

## Abstract

Jet quenching parameter is an important quantity in order to understand energy losses in heavy ion collisions and to get insights into properties of deconfined quark-gluon plasma. Soft Collinear Effective theory (SCET) provides framework to define momentum broadening of probing quark/gluon and thus define jet quenching parameter as the expectation value of two space-like separated light-like Wilson lines which can be evaluated for the desired medium. In this work we evaluate transverse momentum broadening distribution function at weak coupling for quark-gluon plasma in thermal equilibrium using Hard Thermal Loop (HTL) resummed effective thermal field theory and build the framework for going beyond HTL approximation.

## Hard Thermal Loops

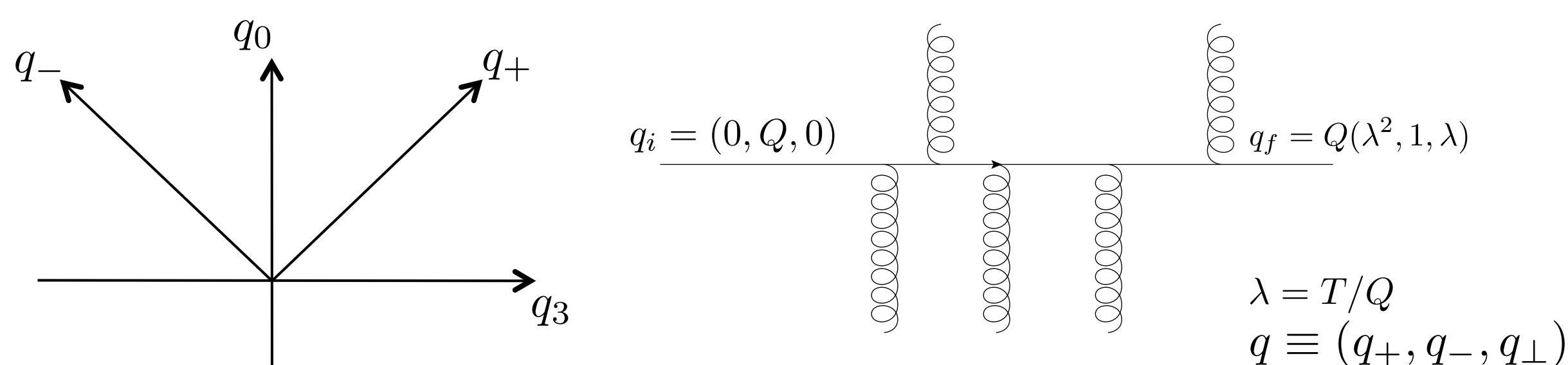
- When using (bare) thermal field theory, loop corrections to self-energy are of order  $(gT)^2/P^2$ .
- When external momentum is "soft" ( $P^2 \sim (gT)^2$ ), loop expansion breaks down ( $g \ll 1$ ).
- Hard Thermal Loops (HTL) is an effective resummed field theory, which enables loop expansion when external momentum is soft.
- Nevertheless, HTL approximation is not valid when external momentum is "hard" ( $P^2 \sim T^2$ ).

## Problem set up

- We are interested in evaluating the jet quenching parameter, defined by

$$\hat{q} = \frac{1}{L} \int \frac{d^2 q_{\perp}}{(2\pi)^2} q_{\perp}^2 P(q_{\perp}),$$

where  $L$  is the longitudinal length of the thermal medium in the direction in which the probing quark(gluon) is propagating.



Definition of light-cone coordinates(left) and depiction of hard (Collinear) probe interacting with medium via thermal (Glauber) gluons which induce transverse momentum broadening.

- Transverse momentum broadening distribution function of quark(gluon) traveling through general medium is given by [1,2,3]:

$$P(q_{\perp}) = \int d^2 x_{\perp} e^{-iq_{\perp} \cdot x_{\perp}} \mathcal{W}_{\mathcal{R}}(x_{\perp}), \text{ normalized as } \int \frac{d^2 q_{\perp}}{(2\pi)^2} P(q_{\perp}) = 1$$

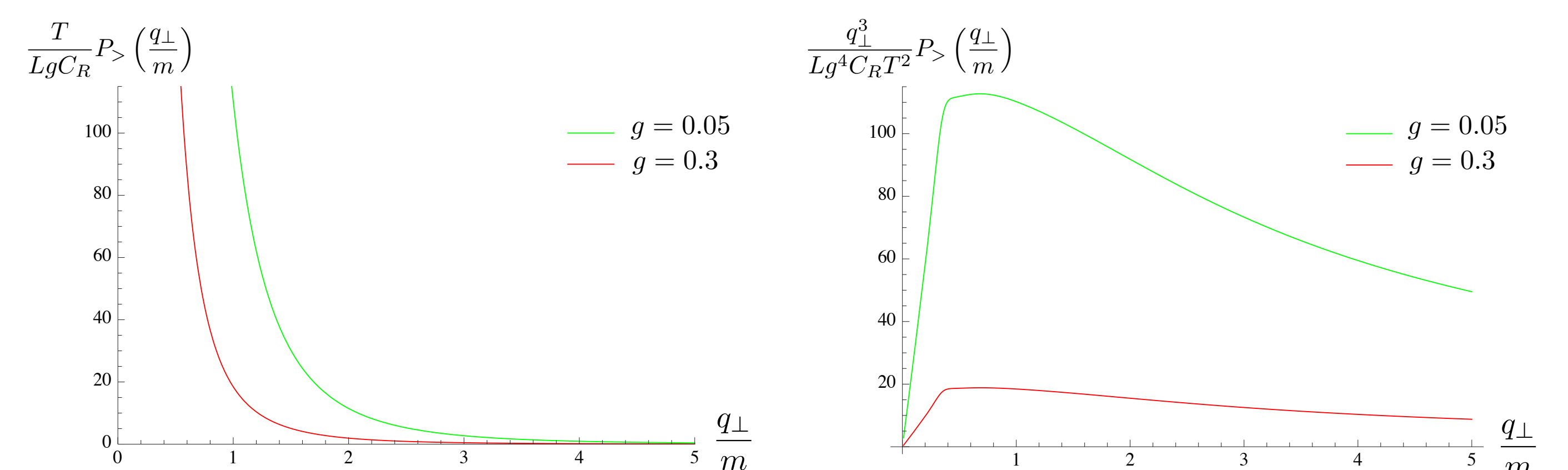
$$\text{and } \mathcal{W}_{\mathcal{R}}(x_{\perp}) = \frac{1}{d(\mathcal{R})} \left\langle \text{Tr} \left[ W_{\mathcal{R}}^{\dagger}[x^+ = 0, x_{\perp}] W_{\mathcal{R}}[x^+ = 0, 0] \right] \right\rangle$$

$W_{\mathcal{R}}$  is a Wilson line along  $x^-$  direction of length  $L^-$ ;  $\mathcal{R}$  is the  $SU(N)$  representation to which the collinear particle belongs.  $\langle \rangle$  in this case is an average over thermal medium in equilibrium.

## Expression for $P(q_{\perp})$

- Using HTL approximation transverse momentum broadening distribution function takes the form:

$$P_{>}(q_{\perp}) = g^2 C_R \sqrt{2} L \int_{-\infty}^{\infty} \frac{dq_{-}}{2\pi} q_{\perp} \coth\left(\frac{q_{-}}{2\sqrt{2}T}\right) \rho_{HTL}^{++}(0, q_{-}, q_{\perp})$$

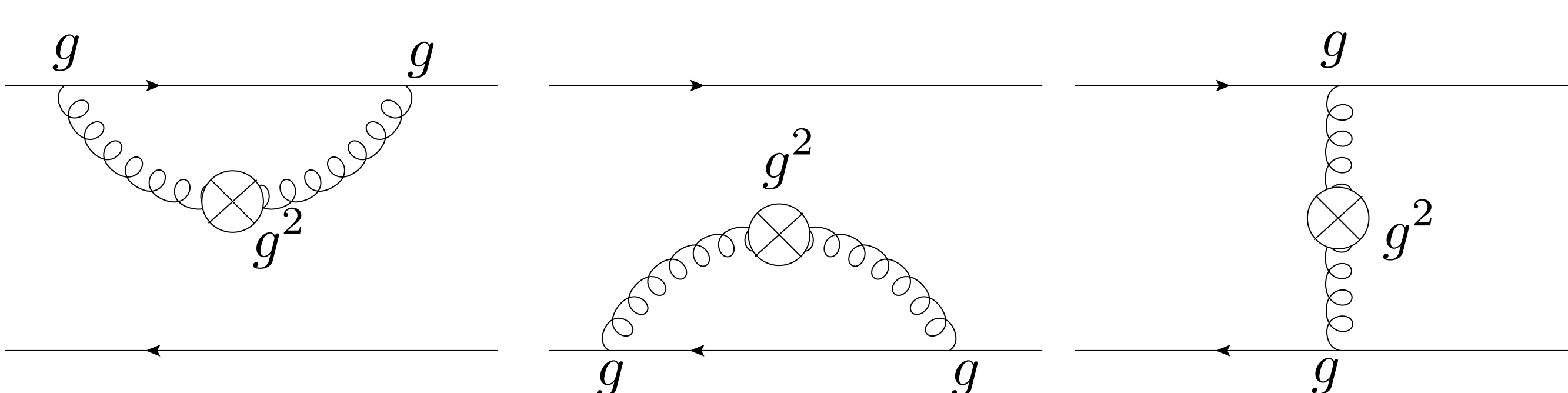


Transverse momentum broadening distribution (left) and distribution multiplied by  $q_{\perp}^3$  (right) which is used when evaluating  $\hat{q}$ ;  $m = m_{QCD} = \frac{\sqrt{3}}{2} gT$  is thermal mass.

- It is ensured that expression for  $\hat{q}$  is IR finite and gauge fixing parameter independent.
- For  $q_{\perp} \lesssim gT$  the contribution of  $q_{-}$  from outside of "soft" region when evaluating  $P(q_{\perp})$  is estimated to be small. So non-HTL corrections to  $P(q_{\perp} \lesssim gT)$  are assumed to be small.
- For  $q_{\perp} \gtrsim T$ , HTL is not valid and appropriate expression for spectral function  $\rho^{++}$  needs to be used.

## Estimating contributions

- Tree level diagrams involving bare propagators or vertices vanish.
- Leading order contribution to  $P(q_{\perp})$  is of order  $g^4$ .



Leading order diagrams: parton self-energy (left and middle), gluon exchange (right)

- General form  $P(q_{\perp}) = g(q_{\perp}^0) \delta^2(q_{\perp}) + P_{>}(q_{\perp})$ , where  $P_{>}(q_{\perp})$  corresponds to the gluon exchange diagram and term proportional  $\delta^2(q_{\perp})$  originates from the first two diagrams and does not contribute to  $\hat{q}$ .

## Comments and Outlook

- Wilson line expression for momentum broadening obtained using SCET is taken one step further by evaluating it in equilibrium quark-gluon plasma with the help of HTL.
- We obtained expression for transverse momentum broadening distribution function  $P(q_{\perp})$  for "soft" momentum.
- The result is general and it is one of the first attempts to calculate the jet quenching parameter using full field theoretical methods.
- The work of including (bare) thermal field theory contributions to  $P(q_{\perp})$  is under progress.

## References

- [1] F. D'Eramo, H. Liu and K. Rajagopal, arXiv:1006.1367 [hep-ph].
- [2] J. Casalderrey-Solana and C. A. Salgado, Acta Phys. Polon. B 38, 3731 (2007) [arXiv:0712.3443 [hep-ph]].
- [3] Z. T. Liang, X. N. Wang and J. Zhou, Phys. Rev. D 77, 125010 (2008) [arXiv:0801.0434 [hep-ph]].