FLOW FLUCTUATIONS IN HEAVY-ION COLLISIONS

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FLOW FLUCTUATIONS

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TAKE-AWAY MESSAGES

- Flow encompasses more phenomena than previously realized.
- New flow observables will tightly constrain models.
- 1 LONG-RANGE TWO-PARTICLE CORRELATIONS
- **2** FLOW FLUCTUATIONS
- 3 NEW FLOW OBSERVABLES



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Azimuthal distribution of emitted particles :



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Azimuthal distribution of emitted particles :

$$rac{dN}{d\phi} \propto 1 + 2v_2\cos 2\phi + 2v_4\cos 4\phi + \dots$$



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(Elliptic) Flow c. QM09

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 $\left<\left< e^{i2(\phi_1 - \phi_2)} \right> \right>$



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$$\left\langle \frac{dN_{\text{pairs}}}{d\Delta\phi} \right\rangle^{\text{(flow)}} 1 + \left\langle v_2^2 \right\rangle \cos 2(\Delta\phi) + \left\langle v_4^2 \right\rangle \cos 4(\Delta\phi) + \dots$$

TWO-PARTICLE CORRELATIONS



⁽PHOBOS, Phys. Rev. C75(2007)054913)

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(STAR, arXiv:1010.0690)

(PHOBOS, Phys.Rev. C81 (2010) 024904)



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... can be generated by purely collective flow.

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$$rac{dN}{d\phi} \propto 1\!+\!2v_1\cos\phi$$

$$+2v_2\cos 2\phi$$

$$+2v_3\cos 3\phi+\ldots$$



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$\frac{dN}{d\phi} \propto 1 + 2v_1 \cos \phi + 2v_1^s \sin \phi + 2v_2 \cos 2\phi + 2v_2^s \sin 2\phi + 2v_3 \cos 3\phi + \dots$



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$$\frac{dN}{d\phi} \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos n(\phi - \psi_n)$$



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$$\frac{dN}{d\phi} \propto 1 + \sum_{n=1}^{\infty} 2v_n \cos n(\phi - \psi_n)$$
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$$v_2 e^{2i\psi_2} \propto \varepsilon_2 e^{2i\Phi_2} \equiv -\frac{\{r^2 e^{2i\phi}\}}{\{r^2\}}$$

(Holopainen, Niemi, Eskola, Phys.Rev.C83, 034901 (2011))



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$$v_3 e^{3i\psi_3} \propto \varepsilon_3 e^{3i\Phi_3} \equiv -\frac{\{r^3 e^{3i\phi}\}}{\{r^3\}}$$

(Qin, Petersen, Bass, Muller, Phys. Rev. C 82, 064903 (2010); Qiu, Heinz, arXiv:1104.0650)



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$$v_1 e^{i\psi_1} \propto \varepsilon_1 e^{i\Phi_1} \equiv -\frac{\{r^3 e^{i\phi}\}}{\{r^3\}}$$

(Teaney & Yan, arXiv:1010.1876; Gardim, Grassi, Hama, Luzum, Ollitrault, arXiv:1103.4605)



FLOW FLUCTUATIONS

- Centrality dependence, size, of v₃ and v₂ (Alver & Roland, Phys.Rev. C81 (2010) 054905)
- *p_t*-dependence and orientation with respect to event plane (Luzum, Phys.Lett. B696 (2011) 499-504)
 (Luzum & Ollitrault, Phys.Rev.Lett. 106 (2011) 102301)
- Centrality dependence of "ridge amplitude" (Sorensen, Bolliet, Mocsy, Pandit, Pruthi, arXiv:1102.1403)
- Factorization, mass dependence, ... (See plenary/parallel talks from ALICE, ATLAS, CMS, ...)

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Quantitative evidence of flow hypothesis:

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Current long-range 2-particle data can be explained by flow alone*.

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$$^{*}\left\langle \cos\Delta\phi
ight
angle = extsf{v}_{1}^{(t)} extsf{v}_{1}^{(a)} - rac{ extsf{p}_{t}^{(t)} extsf{p}_{t}^{(a)}}{\left\langle\sum extsf{p}_{t}^{2}
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Current long-range 2-particle data can be explained by flow alone*.

 \implies can accurately measure many new flow observables with little non-flow contamination

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CMS preliminary, Velkovska plenary



LESSONS:

v_3 is a more sensitive probe of η/s

PHENIX, arXiv:1105.3928



LESSONS:

Combining v_2 and v_3 can rule out IC models

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PHENIX, arXiv:1105.3928



LESSONS:

Combining v_2 and v_3 can rule out IC models (CGC is *not* ruled out)

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ALICE, arXiv:1105.3865



LESSONS:

Glauber may not work either

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ATLAS, Jia plenary



LESSONS:

Higher coefficients are measurable and add more constraints

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STAR preliminary, Sorensen plenary





Stay tuned (see Paul Sorensen's talk later this session)

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Many other independent flow observables can be measured:

 $v\{n_1, n_2, \ldots, n_k\} \equiv \langle \cos(n_1\phi_1 + \ldots + n_k\phi_k) \rangle$

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$$\begin{aligned} \mathbf{v}_{n}\{2\}^{2} &\equiv \mathbf{v}\{n, -n\} &= \left\langle \mathbf{v}_{n}^{2} \right\rangle \\ 2\mathbf{v}_{n}\{2\}^{4} - \mathbf{v}_{n}\{4\}^{4} &\equiv \mathbf{v}\{n, n, -n, -n\} &= \left\langle \mathbf{v}_{n}^{4} \right\rangle \\ \mathbf{v}_{24} &\equiv \mathbf{v}\{2, 2, -4\} &= \left\langle \mathbf{v}_{2}^{2}\mathbf{v}_{4}\cos 4(\psi_{2} - \psi_{4}) \right\rangle \\ \mathbf{v}_{23} &\equiv \mathbf{v}\{2, 2, 2, -3, -3\} &= \left\langle \mathbf{v}_{2}^{3}\mathbf{v}_{3}^{2}\cos 6(\psi_{2} - \psi_{3}) \right\rangle \\ \mathbf{v}_{12} &\equiv \mathbf{v}\{1, 1, -2\} &= \left\langle \mathbf{v}_{1}^{2}\mathbf{v}_{2}\cos 2(\psi_{1} - \psi_{2}) \right\rangle \\ \mathbf{v}_{13} &\equiv \mathbf{v}\{1, 1, 1, -3\} &= \left\langle \mathbf{v}_{1}^{3}\mathbf{v}_{3}\cos 3(\psi_{1} - \psi_{3}) \right\rangle \\ \mathbf{v}_{123} &\equiv \mathbf{v}\{1, 2, -3\} &= \left\langle \mathbf{v}_{1}\mathbf{v}_{2}\mathbf{v}_{3}\cos(\psi_{1} + 2\psi_{2} - 3\psi_{3}) \right\rangle \end{aligned}$$

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QUARK MATTER 2011 12 / 13

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$$v\{n_1, n_2, \ldots, n_k\} \equiv \langle \cos(n_1\phi_1 + \ldots + n_k\phi_k) \rangle$$

$$v_{n}\{2\}^{2} \equiv v\{n, -n\} = \left\langle v_{n}^{2} \right\rangle$$

$$2v_{n}\{2\}^{4} - v_{n}\{4\}^{4} \equiv v\{n, n, -n, -n\} = \left\langle v_{n}^{4} \right\rangle$$

$$v_{24} \equiv v\{2, 2, -4\} = \left\langle v_{2}^{2}v_{4}\cos 4(\psi_{2} - \psi_{4}) \right\rangle$$

$$v_{23} \equiv v\{2, 2, 2, -3, -3\} = \left\langle v_{2}^{3}v_{3}^{2}\cos 6(\psi_{2} - \psi_{3}) \right\rangle$$

$$v_{12} \equiv v\{1, 1, -2\} = \left\langle v_{1}^{2}v_{2}\cos 2(\psi_{1} - \psi_{2}) \right\rangle$$

$$v_{13} \equiv v\{1, 1, 1, -3\} = \left\langle v_{1}^{3}v_{3}\cos 3(\psi_{1} - \psi_{3}) \right\rangle$$

$$v_{123} \equiv v\{1, 2, -3\} = \left\langle v_{1}v_{2}v_{3}\cos(\psi_{1} + 2\psi_{2} - 3\psi_{3}) \right\rangle$$

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Many other independent flow observables can be measured:

$$v\{n_1, n_2, \ldots, n_k\} \equiv \langle \cos(n_1\phi_1 + \ldots + n_k\phi_k) \rangle$$

$$\begin{array}{ll} v_n\{2\}^2 \equiv v\{n, -n\} &= \left\langle v_n^2 \right\rangle \\ 2v_n\{2\}^4 - v_n\{4\}^4 \equiv v\{n, n, -n, -n\} &= \left\langle v_n^4 \right\rangle \\ v_{24} \equiv v\{2, 2, -4\} &= \left\langle v_2^2 v_4 \cos 4(\psi_2 - \psi_4) \right\rangle \\ v_{23} \equiv v\{2, 2, 2, -3, -3\} &= \left\langle v_2^3 v_3^2 \cos 6(\psi_2 - \psi_3) \right\rangle \\ v_{12} \equiv v\{1, 1, -2\} &= \left\langle v_1^2 v_2 \cos 2(\psi_1 - \psi_2) \right\rangle \\ v_{13} \equiv v\{1, 1, 1, -3\} &= \left\langle v_1^3 v_3 \cos 3(\psi_1 - \psi_3) \right\rangle \\ v_{123} \equiv v\{1, 2, -3\} &= \left\langle v_1 v_2 v_3 \cos(\psi_1 + 2\psi_2 - 3\psi_3) \right\rangle \end{array}$$

• Flow fluctuations have important, non-negligible effects on measured correlations

- ullet \Longrightarrow flow has further-reaching effects than was previously realized
- With this understanding comes many new possible independent flow measurements
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CMS PP RIDGE

(d) CMS N \geq 110, 1.0GeV/c<p_<3.0GeV/c



NEW FLOW OBSERVABLES: TRIANGULAR FLOW

Predictions:



MATT LUZUM (IPHT)

FLOW FLUCTUATIONS

Still missing: new v1 from fluctuations (Teaney & Yan, arXiv:1010.1876)

- $v_1 = v_1^a + v_1^s$
- $v_1^a(\eta) = -v_1^a(-\eta)$ = usual directed flow
- $v_1^s(\eta) = v_1^s(-\eta)$ = new "directed flow at midrapidity"
- To measure from a 2-particle correlation, must remove "momentum conservation" correlation (Luzum, Ollitrault, Phys.Rev.Lett.106:102301,2011)

New flow observables: directed flow at midrapidity

$$\langle \cos \Delta \phi \rangle = v_1^{(t)} v_1^{(a)} - \frac{p_t^{(t)} p_t^{(a)}}{\langle \sum p_t^2 \rangle}$$

(Teaney & Yan, arXiv:1010.1876) <u>MATT LUZUM (IPHT)</u>

FLOW FLUCTUATIONS

QUARK MATTER 2011 17 / 13

(Gardim, Grassi, Hama, Luzum, Ollitrault, arXiv:1103.4605)