Higgs Physics and Beyond the Standard Model (mainly SUSY)

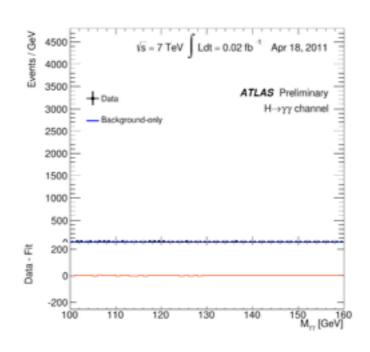
Koichi Hamaguchi (University of Tokyo)

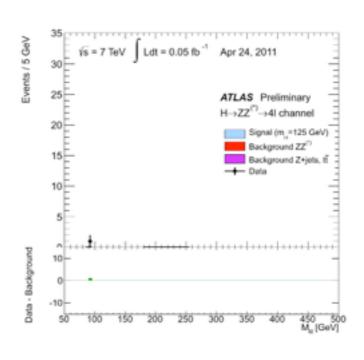
@AEPSHEP2014, Puri, November 11–16, 2014

Discovery of the Higgs boson

2012. July 4



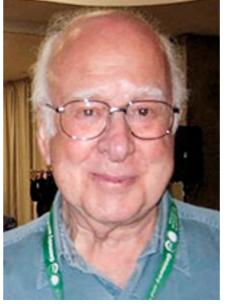


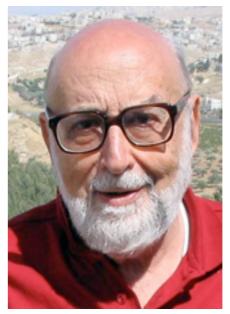


2013. October 8

François Englert
Peter W. Higgs



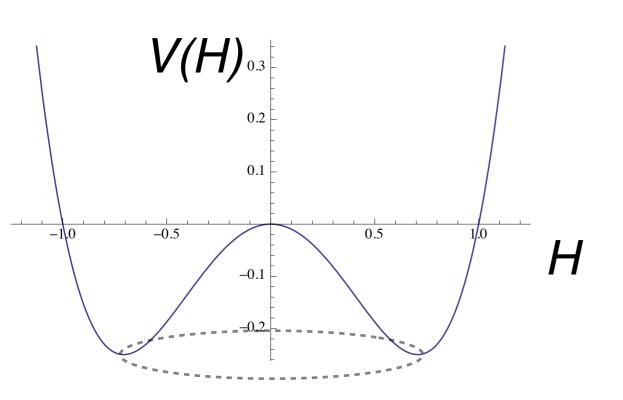






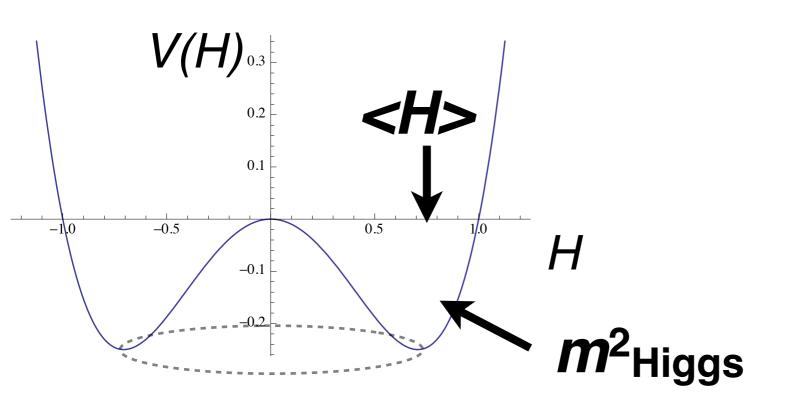
R. Brout (1928-2011)

$$V(H) = -m^2(H^{\dagger}H) + \lambda_{H}(H^{\dagger}H)^2$$



$$V(H) = -m^2(H^{\dagger}H) + \lambda_{\mathbf{H}}(H^{\dagger}H)^2$$

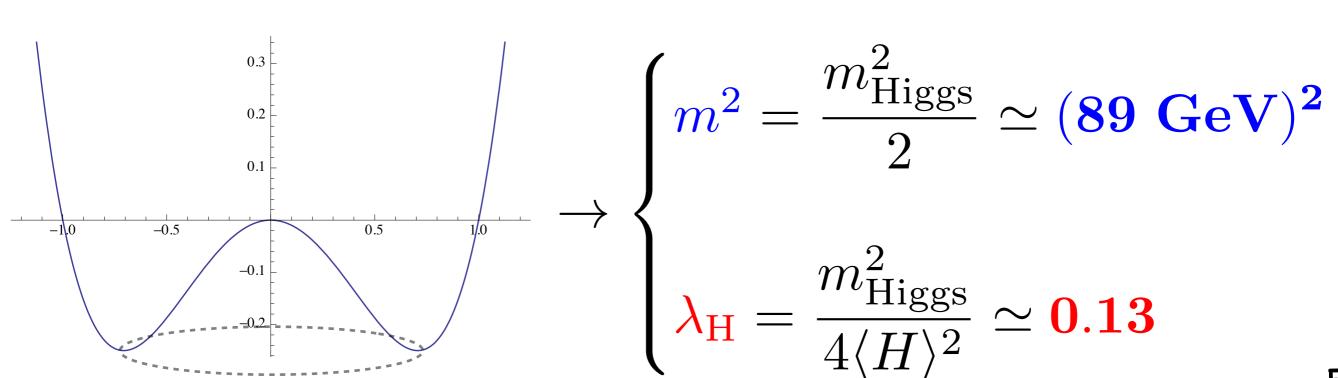
$$\rightarrow \begin{cases} \langle H \rangle^2 = \frac{m^2}{2 \; \lambda_{\rm H}} \quad \overset{\text{We knew-1}}{=} \frac{1}{2 \sqrt{2} \, G_F} \simeq (174 \; {\rm GeV})^2 \\ m_{\rm Higgs}^2 = 2 \; m^2 \quad \overset{\text{Now we also know}}{\simeq} \left(126 \; {\rm GeV} \right)^2 \end{cases}$$



$$V(H) = -m^2(H^{\dagger}H) + \lambda_{\mathbf{H}}(H^{\dagger}H)^2$$

$$\rightarrow \begin{cases} \langle H \rangle^2 = \frac{m^2}{2 \lambda_{\rm H}} & \stackrel{\text{We knew.}}{=} \frac{1}{2\sqrt{2} G_F} \simeq (174 \text{ GeV})^2 \\ m_{\rm Higgs}^2 = 2 m^2 & \stackrel{\text{Now we also know}}{\simeq} \left(126 \text{ GeV} \right)^2 \end{cases}$$

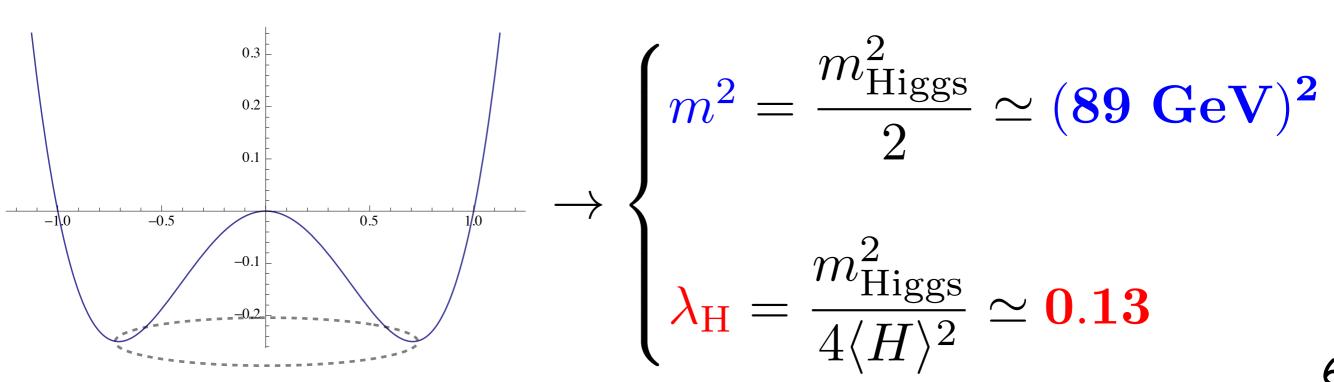
$$m_{\mathrm{Higgs}}^2 = 2 \, m^2 \simeq (\mathbf{126} \, \mathbf{GeV})^2$$



$$V(H) = -m^{2}(H^{\dagger}H) + \lambda_{H}(H^{\dagger}H)^{2}$$

$$(89 \text{ GeV})^{2} \qquad \mathbf{0.13}$$

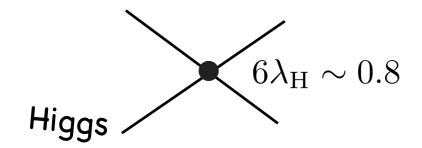
completely determined!



$$V(H) = -m^{2}(H^{\dagger}H) + \lambda_{H}(H^{\dagger}H)^{2}$$

$$(89 \text{ GeV})^{2} \qquad \mathbf{0.13}$$

It seems... Higgs sector is also described by weakly coupled, perturbative QFT. (at least no sign of strong interaction etc, so far...)



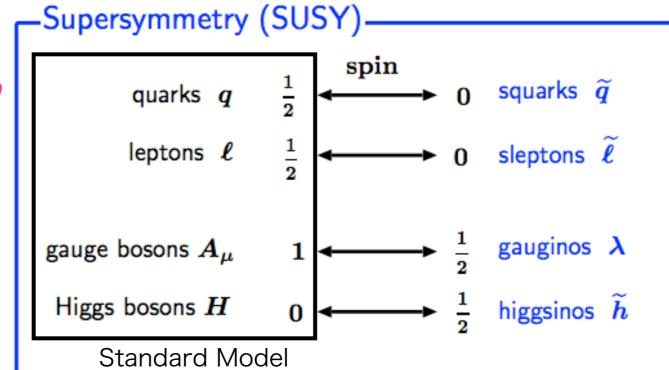
Imlications for BSM (in my opinion....)

This is compatible with....

- GUT and coupling unification in perturbative QFT. 👉 §2
- heavy right-handed neutrinos (Seesaw + Leptogenesis)
- # §2 and Lectures by Prof. Rubakov and Prof. Xing.
- Supersymmetry # §3

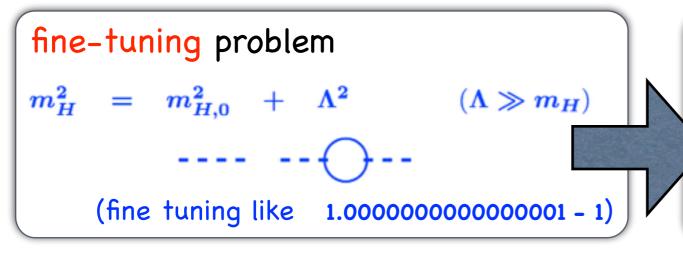
Supersymmetry

boson ⇔ fermion





naturalness

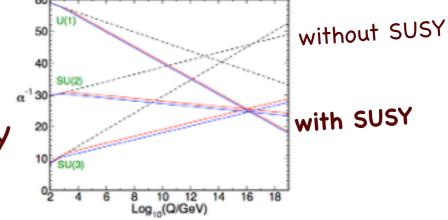


 $m_H^2 = m_{H,0}^2 + (\Lambda^2 - \Lambda^2)$ fermion boson



coupling unification

Grand Unified Theory





Dark Matter = Lightest SUSY particle

OK, then,....

What's the implications of 126 GeV Higgs for Supersymmetry (SUSY) ??

$$V(H) = -m^2(H^\dagger H) + \lambda_H(H^\dagger H)^2 \ (89~{
m GeV})^2 \qquad 0.13 \ = \lambda_{H_{\wedge}}^{
m tree} + \delta \lambda_H^{
m loop} \ rac{g^2 \cos^2 2 eta}{8 \cos^2 heta_W} \simeq {
m 0.069} \cos^2 2 eta \ {
m too~small...}$$

(known)

$$\begin{array}{c} V(H) = -m^2(H^\dagger H) + \lambda_{\rm H}(H^\dagger H)^2 \\ {}_{(89~{\rm GeV})^2} & {\bf 0.13} \\ {}_{\rm in~SUSY...} & = \lambda_{H_{\rm A}}^{\rm tree} + \delta \lambda_{H_{\rm A}}^{\rm loop} \end{array}$$

$$\left(\frac{g^2\cos^2 2\beta}{8\cos^2 \theta_W} \simeq \mathbf{0.069} \cos^2 2\beta\right)$$

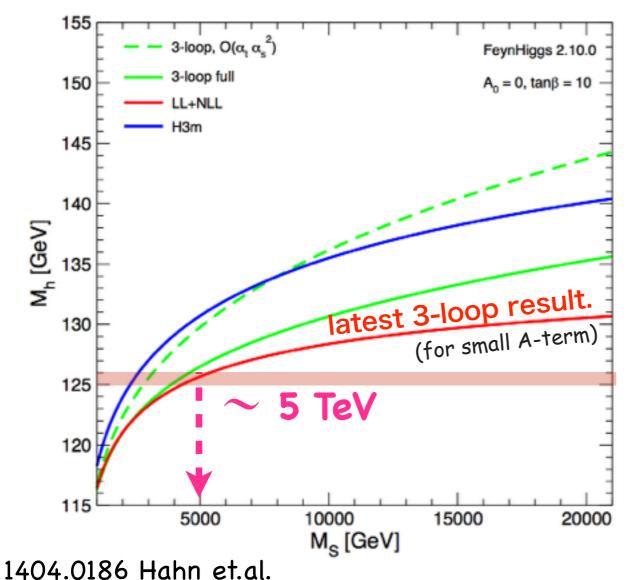
$$\frac{3y_t^4}{16\pi^2} \left(\log \left(\frac{m_{\text{stop}}^2}{m_t^2} \right) + \alpha^2 - \frac{\alpha^4}{12} \right) + \cdots$$
 for large tan β . $(\alpha \simeq A_t/m_{\text{stop}})$

...requires heavy stop and/or large A-term

$$V(H) = -m^{2}(H^{\dagger}H) + \lambda_{H}(H^{\dagger}H)^{2}$$

$$(89 \text{ GeV})^{2} \qquad 0.13$$

in SUSY...



$$=\lambda_{H_{\Lambda}}^{\text{tree}} + \delta\lambda_{H_{\Lambda}}^{\text{loop}}$$

$$\frac{g^2 \cos^2 2\beta}{8 \cos^2 \theta_W} \simeq \mathbf{0.069} \cos^2 2\beta$$

$$\frac{3y_t^4}{16\pi^2} \left(\log \left(\frac{m_{\text{stop}}^2}{m_t^2} \right) + \alpha^2 - \frac{\alpha^4}{12} \right) + \cdots$$
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$$V(H) = -m^2(H^\dagger H) + \lambda_H(H^\dagger H)^2$$
 on the other hand
$$\begin{array}{c} 89 \, \mathrm{GeV})^2 & \mathbf{0.13} \\ = \lambda_{H_A}^{\mathrm{tree}} + \delta \lambda_{H_A}^{\mathrm{loop}} \end{array}$$

$$-m^2 \simeq |\mu|^2 + m_{H_u}^{2\, ({
m tree})} + \delta m_{H_u}^{2\, ({
m loop})}$$
 up to $\mathcal{O}\left(rac{1}{ an^2eta}
ight)$ Higgsino mass

soft mass for up-type Higgs

$$\frac{g^2 \cos^2 2\beta}{8 \cos^2 \theta_W} \simeq \mathbf{0.069} \cos^2 2\beta$$

$$\frac{3y_t^4}{16\pi^2} \left(\log \left(\frac{m_{\text{stop}}^2}{m_t^2} \right) + \alpha^2 - \frac{\alpha^4}{12} \right) + \cdots$$
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$$V(H) = -m^2(H^\dagger H) + \lambda_H(H^\dagger H)^2$$
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$$\begin{array}{c} 89 \, \mathrm{GeV})^2 & \mathbf{0.13} \\ = \lambda_{H_A}^{\mathrm{tree}} + \delta \lambda_{H_A}^{\mathrm{loop}} \end{array}$$

$$-m^2 \simeq |\mu|^2 + m_{H_u}^{2 \text{ (tree)}} + \delta m_{H_u}^{2 \text{ (loop)}}$$

$$\text{up to } \mathcal{O}\left(\frac{1}{\tan^2 \beta}\right)$$

large μ ----> fine-tuning.

e.g.,
$$\simeq (1000~{\rm GeV})^2 - (1004~{\rm GeV})^2$$
 for $|\mu| \simeq 1~{\rm TeV}$

requires Light Higgsino to avoid a fine-tuning.

$$\frac{g^2 \cos^2 2\beta}{8 \cos^2 \theta_W} \simeq \mathbf{0.069} \cos^2 2\beta$$

$$\frac{3y_t^4}{16\pi^2} \left(\log \left(\frac{m_{\text{stop}}^2}{m_t^2} \right) + \alpha^2 - \frac{\alpha^4}{12} \right) + \cdots$$
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... requires heavy stop and/or large A-term

$$V(H) = -m^2(H^\dagger H) + \lambda_H(H^\dagger H)^2$$
 on the other hand $(89 \, \mathrm{GeV})^2$ 0.13 $= \lambda_{H_A}^{\mathrm{tree}} + \delta \lambda_{H_A}^{\mathrm{loop}}$

$$-m^2 \simeq |\mu|^2 + m_{H_u}^{2 \text{ (tree)}} + \delta m_{H_u}^{2 \text{ (loop)}}$$

Moreover,

$$\delta m_{H_u}^{2 \text{ (loop)}} \sim \frac{-3y_t^2}{8\pi^2} \left(m_{\widetilde{t_L}}^2 + m_{\widetilde{t_R}}^2 + |A_t|^2 \right) \log \left(\frac{M_{\text{mess}}}{m_{\sim}} \right) + \cdots$$

 $\frac{g^2 \cos^2 2\beta}{8 \cos^2 \theta_W} \simeq \mathbf{0.069} \cos^2 2\beta$

$$-|A_t|^2 \log\left(\frac{M_{ ext{mess}}}{m^2}\right) + \cdots$$
 $-\frac{3y_t^4}{r^2} \left(\log\left(\frac{m_{ ext{stop}}^2}{m_t^2}\right) + \alpha^2 - \frac{\alpha^4}{12}\right) + \cdots$ for large $\tan \beta$. $(\alpha \simeq A_t/m_{ ext{stop}})$

for large tan β . $(\alpha \simeq A_t/m_{\rm stop})$

requires Light stop and

small A-term

to avoid a fine-tuning.

... requires heavy stop and/or large A-term

Fine-tuning worse than 1% seems unavoidable in MSSM.

(MSSM = Minimal SUSY Standard Model)

What does it imply ??

- 1. No SUSY?
- 2. (It's anyway fine-tuned, then....)

Very heavy SUSY? (10-100 TeV, or even higher...)

- **3.** (still.....)
 - O(0.1-1) TeV SUSY? (fine-tuned, but less than 2 and 3...)

Plan

- O. Introduction (... done)
- 1. Higgs
 - 2. Beyond the Standard Model
 - 3. SUSY
 - 4. SUSY after Higgs discovery

Please interrupt and ask questions at any time!

Any questions so far?

he last day. again

Now we know the Higgs properties very well.

For a review, see

Review of Particle Physics

(K.A. Olive et al. (PDG), Chin. Phys. C38, 090001 (2014), http://pdg.lbl.gov)

"11. Status of Higgs Boson Physics,"

by M. Carena, C. Grojean, M. Kado, and V. Sharma

Now we know the Higgs properties very well.

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eview of fairfield frigulation		
11. STA	11. Status of Higgs boson physics 1 FUS OF HIGGS BOSON PHYSICS)
Written November : the University of C	III. The discovery of a Higgs boson	
Barcelona), M. Kad	III.1. The discovery channels	
V. Sharma (Universi	III. IV. Properties and nature of the new bosonic resonance	
. Introduction .	III.2 IV.1. Theoretical framework	
II. The Standard Mc	V. New physics models of EWSB in the light of the Higgs boson dis	covery
II.1. The SM Higgs	V.1. Higgs bosons in the Minimal Supersymmetric Standard Mode	l (MSS
II.2. The SM custo	III. V.1.1. The MSSM Higgs boson masses	`
II.3. Stability of the	V.1.2. MSSM Higgs boson couplings	
	III.5	

Before the Higgs discovery, we already knew its couplings very well (assuming that there is no deviation from the SM). Only the mass was unknown.

(mass of a SM particle \mathbf{X}) = (its coupling to Higgs) × (Higgs VEV)

$$\implies (\mathbf{X}\text{'s coupling with Higgs}) = \frac{(\mathbf{X}\text{'s mass})}{(\mathrm{Higgs VEV})} \overset{\text{known}}{\longleftarrow}$$

Before the Higgs discovery, we already knew its couplings very well (**assuming** that there is no deviation from the SM). Only the mass was unknown.

(from RPP, "11. Status of Higgs Boson Physics.")

the square of the boson masses. The SM Higgs boson couplings to gauge bosons, Higgs bosons and fermions are summarized in the following Lagrangian:

$$\mathcal{L} = -g_{Hff}\bar{f}fH + \frac{g_{HHH}}{6}H^3 + \frac{g_{HHHH}}{24}H^4 + \delta_V V_\mu V^\mu \left(g_{HVV}H + \frac{g_{HHVV}}{2}H^2\right)$$
(11.9)

with

known

$$g_{Hf\bar{f}} = \frac{m_f}{v}, \quad g_{HVV} = \frac{2m_V^2}{v}, \quad g_{HHVV} = \frac{2m_V^2}{v^2}$$
 (11.10)

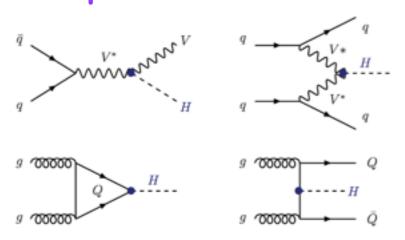
$$g_{HHH} = \frac{3m_H^2}{v}, \quad g_{HHHH} = \frac{3m_H^2}{v^2}$$
 (11.11)

where $V = W^{\pm}$ or Z and $\delta_W = 1, \delta_Z = 1/2$. As a result, the dominant mechanisms for

Before the Higgs discovery, we already knew its couplings very well (assuming that there is no deviation from the SM). Only the mass was unknown.

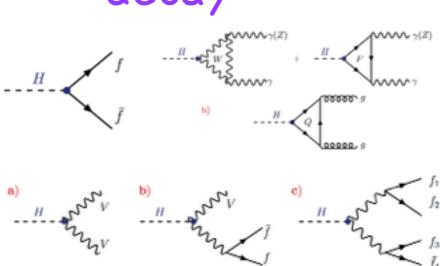
Thus, both the **production cross section** and the **decay rates** of the Higgs were also calculated (as a function of Higgs mass) very precisely.

production



+ higher order diagrams...

decay



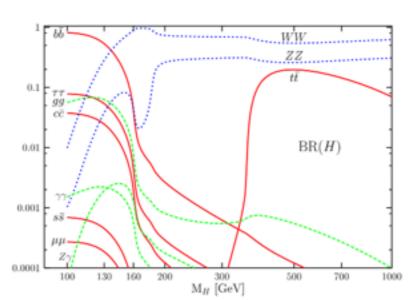


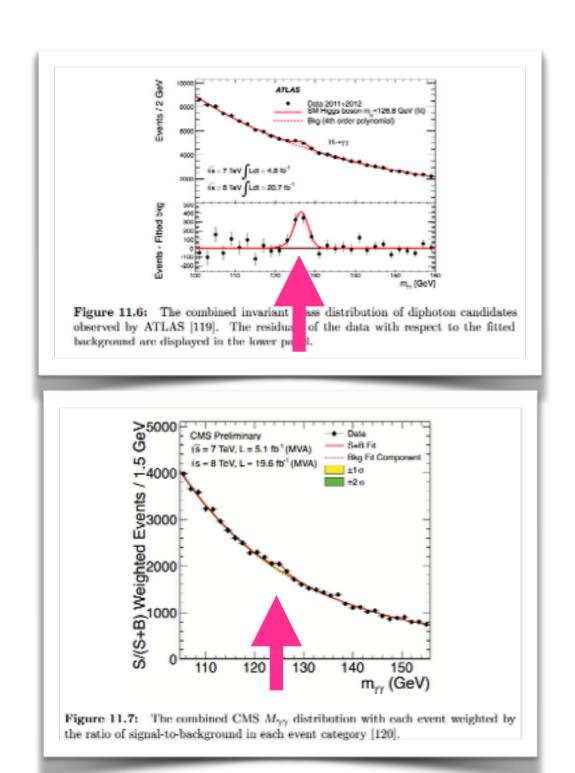
Figure 2.25: The SM Higgs boson decay branching ratios as a function of M_H

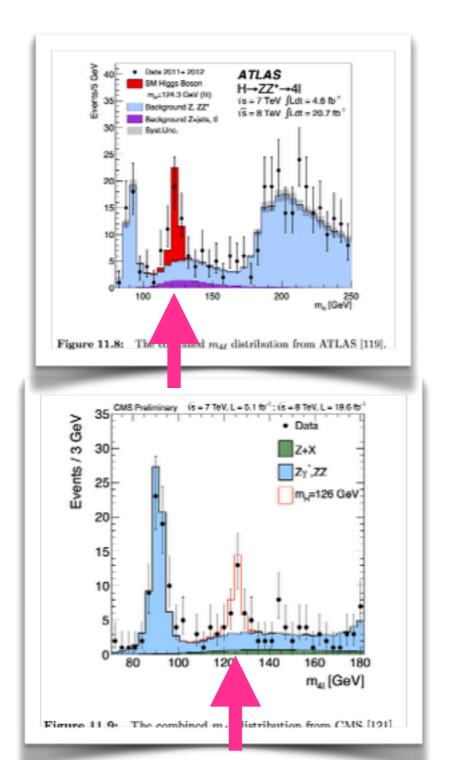
(figures are A.Djouadi, hep-ph/0503172)

+ higher order diagrams...

(figures are from RPP, "11. Status of Higgs Boson Physics.")

Now the Higgs mass is also well measured.





Now the Higgs mass is also well measured.

(from RPP, Particle Listings, "Gauge and Higgs Bosons")

(from RPP,

"11. Status of Higgs Boson Physics.")

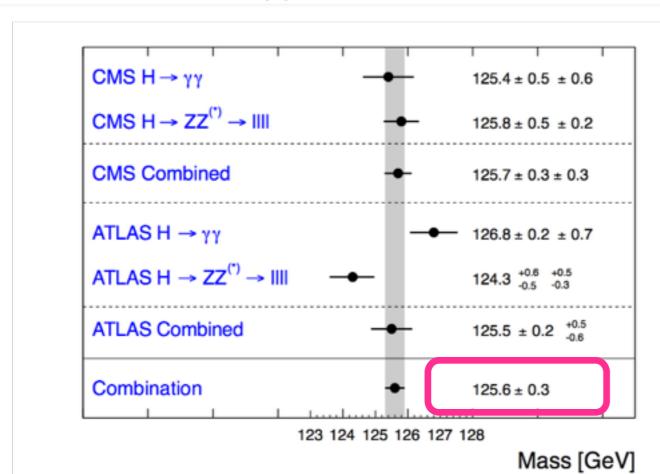


Figure 11.10: A compilation of the CMS and ATLAS mass measurements in the $\gamma\gamma$ and ZZ channels, the combined result from each experiment and our average of the combinations.

Citation: K.A. Olive et al. (Particle Data Group), Chin. Phys. C38, 090001 (2014) (URL: http://pdg.lbl.gov)



$$J = 0$$

In the following H^0 refers to the signal that has been discovered in the Higgs searches. Whereas the observed signal is labeled as a spin 0 particle and is called a Higgs Boson, the detailed properties of H^0 and its role in the context of electroweak symmetry breaking need to be further clarified. These issues are addressed by the measurements listed below.

Concerning mass limits and cross section limits that have been obtained in the searches for neutral and charged Higgs bosons, see the sections "Searches for Neutral Higgs Bosons" and "Searches for Charged Higgs Bosons (H^{\pm} and $H^{\pm\pm}$)", respectively.

HO MASS

A combination of the results from ATLAS and CMS, where a recent unpublished result from CMS is used, yields an average value of 125.6 ± 0.3 GeV, see the review on "Status of Higgs Boson Physics."

125.7±0.4 OUR AVERAGE

125.5±0.2 $^{+0.5}_{-0.6}$ 1,2 AAD 13AK ATLS pp, 7 and 8 TeV 125.8±0.4±0.4 1,3 CHATRCHYAN 13J CMS pp, 7 and 8 TeV • • • We do not use the following data for averages, fits, limits, etc. • • • 126.8±0.2±0.7 2 AAD 13AK ATLS pp, 7 and 8 TeV, $\gamma\gamma$ 124.3 $^{+0.6}_{-0.5}$ 2 AAD 13AK ATLS pp, 7 and 8 TeV, $ZZ^* \rightarrow 4\ell$ 126.2±0.6±0.2 3 CHATRCHYAN 13J CMS pp, 7, 8 TeV, $ZZ^* \rightarrow 4\ell$ 126.0±0.4±0.4 1,4 AAD 12AI ATLS pp, 7 and 8 TeV 125.3±0.4±0.5 1,5 CHATRCHYAN 12N CMS pp, 7 and 8 TeV

- ¹Combined value from $\gamma\gamma$ and $ZZ^* \rightarrow 4\ell$ final states.
- 2 AAD 13AK use 4.7 fb $^{-1}$ of pp collisions at $E_{\rm cm}$ =7 TeV and 20.7 fb $^{-1}$ at $E_{\rm cm}$ =8 TeV.
- 3 CHATRCHYAN 13J use 5.1 fb $^{-1}$ of pp collisions at $E_{\rm cm}=7$ TeV and 12.2 fb $^{-1}$ at $E_{\rm cm}=8$ TeV.
- ⁴ AAD 12N obtain results based on 4.6–4.8 fb⁻¹ of pp collisions at $E_{\rm cm}=7$ TeV and 5.8–5.9 fb⁻¹ at $E_{\rm cm}=8$ TeV. An excess of events over background with a local significance of 5.9 σ is observed at $m_{H^0}=126$ GeV. See also AAD 12DA.
- ⁵ CHATRCHYAN 12N obtain results based on 4.9–5.1 fb⁻¹ of pp collisions at $E_{\rm cm}=7$ TeV and 5.1–5.3 fb⁻¹ at $E_{\rm cm}=8$ TeV. An excess of events over background with a local significance of 5.0 σ is observed at about $m_{H^0}=125$ GeV. See also CHATRCHYAN 12BY and CHATRCHYAN 13Y.

24

Now the Higgs mass is also well measured.

==> No free parameter within the SM.

==> One can compare the predictions and measurements to see if there is any deviation.

Now the Higgs mass is also well measured.

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(from RPP, "11. Status of Higgs Boson Physics.")

$$\mu = \frac{(\sigma \cdot \mathcal{B})_{\text{observed}}}{(\sigma \cdot \mathcal{B})_{\text{SM prediction}}}$$

 σ : Higgs production cross section

 \mathcal{B} : Higgs decay branching ratio

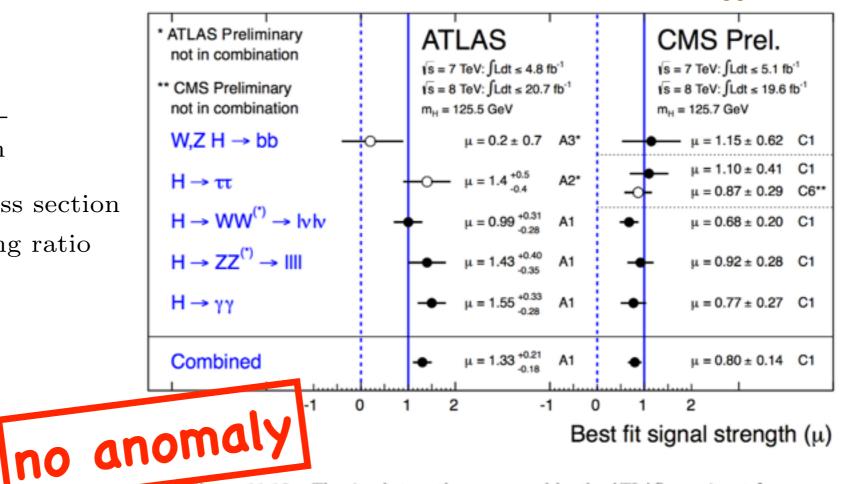


Figure 11.16: The signal strengths μ measured by the ATLAS experiment from Refs. A1 [119], A2 [133] and A3 [138], and CMS experiment from Ref. C1 [124] and C6 [132] in the five principal channels and their combination. It should be noted that the ATLAS combination only includes the bosonic $\gamma\gamma$, ZZ and WW channels.

Now the Higgs mass is also well measured.

==> No free parameter within the SM.

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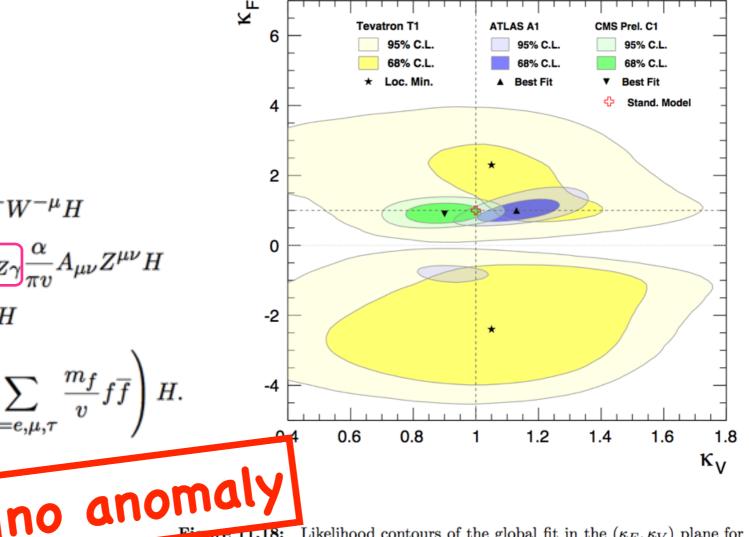
Using effective Lagrangian...

$$\mathcal{L} = \kappa_{3} \frac{m_{H}^{2}}{2v} H^{3} + \kappa_{Z} \frac{m_{Z}^{2}}{v} Z_{\mu} Z^{\mu} H + \kappa_{W} \frac{2m_{W}^{2}}{v} W_{\mu}^{+} W^{-\mu} H$$

$$+ \kappa_{g} \frac{\alpha_{s}}{12\pi v} G_{\mu\nu}^{a} G^{a\mu\nu} H + \kappa_{\gamma} \frac{\alpha}{2\pi v} A_{\mu\nu} A^{\mu\nu} H + \kappa_{Z\gamma} \frac{\alpha}{\pi v} A_{\mu\nu} Z^{\mu\nu} H$$

$$+ \kappa_{VV} \frac{\alpha}{2\pi v} \left(\cos^{2}\theta_{W} Z_{\mu\nu} Z^{\mu\nu} + 2W_{\mu\nu}^{+} W^{-\mu\nu} \right) H$$

$$- \left(\kappa_{t} \sum_{f=u,c,t} \frac{m_{f}}{v} f \overline{f} + \kappa_{b} \sum_{f=d,s,b} \frac{m_{f}}{v} f \overline{f} + \kappa_{\tau} \sum_{f=e,\mu,\tau} \frac{m_{f}}{v} f \overline{f} \right) H.$$



 $\kappa_X = 1 \text{ in SM }_{(\kappa_{Z\gamma} = 0)}$

 $\implies \kappa_X \neq 0$ (modified "by hand")

Figure 11.18: Likelihood contours of the global fit in the (κ_F, κ_V) plane for the ATLAS A1 [119], the CMS C1 [120] and the D0 and CDF combined T1 [108] results.

Now the Higgs mass is also well measured.

==> No free parameter within the SM.

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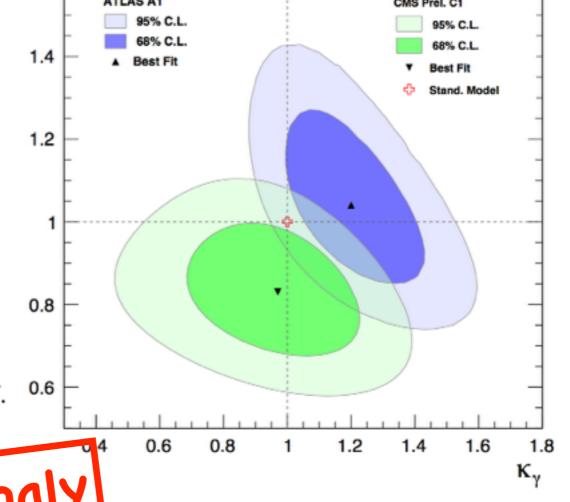
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$$+ \kappa_{VV} \frac{\alpha}{2\pi v} \left(\cos^{2}\theta_{W} Z_{\mu\nu} Z^{\mu\nu} + 2W_{\mu\nu}^{+} W^{-\mu\nu} \right) H$$

$$- \left(\kappa_{t} \sum_{f=u,c,t} \frac{m_{f}}{v} f \overline{f} + \kappa_{b} \sum_{f=d,s,b} \frac{m_{f}}{v} f \overline{f} + \kappa_{\tau} \sum_{f=e,\mu,\tau} \frac{m_{f}}{v} f \overline{f} \right) H. \quad 0.6$$



 $\kappa_X = 1 \text{ in SM }_{(\kappa_{Z\gamma} = 0)}$

 $\implies \kappa_X \neq 0$ (modified "by hand")

Figure 11.19: Likelihood contours of the global fit in the $(\kappa_g, \kappa_{\gamma})$ plane for the ATLAS experiment A1 [119] and the CMS experiment C1 [120] results.

SUMMARY

Higgs was discovered, and everything seems consistent with the SM.

..... any hint of physics beyond the SM ??



Plan

- O. Introduction
- 1. Higgs done
- 2. Beyond the Standard Model
- 3. SUSY
- 4. SUSY after Higgs discovery

Any questions so far?

Plan

- O. Introduction
- 1. Higgs
- 2. Beyond the Standard Model
 - 2.1. puzzles in SM = hints of BSM.
 - 2.2. renormalization and naturalness

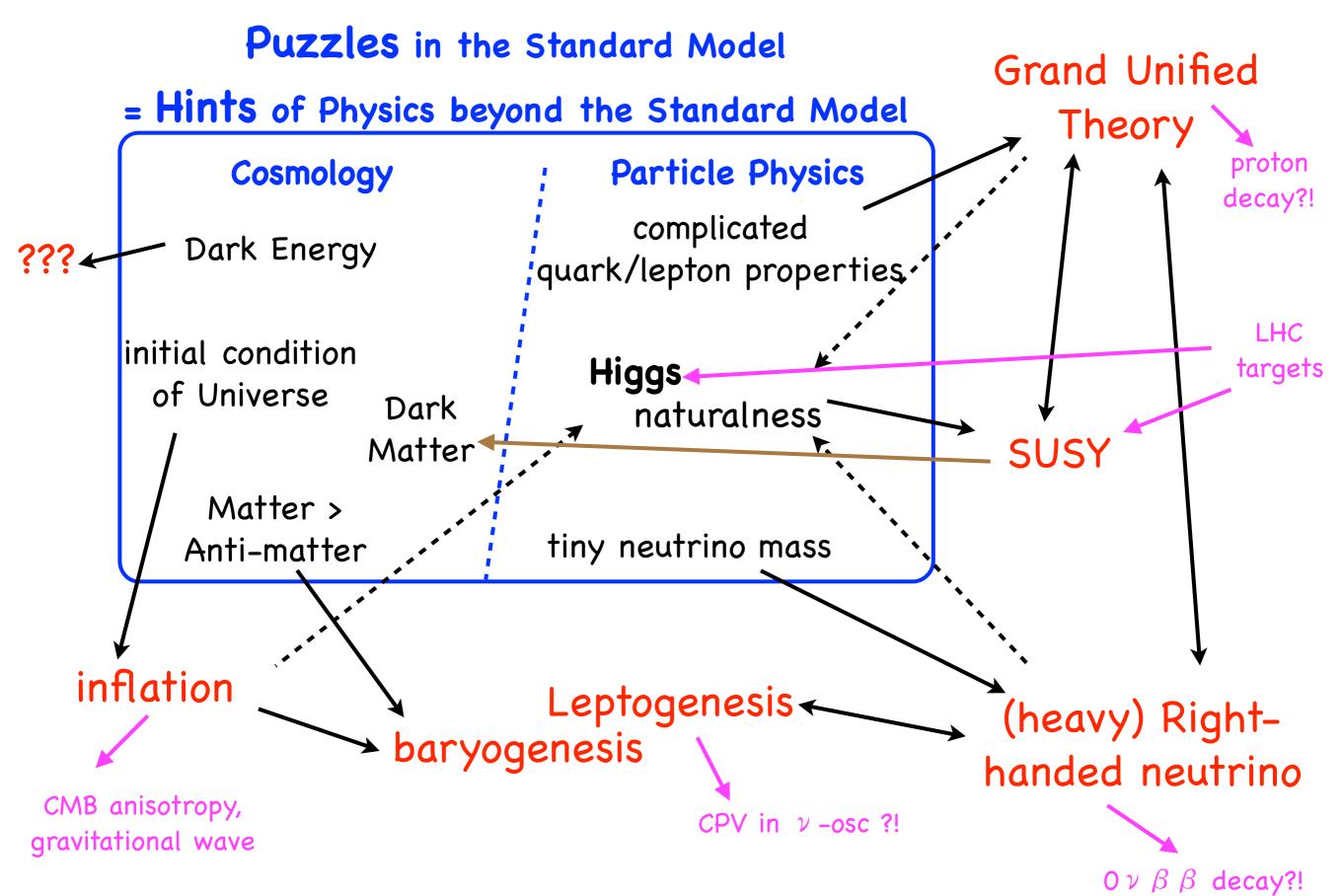
The purpose of this section is:

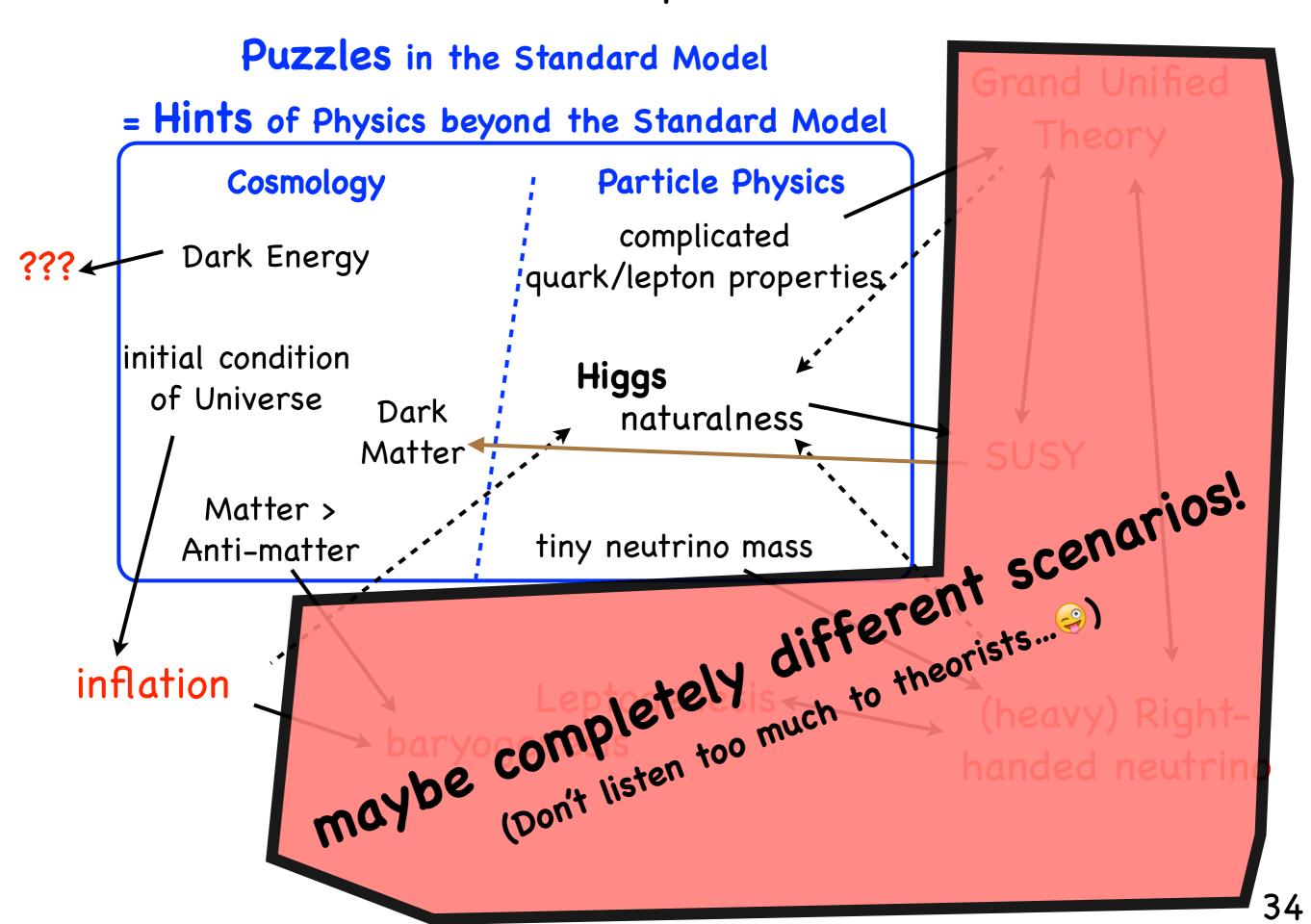
to introduce some of

Motivations for BSM physics and

Candidates for BSM physics.

(NOTE: biased!)





comment:

perturbative 4d QFT up to high scale?

YES

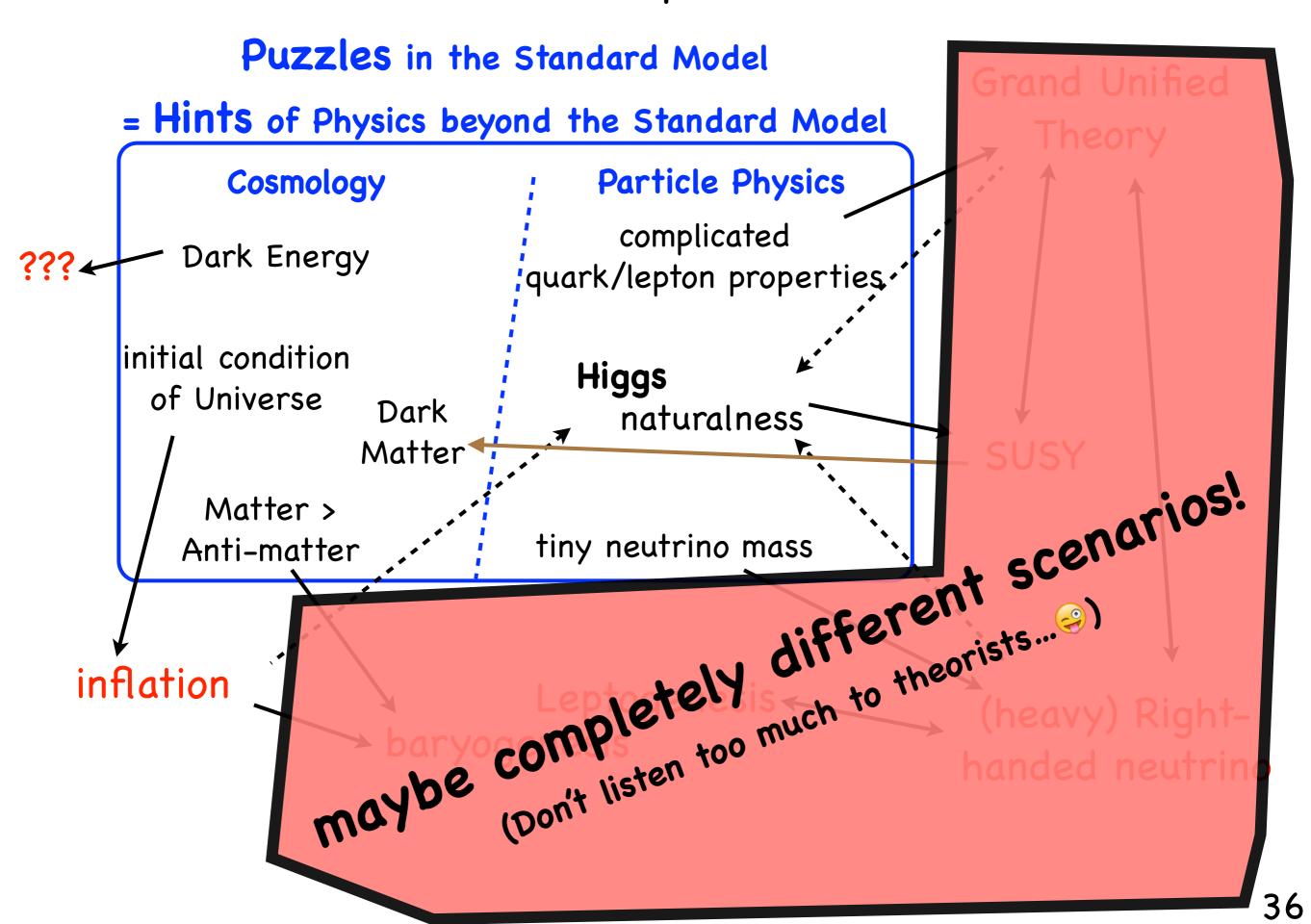
- compatible with
- perturbative GUT,
- seesaw and leptogenesis with heavy right-handed neutrino.
 - cosmology can be discussed within perturbative QFT. (inflation, baryogenesis,...)

NO

example

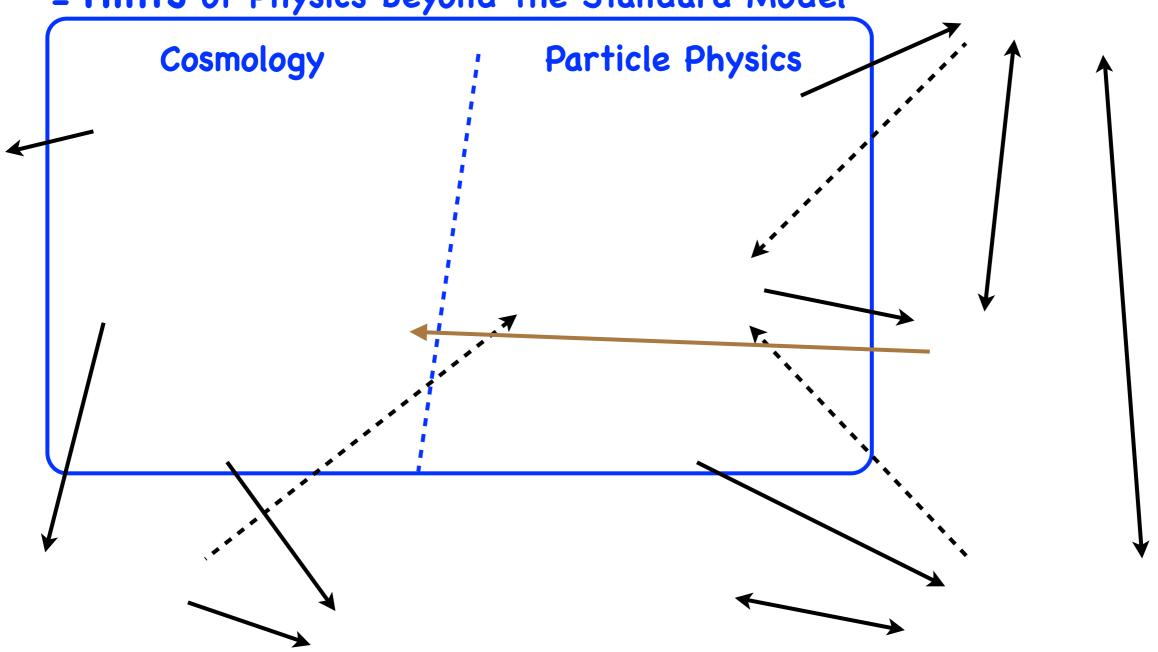
- extra-dimension (large, warped,...)
- composite Higgs
-
- => maybe we can discuss in discussion session.....

this lecture



Puzzles in the Standard Model

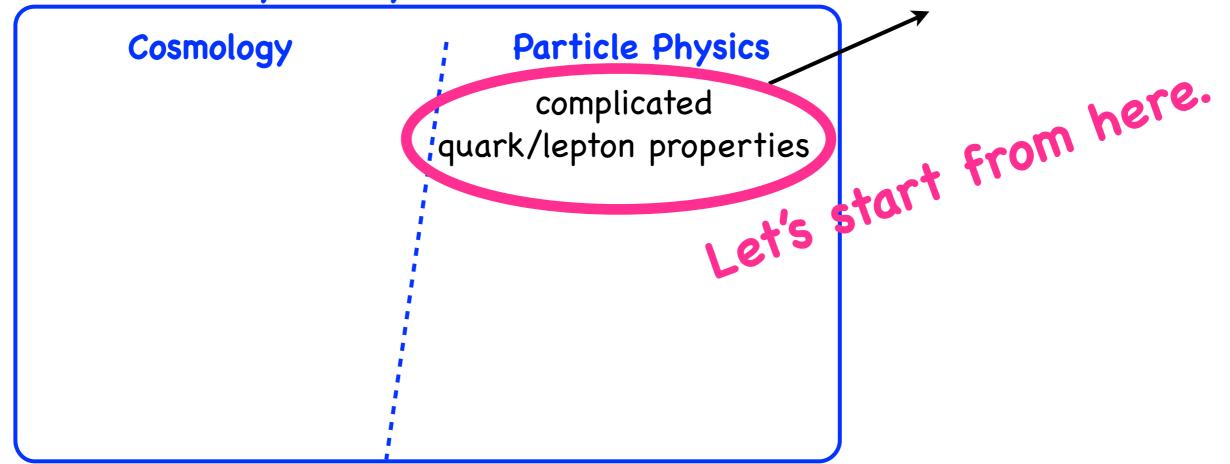
= Hints of Physics beyond the Standard Model



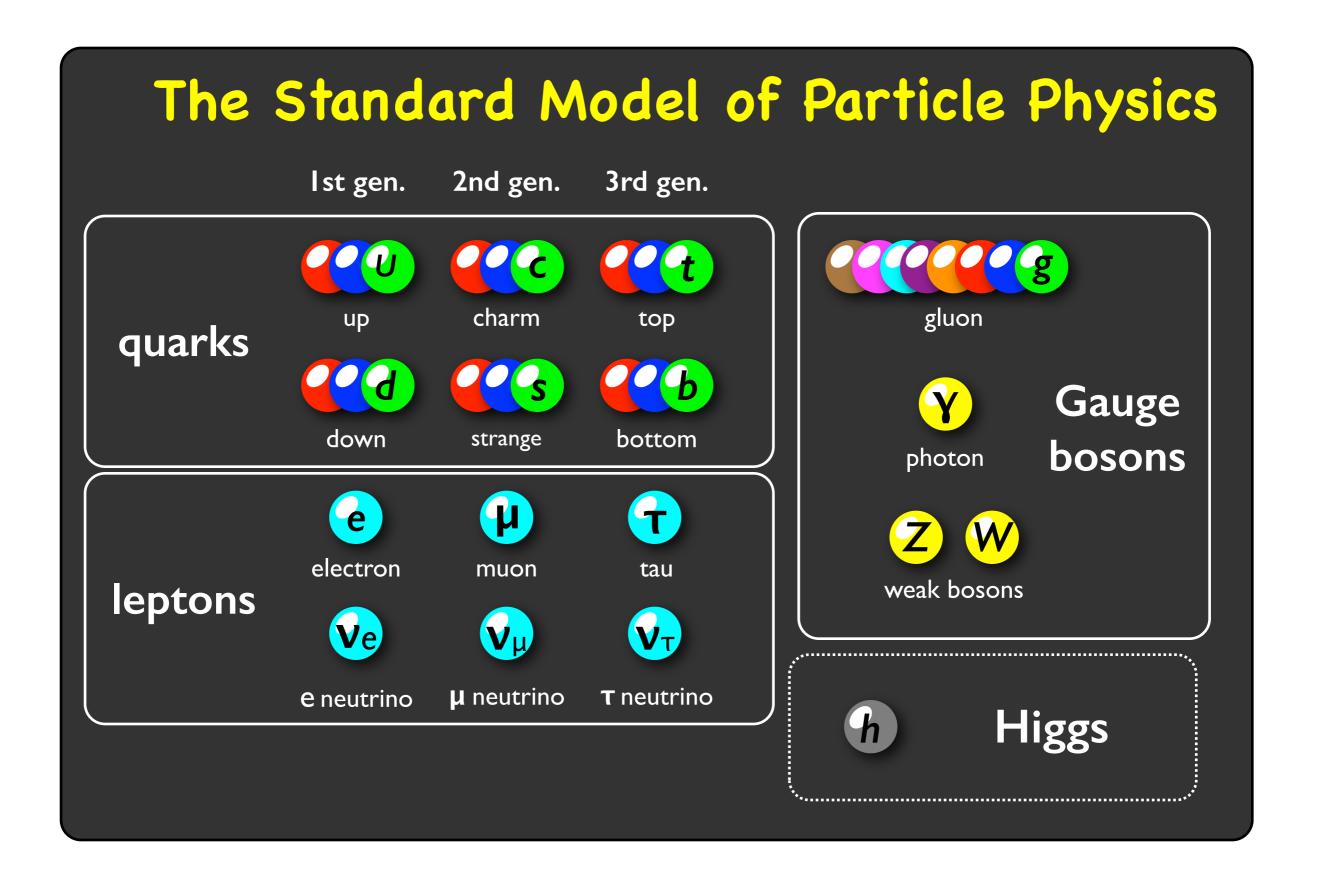
Where shall we start...?

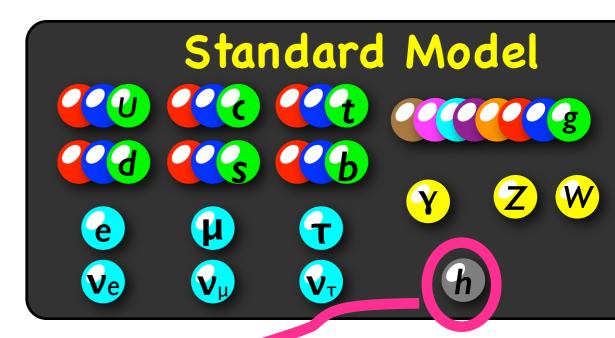
Puzzles in the Standard Model

= Hints of Physics beyond the Standard Model



Where shall we start ...?

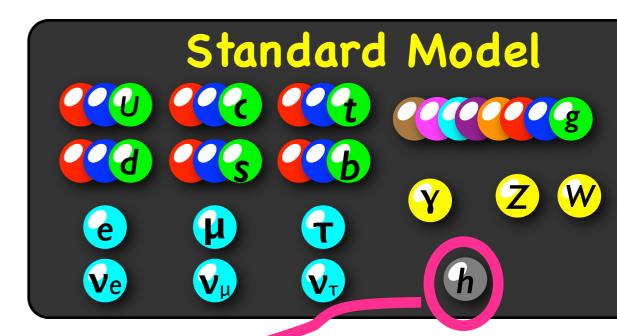




First of all, Let's recall how Higgs gives masses to fermions.

cf. Lecture by Prof. Godbole.

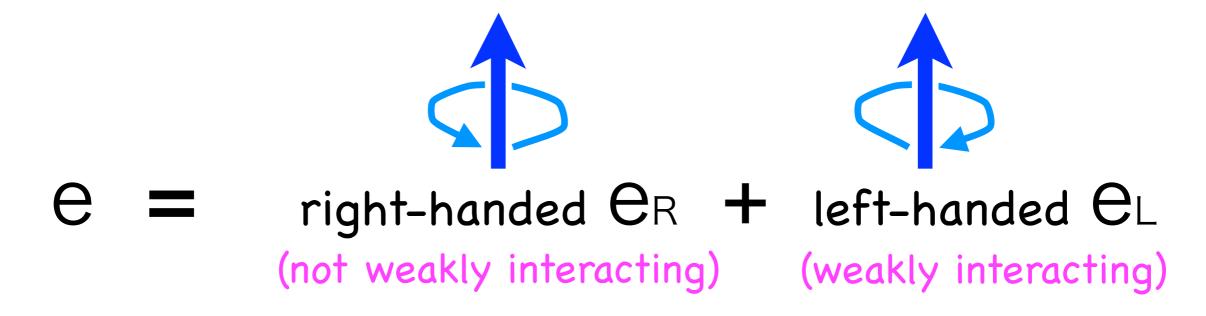
e.g. electron:

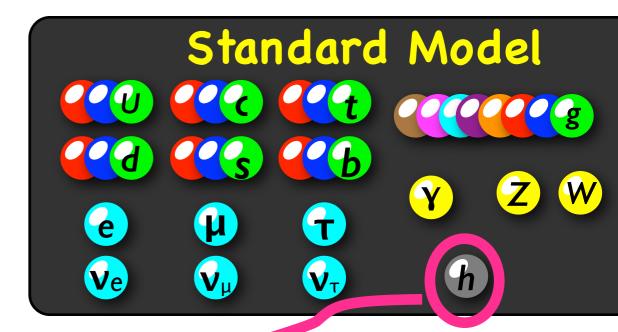


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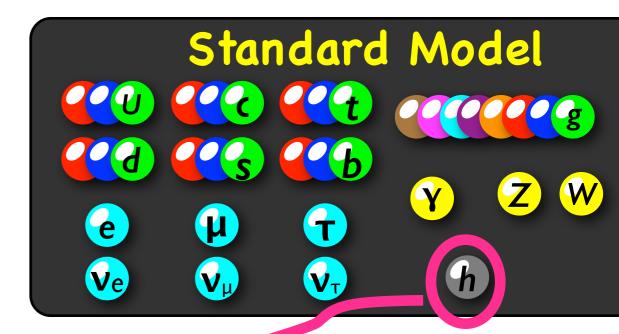
cf. Lecture by Prof. Godbole.

e.g. electron: without Higgs,.....

(not weakly interacting)
right-handed CR
left-handed CL
(weakly interacting)

different particles !! zero masses !

(moving with a speed of light)

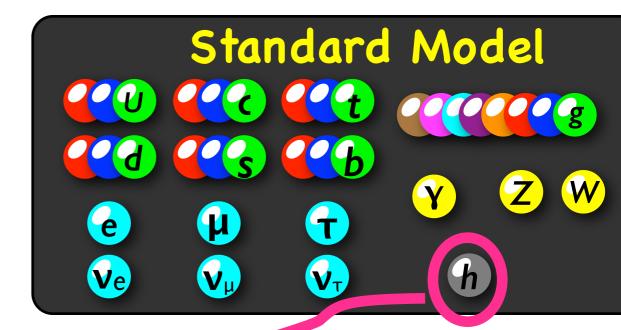


First of all, Let's recall how Higgs gives masses to fermions.

cf. Lecture by Prof. Godbole.

e.g. electron: the Higgs connects the two components.

(not weakly interacting)
right-handed Θ_R left-handed Θ_L Yukawa
interaction



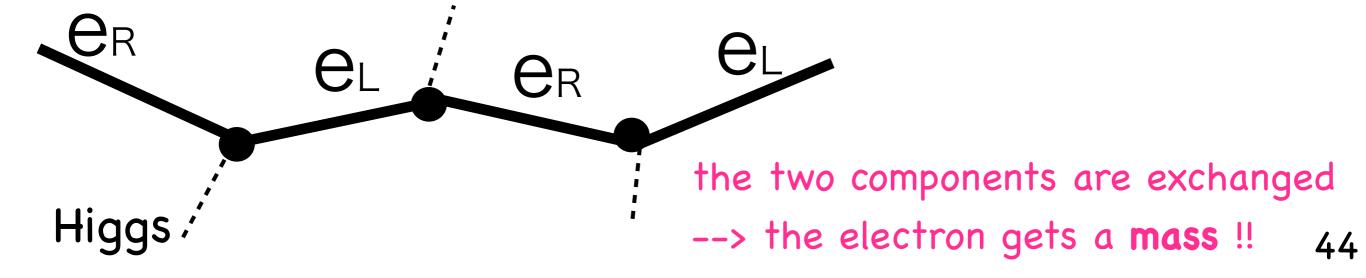
First of all,

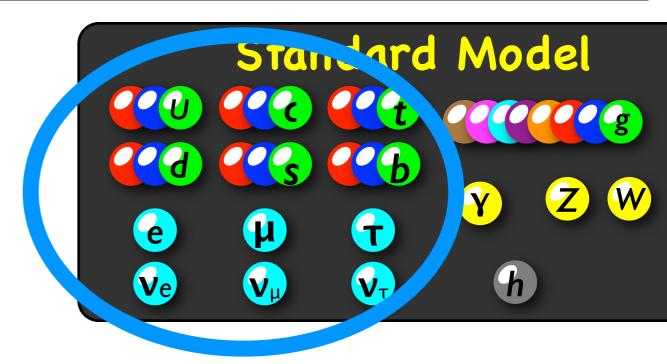
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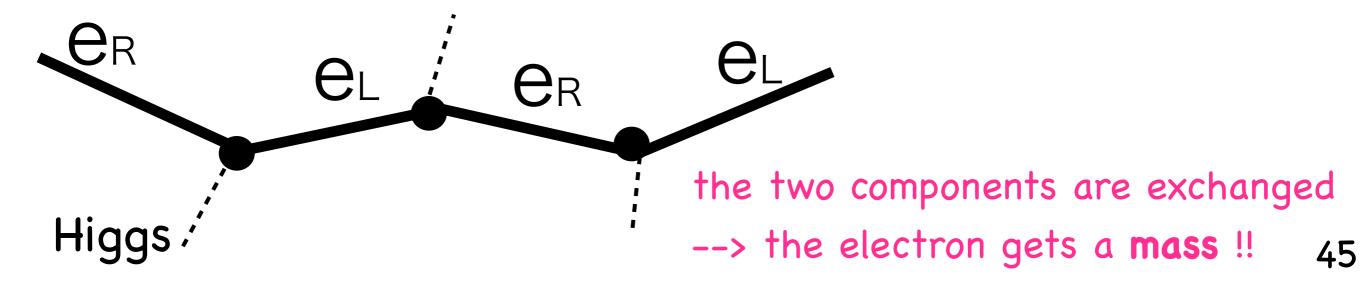
e.g. electron:

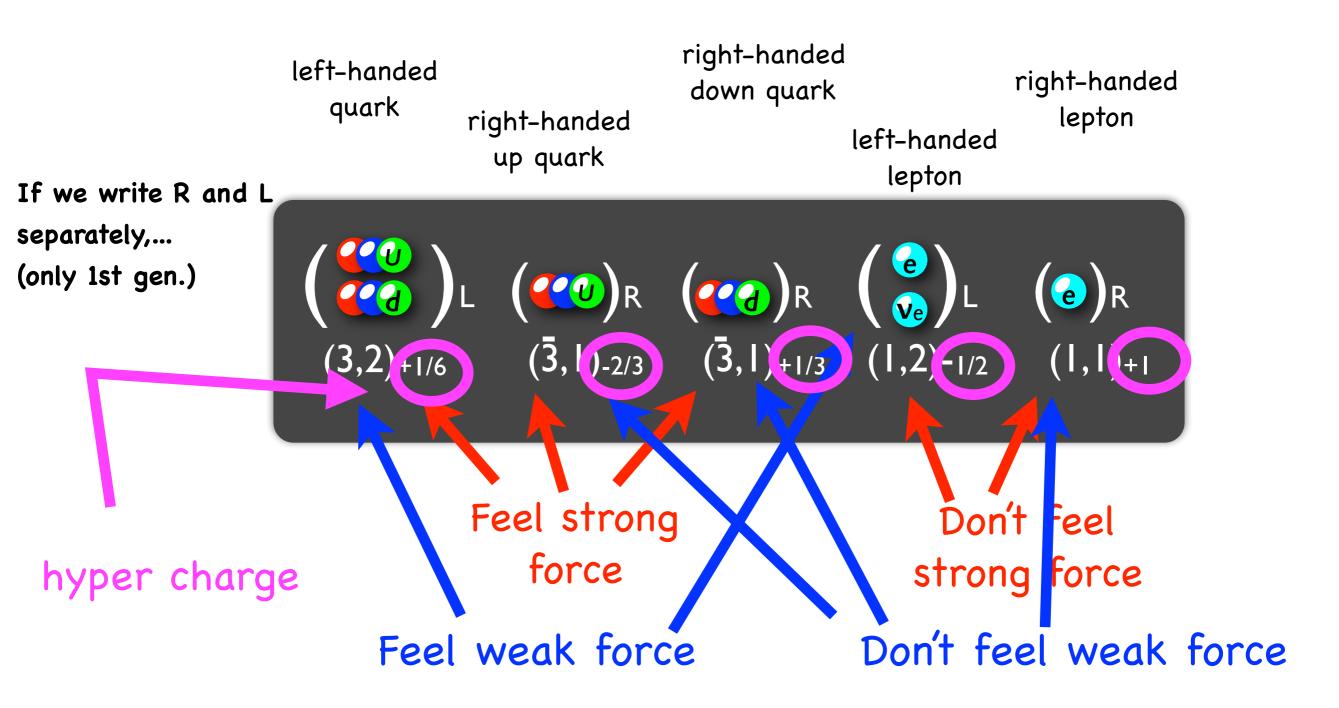
thanks to the "Higgs condensation" in vacuum....





Other quarks and leptons get masses in similar ways. "Left + Right"

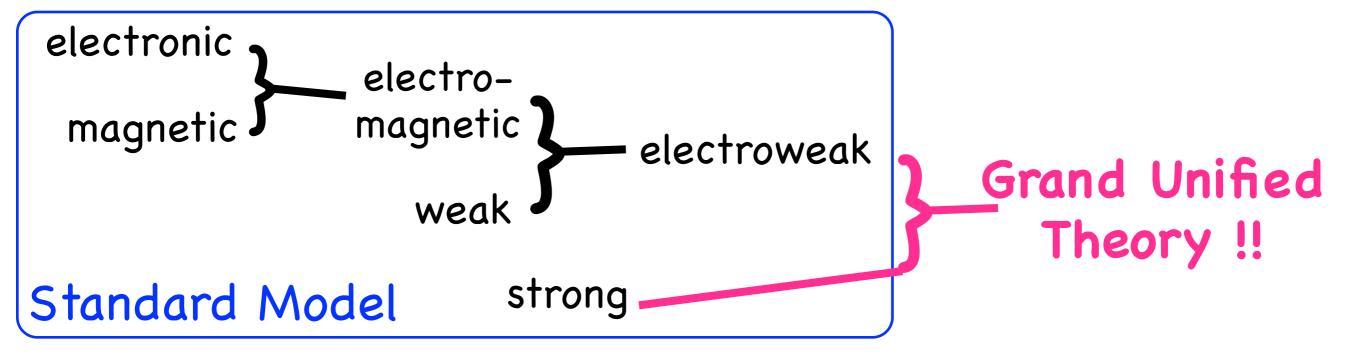




... very complicated !!

Q: any simple, unified theory to explain it?

Q: any simple, unified theory to explain it?



Standard Model

= hir

L

R

2.1. put

= hir

Ve

L

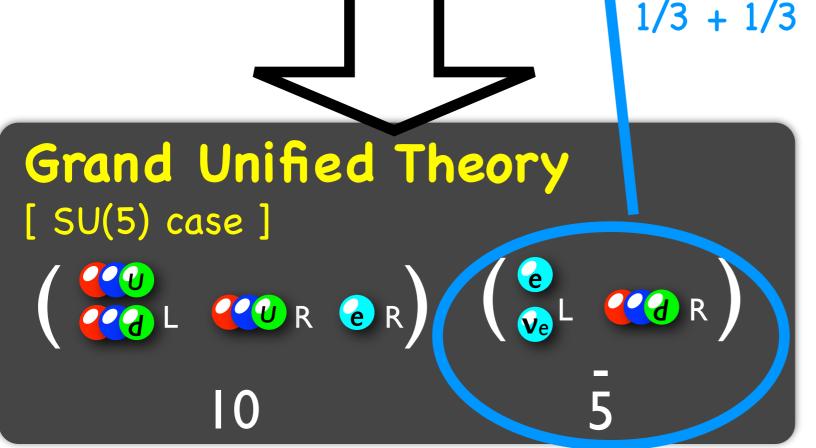
P

R

 $(\overline{3},1)_{+1/3}$ $(1,2)_{-1/2}$

2.1. puzzles in SM= hints of BSM.

Q: any simple, unified theory to explain it?



 $(\bar{3},1)_{-2/3}$

(3,2)+1/6

1/3 + 1/3 + 1/3 - 1/2 - 1/2 = 0

Complicated numbers are naturally explained!

... beautifully unified into simple SU(5) representations!

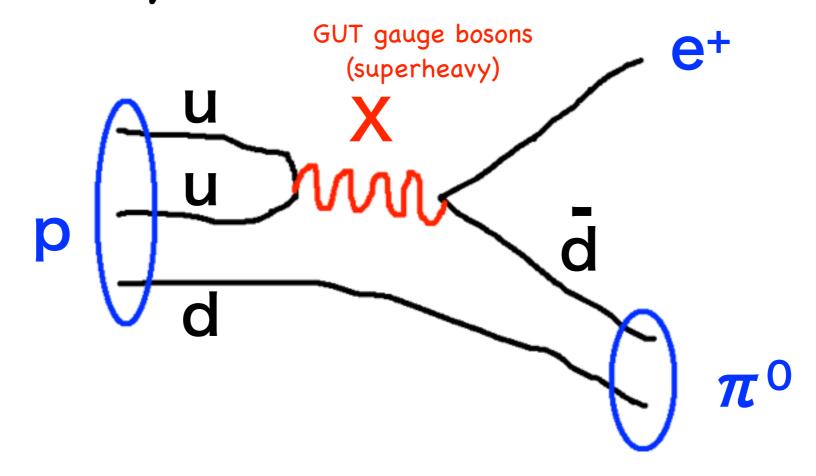
[Georgi, Glashow 1974]

Puzzles in the Standard Model Grand Unified = Hints of Physics beyond the Standard Model Theory Particle Physics Cosmology complicated quark/lepton properties

Any prediction by Grand Unified Theorie (GUT)?

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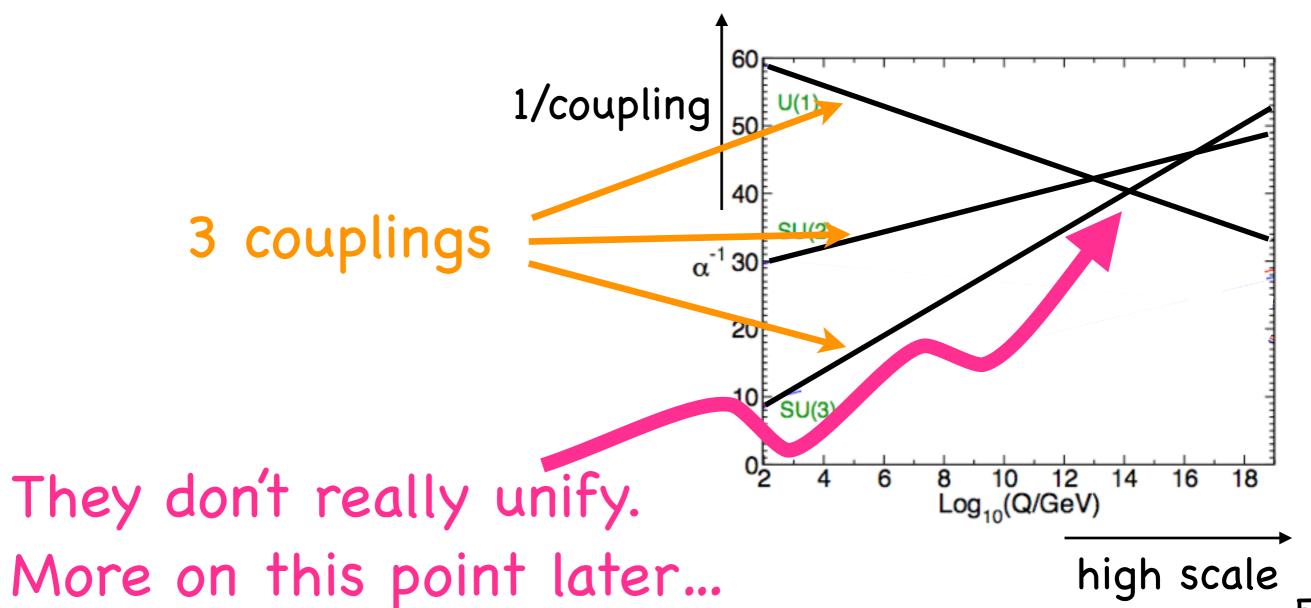
prediction 1. If GUT is correct,..... protons decay!



Any prediction by Grand Unified Theorie (GUT)?

prediction 2. If GUT is correct,.....

the three gauge couplings unify at high scale.



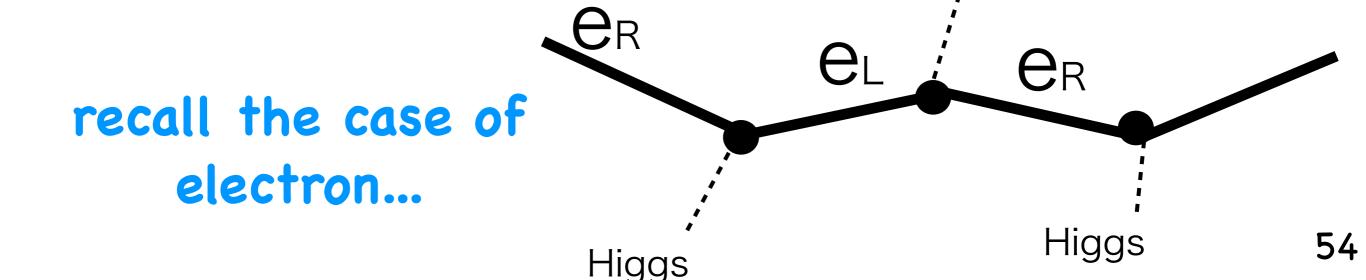
Puzzles in the Standard Model Grand Unified = Hints of Physics beyond the Standard Model Theory proton Particle Physics Cosmology decay?! complicated quark/lepton properties tiny neutrino mass

Next,...

puzzle: neutrino masses of cf. Lecture by Prof. Xing.

```
left-handed quark right-handed up quark right-handed lepton right-handed lepton
```

The neutrino has only left-handed component.



puzzle: neutrino masses of cf. Lecture by Prof. Xing.

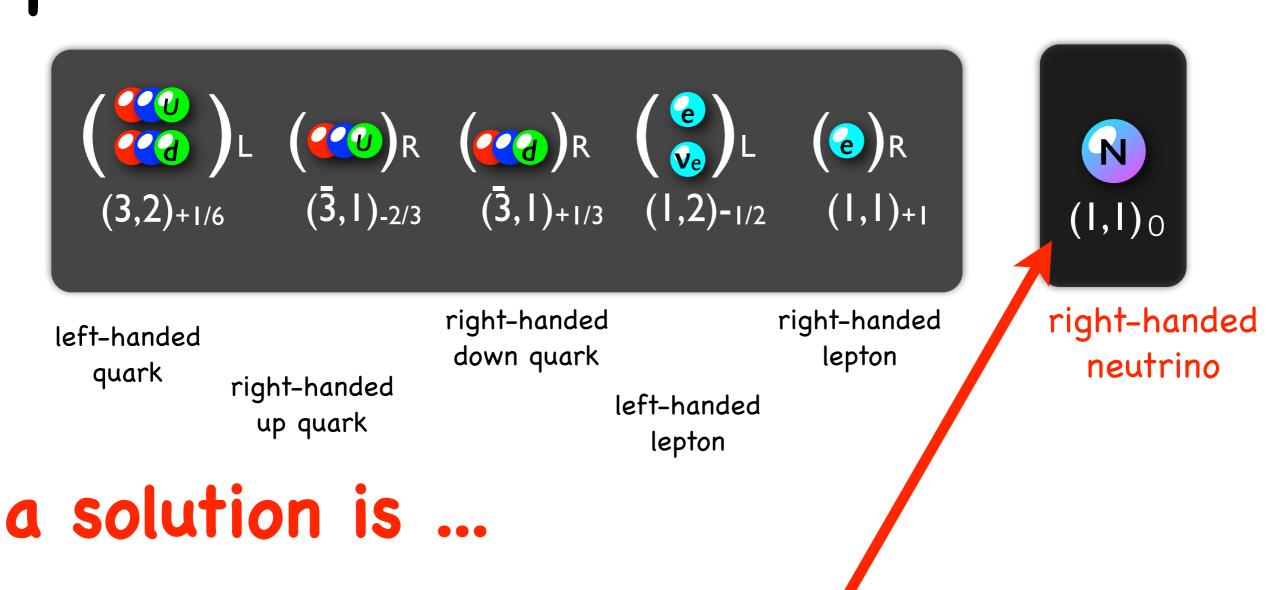
```
left-handed quark right-handed up quark left-handed lepton
```

The neutrino has only left-handed component.

===> Massless !!

But the neutrino masses are confirmed by neutrino oscillations!

puzzle: neutrino masses of cf. Lecture by Prof. Xing.



To add a right-handed neutrino!!



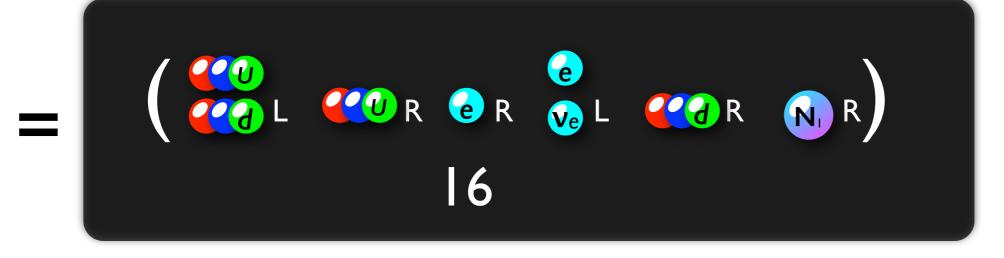
1 quarks and leptons completely unified

```
SU(5)
GUT

10

SU(5)
GUT

(1,1)0
```



SO(10) GUT

All quarks and leptons unified!

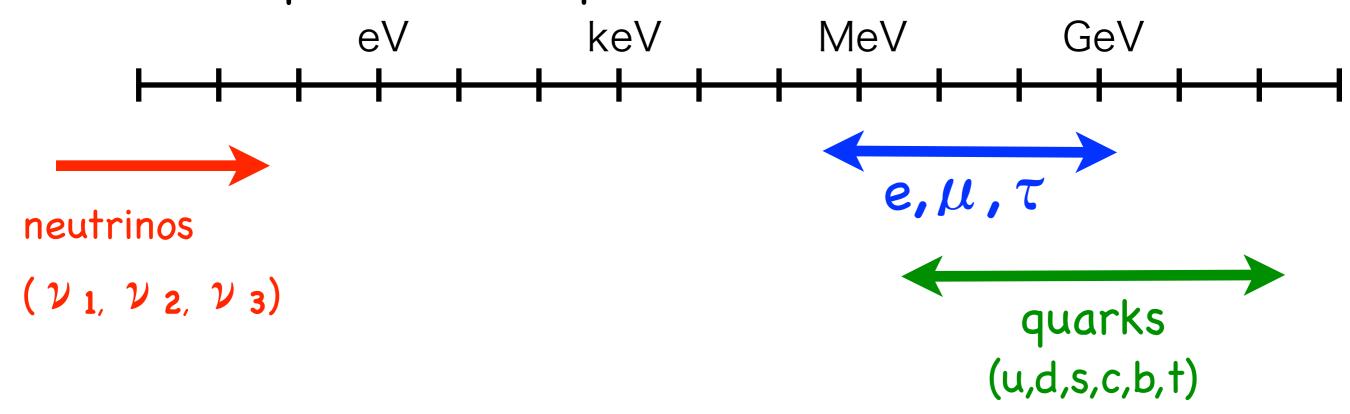
The right-handed neutrino open plays a triple role.

(1) quarks and leptons completely unitarity

```
(\uparrow\downarrow\uparrow\uparrow\downarrow\downarrow)
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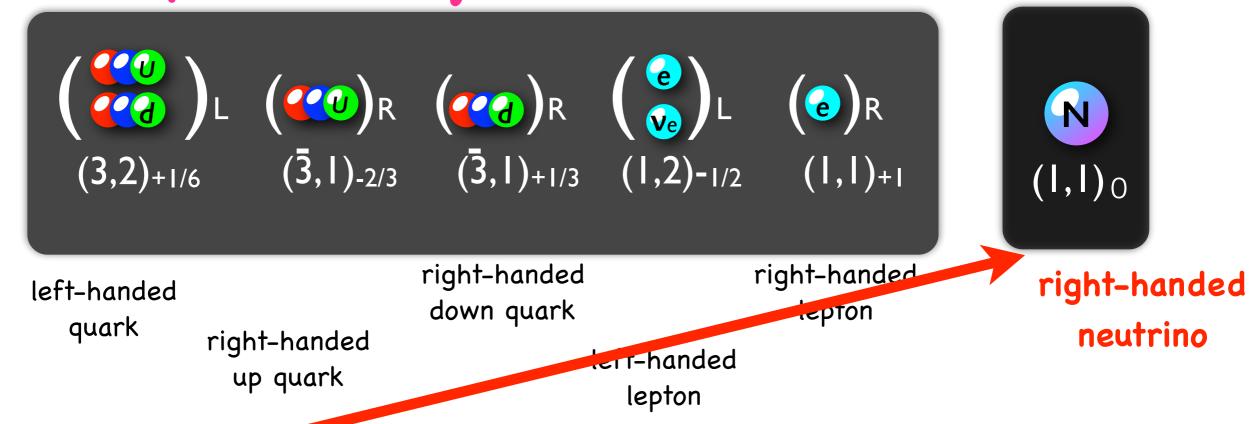
2 explains tiny neutrino mass

masses of quarks and leptons



· · · why neutrino masses are so small??

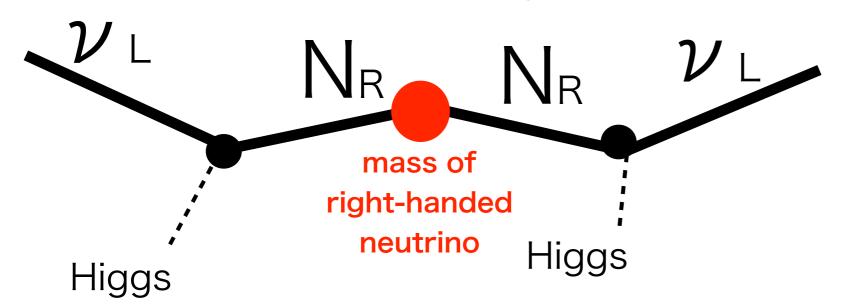
2 explains tiny neutrino mass



This guy is special singlet (feels none of three (EM, weak, and strong) forces.)

- \rightarrow it has no charge.
- \rightarrow it can be its own anti-particle.
- \rightarrow it can have a mass without Higgs.

2 explains tiny neutrino mass



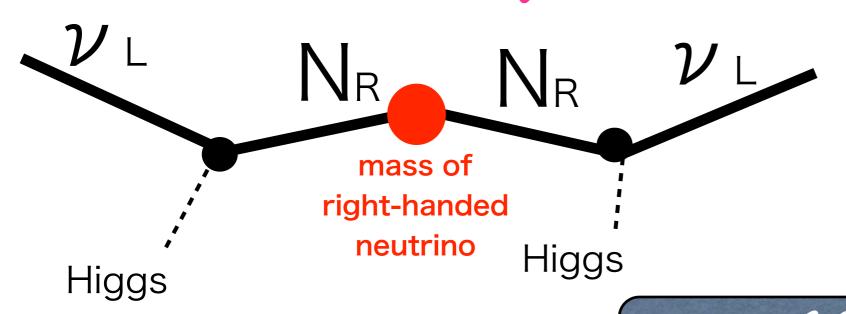
Neutrino mass = (seen e.g., by oscillation exp.)

(~ other quark/lepton mass)^2

mass of right-handed neutrino

heavy R.H. $\nu \rightarrow$ small neutrino masses ("see-saw mechanism")

2 explains tiny neutrino mass



Neutrino mass

0.01 eV

e.g. 100 GeV (~ otner quark/lepton mass)^2

mass of right-handed neutrino e.g. 10¹⁵ GeV

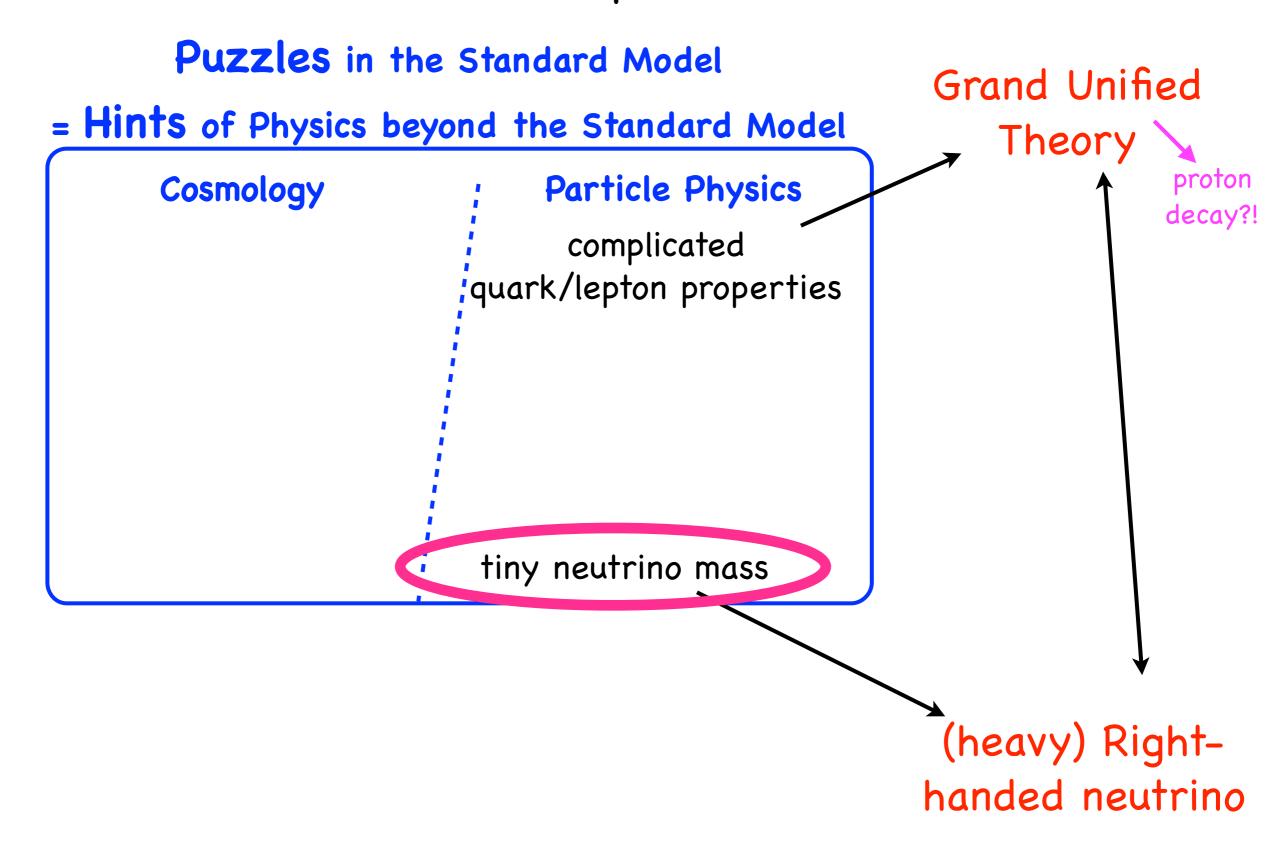
heavy R.H. $\nu \rightarrow$ small neutrino masses ("See-SaW mechanism")

3 explains matter > anti-matter

asymmetry of the universe.

----> Leptogenesis!!

more on this later ...



• to be continued...