



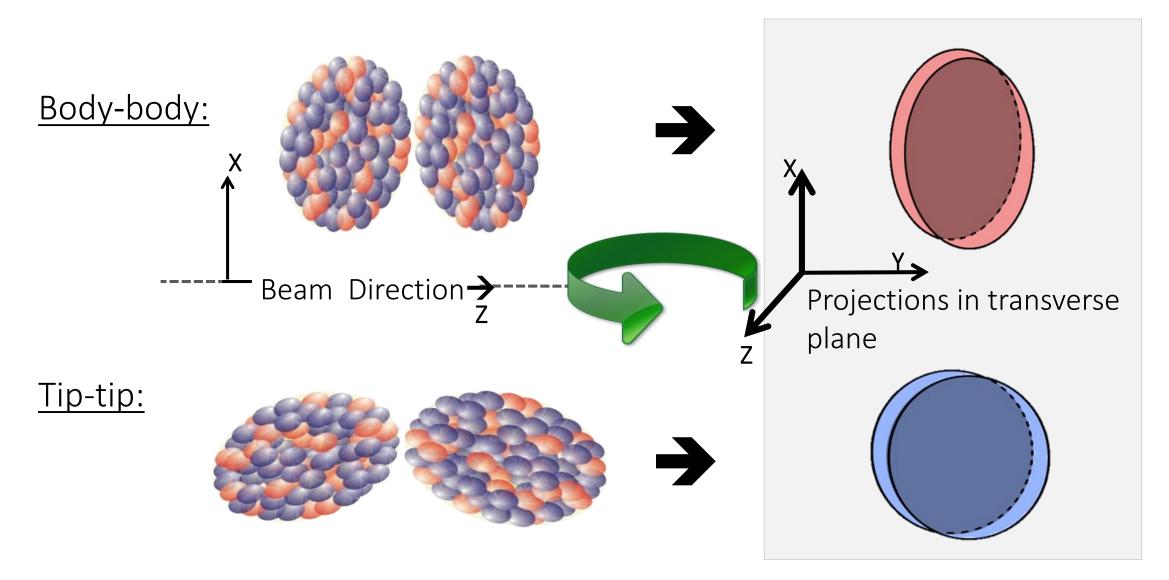
# Orientation Resolution in Uranium-Uranium Collisions

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# Definitions for desired collision configurations





## Motivation



Selection of tip-tip and body-body collision configurations opens up interesting features in the RHIC uranium data.

- Body-body collisions can be used to study the azimuthal behavior of the path length dependence in jet quenching.
- Tip-tip collisions provide the highest density collisions systems.
- For more, see papers by Heinz and Kuhlman: PRL 94 (2005) 132301
   PRC 72 (2005) 037901



## Introduction



#### 1. Initial condition model

 We use a two parameter MC-Glauber model with fluctuations to study deformed Uranium nuclei

### 2. Selecting collision orientation

- We discuss the experimental controls available for selecting orientations
- We address a model derived "knee" structure in the centrality dependent ellipticity

### 3. Zero degree calorimeters

- We use the model analogue of ZDCs to select on fully overlapping collisions
- We construct an explicit definition for tip-tip and body-body collisions to study the configuration selection efficiency
- We use forward and backward ZDC correlations to remove noise in tip-tip selection generated by asymmetric ZDC signals

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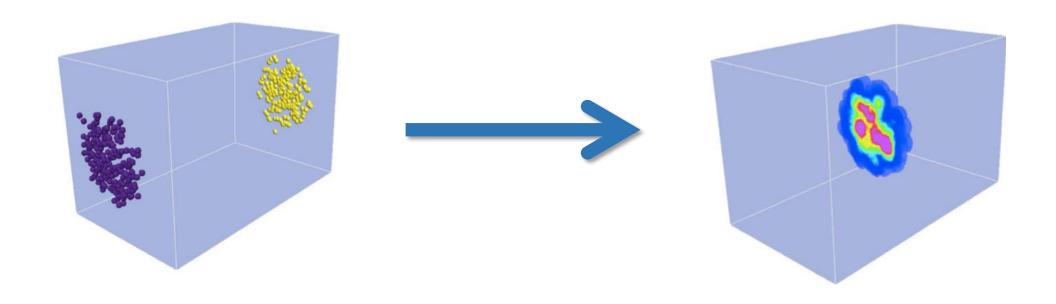
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# Initial condition model

## Deformed Woods Saxon Potential



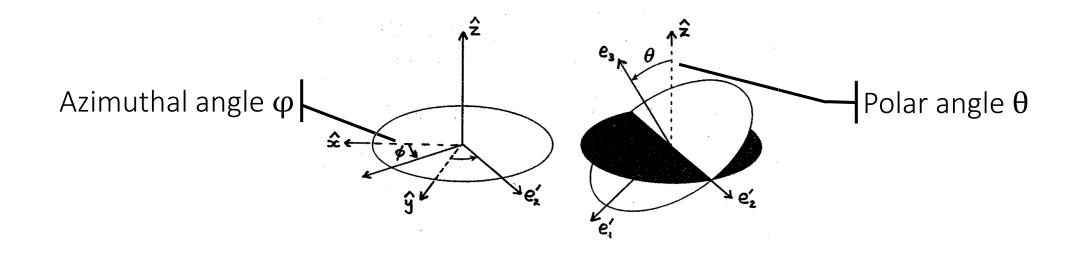
Uranium has a prolate shape given according to the deformed Woods-Saxon potential

$$\rho(r,\theta) = \frac{\rho_0}{1 + e^{(r-R[1+\beta_2 Y_2^0(\theta)+\beta_4 Y_4^0(\theta)])/d}}$$

$$\beta_2 = 0.28 \qquad d = 0.44$$

$$\beta_4 = 0.093 \qquad R = 6.86$$

$$\rho_0 = 0.166$$



# Monte Carlo Glauber (1/2)



The two-component Monte Carlo Glauber model is used to determine the entropy production from each collision event

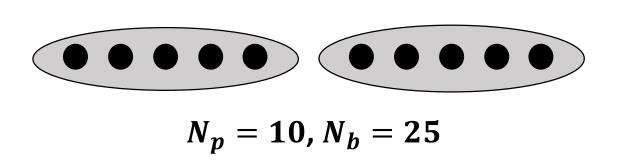
$$\frac{dS}{dy} = \frac{(1-\alpha)N_p}{2} + \alpha N_b \qquad \alpha = 0.12$$

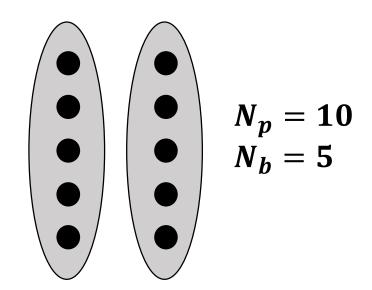
MC-Glauber is motivated by an observed nonlinear dependence of multiplicity on  $N_p$  in collisions between spherical nuclei.

# Monte Carlo Glauber (2/2)



Binary collisions  $N_b$  allow for events with the same number of participants  $N_p$  to contribute very different initial integrated entropy densities.



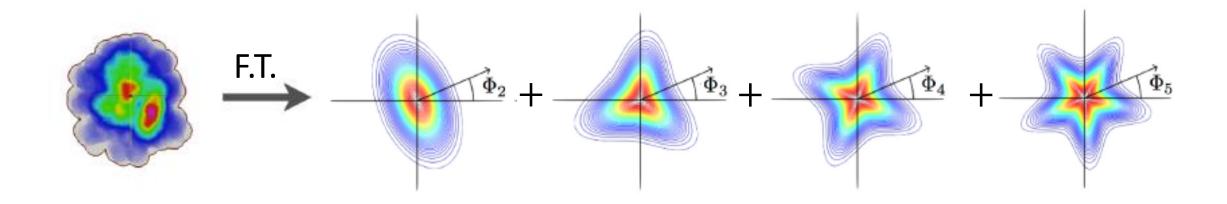


## **Eccentricities**



The shape of the energy distribution e(x,y) in the transverse plane can be described by  $r^2$ -weighted eccentricity coefficients

$$\varepsilon_n e^{in\Phi_n} = -\frac{\int dx dy r^2 e^{in\varphi} e(x,y)}{\int dx dy r^2 e(x,y)}$$



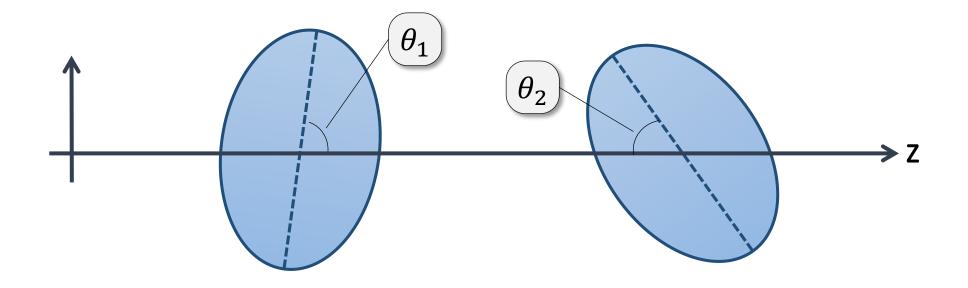
# Multiplicity Fluctuations



We implement multiplicity fluctuations as a weight at each nucleon using a  $\Gamma$  distribution

$$\Gamma(u,\kappa) = \frac{u^{\kappa-1}\kappa^{\kappa}e^{-\kappa u}}{\Gamma(\kappa)}, u = \frac{e_{max}}{\langle e_{max} \rangle} \in [0,\infty)$$

We use  $1/\kappa = 0.9175$  for the scale parameter of this distribution which has a variance of  $\sigma^2 = 1/\kappa$  and average  $\langle u \rangle = 1$ .



# Selecting collision orientations

# Experimental observables



Experiment ↔ Initial Condition

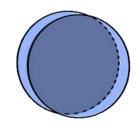
$$v_2 \leftrightarrow \varepsilon_2$$
 $dN/dy \leftrightarrow dS/dy$ 

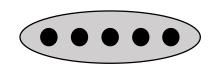
The experimental elliptic flow  $v_2$  and integrated particle distribution dN/dy are used to select orientations—these map linearly to their initial condition analogues  $\varepsilon_2$  and dS/dy.

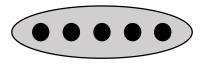
# Definitions for desired collision configurations using experimental observables



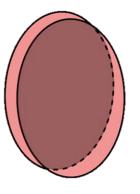
Tip-tip events have small  $\varepsilon_2$  and large  $N_b$ 

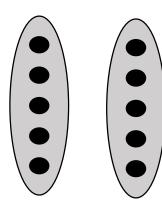






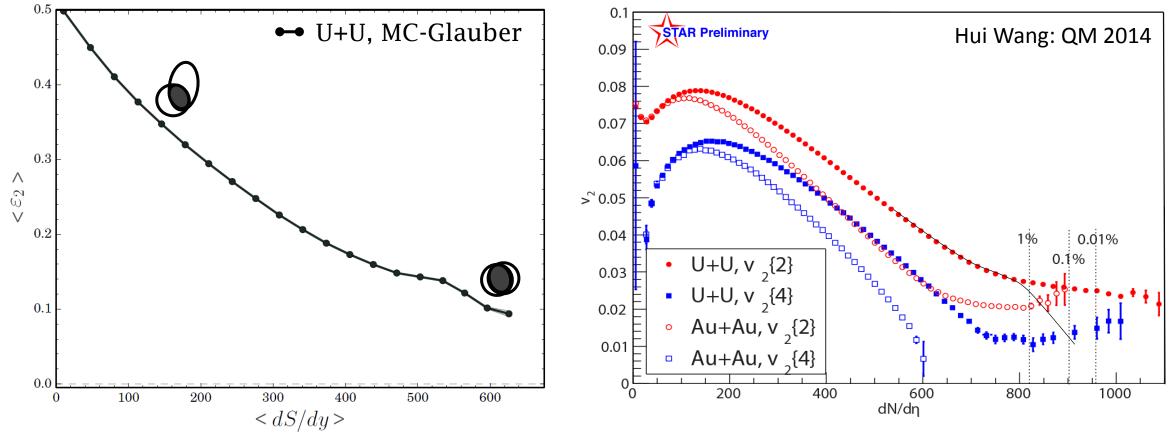
Body-body events have large  $\mathcal{E}_2$  and small  $N_b$ 





# A model dependent "knee" structure

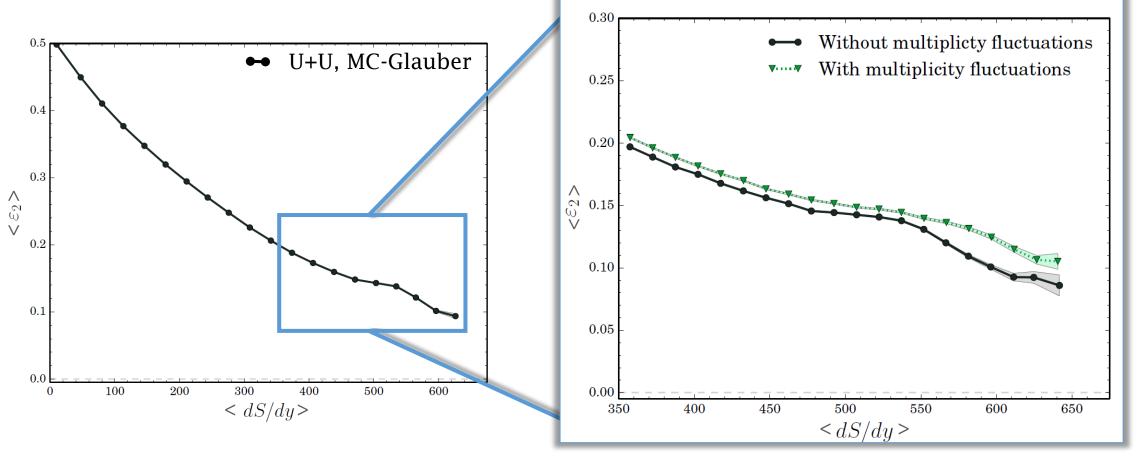




Within the two parameters mixed model, a preferential selection of tip-tip events at high dS/dy results in a "knee" structure in the centrality dependent ellipticity. The knee structure is not seen in STAR results.

# "Knee" structure with multiplicity fluctuations

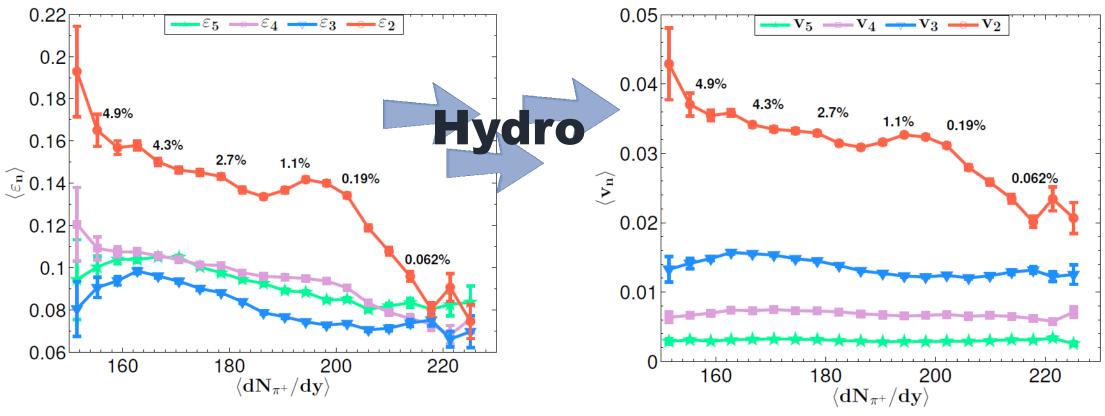




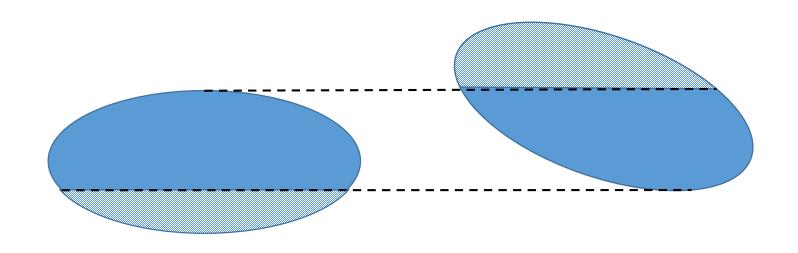
Adding multiplicity fluctuations shifts the "knee" to larger multiplicity and slightly diminishes the effect.

# "Knee" structure with hydrodynamics





The "knee" is seen in  $\varepsilon_2$  and weakly in  $\varepsilon_4$ . Event-by-event viscous hydrodynamics are run for a subset of the data for direct comparison to experiment. We find the "knee" is preserved in the final elliptic flow  $v_2$ .



# Zero Degree Calorimeters:

Tightening the collision geometry by using information about spectators

## Zero Degree Calorimeters



$$v_2 \leftrightarrow \varepsilon_2$$
 $dN/dy \leftrightarrow dS/dy$ 
 $%2DC \leftrightarrow %0N_S$ 

ZDCs catch neutrons that pass through a collision without participating, giving an idea of the collision overlap—the model analogue is the number of nucleon spectators  $N_s = 2A - N_p$ .

We use MC-Glauber initial conditions with multiplicity fluctuations and select on 1% of events that have the smallest  $N_s$ .

# Definitions for desired collision configurations using Euler angles



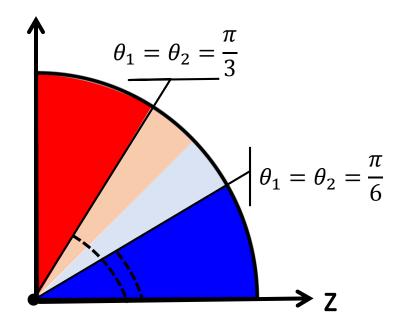
### Tip-tip

$$\sqrt{\frac{\cos^2(\theta_1) + \cos^2(\theta_2)}{2}} > \frac{\sqrt{3}}{2}$$

### Body-body

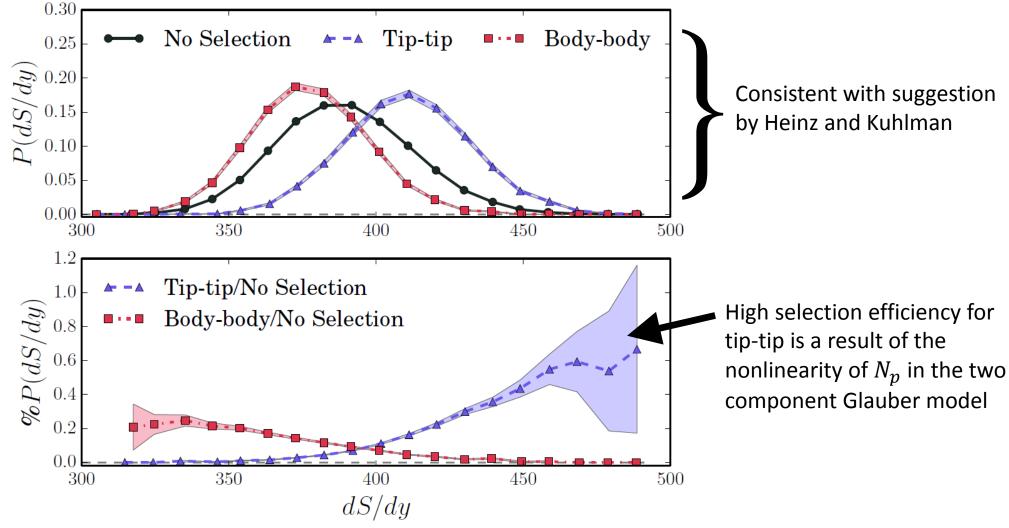
$$\sqrt{\frac{\cos^{2}(\theta_{1}) + \cos^{2}(\theta_{2})}{2}} < \frac{1}{2}$$

$$|\varphi_{1} - \varphi_{2}| < \pi/10$$



# Multiplicity distributions for different collision configurations for 1% ZDC selected events

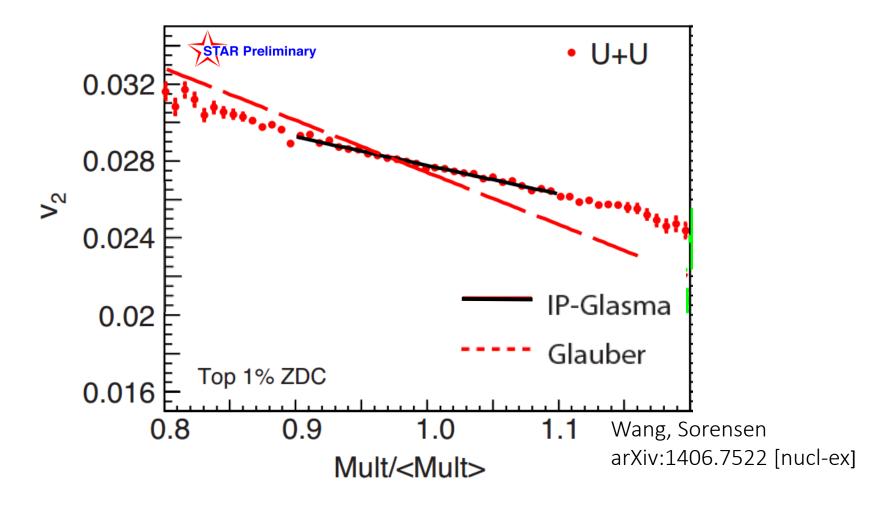




Lower dS/dy have more body-body. Higher dS/dy have more tip-tip.

# Model differences in elliptic flow vs multiplicity

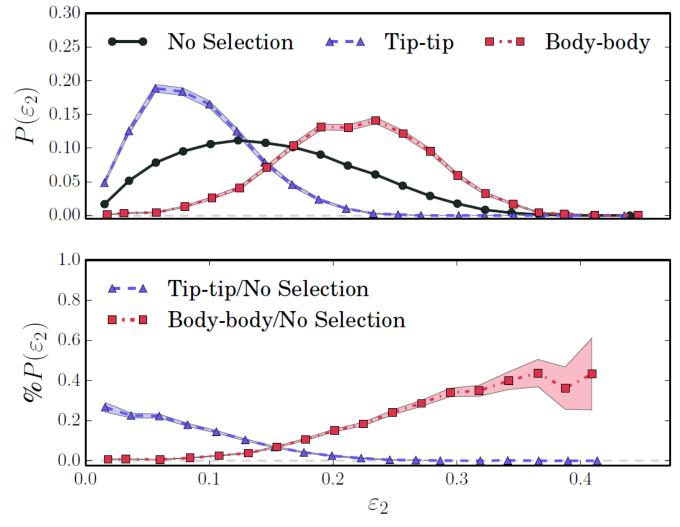




The nonlinear dependence of  $N_p$  on  $N_b$  in the 2-component Glauber model results in a steeper slope than in the STAR data and in the IP-Glasma model.

# Ellipticity distributions for different collision configurations for 1% ZDC selected events



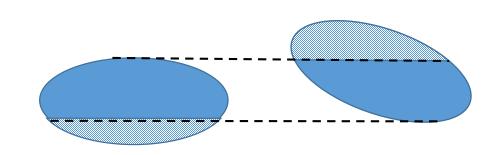


Lower  $\varepsilon_2$  have more tip-tip. Higher  $\varepsilon_2$  have more body-body.

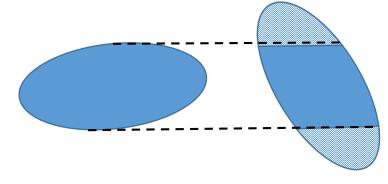
# Correlating ZDCs for additional constraints on collision geometry (1/2)



Events with symmetric ZDC signature are better described by the tip-tip collision criteria than asymmetric events (spectator regions are shaded).



Symmetric ZDC signature contributes more tip-tip

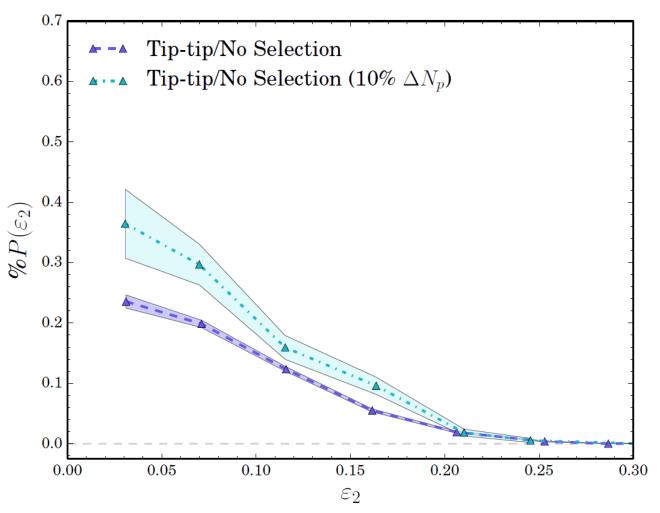


Asymmetric ZDC signature contribute less tip-tip

We define  $\Delta N_s = |N_{s,1} - N_{s,2}|$  and select events from the bottom 10% of  $\Delta N_s$  to diminish the contribution of asymmetric collisions and increase tiptip selection efficiency.

# Correlating ZDCs for additional constraints on collision geometry (2/2)





Selection efficiency of tip-tip collisions increases by removing asymmetric noise.

# Conclusions, future study



#### **Conclusions**

The MC-Glauber "knee" structure in the multiplicity dependent ellipticity is preserved after hydrodynamic evolution. It is weakened by multiplicity fluctuations.

For events cut on 1% ZDC, selection efficiency for tip-tip by cutting additionally on large dS/dy is ~50%, for body-body by cutting on low dS/dy it is never above 20%. These numbers are model dependent and correlate with the "knee" structure.

Selection efficiency of  $\varepsilon_2$  is ~25% for tip-tip and ~40% for body-body. Tip-tip results can be improved to ~35% using ZDC correlation.

#### **Future study**

In the process of running a hydrodynamic analysis of collision events that includes multiplicity fluctuations in the initial conditions.

Will use correlated forward and backward ZDC to study tip-body collisions.

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