

# Direct Photons in PHENIX at RHIC

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## **Abstract.**

PHENIX results on direct photon production in  $p+p$ ,  $d+Au$ ,  $Cu+Cu$  and  $Au+Au$  collisions are presented. Direct photon yield at high  $p_T$  is consistent with the dominance of the LO processes (inverse Compton scattering and quark-antiquark annihilation). The low  $p_T$  direct photons exhibit an order of magnitude exponential enhancement consistent with thermal emission from QGP. The thermal nature of this enhancement is further confirmed by large elliptic flow of direct photons below  $p_T \sim 5 \text{ GeV}/c$ .

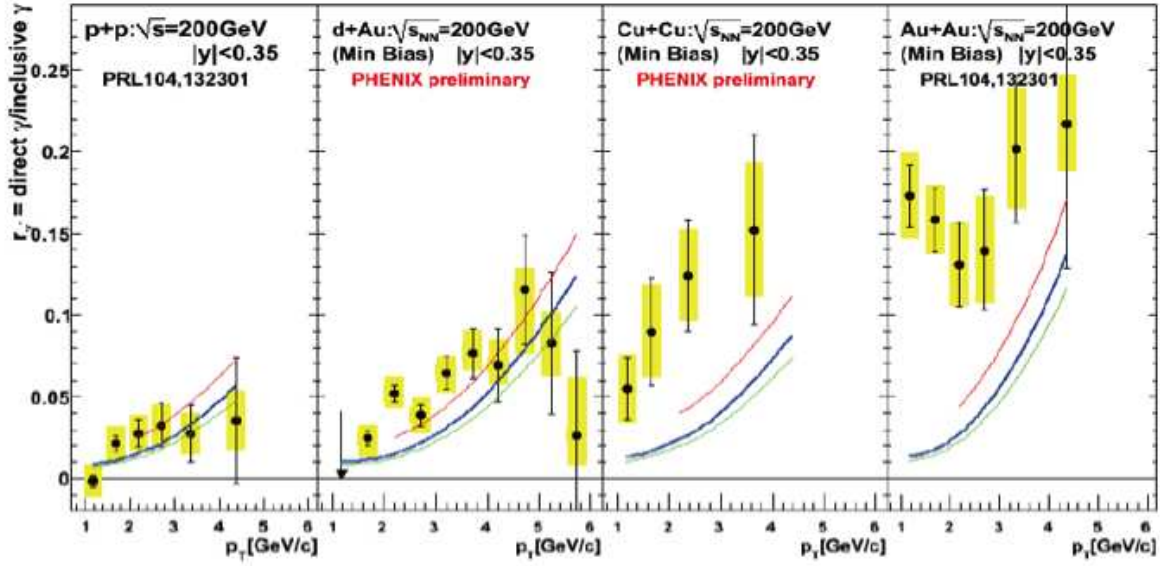
## **1. Introduction**

Among the observables used to probe the high temperature and high density phase of heavy nucleus collisions direct photons are considered of particular interest. Nuclear modifications to the yield of direct photons reflects media effects on contributing production mechanisms. The anisotropy of high  $p_T$  photon production reflects the geometry of the medium. The emission from expanding media influences both the rate and azimuthal asymmetries of photon production at low to medium  $p_T$  (1-5 GeV/c) [1, 2].

## **2. Direct photon measurements**

The data presented in this paper are from  $p+p$ ,  $d+Au$ ,  $Cu+Cu$  and  $Au+Au$  data sets at  $\sqrt{s_{NN}}=200$  GeV taken with the PHENIX detector [3] in 2004-2008. The PHENIX central arms, each covering  $\pm 0.35$  units of pseudorapidity around midrapidity and  $90^\circ$  in azimuth, contain charged-particle tracking chambers and two kinds of electromagnetic calorimeters. The BBC and Zero-Degree Calorimeters (ZDC) were used for minimum bias event selection and centrality determination (when appropriate).

Direct photon yields, their flow and gamma-jet correlations are measured by statistical subtraction of the estimated meson (mainly  $\pi^0$ ) decay photon contribution from the inclusive photon and  $\gamma-h$  samples. An alternative method for measuring direct photons is provided by the low-mass lepton  $e^+e^-$  pairs produced by a higher order QED correction to the real photon emission process. Any source of real photons must also emit virtual photons and their yield is related to that of real photons. In the low mass



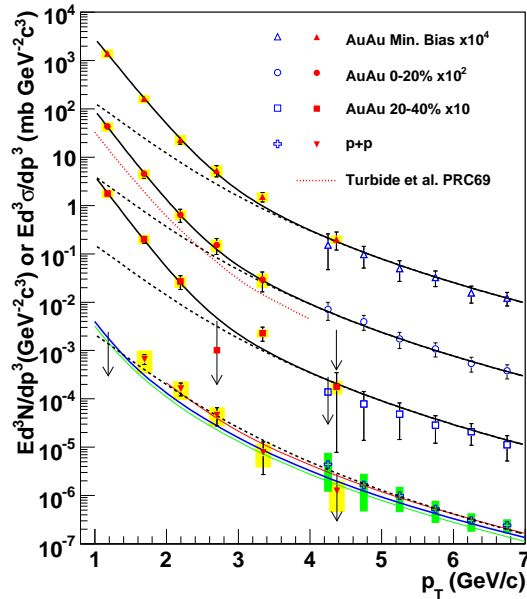
**Figure 1.** (color online) The fraction  $r_\gamma$  of direct photons in the inclusive photon yield as a function of  $p_T$  in  $p+p$ ,  $d+Au$ ,  $Cu+Cu$  and  $Au+Au$  (min. bias) collisions [1]. The error bars and the boxes represent statistical and systematic uncertainties, respectively. The curves are from an NLO pQCD calculation.

region, where the  $p_T$  of the  $e^+e^-$  pair is much greater than its mass ( $m_{ee} \ll p_T$ ), the yield of the intermediate virtual photons is approximately the same as that of real photons. Therefore, in this quasi-real virtual photon region, the production of real direct photons can be deduced from measurements of  $e^+e^-$  pairs.

Figure 1 shows the fraction  $r_\gamma = \frac{\text{direct } \gamma}{\text{inclusive } \gamma}$  of the direct photon component in  $p+p$ ,  $d+Au$ ,  $Cu+Cu$  and  $Au+Au$  collisions, respectively. The curves represent the expectations from a NLO pQCD calculation [4]. The three curves correspond (from top to bottom) to the theoretical scales set to  $\mu = 0.5 p_T$ ,  $p_T$ , and  $2 p_T$ , respectively. While the fraction  $r_\gamma$  is consistent with the NLO pQCD calculation in  $p+p$  and  $d+Au$ , it is larger than the calculation in both  $Cu+Cu$  and  $Au+Au$  for  $p_T < 5$  GeV/c.

The direct photon yield computed using the relation  $dN_\gamma^{dir}(p_T) = r_\gamma \times dN_\gamma^{incl}(p_T)$  is compared in Figure 2 to the direct photon data from [5, 6] and NLO pQCD calculations. In central  $Au+Au$  it shows excess over collision scaled  $p+p$  data, and the shape of the excess is well described by an exponential with inverse slope  $T \simeq 220$  MeV. If the direct photons in  $Au+Au$  collisions are of thermal origin, the inverse slope  $T$  is related to the initial temperature  $T_{init}$  of the dense matter. In hydrodynamical models,  $T_{init}$  is 1.5 to 3 times  $T$  due to the space-time evolution [7].

If photons are radiated inside an expanding matter having azimuthally anisotropic momentum distribution, their momenta add or subtract for radiation along or opposite to the motion. Neglecting pion mass, thermal photons must have the same or greater  $v_2$  as pions [8]. An earlier low  $p_T$  measurements of photon flow ( $v_2$ ) for  $\pi^0$  and inclusive

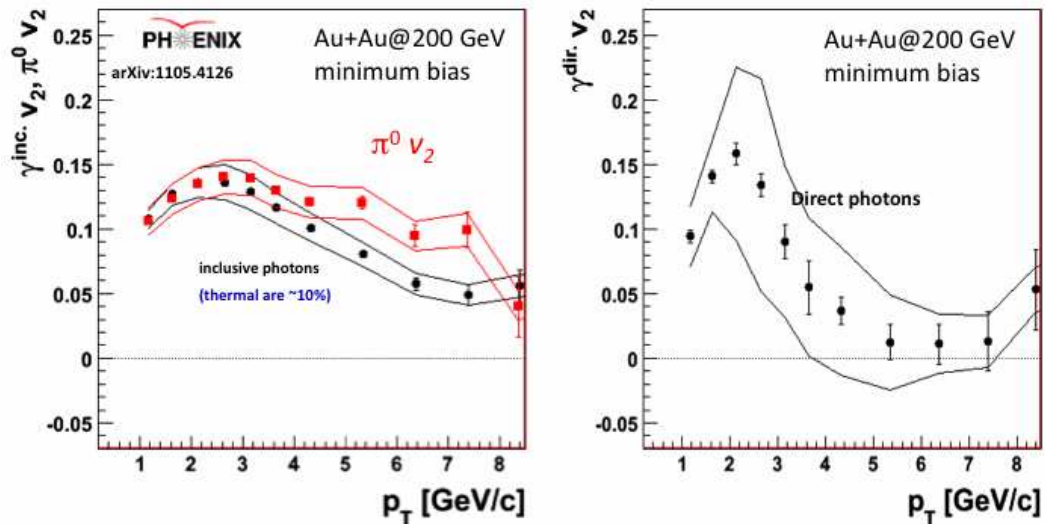


**Figure 2.** (color online) Invariant cross section ( $p+p$ ) and invariant yield ( $Au+Au$ ) of direct photons as a function of  $p_T$  [1, 5, 6]. The solid curves on the  $p+p$  data represent NLO pQCD direct photon calculations [4]. The dashed curves on the  $Au+Au$  data show the  $p+p$  fit scaled by  $T_{AA}$ . The solid curves on the  $Au+Au$  data are an exponential fit plus the  $T_{AA}$  scaled  $p+p$  fit.

photons) has been published in [9]. Using the high statistics  $p+p$  and  $Au+Au$  data from 2007run PHENIX extended  $p_T$  range for  $v_2$  measurements to 15 GeV/c and dramatically improved  $v_2$  precision in the low to medium  $p_T$  range.

To correct for a large contribution from hadron decays, predominantly from  $\pi^0$  (80%) and  $\eta$  (15%), and cluster merging for high  $p_T$   $\pi^0$  the elliptic  $v_2$  flow of the direct photons was calculated as  $v_2^{\gamma,dir} = \frac{R_\gamma(p_T) \times v_2^{\gamma,inc} - v_2^{\gamma,bg}}{R_\gamma(p_T) - 1}$ , where  $N^{inc} = N^{meas} - N^{hadr}$  is for inclusive photons and  $R_\gamma(p_T) = N^{inc}(p_T)/N^{bg}(p_T)$  is the direct photon excess ratio. The  $v_2$  values for  $\pi^0$  in minimum bias  $Au+Au$  events are compared to similar data for inclusive photons in Figure 3 left panel [2]. The  $v_2$  data for direct photons in the most central  $Au+Au$  events are shown in the right panel in the same Figure.

Two sets of points (inclusive photons and  $\pi^0$ ) are barely different in the thermal  $p_T$  range indicating the dominance of the component with flow values close to that of hadrons. Flow of photons drops at high  $p_T$  consistent with a growing contribution from hard scattering. Qualitative conclusions of the left panel are confirmed quantitatively in the right panel displaying subtraction data for direct photons in the most central  $Au+Au$  collisions. The measured  $v_2$  has value comparable to that of hadrons below  $p_T \sim 5 \text{ GeV}/c$  and drops to zero at higher  $p_T$  as expected if LO processes (reverse Compton scattering and quark-antiquark annihilation) dominate in that  $p_T$  range.



**Figure 3.** (color online) Left panel: Comparison between elliptic flow ( $v_2$ ) of  $\pi^0$  and inclusive photons in mim. bias  $Au+Au$  collisions. Right panel: Direct photon flow in the most central  $Au+Au$  collision computed using a subtraction procedure.

### 3. Summary

The direct photon production data accumulated by PHENIX confirm the unique role of direct photons in probing the prehistory of the sQGP and its properties. Consistency between high  $p_T$  yields of direct photons in  $Au+Au$  collisions and collision scaled  $p+p$  yield confirms LO dominance at high  $p_T$  and validity of collision scaling. The low  $p_T$  behavior of direct photons in central  $Au+Au$  collisions is dramatically different from all other particles exhibiting an order of magnitude exponential enhancement as  $p_T \rightarrow 0$ , suggestive of thermal emission from the sQGP. The assumed thermal nature of low to medium  $p_T$  photons is further confirmed by the presence of a large flow of direct photons in the  $p_T$  range below 5 GeV/c.

### References

- [1] A. Adare et al., Phys. Rev. C **81**, 034911 (2010)
- [2] arXiv:1105.4126v1 [nucl-ex] 20 May 2011
- [3] K. Adcox et al. (PHENIX), Nucl. Instrum. Meth., **A499** 469 (2003)
- [4] L. E. Gordon and W. Vogelsang, Phys. Rev. D **48**, 3136 (1993) and W. Vogelsang, private communication (2008).
- [5] S.S. Adler et al., Phys. Rev. Lett. **94**, 232301 (2005).
- [6] S.S. Adler et al., Phys. Rev. Lett. **98**, 012002 (2007).
- [7] D. d'Enterria and D. Peressounko, Eur. Phys. J. **C46**, 451 (2006).
- [8] B. Kopeliovich, Private communication (May 2011)
- [9] S. Adler et al., Phys. Rev. Lett. **96**, 032302 (2006)