# **Directed flow and early thermalization**

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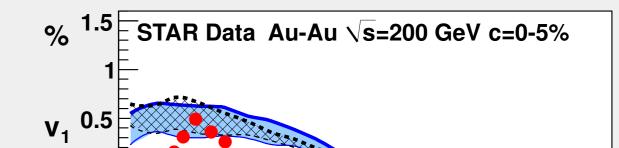
#### Abstract

The generation of the directed flow of particles emitted from the fireball created in heavy-ion collisions at RHIC is described using a 3+1Dhydrodynamical model. The initial density of the tilted fireball is constructed as a sum of contributions from forward and backward going participants. Our model reproduces the experimentally observed negative directed flow in a wide range of central pseudorapidities and reproduces correctly the scaling of the directed flow when going from Au-Au to Cu-Cu systems [1]. The directed flow is a very sensitive measure of the pressure equilibration in the first  $\mathbf{1}$  fm/c of the evolution, it is strongly reduced in the presence of even a very short pressure anisotropy. Our calculations show that the system must thermalize fast (< 0.25 fm/c). This suggests that the matter behaves as a strongly coupled system already at the first stages [2]

## **Collective flow**

For non-central collisions the interaction region is azimuthally asymmetric and asymmetric emissions of particles takes place as a results of the collective expansions of matter. The effect can be quantified in therms of Fourier coefficients in the expansion of the measured particle spectra

#### **Directed flow**

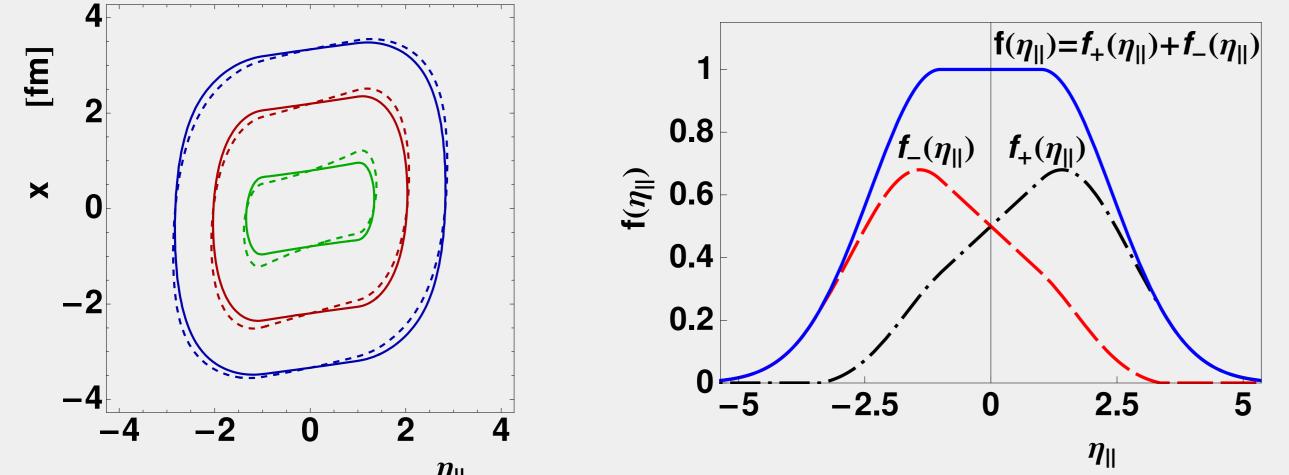


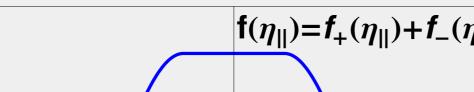
 $\frac{\mathrm{dN}}{\mathrm{d}^2 \mathbf{p}_{\perp} \mathrm{d}\eta} = \frac{\mathrm{dN}}{2\pi \mathbf{p}_{\perp} \mathrm{d}\mathbf{p}_{\perp} \mathrm{d}\eta} (1 + 2\mathbf{v}_1 \cos(\phi) + 2\mathbf{v}_2 \cos(2\phi) + ...)$ 

The coefficient  $v_1$  of the directed flow is measured at RHIC energies. It is zero at zero rapidity for collisions for symmetric nuclei, and becomes negative (positive) when moving to forward (backward) pseudorapidities  $\eta$ .

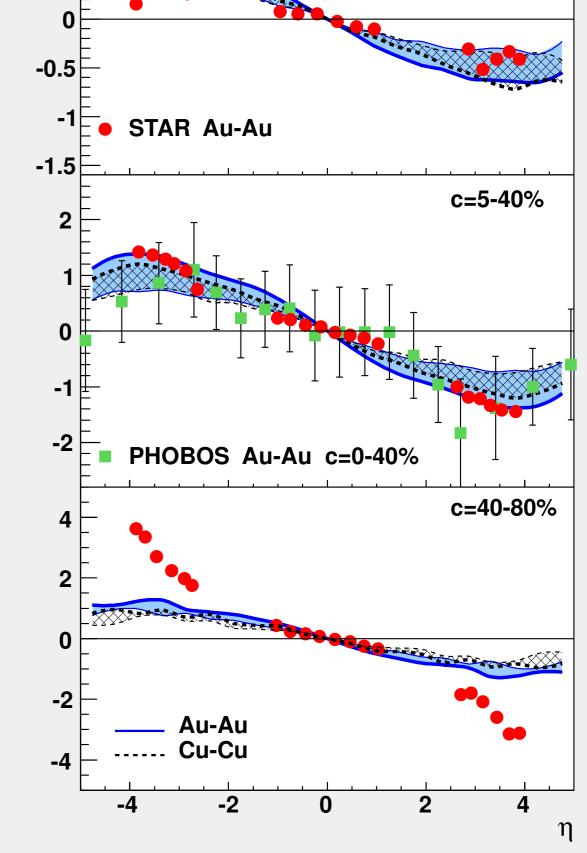
### Initial conditions

The formation of the directed flow in ultrarelativistic collisions requires a mechanism that breaks the symmetry with respect to the collisions axis and some effective transverse and longitudinal acceleration of the fluid elements. The asymmetric emission in space-time rapidity from forward and backward going wounded nucleons [3] results in a tilt of the initial fireball away from the collisions axis [left panel]. Therefore we take the initial profile as a sum of two contribution  $f_+$  and  $f_-$  representing the emissions from forward and backward going participant nucleons[right panel].



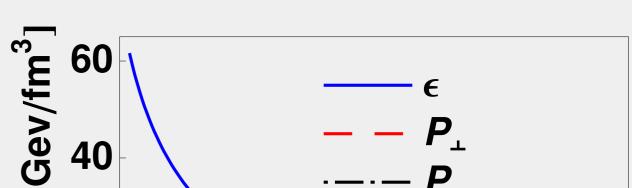


Using our 3+1 dimensional hydrodynamic model with a tilted source we have calculated the directed flow of charged particles. The thick solid lines represent the results for Au-Au collisions and dashed lines correspond to Cu-Cu collisions. The experimental data [4-5] are reproduced in the central rapidity region in a satisfactory way. The shaded bands in the figure represent the uncertainty of the model.



## 3+1D expansion with off-equilibrium pressure

We study the dynamics of the system with anisotropic pressures. A phenomenological correction is added to the ideal fluid energy-momentum ten-



[fm/c]

The mechanism generating the directed collective flow from a tilted initial source can be understood when considering the acceleration equations of relativistic hydrodynamics. For small initial times the accelerations in the transverse **x** and longitudinal (space-time rapidity)  $\eta_{\parallel}$  directions are

$$\partial_{\tau} \mathbf{v}_{\mathbf{x}} = -\frac{\partial_{\mathbf{x}} \mathbf{P}_{\perp}}{\epsilon + \mathbf{P}}$$
  
 $\partial_{\tau} \mathbf{Y} = -\frac{\partial_{\eta_{\parallel}} \mathbf{P}_{\parallel}}{\tau(\epsilon + \mathbf{P})}$ 

According to the above equations, the acceleration of the fluid elements in the tilted source generates anti-flow ( $v_1 < 0$  for  $\eta > 0$ ). It should be stressed that generating the directed flow requires non-zero gradients of both the transverse  $\mathbf{P}_{\perp}$  and longitudinal  $\mathbf{P}_{\parallel}$  pressures.

sor resulting in anizotropic pressures

$$\mathsf{P}_{\perp} - \mathsf{P}_{\parallel} = rac{3}{2} \mathsf{P}_{\mathrm{eq}}( au_0) \exp\left(-rac{ au - au_0}{ au_{\mathrm{ico}}}
ight)$$

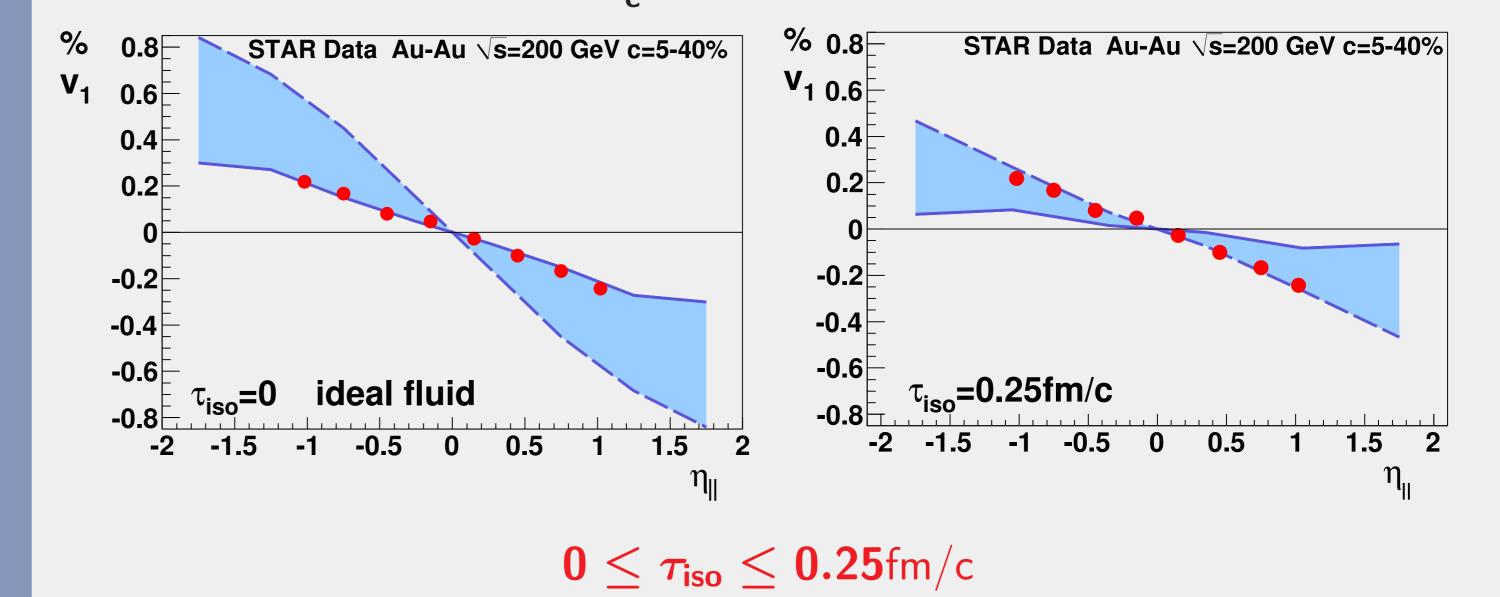
It means that initially the longitudinal pressures is zero and that the pressure anisotropy decreases with a relaxation time  $au_{iso}$ .

## Early thermalization

To estimate the thermalization time  $\tau_{iso}$  we use two extreme assumptions for the value of the source tilt within the Glauber model. For the smaller tilt the experimental data are described using  $\tau_{iso} = 0$ . The expansion of the source with the larger tilt is compatible with the data if the longitudinal pressure is retarded by  $0.25\frac{\text{fm}}{c}$  with respect to the transverse pressure.

ر **م**<sup>=</sup> 20

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#### References

[1] P. Bozek and I. Wyskiel *Directed flow in ultrarelativistic heavy-ion collisions*, Phys. Rev. **C 81** 054902 (2010)

[2] P. Bozek and I. Wyskiel-Piekarska Indications of early thermalization in relativistic *heavy-ion collisions*, Phys. Rev. C 83 024910 (2011) [3] A. Białas and W. Czyż, Acta Phys. Polon. **B36**, 905 (2005) [4] B. B. Back *et al.* (PHOBOS), Phys. Rev. Lett. **97**, 012301 (2006)

[5] B. I. Abelev et al. (STAR), Phys. Rev. Lett. **101**, 252301 (2008)

The conclusion of this analysis is that the isotropization time of the pressure is smaller than  $0.25\frac{\text{fm}}{c}$ . Such a small value of the delay for the appearance of the longitudinal pressure indicates that the system is strongly coupled.

#### Conclusions

We demonstrate in explicit hydrodynamic calculation that the directed flow

- ▶ is **formed in the first stage** of the expansion,
- is built trough a simultaneous action of the transverse and longitudinal pressures Using the directed flow of particle we estimate that the thermalization time is smaller than  $0.25\frac{\text{fm}}{2}$