Probing nuclear parton densities and parton energy loss processes through photon + heavy-guark jet production in p-A and A-A collisions Quark Matter 2011

Heavy Quark Energy Loss in $\gamma + Q$ Production

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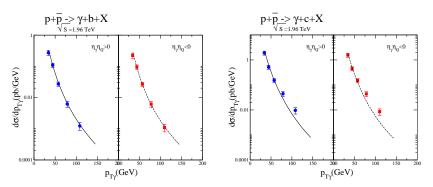
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- ▶ Prompt photons produced in hard scattering or via fragmentation
- ▶ Heavy quark jet charm or bottom jet
- Depending on the collision type this process can be useful in various ways
- ▶ In $p \bar{p}$ collisions (Tevatron) it is sensitive to the charm (bottom) PDF arXiv:0901.3791 useful for testing for intrinsic charm or bottom (IC or IB) (theory data comparison : arXiv:0901.0739).
- ▶ In p-p, p-A collisions (LHC,RHIC) can test for IC/IB again (work in progress) and can be used to constrain the gluon nuclear PDF (nPDF) (arXiv:1012.1178) (dominant subprocess is g-Q initiated)
- Knowing the precise nPDFs is necessary for obtaining reliable predictions in A – A collisions!
- ▶ In A A collisions it can help us obtain a better understanding of the parton energy loss processes in the massive quark sector (work in progress)

Comparison between theory and data @ $p - \bar{p}$

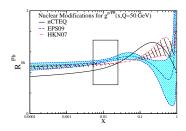
Measurements by DØ Collaboration

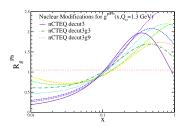


- ▶ Really good agreement between data and theory for $\gamma + b$
- ▶ For $\gamma + c$ data at large $p_{T\gamma}$ is above the theory curve \rightarrow possible explanation existence of intrinsic charm

▶ The nuclear PDFs are far less constrained than the proton PDFs. especially the gluon nuclear PDF. It is largely unconstrained, illustrated by the nuclear modification factor:

$$R_g^{Pb}(x,Q) = g^{p/Pb}(x,Q)/g^p(x,Q)$$

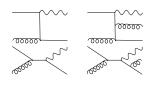




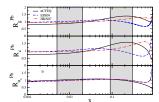
▶ Different nPDF sets (nCTEQ, EPS and HKN) including the errors → differing predictions - need a more precise determination of $g^{p/Pb}(x,Q) \Rightarrow LHC$ data is needed! ◆□ → ◆周 → ◆三 → ◆ □ → ◆ ○ ◆

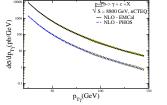
How can $\gamma + Q$ help?

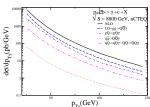
- At LO only one hard scattering subprocess Compton subprocess g − Q initiated
- ▶ Standard approach: HQ PDFs are generated radiatively \Rightarrow $R_{\sigma}^{Pb} \simeq R_{c}^{Pb}$



▶ Direct access to gluon nPDF







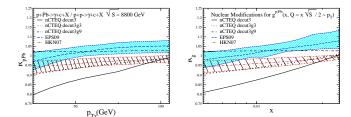
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- ▶ g & Q initiated subprocesses dominate (>80%) ⇒ sensitivity to gluon and HQ PDFs.
- ▶ Using an integrated yearly luminosity of $\mathcal{L} = 10^{-1} pb^{-1}$ a precursory number of events per year at EMCal for $\gamma + c$ is $\mathcal{N}_{\gamma+c}^{pPb} = 11900$ $(\sigma_{\gamma+c}^{PPb}=119nb)$ and for $\gamma+b$ is $\mathcal{N}_{\gamma+b}^{PPb}=2270$ $(\sigma_{\gamma+b}^{PPb}=22.7nb)$

Nuclear Modifications

$$R_{pA}^{\gamma Q} = \frac{\sigma(pA \to \gamma \ Q \ X)}{A \ \sigma(pp \to \gamma \ Q \ X)}$$

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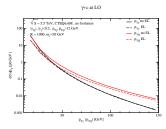
- $ightharpoonup R_{p,A}^{\gamma Q} \simeq R_{\sigma}^{Pb}$ in the x region probed at ALICE
- ▶ Measurements of $\gamma + Q$ with appropriate error bars will allow to distinguish between the different nPDF sets and place useful constraints on the gluon nPDF (arXiv:1012.1178)

$\gamma + Q$ in A - A Collisions

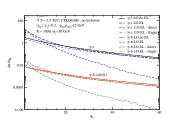
- ▶ Single hadron production R_{AA} not enough to constrain/quantify the amount of energy loss of partons in QGP \Rightarrow need more differential observables, like double inclusive particle production and in particular $\gamma + jet$ correlations (X.-N. Wang, Z. Huang, I. Sarcevic hep-ph/9605213, F. Arleo, P. Aurenche, Z. Belghobsi, J.-P. Guillet hep-ph/041008)
- $ightharpoonup \gamma + Q$ is ideal for probing hot QCD medium
- Q Jet Quenching
- $ightharpoonup \gamma$ is medium insensitive \Rightarrow can gauge HQ's initial energy
- ▶ $\gamma + Q$ can help to clarify the energy loss in the heavy quark sector $(\epsilon_q > \epsilon_c > \epsilon_b)$ (Heavy quark colorimetry of QCD matter -Y. L. Dokshitzer, D.E. Kharzeev)
- ► The two-particle final state further offers a range of observables ⇒ can get a better handle on the energy lost Photon-jet energy asymmetry, Momentum imbalance, Photon-jet pair momentum
- ▶ The energy loss of the heavy quark ϵ_Q is computed on an event by event basis, with the use of the quenching weight obtained perturbatively [Armesto Dainese Salgado Wiedemann 2005 (arXiv:hep-ph/0501225)]
- work in progress T.S., F.Arleo, I. Schienbein



Effects of energy loss on the $\gamma+Q$ cross-section - LO



- ▶ Difference in spectrum vs p_{TQ} in vacuum and in medium ⇒ due to energy loss
- $\frac{d\sigma}{pT\gamma}$ spectrum almost unchanged (small difference between $\frac{d\sigma^{med}}{pT\gamma}$ and $\frac{d\sigma^{vac}}{pT\gamma}$ at low p_T due to experimental cuts)

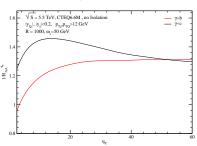


- $ightharpoonup q_{\perp}$ in more detail
- ightharpoonup At LO direct component $q_{\perp} \simeq \epsilon_{\it Q}$
- At LO fragmentation component ε_Q represents the shift of the q_T spectrum in vacuum vs the one in medium

- ▶ Direct and fragmentation components behave very distinctly
- ▶ In medium the direct contribution decreases sharply with increasing q_T \Rightarrow small probability of events with large ϵ_Q
- ▶ In vacuum the direct contribution is non-zero only at $q_T = 0$
- ► Therefore compare only the vacuum and medium fragmentation contributions

q in more detail - Il

LO Fragmentation Contribution



- ▶ $\Delta E_c > \Delta E_b$; as q_T grows the difference disappears, as the quenching weight depends on m/E, which becomes similar for $\gamma + c$ and $\gamma + b$ at large q_T
- Need to compare σ in medium and vacuum at the NLO level, where the particles have a larger kinematic phase-space!

 $\gamma + Q$ production is a very useful process. It can can be employed to constrain the HQ PDFs in hadron-hadron collisions, while measurements in p-A collisions can help constrain the gluon nPDF. In A-A collisions it can be used for an estimate of the HQ energy loss, where it also provides access to the mass hierarchy of parton energy loss.

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