

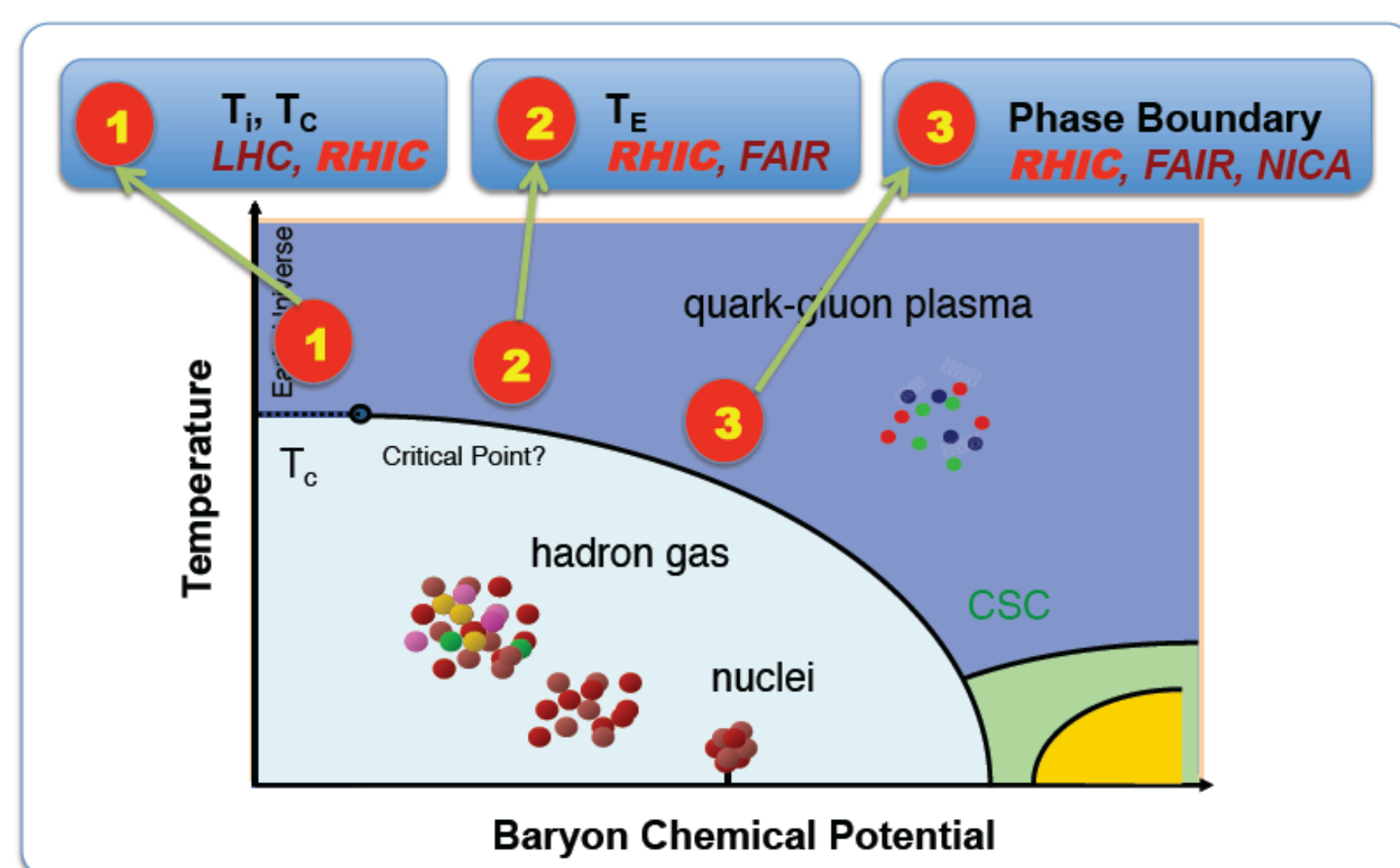
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## ABSTRACT

The production of hypertriton and light nuclei are simulated in a dynamical coalescence model coupled with a multi-phase transport model (AMPT). The beam energy dependence of strangeness population factor,  $S_3 = \frac{^3\text{H}}{^3\text{He} \times \frac{\Lambda}{p}}$ , is calculated to study local baryon-strangeness correlation as a valuable tool to probe the nature of the dense matter created in relativistic heavy ion collisions. We find that AMPT with string melting predicts an increase of  $S_3$  with increasing beam energy, and is consistent with experimental data, while AMPT with only hadronic scattering results in a low  $S_3$  throughout the energy range from AGS to RHIC, and fails to describe the experimental data. And we analyzed coalescence parameters,  $B_2$  and  $B_3$ , based on the production of deuteron, helium-3 and proton. The coalescence parameters of  $B_2$  and  $B_3$  decrease with increasing of beam energy or number of participant. The value of  $B_2$  and  $B_3$  in this model are consistent with the measurement by experiment collaboration in nucleus-nucleus collisions at different beam energy or in different centralities. The freeze-out correlation volume,  $V_f^{A-1}$  (A is atomic mass number), is calculated in AMPT model. The results of coalescence parameter and the freeze-out correlation volume follow the relation of  $B_A \sim V_f^{1-A}$ , which is from coalescence mechanism and observed in experiments.

## QCD Phase transition

The baryon-strangeness local correlation and evolution of collision zone are sensitive to deconfinement phase transition



## Dynamical Coalescence model

The multiplicity of a M-hadron cluster in a heavy ion collision is given by,

$$N_M = G \int d\vec{r}_1 d\vec{q}_1 \dots d\vec{r}_{M-1} d\vec{q}_{M-1} \times \left\langle \sum_{i_1 > i_2 > \dots > i_M} \rho_i^M(\vec{r}_i, \vec{q}_i, \dots, \vec{r}_{M-1}, \vec{q}_{M-1}) \right\rangle$$

Deuteron:  $\rho_d^M(r, \vec{k}) = 8 \left( -\frac{r^2}{\sigma_d^2} \right) \exp[-\vec{k}^2 \sigma_d^2]$

3-hadron cluster:  $\rho_{3h}^M(\rho, \lambda, \vec{k}_p, \vec{k}_n) = 8^2 \left( -\frac{\rho^2 + \lambda^2}{\sigma_{3h}^2} \right) \exp[-(\vec{k}_p^2 + \vec{k}_n^2) \sigma_{3h}^2]$

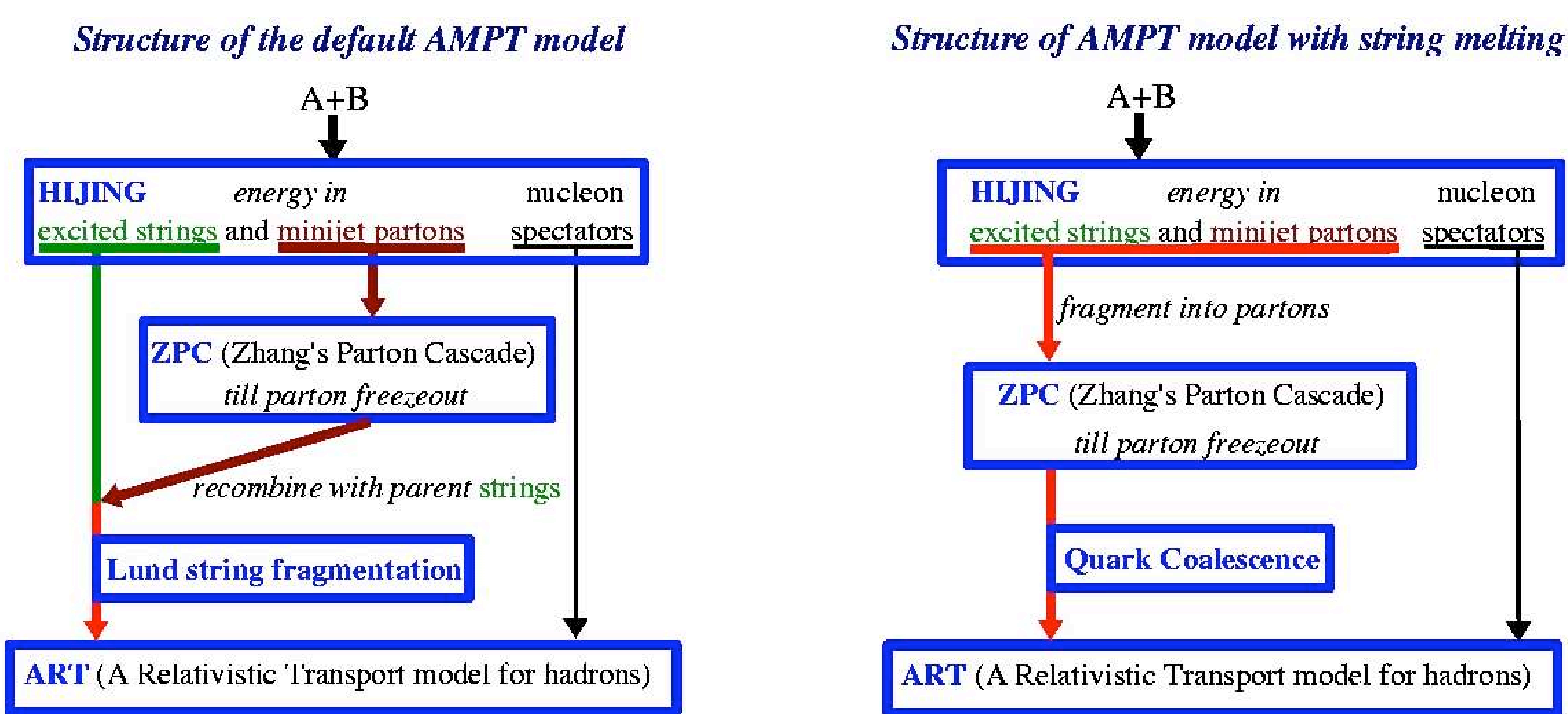
Coalescence picture: d(p,n), t(p,n,n),  $^3\Lambda\text{H}(p,n,\Lambda)$ ,  $^3\text{He}(p,p,n)$

## HTB

CRAB is based on the formula,

$$C(\vec{p}_1, \vec{q}) = 1 + \frac{\int d^4x_1 d^4x_2 S_1(x_1, \vec{p}_1) S_2(x_2, \vec{p}_2) |\phi_{rel}(x'_2 - x'_1)|^2}{\int d^4x_1 d^4x_2 S_1(x_1, \vec{p}_1) S_2(x_2, \vec{p}_2)}$$

## AMPT model

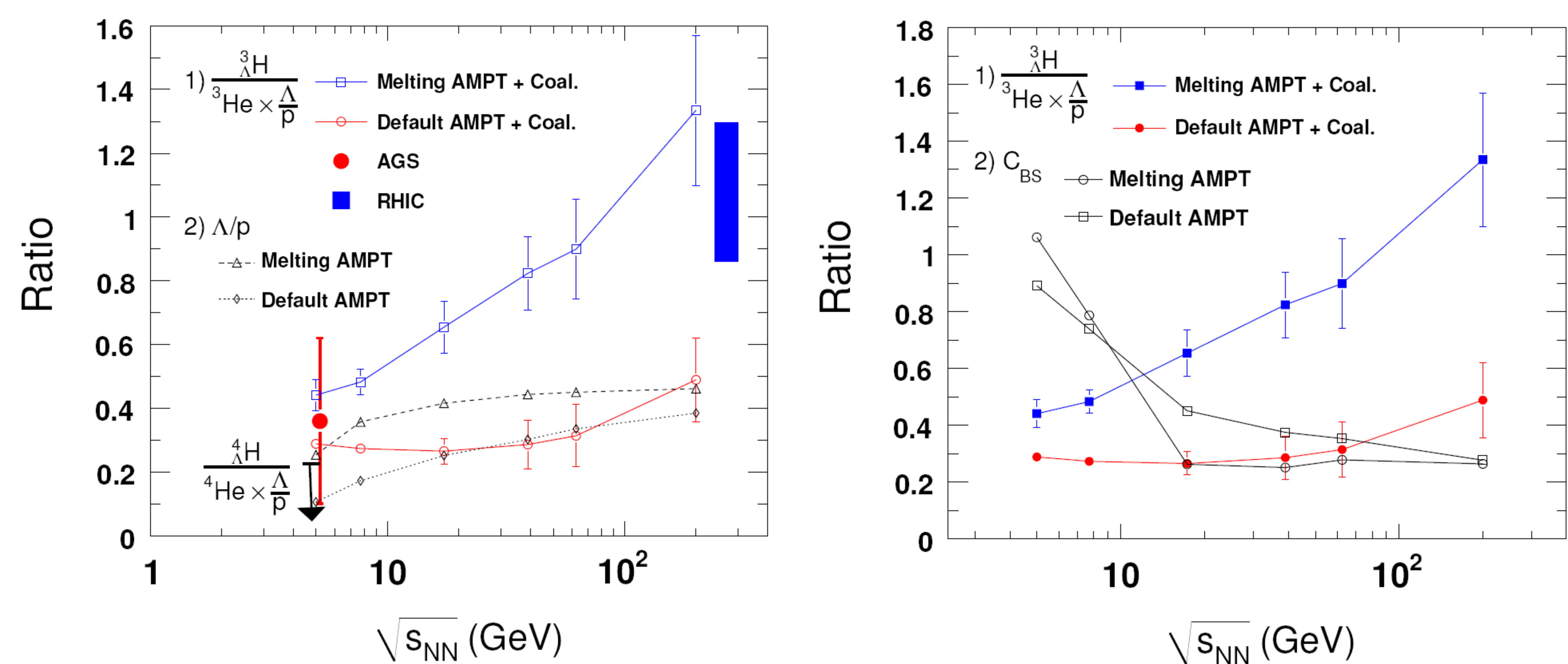


- (1) Initial condition (HIJING); (2) Parton cascade (ZPC);  
(3) Hadronization; (4) Hadronic rescattering (ART)

## Reference

- [1] S. Zhang, J. H. Chen, H. Crawford et al., Phys. Lett. B 684, 224 (2010).  
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[3] B. I. Abelev et al. (STAR Collaboration), arXiv:0909.0566.  
[4] H. H. Gutbrod et al., Phys. Rev. Lett. 37, 667 (1976).  
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## Strangeness population Factor, $S_3 = \frac{^3\text{H}}{^3\text{He} \times \frac{\Lambda}{p}}$

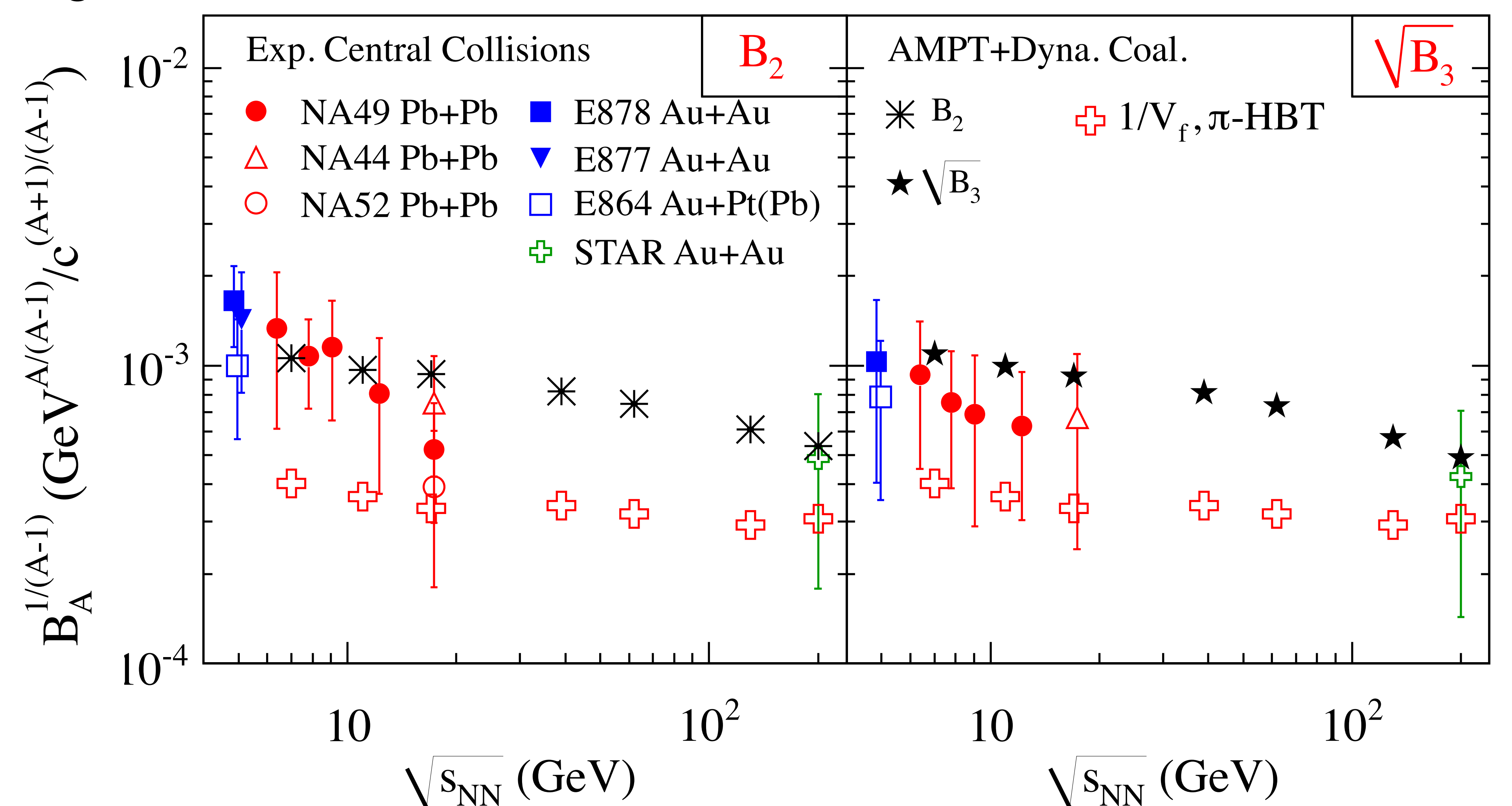


- ◆ Strangeness population factor is sensitive to the freedom degree of the dense medium created in HIC, namely sensitive to the deconfinement;
- ◆ A valuable tool to probe the nature of the dense matter created in HIC

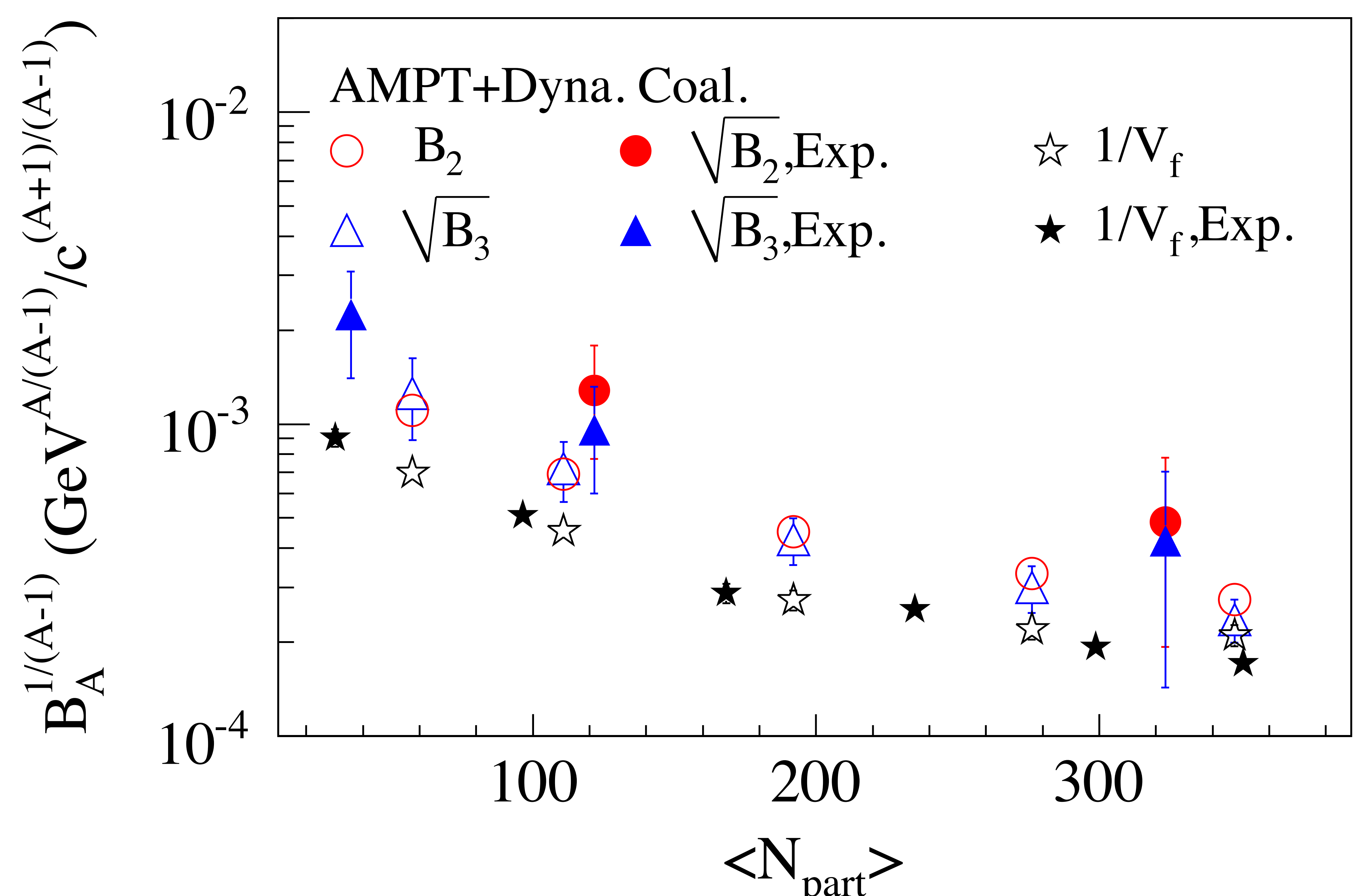
## Coalescence parameter and freeze-out correlation volume

Coalescence parameter:  $B_A = \frac{d^2 N_A / d^3 p_A}{(d^3 N_p / d^3 p_p)^A}$ , assuming neutrons have the same distribution of protons

Freeze-out correlation volume:  $V_f = (2\pi)^{3/2} R_{side}^2 R_{long}$ ,  $R_{side}$  and  $R_{long}$  are the sideward and longitudinal radii from HBT.



- ◆ The coalescence parameters and reverse of freeze-out volume decrease with increasing of beam energy
- ◆ The evolution of collision zone can reach a larger system volume at more high energy collisions



- ◆ The coalescence parameters and reverse of freeze-out volume decrease with increasing of number of participant

## Summary

This beam energy and system size dependences indicate the increase of source size in more high energy collisions and in more central collisions.