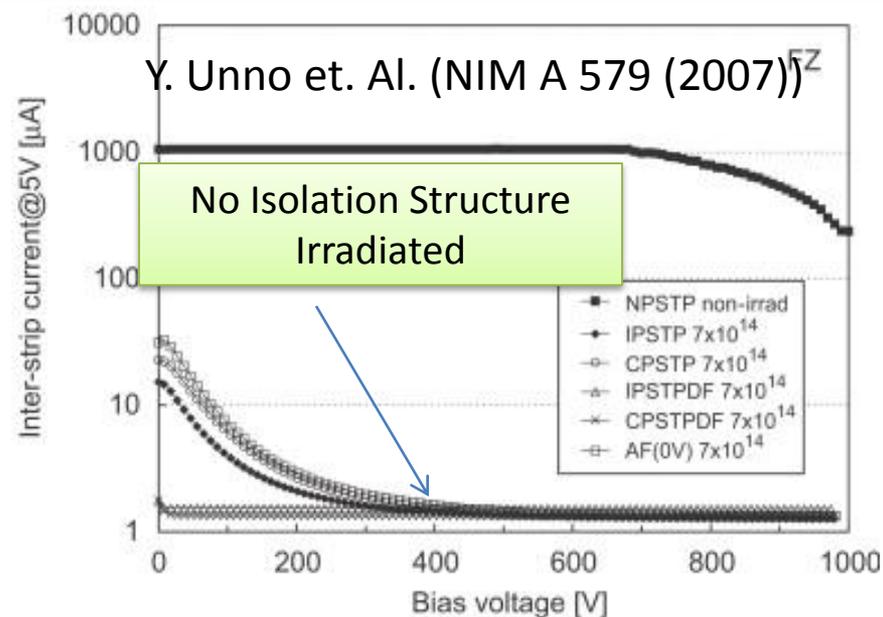
**No interface trap state included!**

**No simulation study which incorporate Bulk + Surface damage simultaneously !**

Simulation using Surface damage only	Simulations using Bulk damage only
1. G. Verzellesi & G.F. Dalla Betta <b>Nucl. Sci. Symp., 2000 IEEE (Vol.-1)</b> Compact modeling of n-side interstrip resistance in p-stop and p-spray isolated double-sided silicon microstrip detectors	1. V. Eremin et al. NIM A 476 (2002) 556–564 The origin of double peak electric field distribution in heavily irradiated silicon detectors
2. P. Claudio (2006) IEEE Trans. ON Nucl. Sci., VOL. 53, NO. 3 Device Simulations of Isolation Techniques for Silicon Microstrip Detectors	2. M. Petasecca <i>et al.</i> NIM A 563 (2006) 192–195 Numerical simulation of radiation damage effects in p-type silicon detectors
3. Y Unno et al. NIM A 636 (2011) S118–S124 Optimization of surface structures in n-in-p silicon sensors using TCAD simulation	3. V. Chiochia et al., IEEE Trans. Nucl. Sci. NS-52 (2005) 1067 Simulation of Heavily Irradiated Silicon Pixel Sensors and Comparison With Test Beam

# Some puzzling observations !

- Very good interstrip resistance values ( $R_{int}$ ) are observed for n-on-p Si sensor with very low p-stop ( $\sim 5 \times 10^{15} \text{cm}^{-3}$ ) and p-spray doping densities ( $\sim 1 \times 10^{15} \text{cm}^{-2}$ ) even after very high proton fluences ( $\sim 1 \times 10^{15} \text{cm}^{-2} \text{neq}$ ) (CMS-HPK tracker phase-II upgrade study)
- Good  $R_{int}$  was reported [see right] for n-on-p strip sensors, irradiate with high proton fluences, having no insulation structures !



Oxide charge density ( $Q_F$ ) appears to be suppressed in the hadron irradiated sensors!

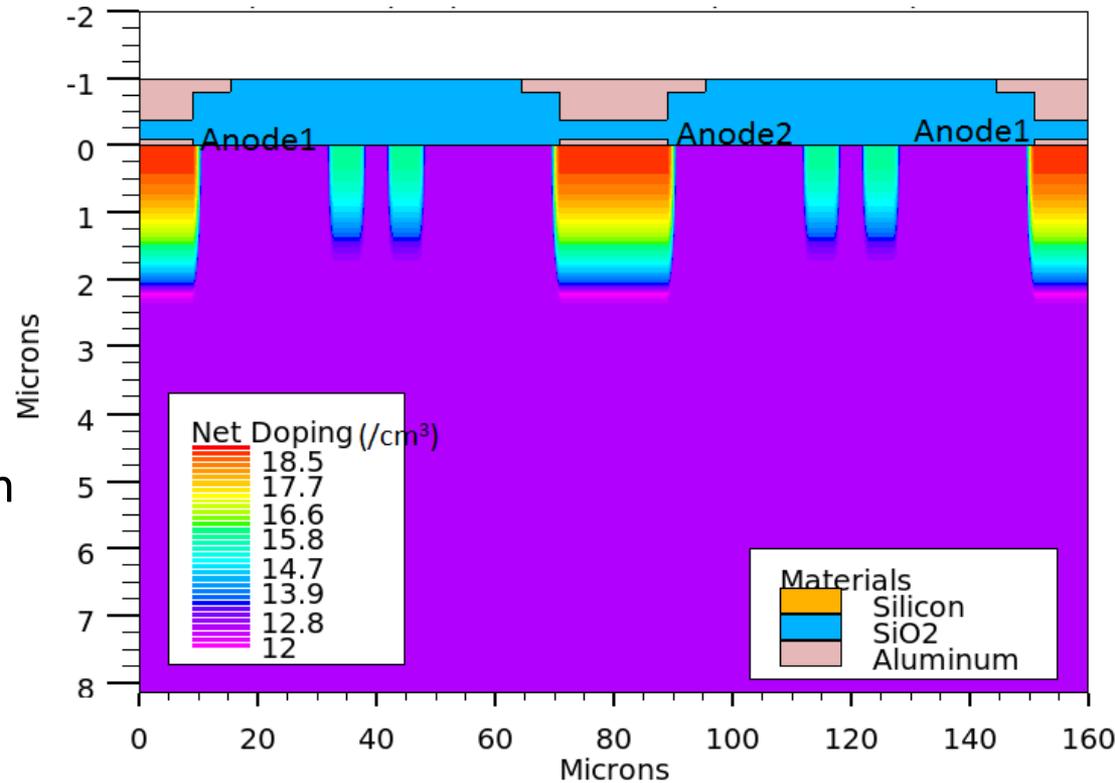
- From eTCT measurements and earlier Bulk damage simulations, higher electric fields are expected near  $n^+$  strips compare to  $p^+$  strip; so n-on-p strip sensors should show lower breakdowns at higher fluences

But p-on-n type strip sensors are found to be much more prone to microdischarge!

These observations (and many others) can be understood as combine effects of Bulk and Surface damage ( $Q_F$  and  $N_{it}$  both) !

# Simulation structure

- Bulk doping =  $3 \times 10^{12} \text{cm}^{-3}$
- 2-D simulations
- Double p-stops
- Each  $4 \mu\text{m}$  wide separated by  $6 \mu\text{m}$
- P-stop doping =  $5 \times 10^{15} \text{cm}^{-2}$   
P-stop doping depth =  $1.6 \mu\text{m}$
- CMS HPK tracker upgrade campaign parameters



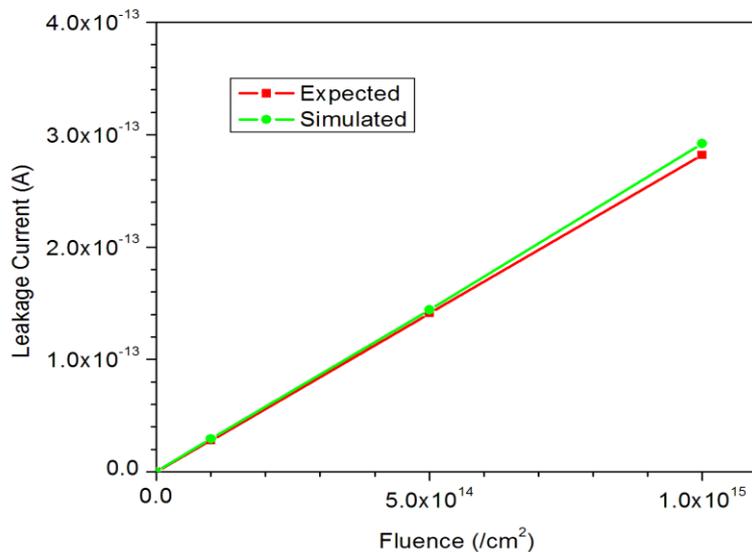
- ❑ Three strips structure was used for  $R_{\text{int}}$  simulations in which bias of 0.2V is given to Central DC Anode while two neighboring Anodes are shorted together. Reverse bias is provided from backside contact (not shown here)
- ❑ Simulations are carried out using Silvaco TCAD tool. (For details on meshing, boundary conditions, models etc., see last vertex talk by Mathieu Benoit)



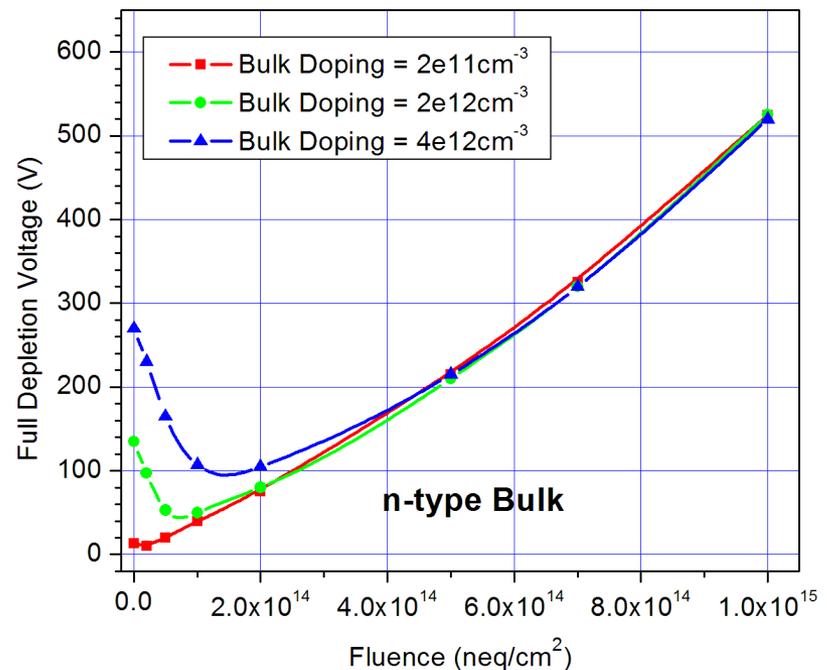
# Bulk + Surface damage model

- Bulk damage model (for proton irradiation)
- Produce experimentally measured currents for irradiated diodes
- Correct full depletion voltages (say, ~500V for  $1e15 \text{ neq/cm}^2$  fluence of proton irradiation)
- Produces electric fields from both sides (See Kramberger's eTCT measurements for charged hadrons, presented in last Vertex)

Trap	Energy Level (eV)	Intro.	$\sigma_e \text{ (cm}^{-2}\text{)}$	$\sigma_h \text{ (cm}^{-2}\text{)}$
Acceptor	$E_C - 0.51\text{eV}$	4	$2 \times 10^{-14}$	$2.6 \times 10^{-14}$
Donor	$E_V + 0.48\text{eV}$	3	$2 \times 10^{-14}$	$2 \times 10^{-14}$



Leakage current comparison (@253K) for  $1 \times 1 \times 300 \mu\text{m}$  diode



Simulated  $V_{FD}$  for different fluences for three n-type different bulk doping

# Double junction effect on $V_{FD}$

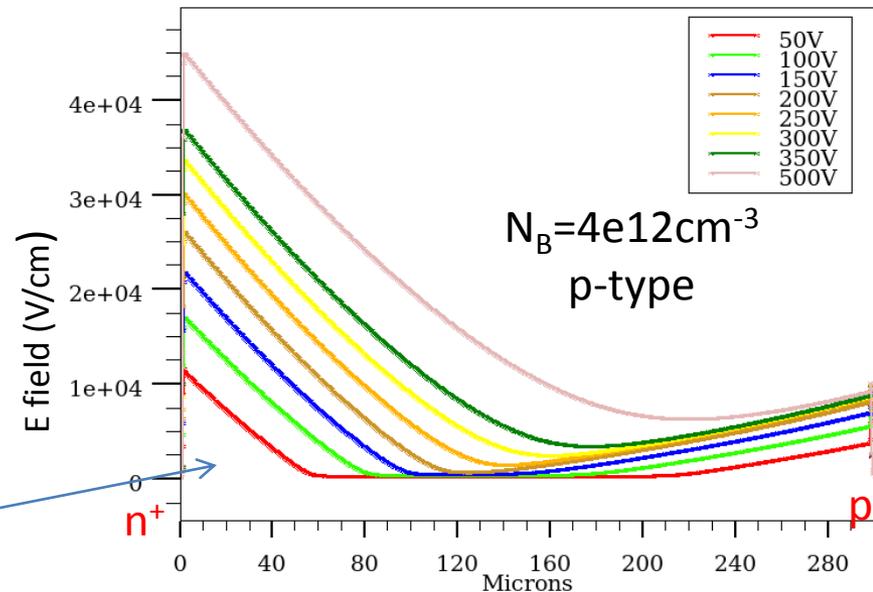
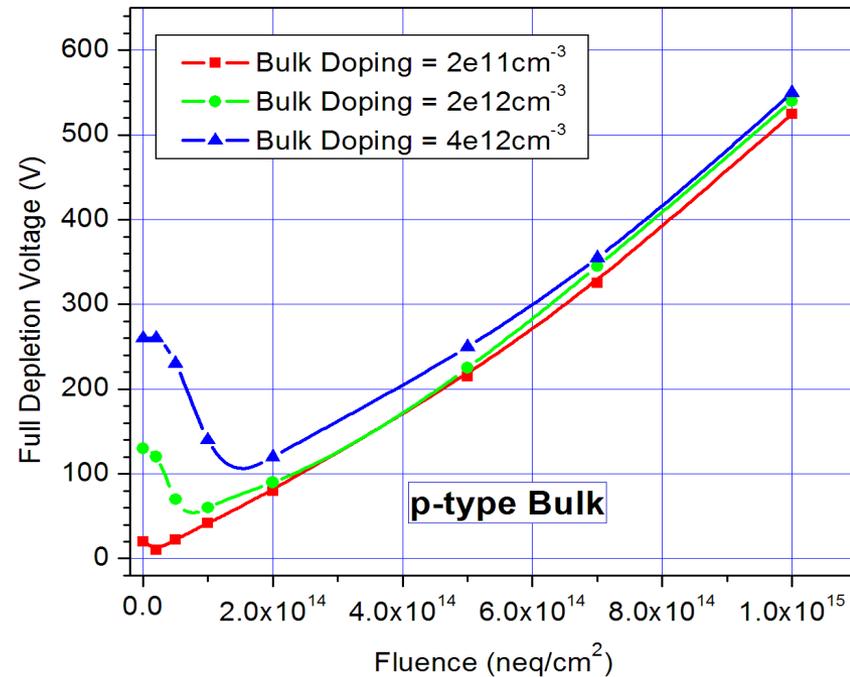
- Initial  $V_{FD}$  drop for all of the three bulk doping
- $V_{FD}$  minimum for bulk doping  $2e11cm^{-3}$  happen at very low fluence but for higher bulk doping,  $V_{FD}$  minimum is at higher fluence

**One is tempted to attribute this effect as “Donor Removal” and “Acceptor Removal” but we have not used any donor/acceptor removal term in the simulations !**

- The initial lowering of  $V_{FD}$  is simply due to the double junction effect in irradiated Si
- Due to the double junction effect, depletion of charge carriers starts from both sides of Si diode
  - Due to depletion from the both ends, Si diodes is depleted at lower  $V_{FD}$  bias , for initial fluences.
  - There may not be any need for the “Donor/Acceptor Removal” terms.

- See my 24<sup>th</sup> RD-50, talk for further discussions!

E field across the diode bulk for different bias for Fluence= $5e14neq/cm^2$

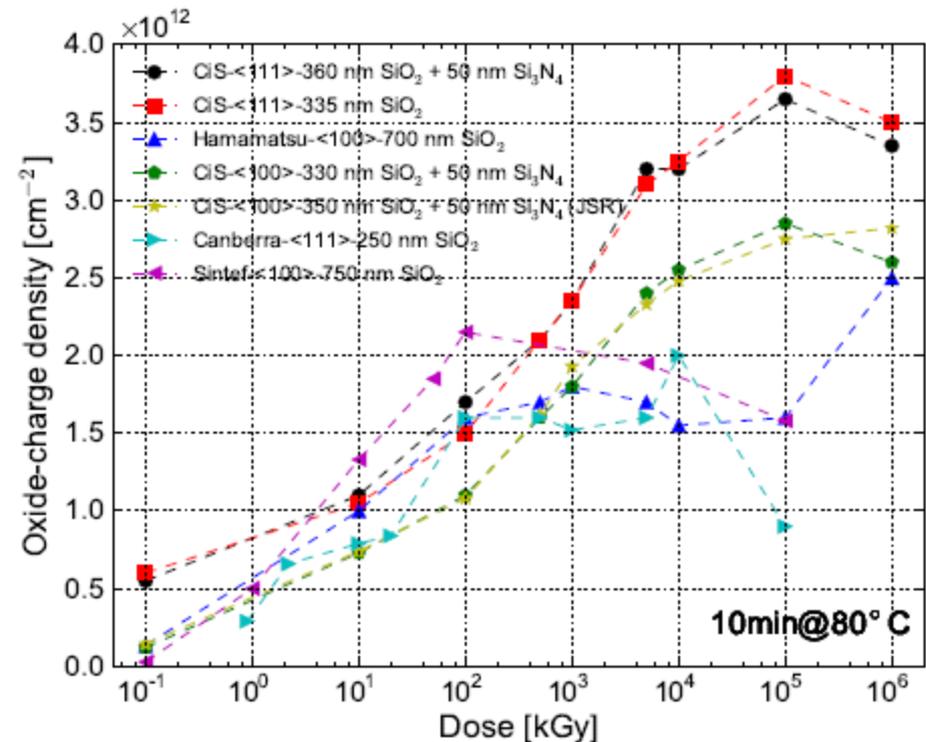


# Surface damage- Oxide charge density ( $Q_F$ )

- Surface damage is incorporated using the Oxide charge density ( $Q_F$ ) and Interface state density ( $N_{it}$ )
- Oxide charge density ( $Q_F$ ) is a complex function of fabrication process, annealing steps, humidity etc.
- Hence, instead of taking one value of  $Q_F$ , Oxide charge density is incorporated in simulations by considering range of  $Q_F$  for a given fluence.

Irradiation fluence (neq/cm <sup>2</sup> )	Range of $Q_F$ (cm <sup>-2</sup> )
0	5e10 to 5e11
1x10 <sup>14</sup>	1e11 to 8e11cm-2
5x10 <sup>14</sup>	5e11 to 1.2e12
1x10 <sup>15</sup>	8e11 to 2e12

**Ranges of Oxide charge density ( $Q_F$ ) used for different fluences**



# Surface damage-Interface state density ( $N_{it}$ )

The interface trap states can play very important role in irradiated Si sensor, because,

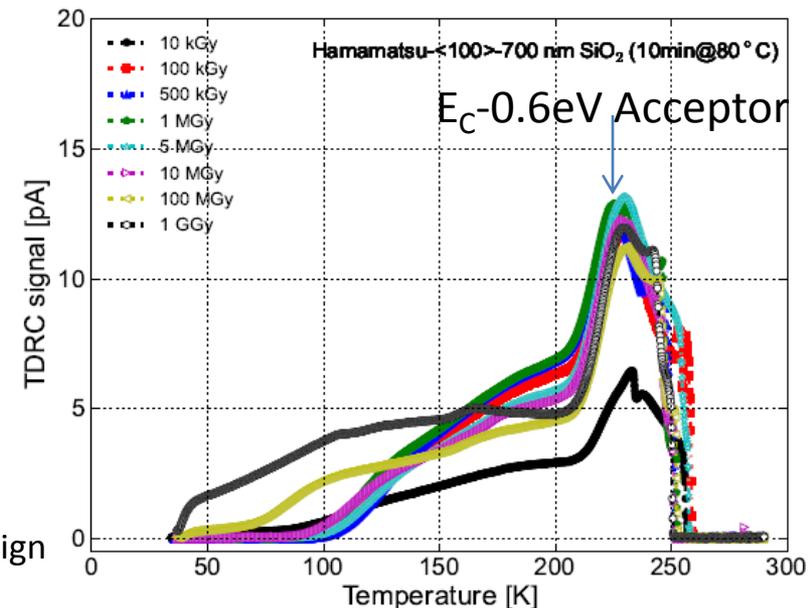
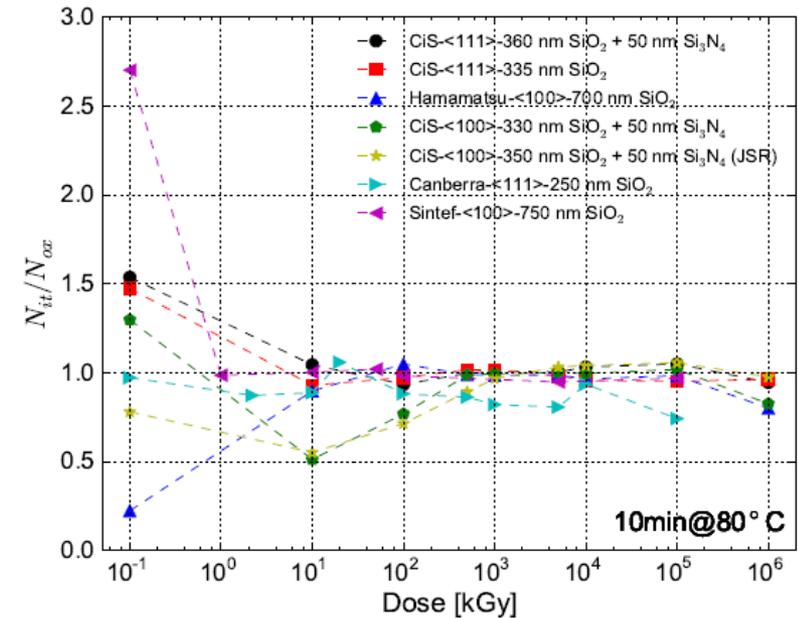
- Its density ( $N_{it}$ ) is comparable to the Oxide charge density ( $Q_F$ ) as shown in the right side plot [1]

- A significant number of  $N_{it}$  states are deep trap states, thus capable of altering the space charge near interface (see the TDRC spectra, right below [1])

- Rint simulations indicates that these interface traps are acceptor type states

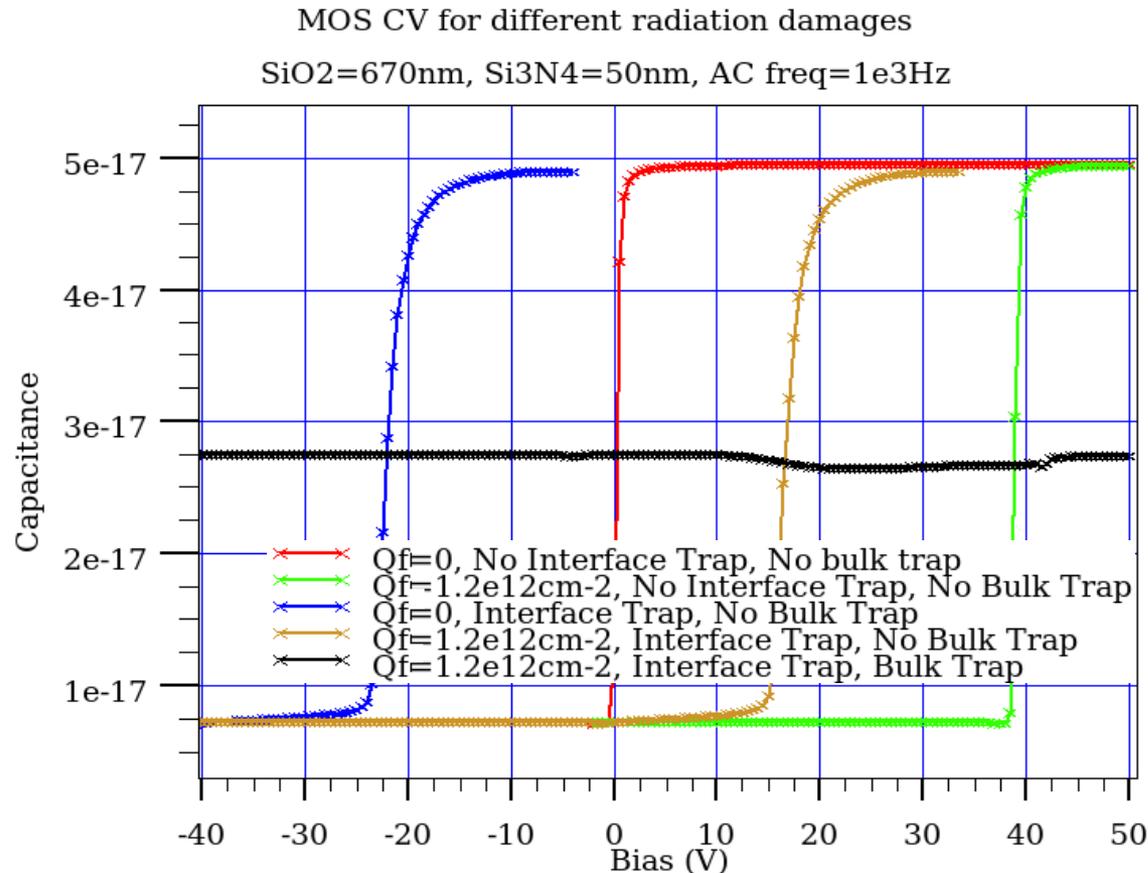
## Assumptions;

1. We have assumed that  $N_{it}$  density is equal to  $Q_F$  density
2. For a given  $N_{it}$ , 60% of the states are deep traps ( $E_C-0.6\text{eV}$ ) and 40% are shallow states ( $E_C-0.39\text{eV}$ ) with  $\sigma_n = \sigma_p = 1\text{e-}15\text{cm}^{-2}$  [1]



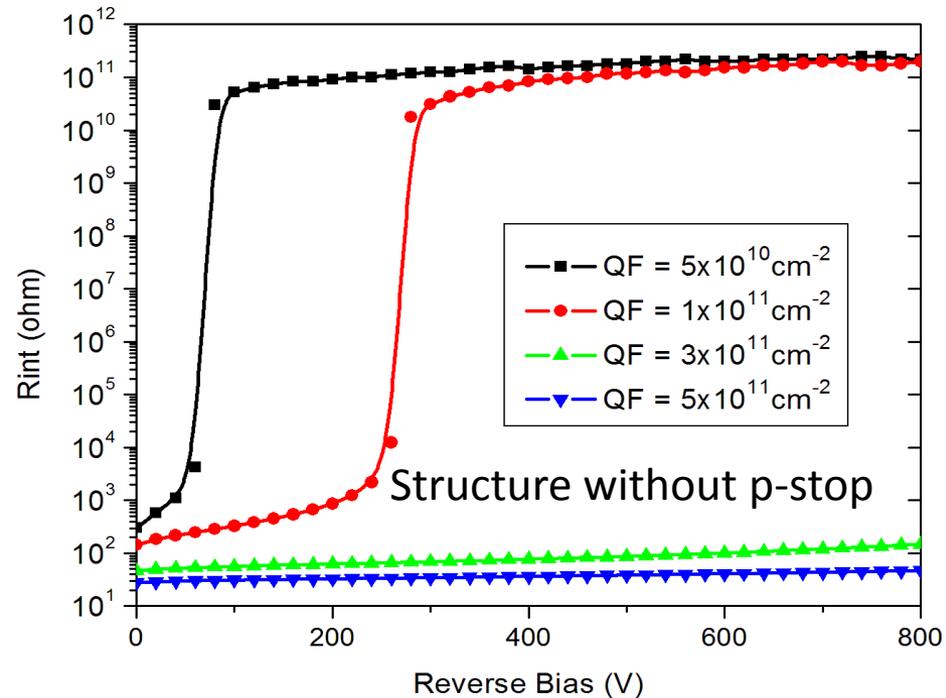
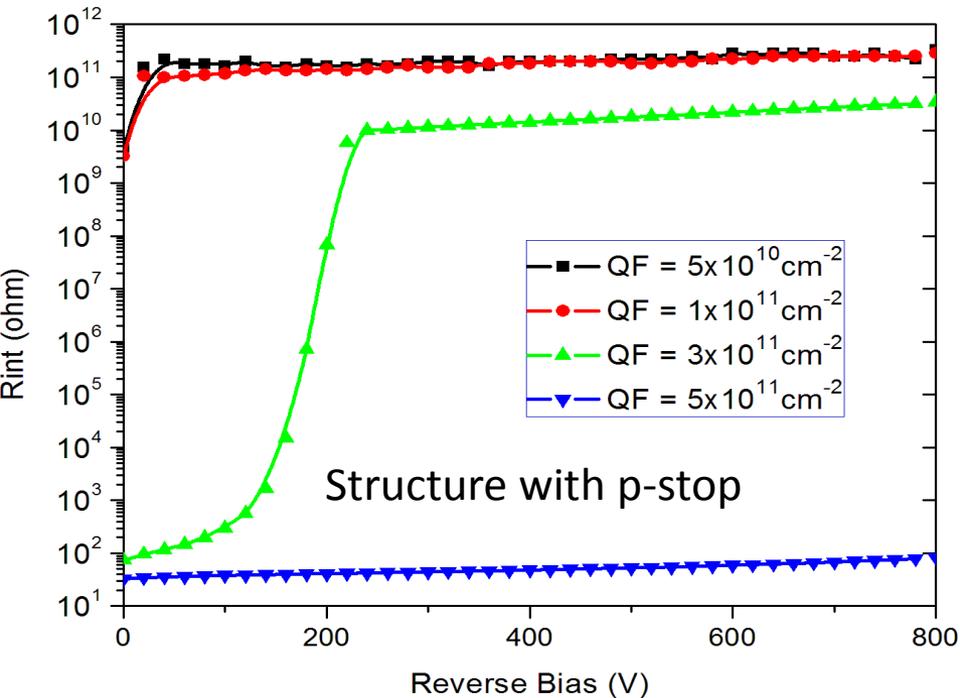
[1] J. Zhang, DESY Thesis-2013, "X-ray radiation damage studies and design of a Si Pixel sensor for different fluences for science at the XFEL"

# Effect of radiation damage - CV for MOS



- Interface states can significantly lower the flatband voltages for MOS
- It is not possible to extract any meaningful flatband voltage for hadron irradiated MOS (CV curve is almost flat for them, as shown by black curve, for fluence=1e15 neq/cm<sup>2</sup>). A similar effect is reported in Maria Bernard-Schwarz's Diploma thesis, Vienna Uni., "Measurements and Irradiation Analysis of Silicon Structures for the CMS Upgrade"

# Simulation of Rint without Bulk damage (Only $Q_F$ variation)



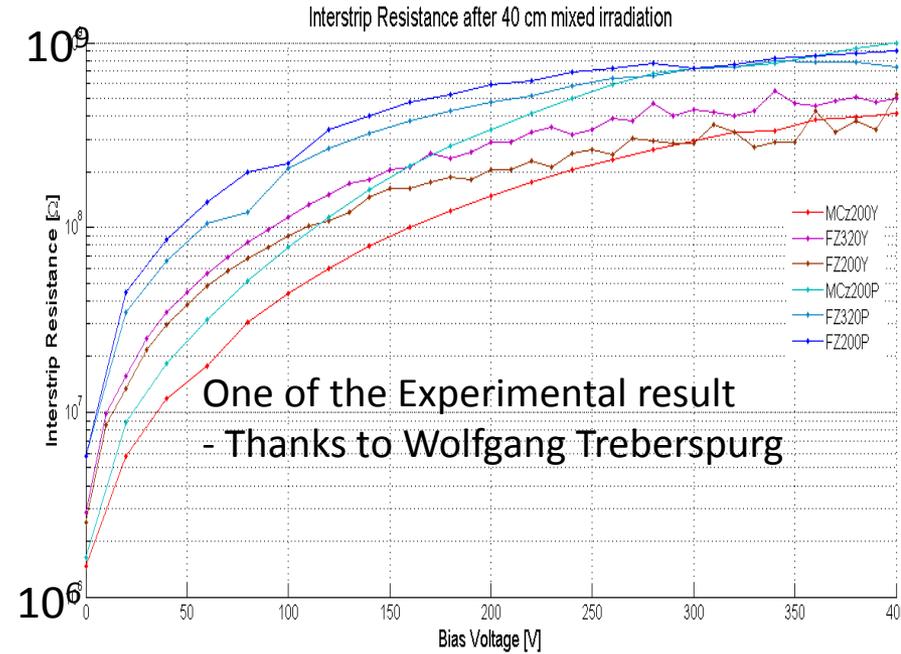
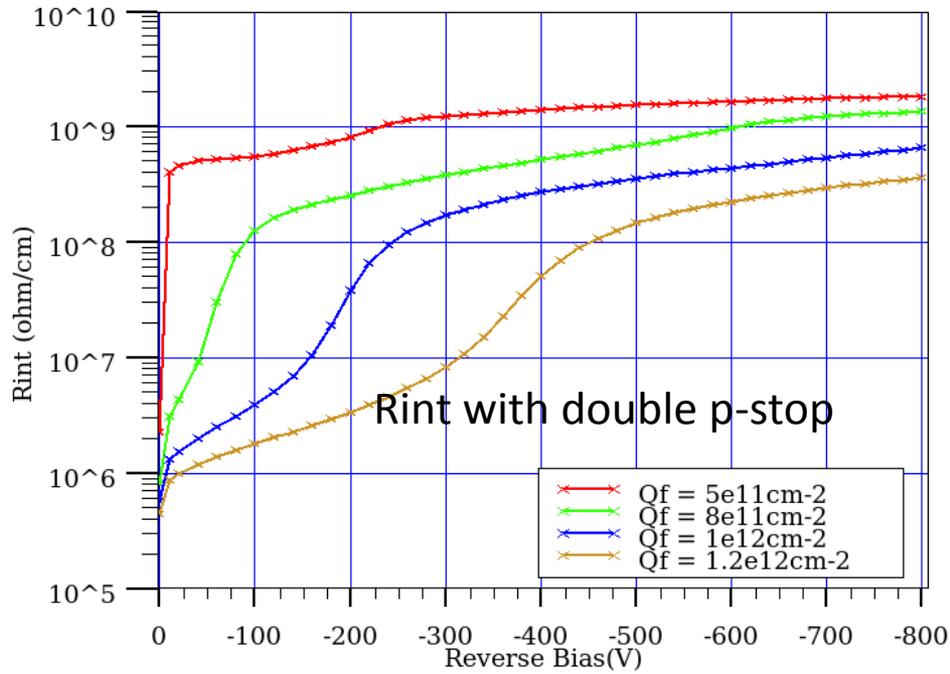
## Three different Rint curves

1. For low values of  $Q_F$ , good strip insulation is obtained even for low bias voltages.
2. For intermediate values of  $Q_F$ , strip insulation is very poor for low voltages, but improves with higher reverse biases, as the electrons from accumulation layer are progressively removed, resulting in a higher Rint.
3. But for higher values of  $Q_F$ ,  $R_{int}$  remains very low up to 800 V.
4. Rint values are lower for structure without pstop.

# Simulations of Rint for Fluence = $5 \times 10^{14} \text{ neq/cm}^2$

Rint for different surface damages for n-on-p strip sensors

Fluence =  $5 \times 10^{14} \text{ neq/cm}^2$ , Temp = 253K, Pitch = 80 micron

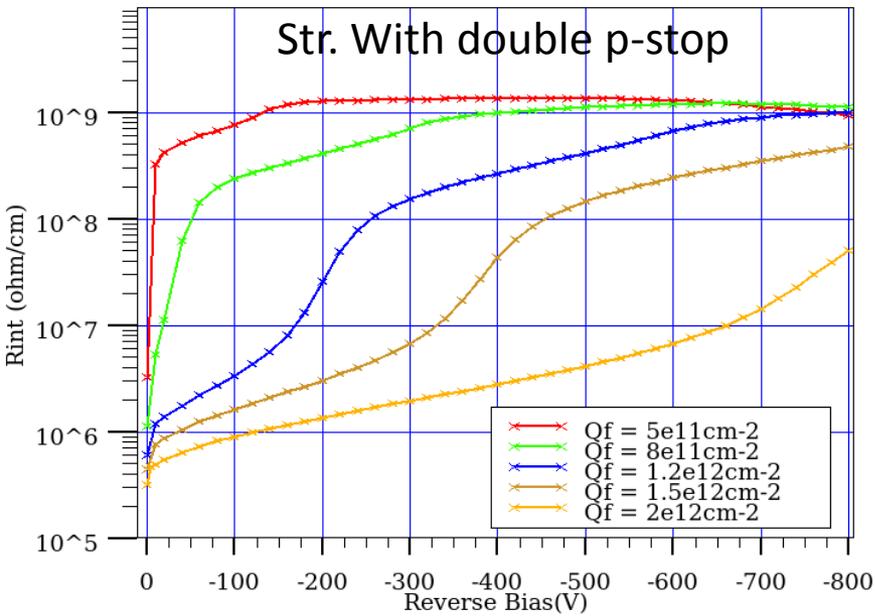


- Both Bulk + Surface damage ( $N_{it} + Q_F$ ) are used
- Rint values of more than 300MΩ is possible for  $Q_F = 1.2 \text{ cm}^{-2}$
- Significant improvement in Rint values, after addition of bulk damage and  $N_{it}$  for higher values of  $Q_F$  (compared to the  $Q_F$  only case, say for  $Q_F = 5 \times 10^{11} \text{ cm}^{-2}$ ,  $n^+$  strips were shorted but after radiation damage, we are having very good Rint)

# Simulations of Rint for Fluence = $1 \times 10^{15} \text{ neq/cm}^2$

Rint for different surface damages for n-on-p strip sensors

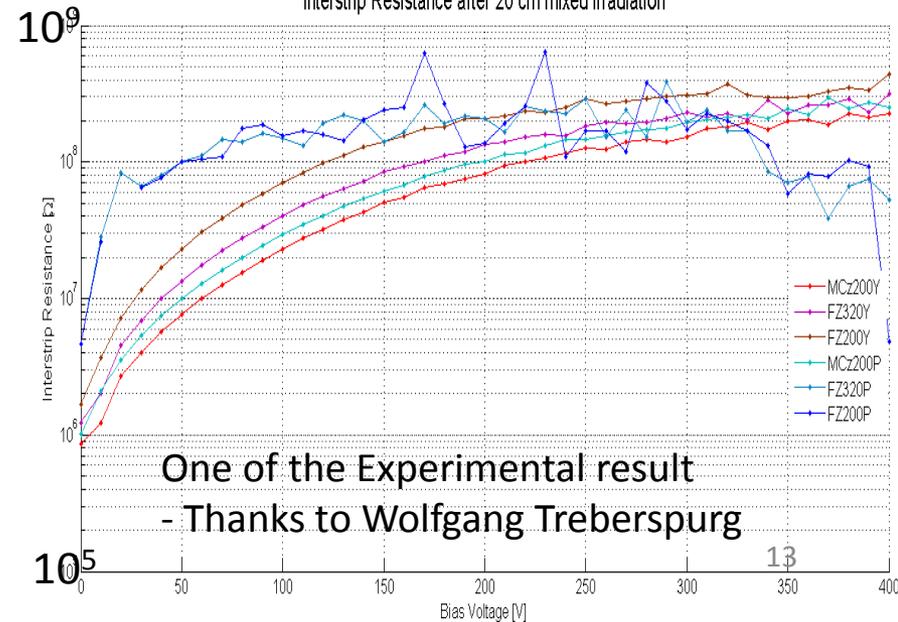
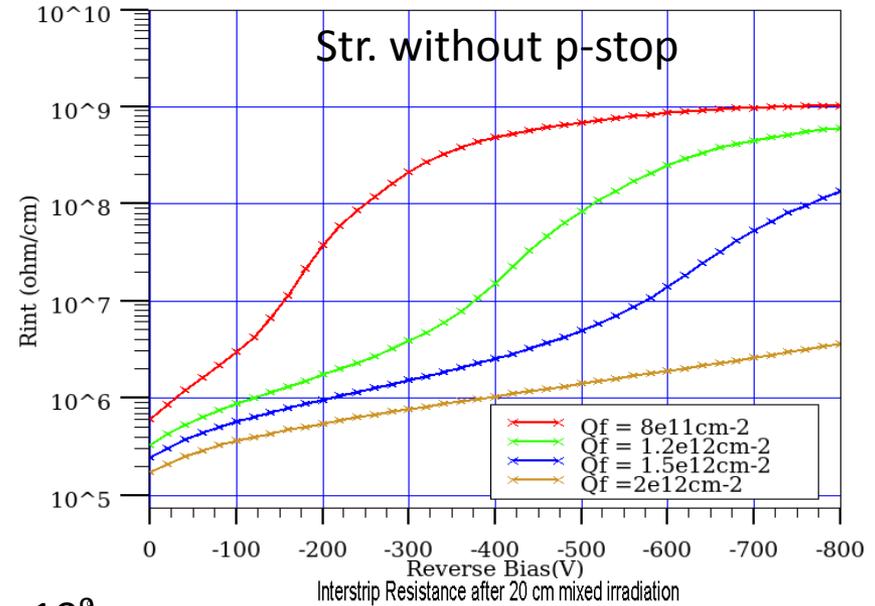
Fluence =  $1 \times 10^{15} \text{ neq/cm}^2$ , Temp = 253K, Pitch = 80 micron



**Str. With double p-stop**

Rint for different surface damages for n-on-p MSSD without Pstops

Fluence =  $1 \times 10^{15} \text{ neq/cm}^2$ , Temp = 253K, Pitch = 80 micron

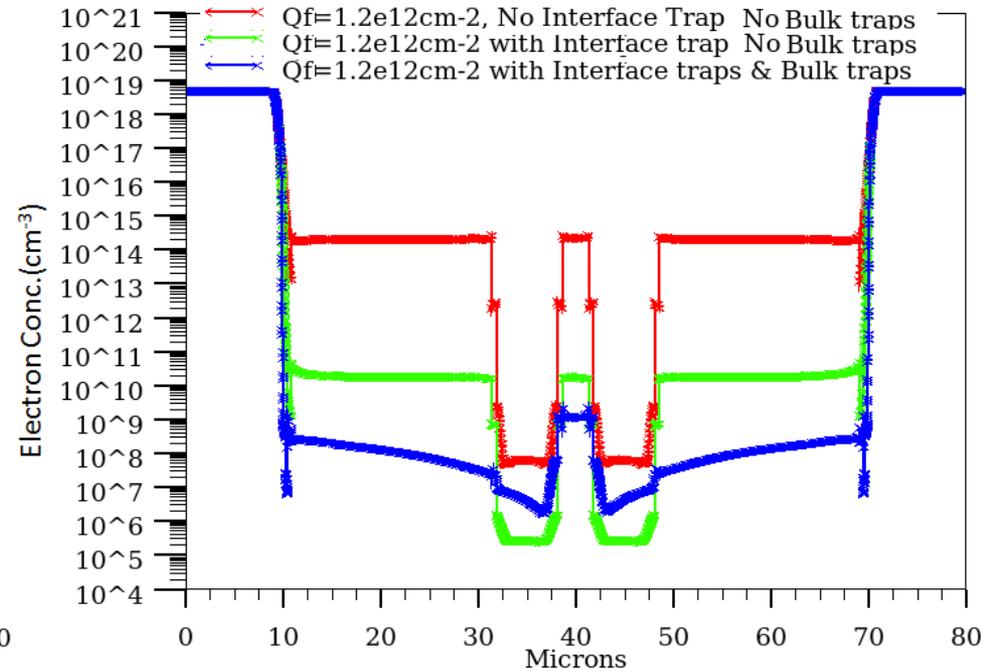
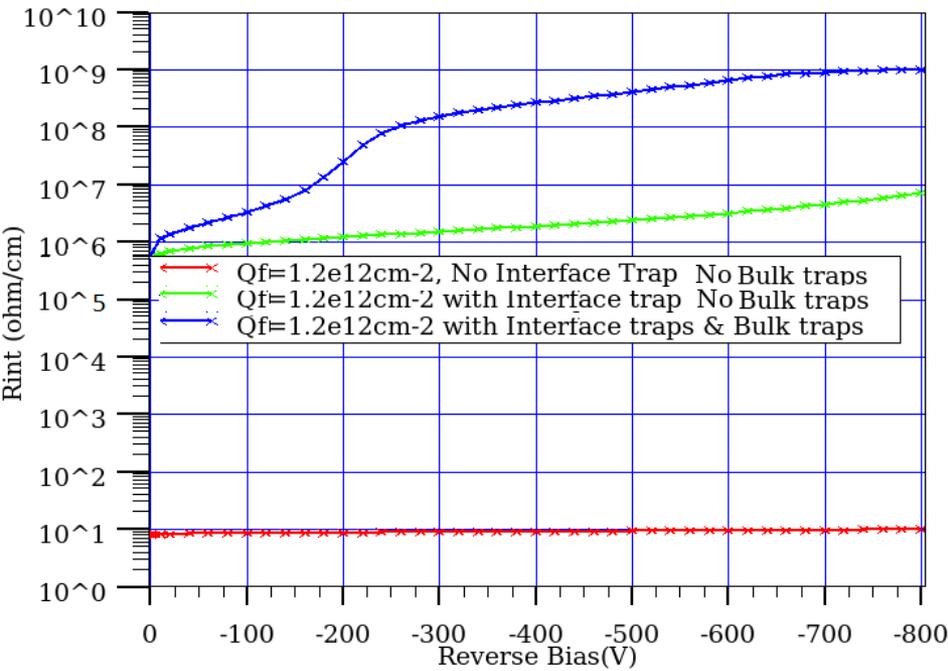


- Rint values of more than 100MΩ is possible for  $Q_f = 1.5 \times 10^{12} \text{ cm}^{-2}$
- Significant improvement in Rint values for higher values of  $Q_f$  (Compared to  $Q_f$  only case)
- Good strip insulation is possible even without any isolation structure! This explain the observed good Rint by Y. Unno et al, (NIM A 579)

# Effect of Bulk and Surface Damages on Rint

Rint for n-on-p MSSD for  $Q_f=1.2e12cm^{-2}$

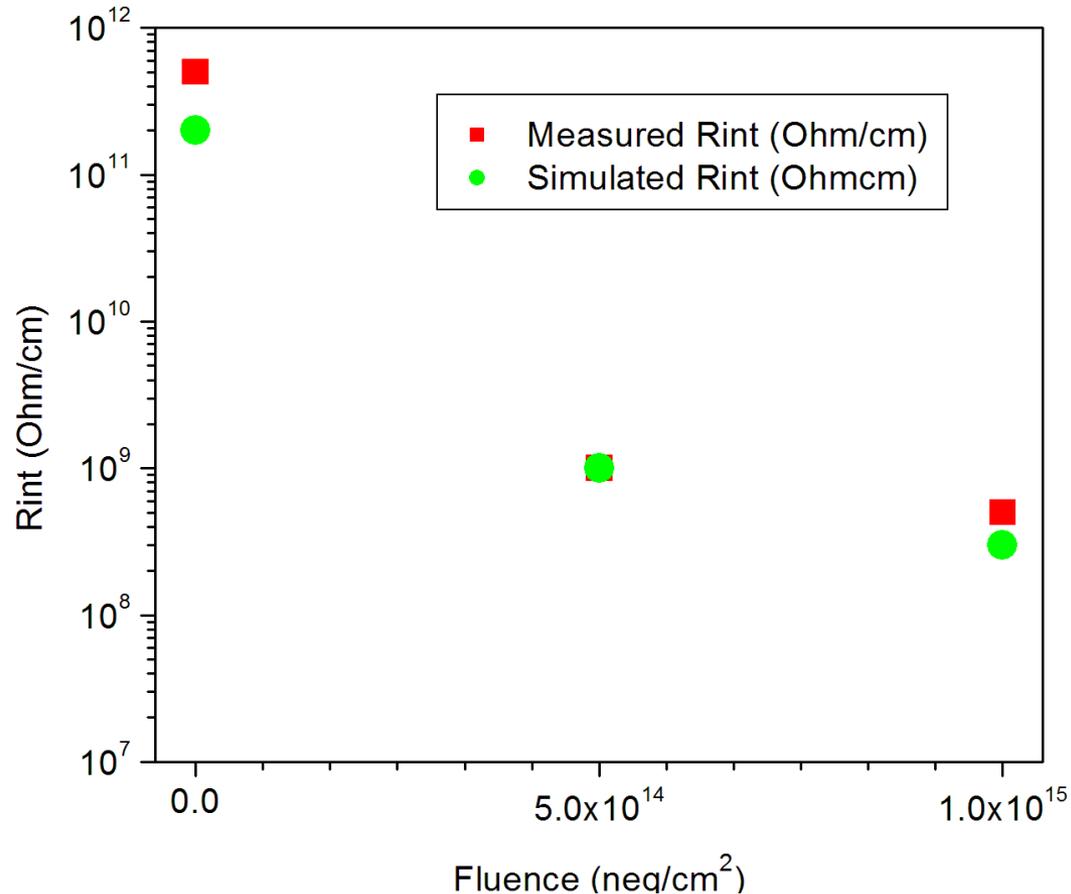
Temp=253K, Pitch= 80 micron



Electron conc. 0.1 $\mu$ m below SiO<sub>2</sub>/Si interface

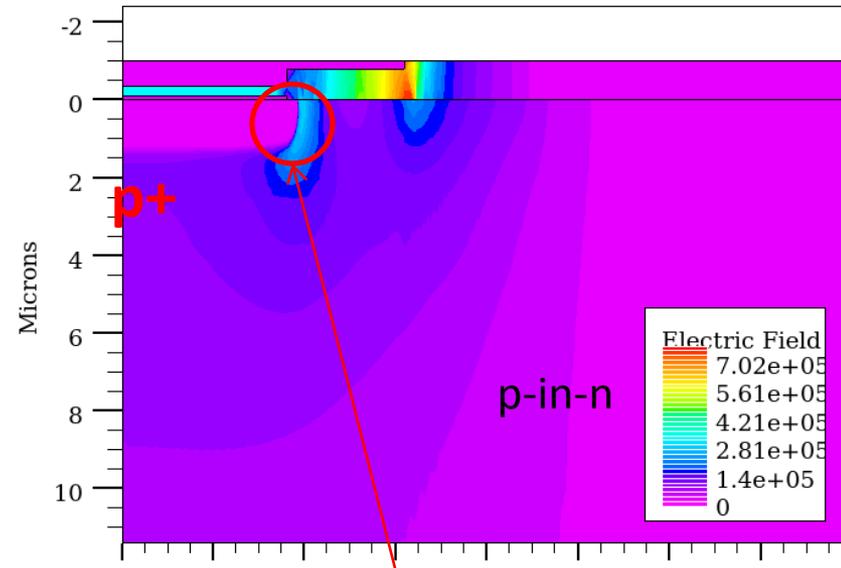
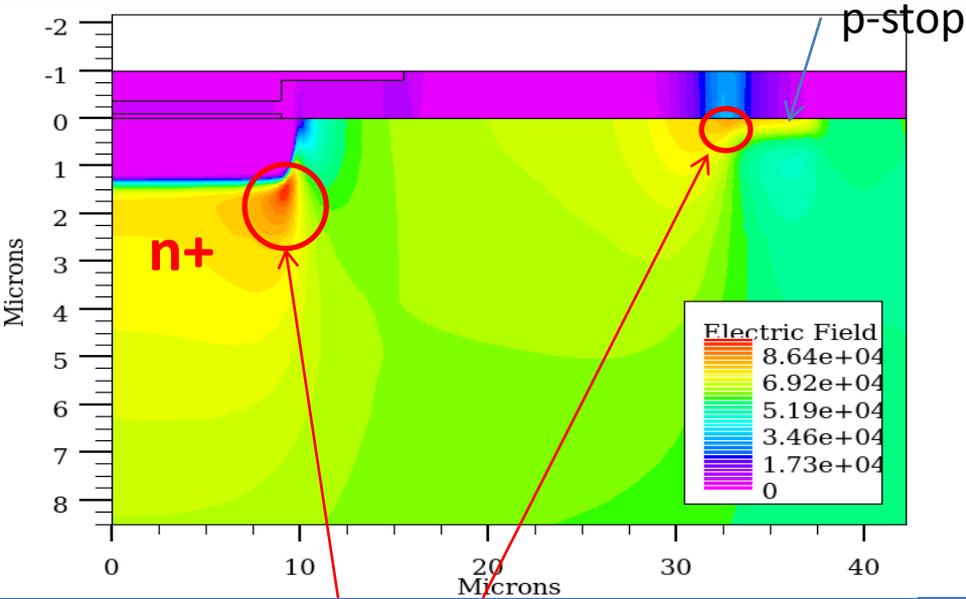
- For a given Oxide charge density value, Rint increases very significantly with inclusion of Interface states
- Rint further increases with inclusion of Bulk traps, which introduces negative space charge near n<sup>+</sup> strip side and further reducing the electron accumulation layer

# Rint simulation and measurements: Comparison



- Good agreement between measured and simulated Rint
- There is variations in measured Rint [4] (for different samples, irradiated with same fluences) and simulated Rint too (assuming different  $Q_F$ s)

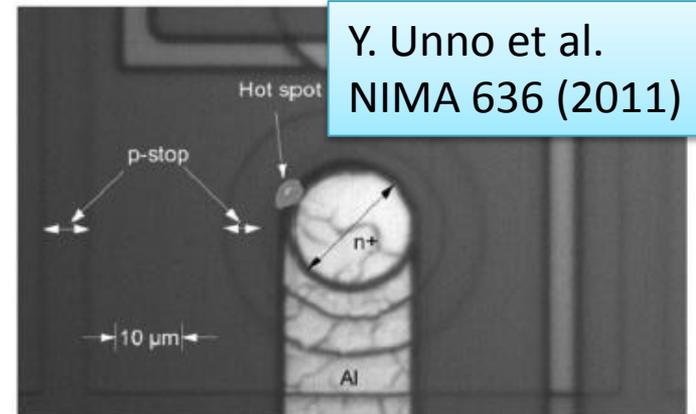
# Maximum E field regions in p-type and n-type sensors



Maximum E field for n-on-p MSSD is near the curvature region of  $n^+$  strip Or just near p-stop, just below  $\text{SiO}_2/\text{Si}$  interface - Shown by cutline  $1.3\mu\text{m}$  below  $\text{SiO}_2/\text{Si}$  interface

Maximum E field for p-on-n MSSD is just below  $\text{SiO}_2/\text{Si}$  interface near  $p^+$  strip - Shown by cutline  $0.1\mu\text{m}$  below  $\text{SiO}_2$

- For p-stop (at least) up to  $2 \times 10^{16} \text{cm}^{-3}$ , highest E field is near  $n^+$  strip curvature (see R. Dalal et al, 23<sup>rd</sup> RD-50)
- For low and intermediate p-stops dopings densities, it is quite possible, that microdischarges are taking place at  $n^+$  curvature.
- Explanation need Bulk + Surface damage. Use of  $Q_F$  only for simulations leads to wrong conclusion!



Y. Unno et al. NIMA 636 (2011)

Fig. 1. Hot spot of microdischarge observed in an n-in-p sensor by an infrared camera. The spot is at the edge of the  $n^+$  electrode.

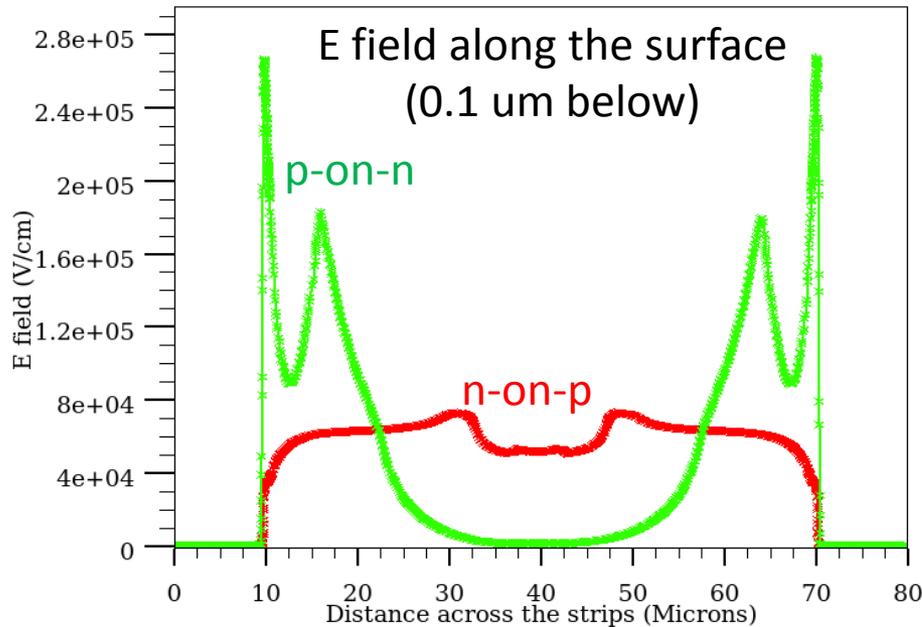
# Effect of very high E field in irradiated sensor

- E field inside irradiated sensors is a strong function of space charge .
- Very high E field in a region can initiate avalanche in that region.
- Once avalanche is started, a lot of free e/h pairs are produced which will compensate the nearby space charge, changing the electric field, thus stopping the further breakdown.

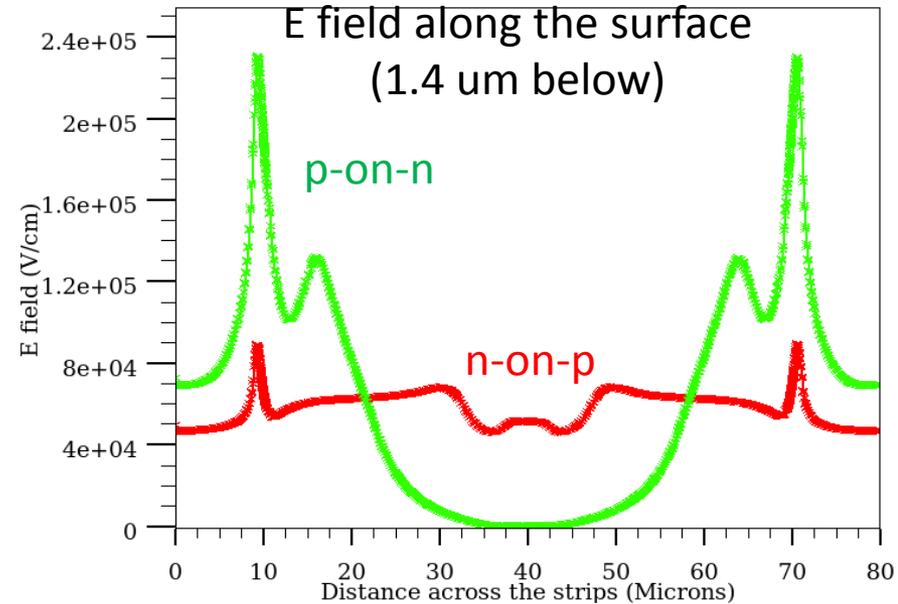
This mechanism may stop the avalanche from turning global and continuous. Thus high E field near the strips can be a reason of non – Gaussian noise events which occur randomly (RGH- Random Gaussian Hits)

# E. Field (Irradiated) comparison : p-in-n & n-in-p sensor

E field for p-on-n and n-on-p type strip sensors for Fluence= $1e15cm^{-2}$   
 $Q_f=1.2e12cm^{-2}$ , Bias=500V, Cutline is 0.1um below SiO<sub>2</sub>/Si



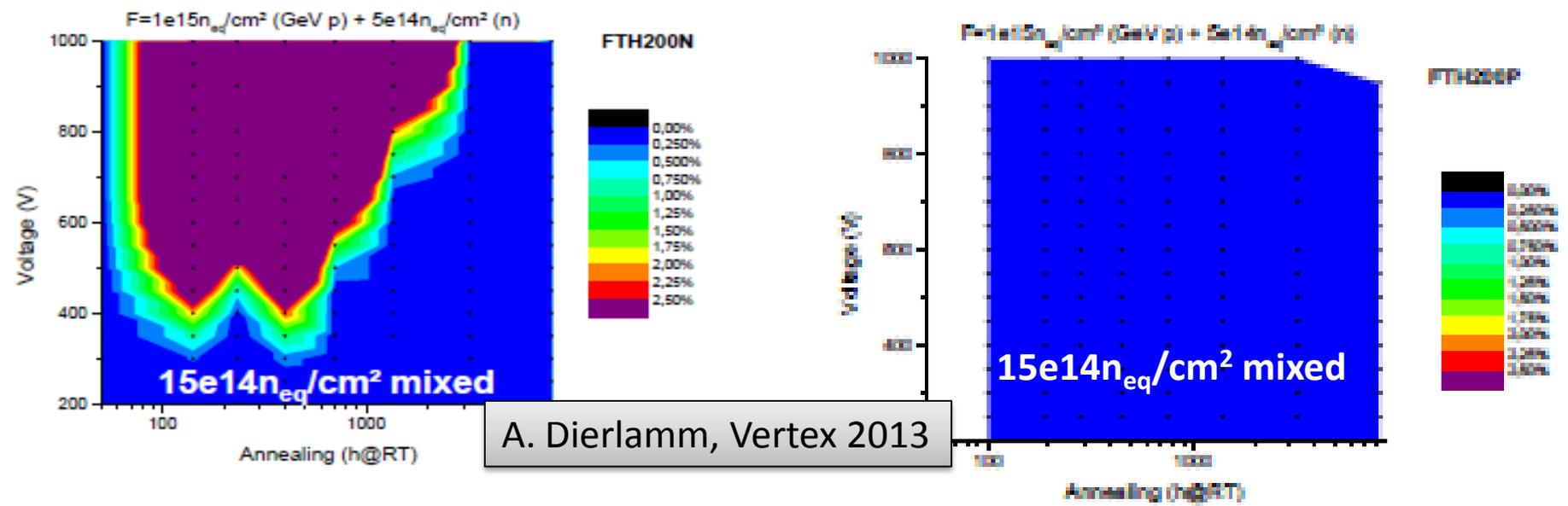
E field for p-on-n and n-on-p type strip sensors for Fluence= $1e15cm^{-2}$   
 $Q_f=1.2e12cm^{-2}$ , Bias=500V, Cutline is 1.4um below SiO<sub>2</sub>/Si



Flux =  $1 \times 10^{15} cm^{-2}$  ;  $Q_f = 1.2 \times 10^{12} cm^{-2}$  ; Bias = 500 V

- Peak electric field is more for p-in-n (n-type) sensor as compared to n-in-p (p-type) sensor for a given bias.
  - Micro-discharge possibility is much more in p-in-n sensors.

# Amount of RGH for irra. n-type and p-type of sensor



N-type sensor

P-type sensor

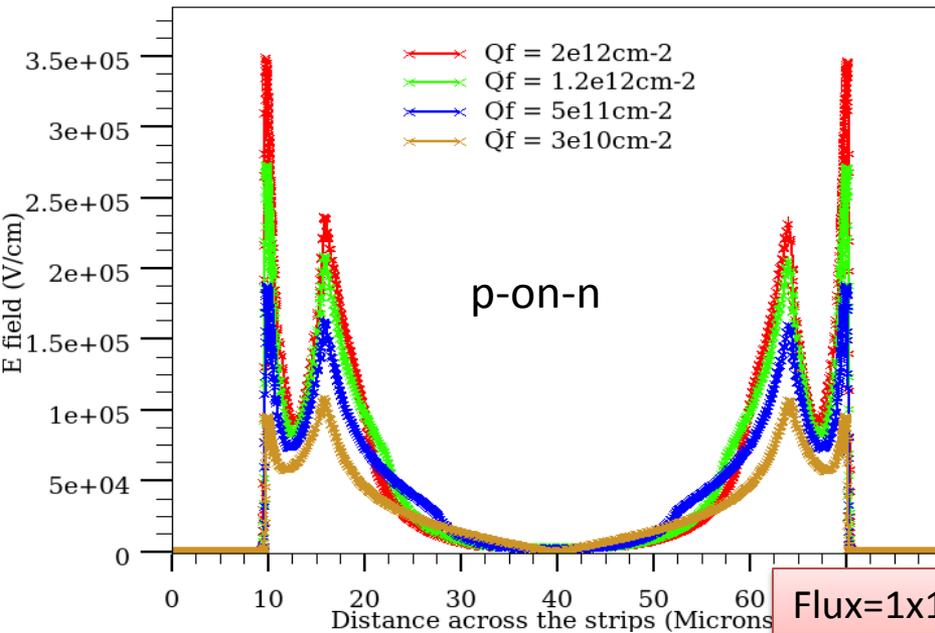
- Significant amount of non-Gaussian noise (RGH) observed in p-in-n sensor (n-type)
- Very less amount of RGH rates observed for n-in-p (p-type) sensor



# E. Field (Irradiated) : Effect of $Q_F$

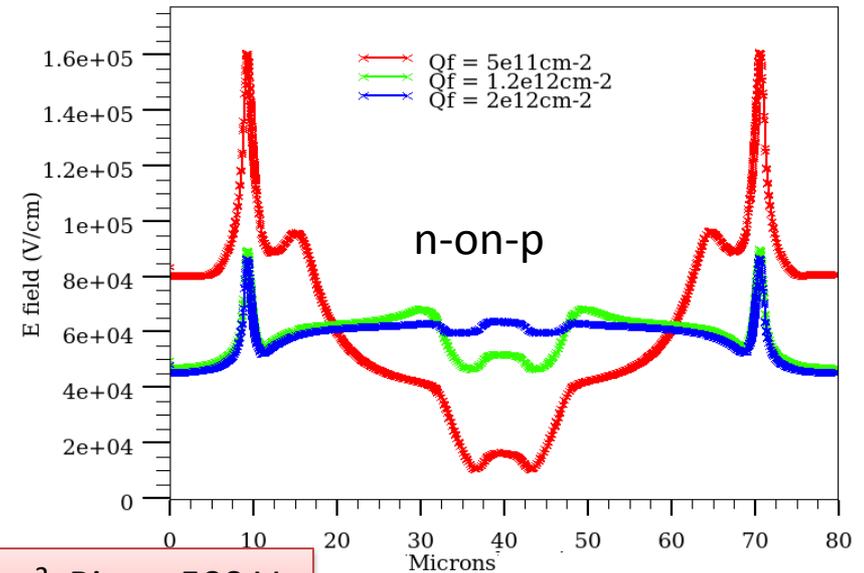
E field for p-on-n strip sensors for different surface damage

Fluence= $1e15$ neq/cm<sup>2</sup>, Bias=500V, Temp=253K



E field for different surface damages for n-on-p strip sensors

Fluence= $1e15$ neq/cm<sup>2</sup>, Temp=253K, Bias=500V

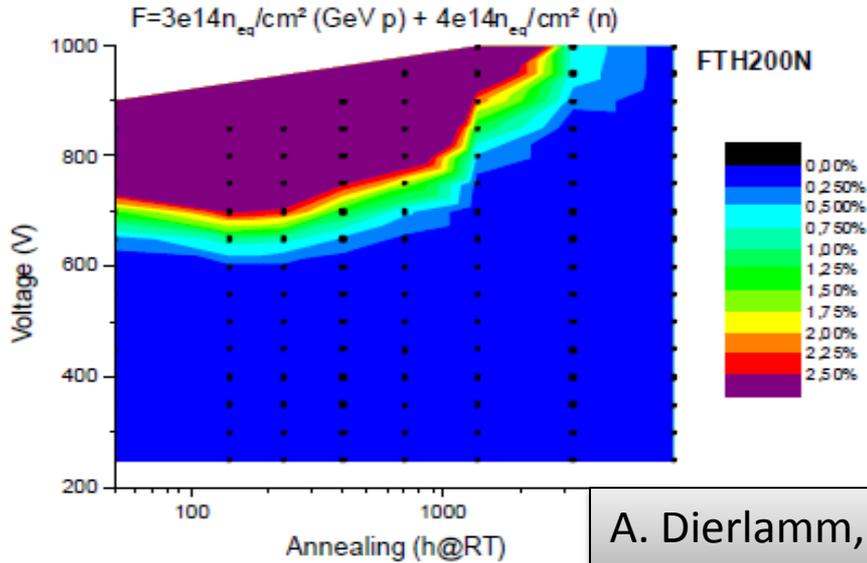


Flux= $1 \times 10^{15}$ cm<sup>-2</sup> Bias = 500 V

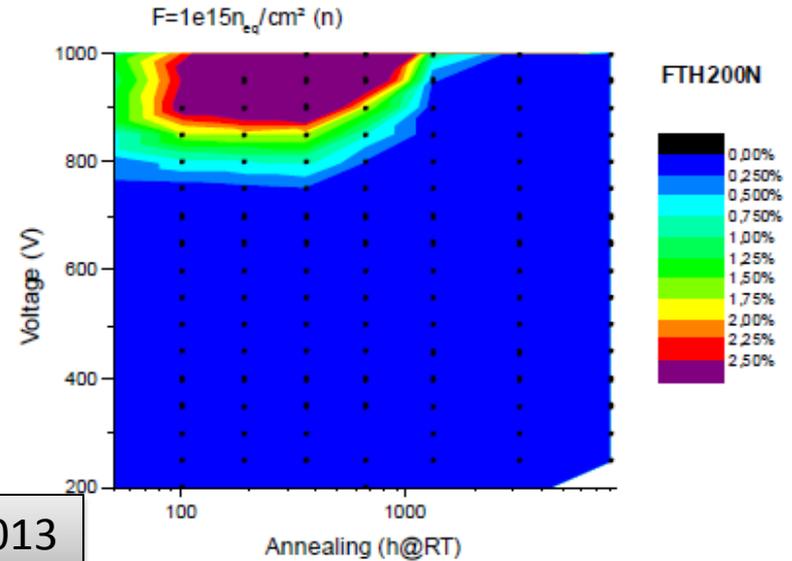
- **p-on-n strip sensors**
  - As  $Q_F$  increases  $\Rightarrow$  Peak Efield increases.
  - Micro-discharge possibility is more for p-on-n strip sensors after proton irradiation or less possibility after neutron irradiation
- **n-on-p strip sensors**
  - Peak field is much less compare to p-in-n sensors
  - As  $Q_F$  increases  $\Rightarrow$  Peak E field decreases.
  - Charge multiplication should be more for neutron irradiated n-on-p strip sensors compare to proton irradiated n-on-p (for same fluence)

# RGH for different type of irradiation (for p-in-n)

$7e14n_{eq}/cm^2$  mixed

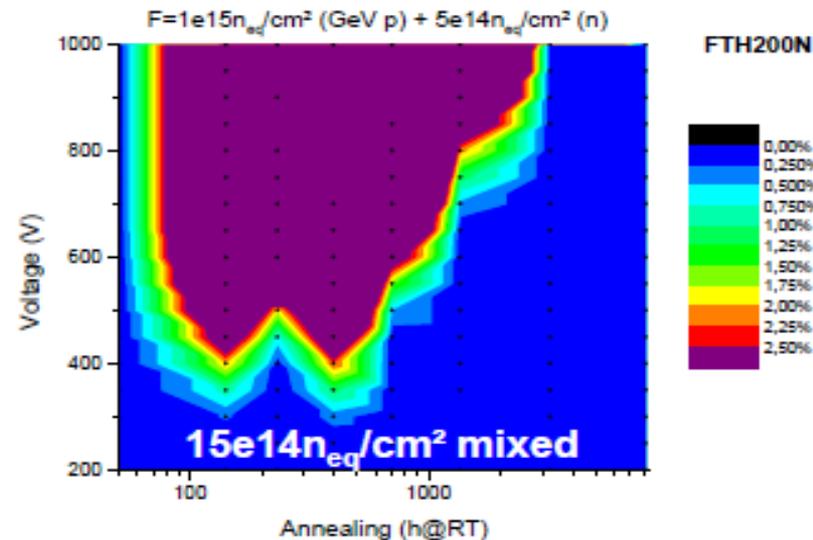


$10e14n_{eq}/cm^2$  neutrons



A. Dierlamm, Vertex 2013

- Much less RGH for neutron irradiation clearly indicate the role of surface damage
- The dependence on ionizing radiation hints toward a combined effect of bulk damage and surface damage
- The  $Q_F$  effect simulation were carried out before these measurements

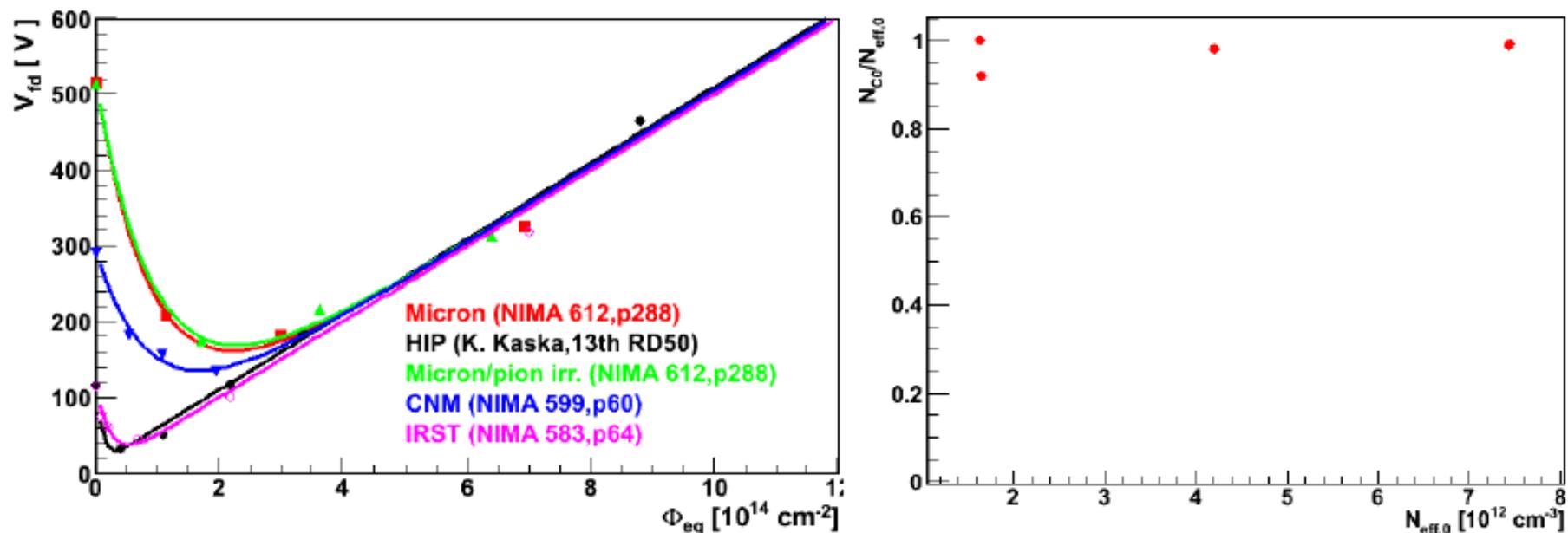


# Summary

- Both bulk and surface damages take place during the hadron irradiation of Si sensors
  - Simultaneous incorporations of bulk damage and surface damage is essential for correct interpretation of observed radiation damage effects
    - Use of only bulk damage or only surface damage may lead to wrong conclusions!
  - Use of interface states, along with Oxide charge density is essential for proper surface damage inclusion in simulations
    - Conc. of  $N_{it}$  and  $Q_F$  (or  $N_{OX}$ ) are comparable
    - Significant fraction of deep traps in  $N_{it}$
  - There may be initial lowering of  $V_{FD}$  with irradiation
    - Double junction effect
    - Acceptor/Donor removal may be avoided
  - Interface trap states and negative space charge, near  $n^+$ , due to bulk damage results in good  $R_{int}$  values for low  $p_{stop}$  doping and even without any  $p_{stop}$
  - Highest electric field in irradiated n-on-p is much lower compared to irradiated p-on-n
- Most of the present work is carried out under RD-50 and CMS Phase-II upgrade
- I am thankful to Michael Moll for lots of suggestions and interesting discussion
- Really thankful to many people for their excellent measurements, which are used to understand the various aspect of radiation damage (like J. Zhang, T. Pohlson, Kramberger, Hadi, Y. Unno, Wolfgang... and many more)
- Finally, many thanks to 23<sup>rd</sup> Vertex Organizers for financial assistance and perfect conference !

*Thanks for your attention!*

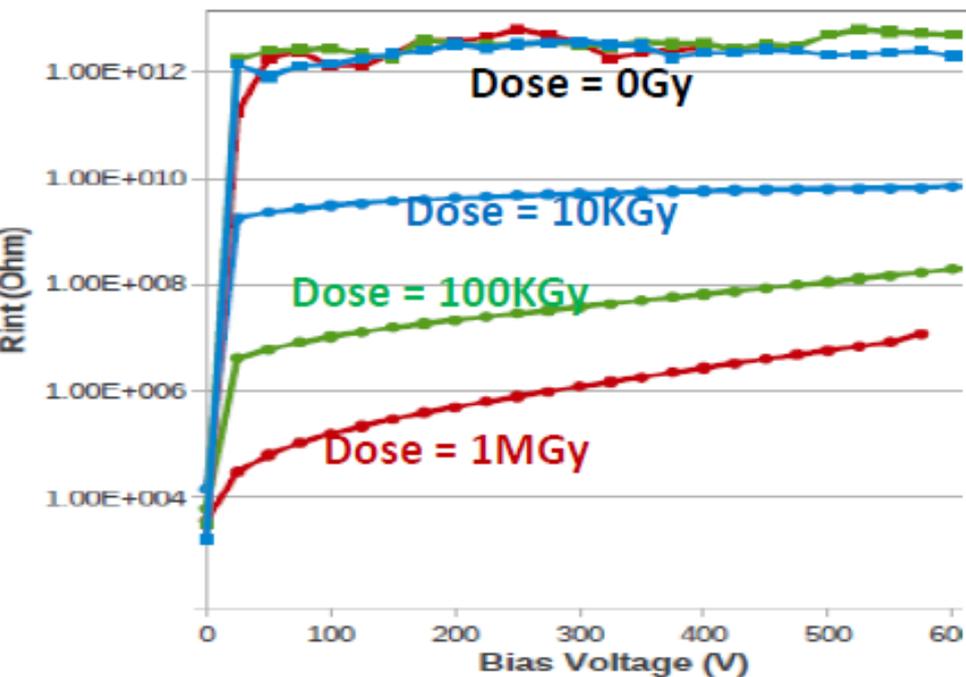
# MCz-p irradiated with charged hadrons



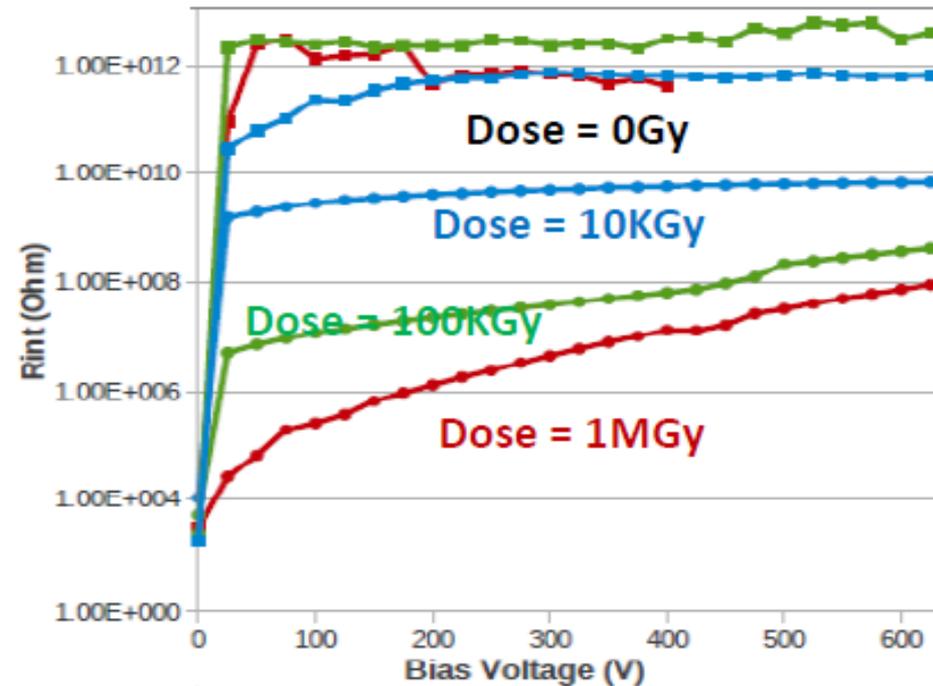
- $g_{eff} = 0.0071$  cm $^{-1}$  (taken from O rich measurements from RD48/50) and seems to be adequate,  $c$  and  $N_{C0}$  were determined from the fit.
- Different producers – **no impact of processing on behavior**
- *Acceptor removal seems to be complete*
- $c \sim 1 \cdot 10^{-14}$  cm $^2$  (seems larger for lower resistivity, but uncertainty is too large for any firm conclusion)

# Rint measurements for irradiated DC-CAP Sensors

- Very high Rint ( $\sim 1 \times 10^{12}$  Ohm for unirradiated DC-CAD test structures)
- Rint decreases with irradiation dose
- For highest dose of 1MGy, Rint is around  $1 \times 10^6$  Ohm, Strip Isolation is effectively lost.



**Pitch = 120 μm**



**Pitch = 80 μm**

Rint measurements by Anna Peisert and Hadi, CERN for different X-ray dose

# What is going inside hadron irradiated sensors !

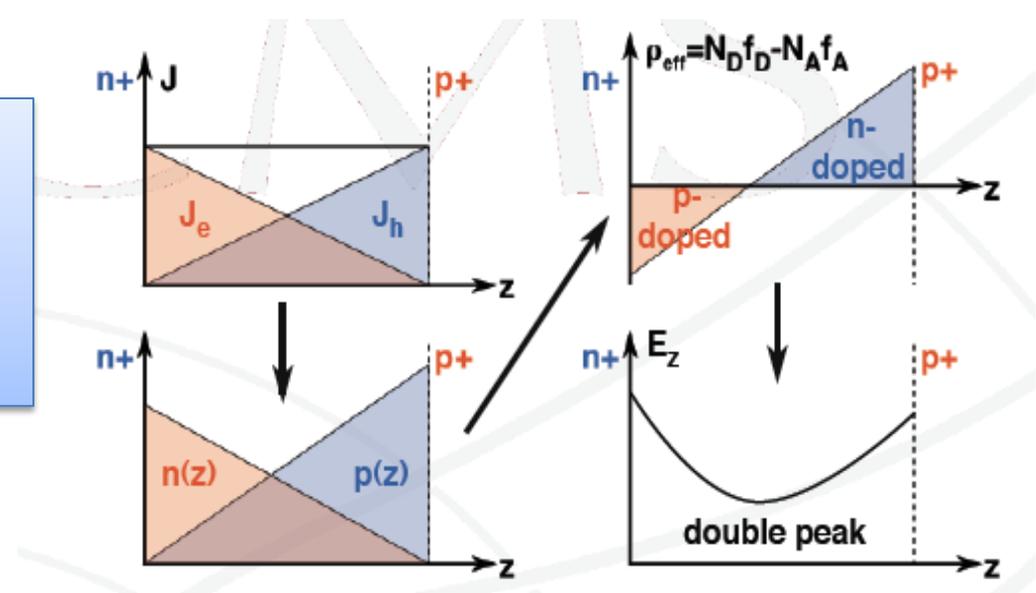
Irradiation of n+p-p+ Si sensor by hadrons :

- Acceptor and Donor traps are created
- Deep traps leads to quite higher leakage current
- Electrons move toward n+ strips while holes move toward p+ backside
- Electron density near n+ is very high leading to filling of Acceptor traps and thus creating negative space charge near n+ strip.
- Similarly, positive space charge is created near p+ by filling of Donor traps by holes

**High negative space charge near n+ strips result in high E-fields near strips (similar for pixel)**

**Can we see this....?**

**Yes by eTCT**



# Two type of irradiation

## ❖ Irradiation with Photons (x-ray and $\gamma$ -ray irradiation)

- Only surface damage is significant, resulting in very high  $Q_f$  (see backup slides)
- Leakage current is very low,  $\alpha$  is at least three orders lower than hadron irradiation and no effect of annealing (very low bulk damage, M.Moll thesis)

For this type irradiation : No High electric field near strips

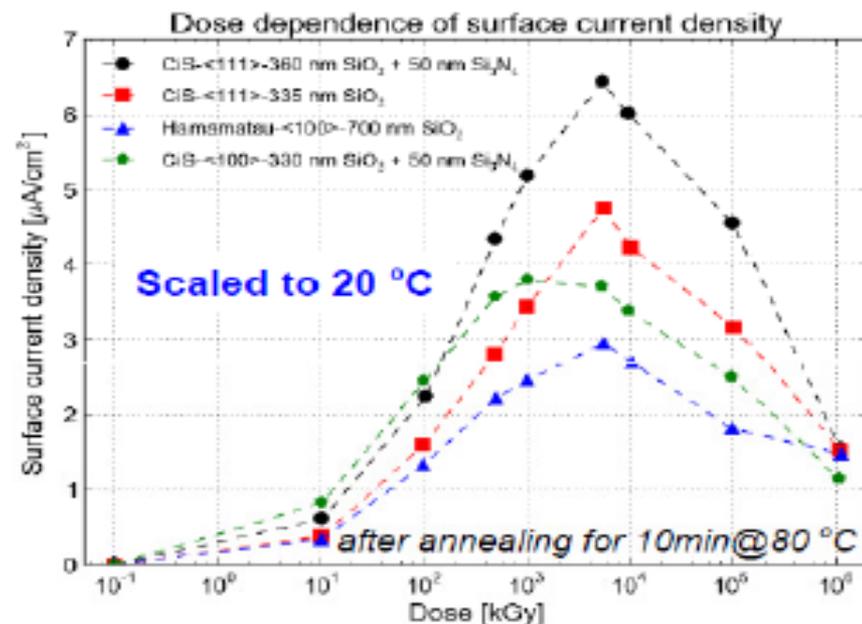
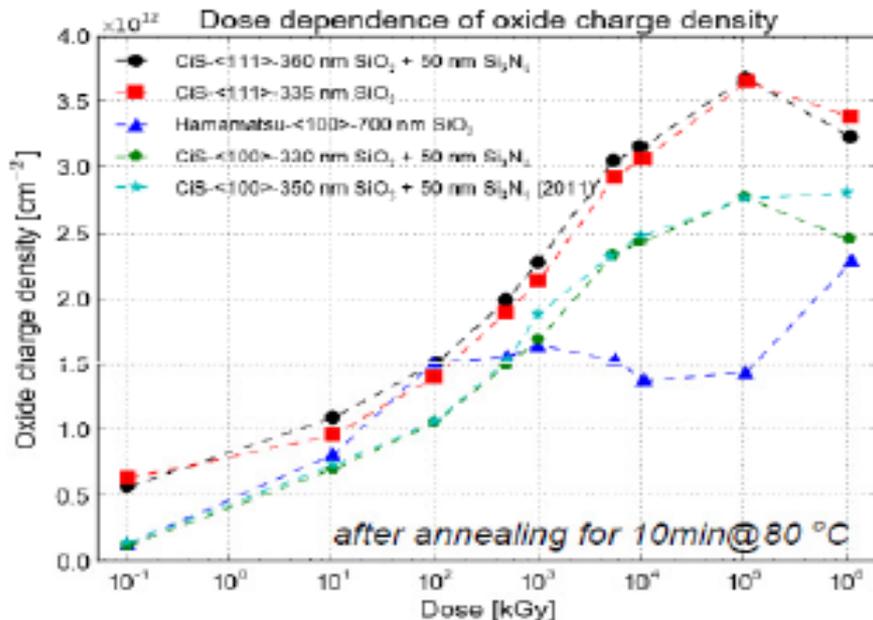
- Oxide charge density  $\sim 2-3 \times 10^{12} \text{ cm}^{-2}$  after irradiation  $\sim 1 \text{ MGy}$  (in MOS as well as in strips and pixel sensors), leading to very serious problems for isolation,  $C_{int}$  & breakdown.

## ❖ Irradiations with Hadrons (p,n or pions irradiations)

- Along with surface damage, significant bulk damage  $\longrightarrow$  very high leakage current
- This leads to high Electric field or high density of negative space charge near the  $n^+$  strips.
- No  $e^-$  accumulation layer formation .... No problems for strip isolation,  $C_{int}$  etc.
- Measurements using MOS will show expected high  $Q_f$  as there is no high leakage current, so, no negative space charge near Si/SiO<sub>2</sub> junction (no suppression of  $Q_f$ ).

### 3. Summary: Dose Dependence of $N_{ox}$ and $J_{surf}$

**Vendors:** CiS, Hamamatsu, Canberra; **Crystal orientations:**  $\langle 111 \rangle$ ,  $\langle 100 \rangle$ ;  
**Insulator:**  $SiO_2$  (335-700 nm), with and without additional 50 nm  $Si_3N_4$



- Results reproducible (after some annealing)
- Spread of about a factor 2
- $N_{ox}$  saturates for ~1 - 10 MGy
- $J_{surf}$  peaks at 1-10 MGy, then decreases

- Equilibrium h-trapping and eh-recombination ?
- E-field effects due to oxide charges ?
- Understanding needs more studies

J.Zhang et al., arXiv:1210.0427(2012)

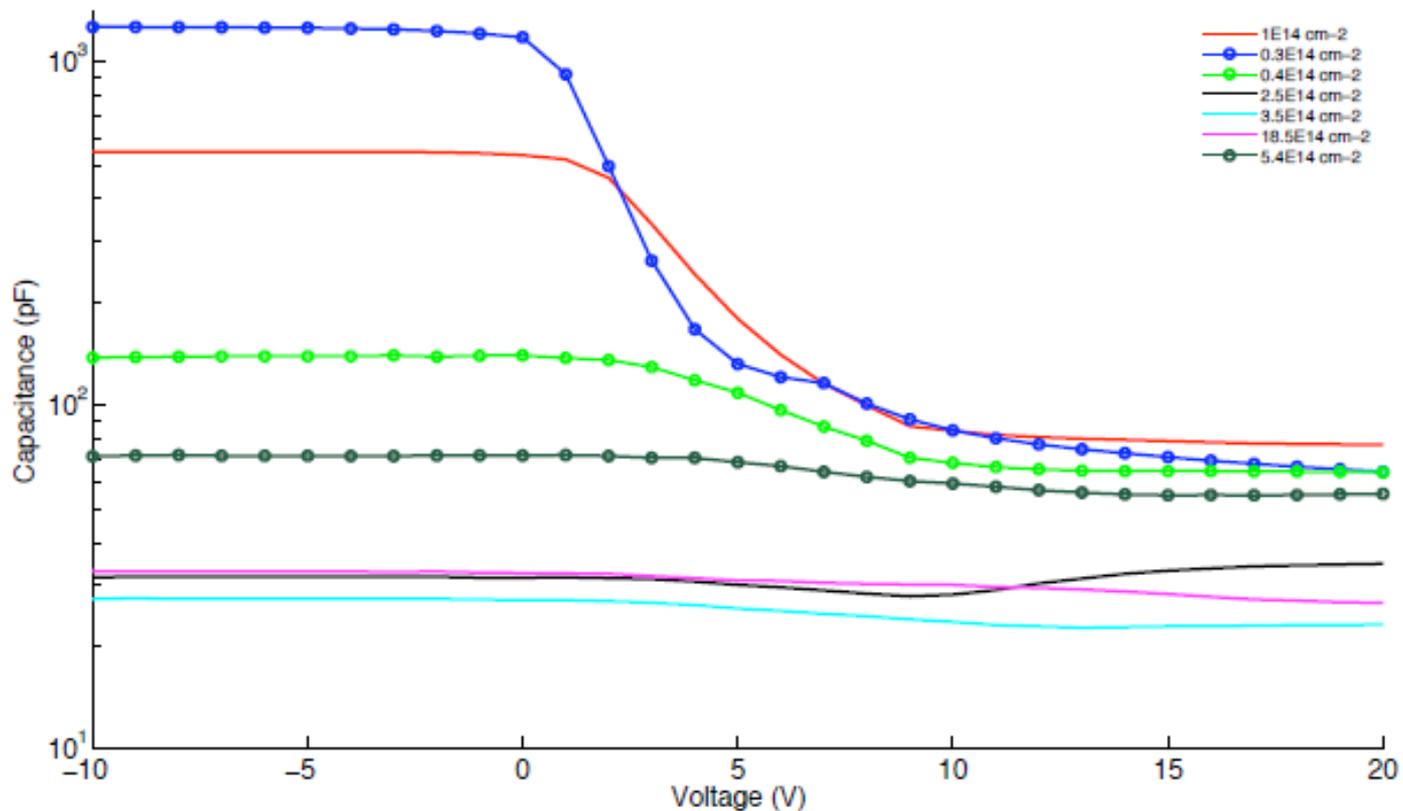
**X-ray radiation damage saturates !!!**



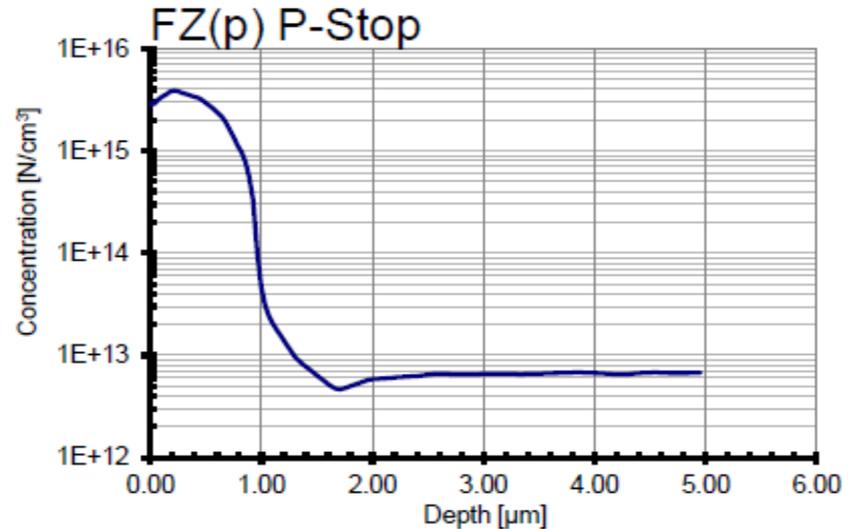
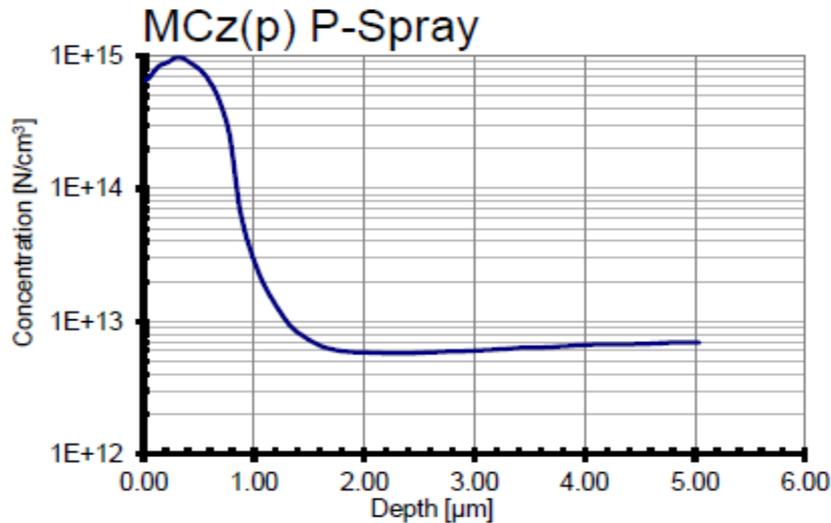
Robert Klanner - Univ. of Hamburg - RESMDD- Firenze - 10 -12. October 2012

# From Maria thesis, MOS measurement

The higher the fluence the more charge traps are introduced. The MOS capacitance decreases with fluence due to trapped charge carriers. The flatband voltage of unirradiated TS is about 1V compared to one of the irradiated result of about 4V. This higher flatband voltage points to additional oxide charges produced by irradiation.



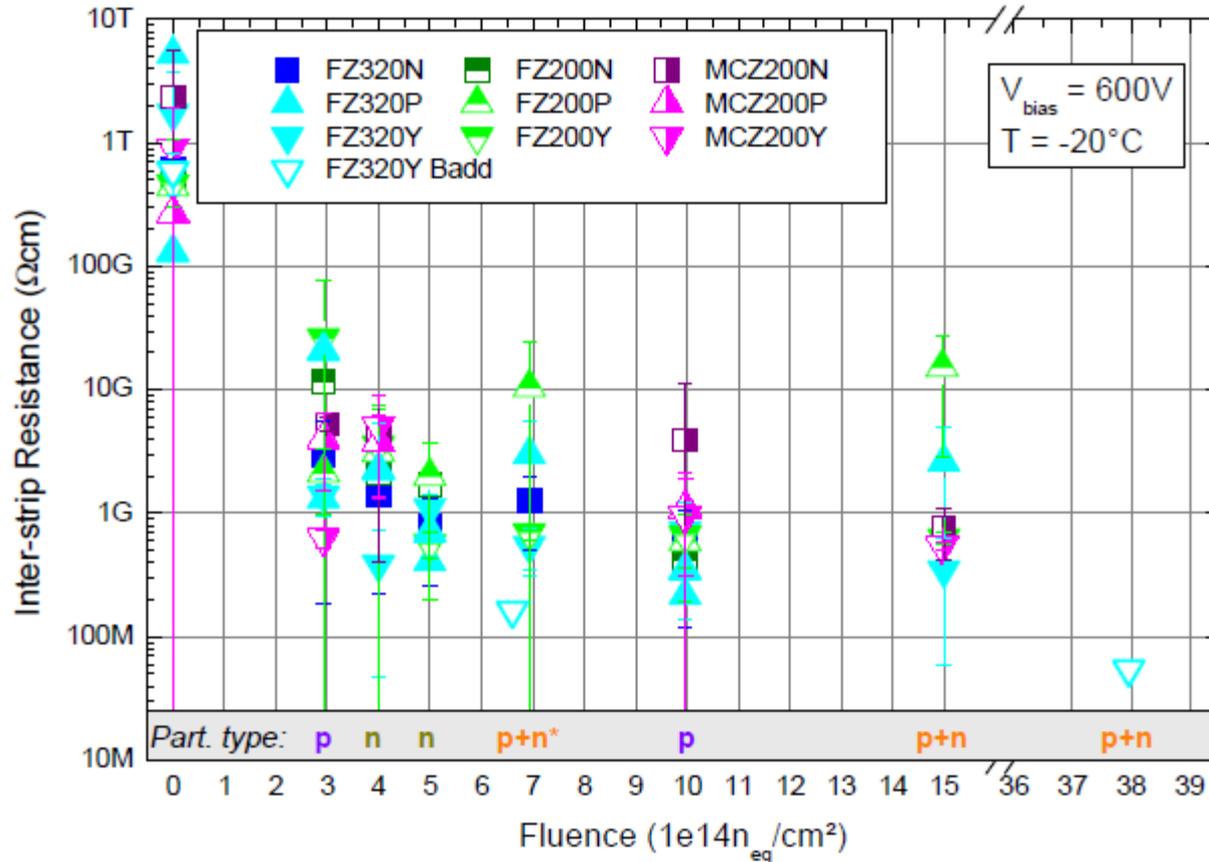
## 2. Doping Profiles: P Spray/Stop Implant



- **P Spray-Implant** in MCz200P: Peak concentration at 1E15, Implant depth app. 1,5 μm, Bulk concentration 8E12
- **P Stop-Implant** in FZ200P: Peak concentration at 4E15, Implant depth app. 1,6 μm, Bulk concentration 8E12
- The P Stop concentration is approximately 4 times higher than the P Spray, the implant depth is almost similar

# Rint variations for different fluence

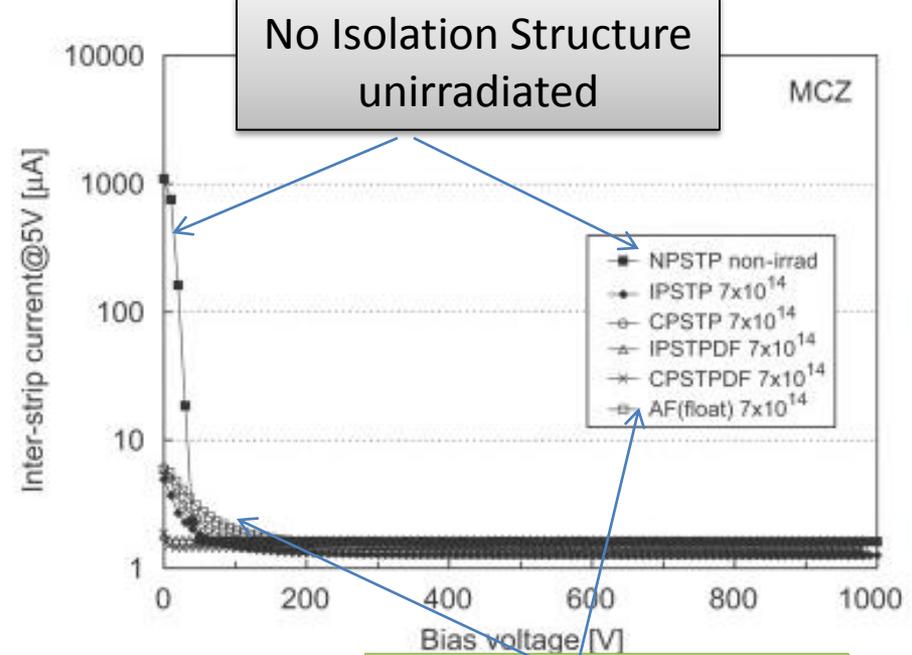
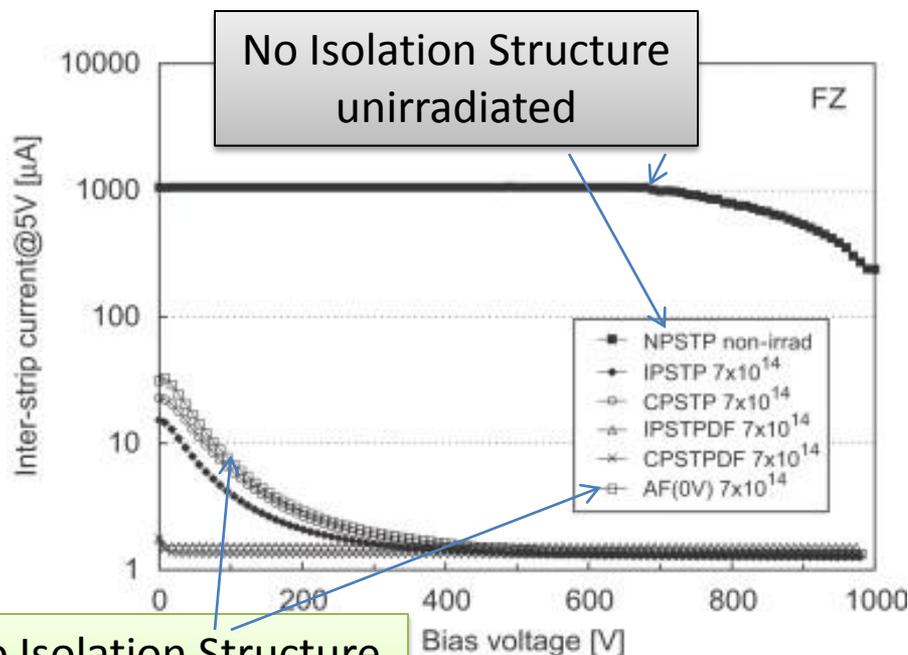
(Alexander D., Vertex-2012, 016 paper)



# Another experimental evidence .....

Y. Unno et. Al. (NIM A 579 (2007))

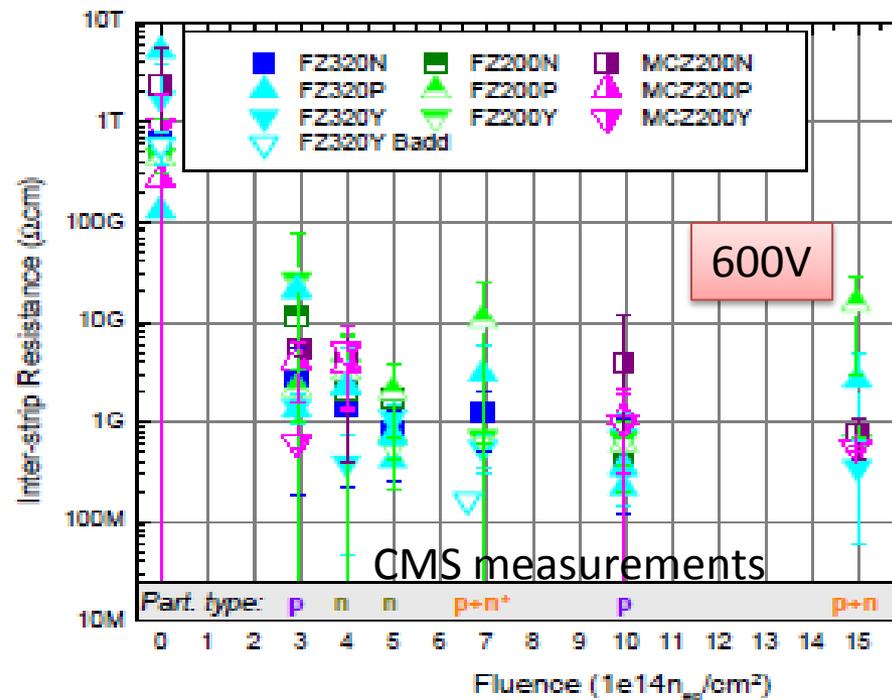
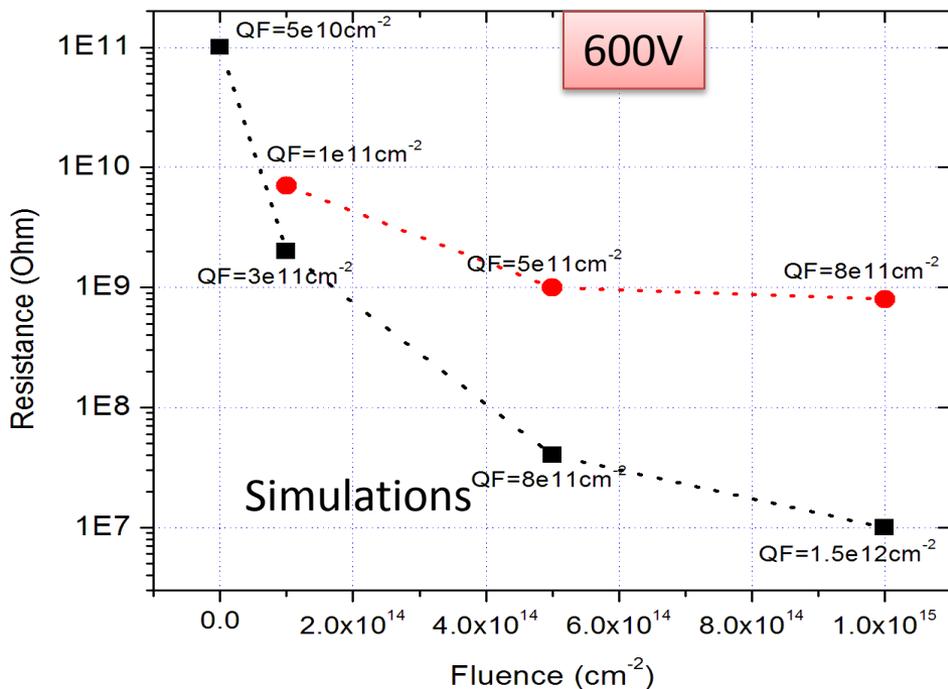
Strip isolation was observed for n+p- sensors without isolation structure after proton irradiation  
- Clear signature of proposed mechanism !



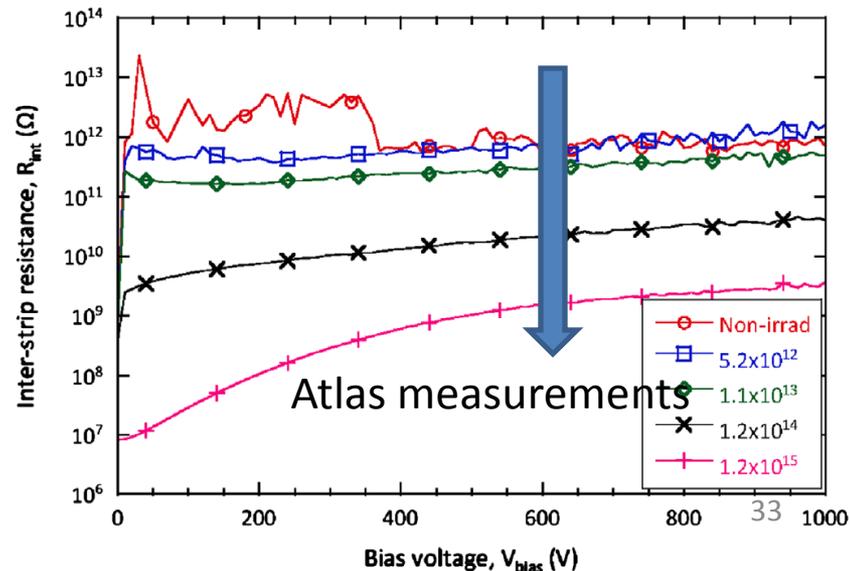
## Plot of Interstrip current vs. applied reverse bias.

- NPSTP – No Isolation structure (non-irradiated).
- AF – No isolation structure (irradiated by flux =  $7 \times 10^{14} \text{cm}^{-2}$ ).
- All other structures are with different layouts of Pstops (Irradiated)
- Voltage difference between two neighboring strips = 5V

# $R_{int}$ variation with Irradiation flux



- $R_{int}$  decreases with increase in fluence
- Similar trends in CMS tracker upgrade measurements (A. Dierlamm, PoS paper, Vertex 2012, 016)
- Similar trends observed in Atlas measurements (Y. Unno et al., NIMA, 2013), <http://dx.doi.org/10.1016/j.nima.2013.04.075>) with different p-stop parameters
- Can not be explained by increase in leakage current



DC external resistance = 1 ohm

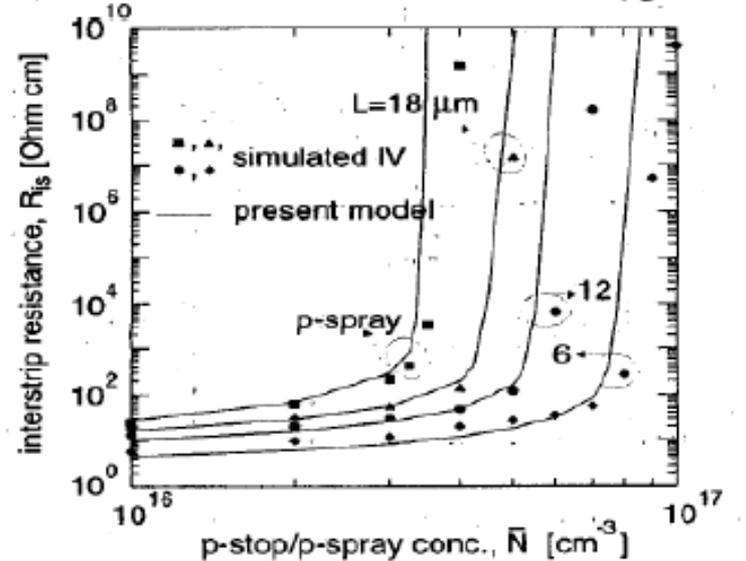
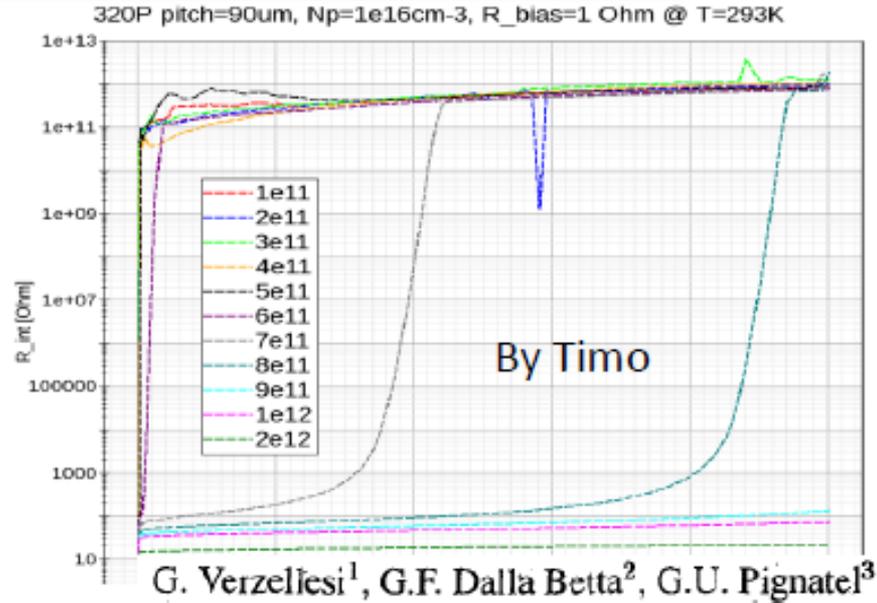
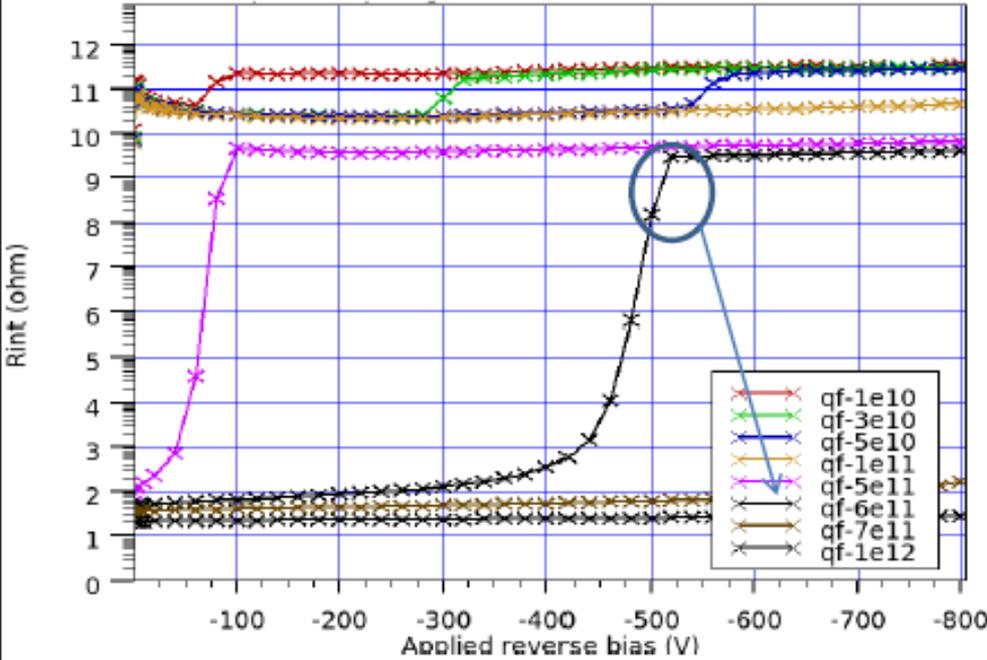


Figure 4: Interstrip resistance values as a function of the p-stop(p-spray) average doping concentration, as obtained from the proposed, analytical model and from simulated  $I_2(V_1)$  curves. A positive charge density of  $2 \times 10^{12} \text{ cm}^{-2}$  is assumed in the oxide.

- Similar, qualitative features for simulation plots
- Slight difference for intermediate values of  $Q_F$