Matrix Element technique at NLO to extract m<sub>top</sub> from dileptonic events

CERN, May 21, 2014 Pierre Artoisenet CP3, UCLouvain

based on ongoing work in collaboration with

Fabio Maltoni

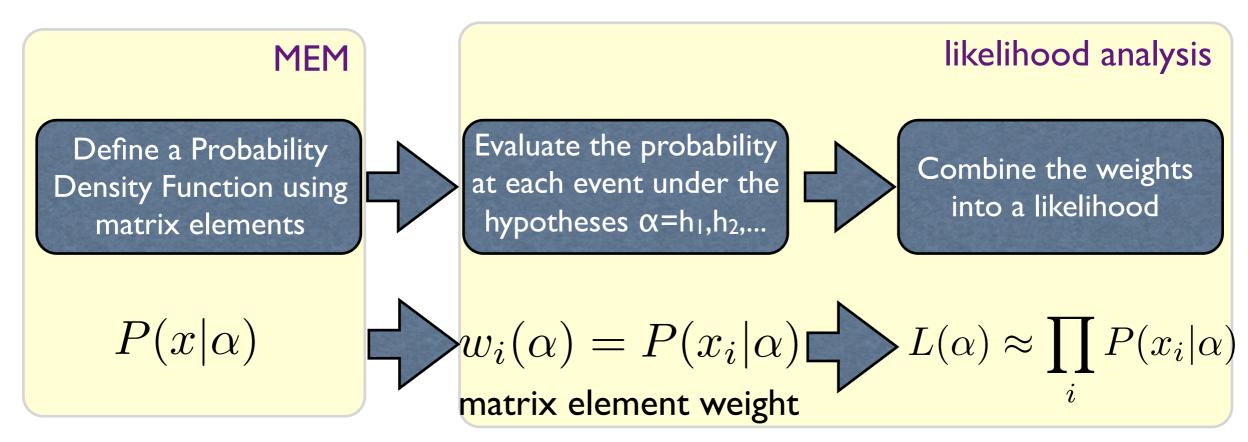
Michele Selvaggi

#### Outline

- Matrix Element Method (MEM)
- $\bullet$  New formulation of the MEM to measure  $m_{top}$
- Results
- Conclusion

# Matrix Element Method

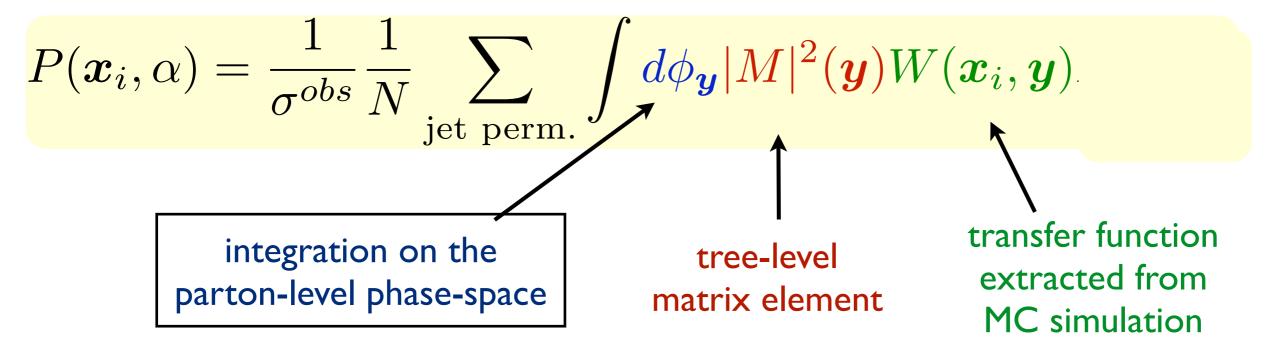
- Construction of the PDF based on hard scattering matrix elements
- Definition of the discriminating variable: likelihood built upon this PDF



- $\boldsymbol{x}$  : kinematics of the reconstructed event
- $\boldsymbol{\alpha}$  : theoretical assumption, in our case the mass of the top quark

#### Definition of the Probability Density Fct in the MEM

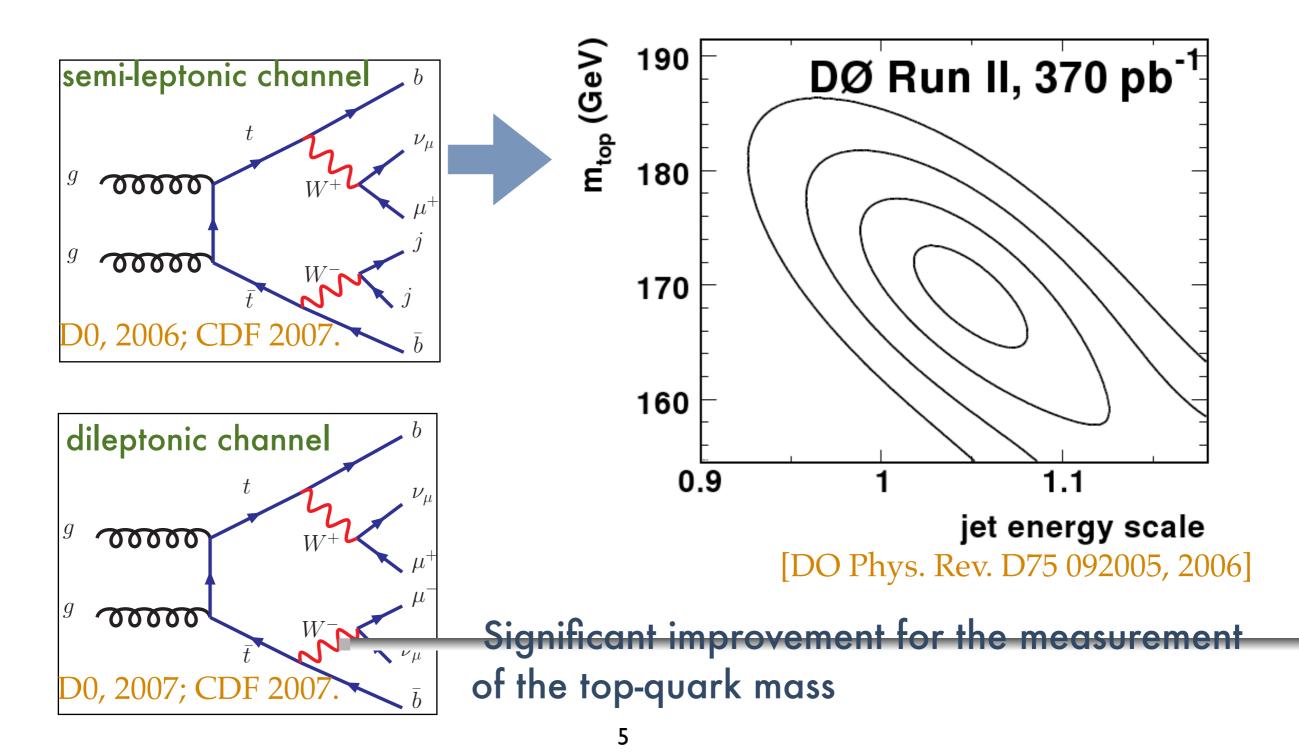
Hard scattering matrix element integrated over unconstrained information and convoluted with the resolution function W for the measured quantities
F. Canelli (D0 collaboration) 2003



In this talk: propose a new formulation of the PS integration that is fast and more convenient for the inclusion of QCD correction and parton shower effects

### Measurement of $m_{top}$ with the MEM at the Tevatron

#### Top-quark mass measurement from $t\overline{t}$ production in hadron collisions



# MEM for the measurement of m<sub>top</sub>: prospects

Improving the measurement of mtop with the MEM is challenging:

- I. Uncertainties in the transfer functions on jet energies
- 2. CPU Expensive method when the statistic is large
- 3. Lack of NLO information in the MEM weights

In this talk we consider the production of top quark pair in the dilepton channel and suggest a new formulation of the method that is fast and flexible enough to account for QCD correction

# General idea

- Consider top quark pair production in the dilepton channel
- Consider only the information from the lepton kinematics P<sub>1</sub>, P<sub>2</sub>



use only best determined quantities in the events

Set the MEM Probability Density Function to the cross section differential in the kinematics of the leptons

$$P(p_{l_1}, p_{l_2}|m_t) = \left(\frac{d\sigma}{[dp_{l_1}][dp_{l_2}]}\right)_{m_{\text{top}}}$$

$$[dp_{l_i}] =$$
  
LIPS measure  
for lepton i

For this specific case, we show that QCD corrections and parton shower effects can be included with no additional CPU costs

Starting point for the evaluation of the PDF: use the Narrow Width Approximation for the tops to factorize the production and decay phase-space measures

$$P(p_{l_1}, p_{l_2}|m_t) = \frac{1}{\sigma_{\text{prod}}} \sum_i w_i \frac{1}{|M_{\text{prod}}(X_i)|^2} \left(\frac{1}{2m_t \Gamma_t(m_t)}\right)^2$$
$$\int \prod_j \left(\frac{d\Phi_{\text{dk}}^j(X_i)}{[dp_{l_j}]} f_{\text{BW}}(p_{W_j}^2, m_W, \Gamma_W)\right) |\tilde{M}_{\text{dk}}(X_i^{\text{dk}})|^2$$

- The sum over i runs over a sample of production events (w/o top quark decay), w<sub>i</sub> = weight of production event i
- The sum over j runs over the decay branches
- $d\Phi^j_{dk}(X_i)$  is the phase-space measure associated with the decay products in branch j,

$$|\tilde{M}_{\rm dk}(X_i^{\rm dk})|^2 = {\rm decayed \, ME,} \quad |M_{\rm prod}(X_i)|^2 = {\rm production \, ME} \ {\rm (whole \, process)} \quad {\rm 8}$$

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covers the phase space for the decay

The sum over i runs over a sample of production events (w/o top quark decay), w<sub>i</sub> = weight of production event i

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Observations:

The phase-space integration associated with each decay branch is 2-dimensional and can be parametrized as follows:

$$\frac{d\Phi_{dk}^{j}(X_{i})}{[dp_{l_{j}}]} = d\phi_{j}dp_{W_{j}}^{2}\sum_{\text{sol}}g_{j}(p_{t_{j}}, p_{l_{j}}, \phi_{j}, p_{W_{j}}^{2})$$

- $p_{W_i}^2$ : squared invariant mass of W boson
  - $\phi_j$ : azimuthal angle of neutrino in the frame where b, v are back-to-back
- The matrix element with the decay (but Breit-Wigner distributions for W pulled out) is independent of the variables  $p_{W_j}^2, \phi_j$

indeed polarized matrix elements for the decay itself read:

$$|\tilde{M}_{\uparrow}(t(p_t) \to l(p_l) + b + \nu)|^2 = g_w^4 (2p_l \cdot p_t^a) (m_t^2 - m_b^2 - 2p_t \cdot p_l)$$
$$|\tilde{M}_{\downarrow}(p_t \to l(p_l) + b + \nu)|^2 = g_w^4 (2p_l \cdot p_t^b) (m_t^2 - m_b^2 - 2p_t \cdot p_l)$$

So ME's need to be evaluated at most once per production event:

$$P(p_{l_1}, p_{l_2}|m_t) = \frac{1}{\sigma_{\text{prod}}} \sum_{i} w_i \frac{|\tilde{M}_{\text{dk}}(X_i^{\text{dk}})|^2}{|M_{\text{prod}}(X_i)|^2} \left(\frac{1}{2m_t \Gamma_t(m_t)}\right)^2$$
$$\int \prod_{j} \left( d\phi_j dp_{W_j}^2 f_{\text{BW}}(p_{W_j}^2, m_W, \Gamma_W) \sum_{\text{sol}} g_j(p_{t_j}, p_{l_j}, \phi_j) \right)$$

So ME's need to be evaluated at most once per production event:

$$\begin{split} P(p_{l_1},p_{l_2}|m_t) &= \frac{1}{\sigma_{\mathrm{prod}}} \sum_i w_i \frac{|\tilde{M}_{\mathrm{dk}}(X_i^{\mathrm{dk}})|^2}{|M_{\mathrm{prod}}(X_i)|^2} \left(\frac{1}{2m_t \Gamma_t(m_t)}\right)^2 \\ &\int \prod_j \left( d\phi_j dp_{W_j}^2 f_{\mathrm{BW}}(p_{W_j}^2,m_W,\Gamma_W) \sum_{\mathrm{sol}} g_j(p_{t_j},p_{l_j},\phi_j) \right) \\ &= \sum_j G(p_{t_j},p_{l_j},m_W,\Gamma_W) \begin{array}{c} \mathrm{can \ be \ evaluated} \\ \mathrm{can \ be \ evaluated} \\ \mathrm{analytically \ !} \end{split}$$

- I. reading a production event,
- 2. computing  $\sum_{i} G(p_{t_j}, p_{l_j}, m_W, \Gamma_W)$  ~overlap |b| the lepton & top momenta,
- 3. multiplying this overlap by the event weight and by the decayed ME over production ME

#### Advantages:

#### **FAST**

Integration over the phase-space for production is optimized by considering a sample of unweighted events

Calculation of one weight (looping over 100k LO production events, 0.5 - 1% accuracy) takes 8 seconds (on average)

#### TH SYSTEMATICS EASILY INCLUDED

Uncertainties from variation of the scales available in the production event file, can be propagated to the weights at no cost

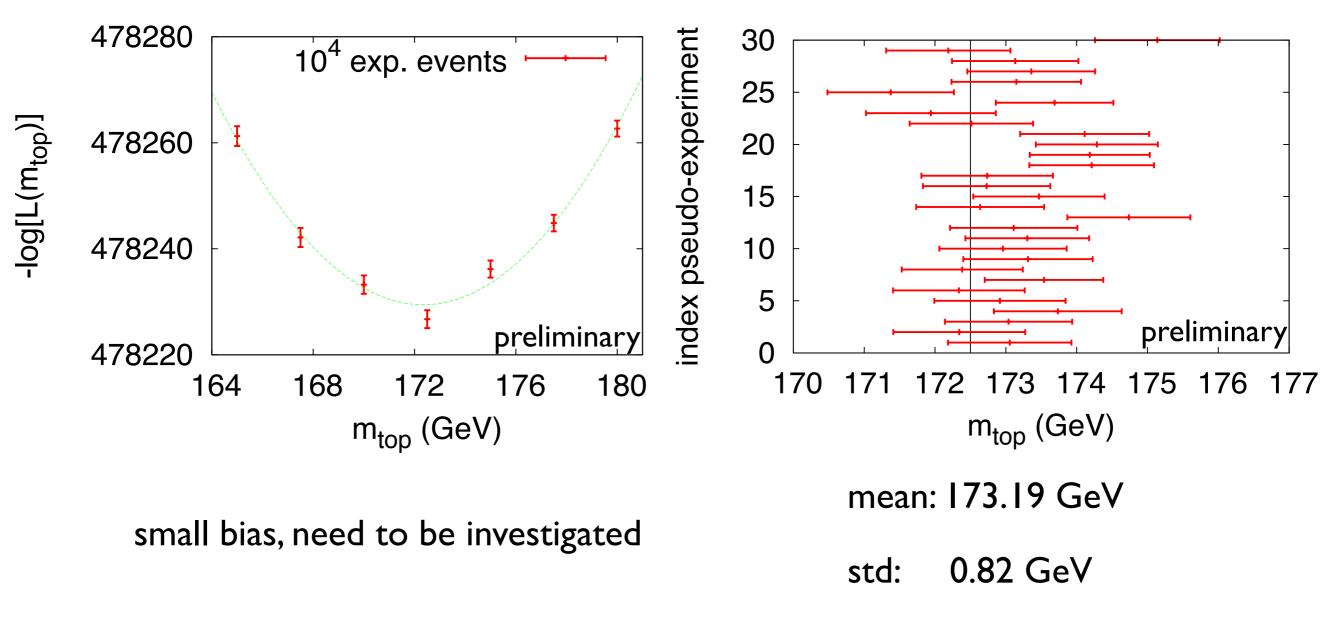
#### • CAN ACCOUNT FOR QCD CORRECTION (in the production)

Simply loop over a sample of NLO production events

# Application (LHC@I3TeV)

30 pseudo-experiment samples of 10K events, no cut,

generated with  $m_{top}$ =172.5 GeV, at LO, no shower



Wednesday 21 May 14

# Including shower effects

Parton shower impacts the kinematics of the top quarks, hence also the kinematics of the leptons

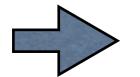
In our setup; such effects can be included by slightly modifying the previous prescription

For each production event,

- I. read the momenta of the top quarks after shower  $p_{t_1}^{
  m sh},\,p_{t_2}^{
  m sh}$
- 2. reconstruct the kinematics of the top quark decays from  $p_{t_1}^{
  m sh},\,p_{t_2}^{
  m sh},\,p_{l_1},\,p_{l_2}$
- 3. boost the kinematics of each branch in a frame where the momenta of the top quarks match those of the "matrix-element event" (= event before the shower)
- 4. Use the matrix-element event to evaluate the ME's

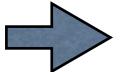
# Ongoing work

Extend the analysis at NLO



loop over a sample of NLO production events

Assess the sensitivity of the method w.r.t different production scheme (LO, LO +PS, LO merged samples with different parton multiplicities, NLO matched to PS)



- loop over different samples of production events
- Assess the impact of cuts on the leptons and on the jets (ex: b-tagging effects)
  - > reconstruction of  $m_{top}$  via a calibration procedure
- Assess the impact of systematic uncertainties
  - **P** 
    - propagate the variation of the scales to the value of the weights and the likelihood

### Conclusion

- In this talk we explored the use of MEM to measure m<sub>top</sub> from dilepton top quark pair events using only the lepton information
- For this very specific case, reweighing events with matrix elements can be achieved by means of a new formulation for the calculation of the weights, that is faster and more appropriate to include QCD effects beyond leading order
- Beside its possible capabilities to improve the mass measurement of the top quark, this analysis represents an interesting case of study to assess the impact of QCD correction on the calculation of the weights
- Whether this option will deliver a competitive measurement of m<sub>top</sub> remains to be established