



THE PICOSEC
MC-NET PROJECT



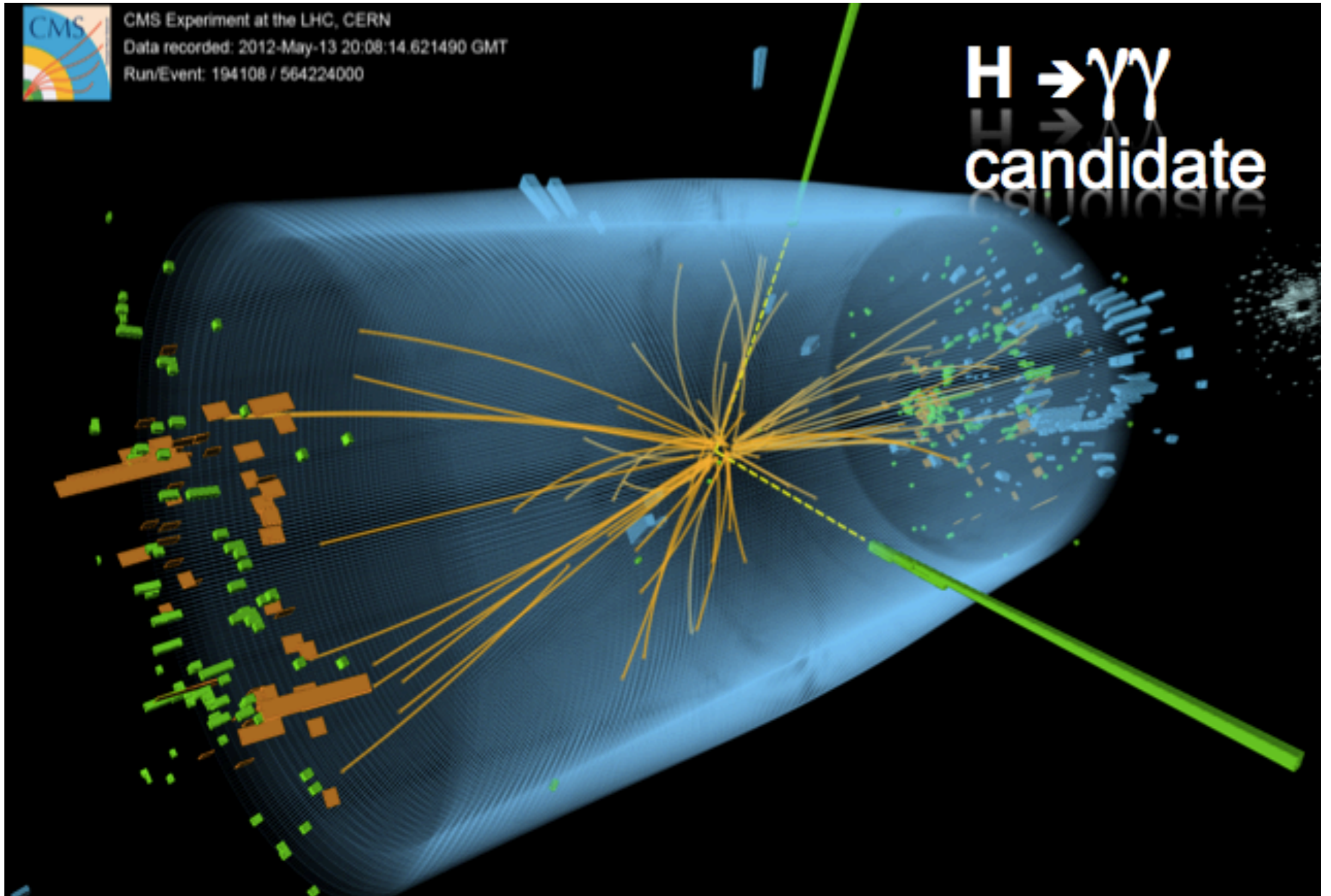
Interaction and detection of particles

Erika Garutti





A “typical” high energy physics event ...





















The Standard Model

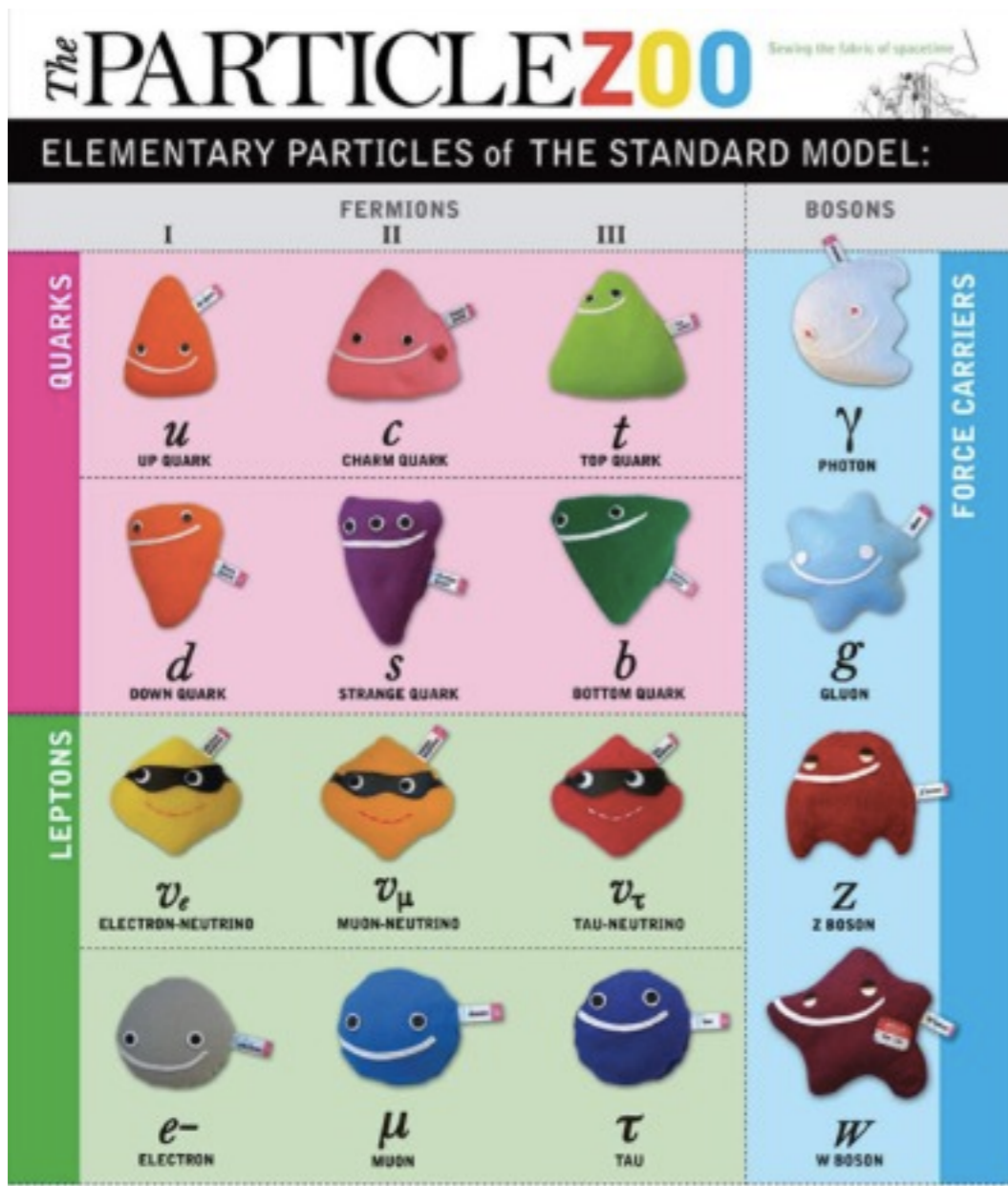
The PARTICLE ZOO Sewing the fabric of spacetime

ELEMENTARY PARTICLES of THE STANDARD MODEL:

	FERMIONS			BOSONS
	I	II	III	
QUARKS	 u UP QUARK	 c CHARM QUARK	 t TOP QUARK	 γ PHOTON
	 d DOWN QUARK	 s STRANGE QUARK	 b BOTTOM QUARK	 g GLUON
LEPTONS	 ν_e ELECTRON-NEUTRINO	 ν_μ MUON-NEUTRINO	 ν_τ TAU-NEUTRINO	 Z Z BOSON
	 e^- ELECTRON	 μ MUON	 τ TAU	 W W BOSON



The Standard Model



Three generations of matter (fermions)

	I	II	III	
mass	$2.4 \text{ MeV}/c^2$	$1.27 \text{ GeV}/c^2$	$171.2 \text{ GeV}/c^2$	0
charge	$2/3$	$2/3$	$2/3$	0
spin	$1/2$	$1/2$	$1/2$	1
name	u up	c charm	t top	γ photon
Quarks	$4.8 \text{ MeV}/c^2$ $-1/3$ $1/2$ d down	$104 \text{ MeV}/c^2$ $-1/3$ $1/2$ s strange	$4.2 \text{ GeV}/c^2$ $-1/3$ $1/2$ b bottom	0 0 1 g gluon
	$< 2.2 \text{ eV}/c^2$ 0 $1/2$ ν_e electron neutrino	$< 0.17 \text{ MeV}/c^2$ 0 $1/2$ ν_μ muon neutrino	$< 15.5 \text{ MeV}/c^2$ 0 $1/2$ ν_τ tau neutrino	$91.2 \text{ GeV}/c^2$ 0 1 Z^0 Z boson
Leptons	$0.511 \text{ MeV}/c^2$ -1 $1/2$ e electron	$105.7 \text{ MeV}/c^2$ -1 $1/2$ μ muon	$1.777 \text{ GeV}/c^2$ -1 $1/2$ τ tau	$80.4 \text{ GeV}/c^2$ ± 1 1 W^\pm W boson
				Gauge bosons



What particles can be detected?

In order for a particle to be detected, it has to interact and deposit energy



heavy charged particles
(with masses $> m_{\text{electron}}$)



electrons/positrons



photons



neutrons



What particles can be detected?

In order for a particle to be detected, it has to interact and deposit energy

stable particles ($c\tau > 500\mu\text{m}$): e^\pm , μ^\pm , π^\pm , K^\pm , p^\pm , K^0 , n , γ



heavy charged particles
(with masses $> m_{\text{electron}}$)



electrons/positrons



photons



neutrons



What particles can be detected?

In order for a particle to be detected, it has to interact and deposit energy

stable particles ($c\tau > 500\mu\text{m}$): e^\pm , μ^\pm , π^\pm , K^\pm , p^\pm , K^0 , n , γ

- Ultimately, the signals are obtained from the interactions of charged particles
- Neutral particles (photons, neutrons) have to transfer their energy to charged particles to be measured calorimeters
- what does **detecting** mean?



To be detected particles need to interact with matter..



To be detected particles need to interact with matter..





What do we measure of a particle?

- A particle is uniquely identified by its **mass and charge**

Baryons qqq and Antibaryons $\bar{q}\bar{q}\bar{q}$					
Symbol	Name	Quark content	Electric charge	Mass GeV/c^2	Spin
p	proton	uud	1	0.938	1/2
\bar{p}	anti-proton	$\bar{u}\bar{u}\bar{d}$	-1	0.938	1/2
n	neutron	udd	0	0.940	1/2
Λ	lambda	uds	0	1.116	1/2
Ω^-	omega	sss	-1	1.672	3/2



What do we measure of a particle?

- A particle is uniquely identified by its **mass and charge**
- most elementary particle have $Z_e = \pm 1$

Baryons qqq and Antibaryons $\bar{q}\bar{q}\bar{q}$					
Symbol	Name	Quark content	Electric charge	Mass GeV/c^2	Spin
p	proton	uud	1	0.938	1/2
\bar{p}	anti-proton	$\bar{u}\bar{u}\bar{d}$	-1	0.938	1/2
n	neutron	udd	0	0.940	1/2
Λ	lambda	uds	0	1.116	1/2
Ω^-	omega	sss	-1	1.672	3/2



What do we measure of a particle?

- A particle is uniquely identified by its

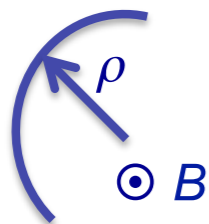
mass and charge

- most elementary particles have $Ze = \pm 1$
- particle mass determination

Baryons qqq and Antibaryons $\bar{q}\bar{q}\bar{q}$					
Symbol	Name	Quark content	Electric charge	Mass GeV/c^2	Spin
p	proton	uud	1	0.938	1/2
\bar{p}	anti-proton	$\bar{u}\bar{u}\bar{d}$	-1	0.938	1/2
n	neutron	udd	0	0.940	1/2
Λ	lambda	uds	0	1.116	1/2
Ω^-	omega	sss	-1	1.672	3/2

$$m^2 c^2 = E^2 - \vec{p}^2$$

requires **energy and momentum**

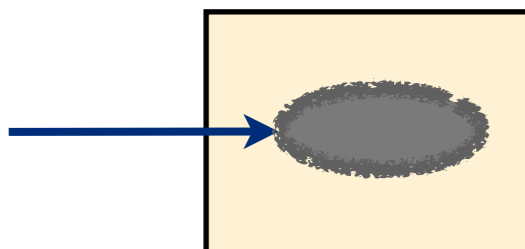


momentum determination via Lorenz relation

$$p = \rho ZeB$$

energy determined via calorimetric measurements

$$E = \gamma mc^2$$





What do we measure of a particle?

$$m^2 c^2 = E^2 - \vec{p}^2$$

E = Energy

p = Momentum



What do we measure of a particle?

$$m^2 c^2 = E^2 - \vec{p}^2$$

E = Energy

Units

[GeV]

p = Momentum [GeV/c]

Note: 1 GeV = 1.000 MeV = 1.000.000 keV



What do we measure of a particle?

$$m^2 c^2 = E^2 - \vec{p}^2$$

E = Energy

Units

[GeV]

p = Momentum

[GeV/c]

Note: 1 GeV = 1.000 MeV = 1.000.000 keV

why such a high energy?

think about the microscope resolution:



What do we measure of a particle?

$$m^2 c^2 = E^2 - \vec{p}^2$$

E = Energy

p = Momentum

Units

[GeV]

[GeV/c]

Note: 1 GeV = 1.000 MeV = 1.000.000 keV

why such a high energy?

think about the microscope resolution:

$$\Delta \sim \frac{\lambda}{2\pi}$$



What do we measure of a particle?

$$m^2 c^2 = E^2 - \vec{p}^2$$

E = Energy

Units

[GeV]

p = Momentum

[GeV/c]

Note: 1 GeV = 1.000 MeV = 1.000.000 keV

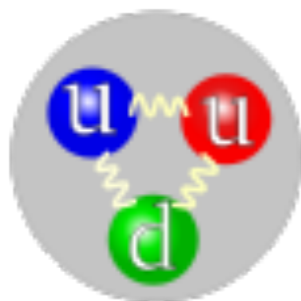
why such a high energy?

think about the microscope resolution:

$$\Delta \sim \frac{\lambda}{2\pi} = \frac{\hbar c}{E}$$

$$\frac{200 \text{ MeV fm}}{200 \text{ GeV}} = 10^{-3} \text{ fm} = 10^{-18} \text{ m}$$

Proton radius = 1 fm = 10^{-15} m (need $E > 200$ MeV)



Quark radius $< 10^{-19}$ m



How do particles interact with matter?

stable particles: e^\pm , μ^\pm , π^\pm , K^\pm , p^\pm , K^0 , n , γ

- 1) photons
- 2) heavy charged particles ($m > m_e$)
- 3) electron / positron
- 4) neutrons

Jets: collimated bunch of (mainly heavy charged) particles from the hadronization of a quark

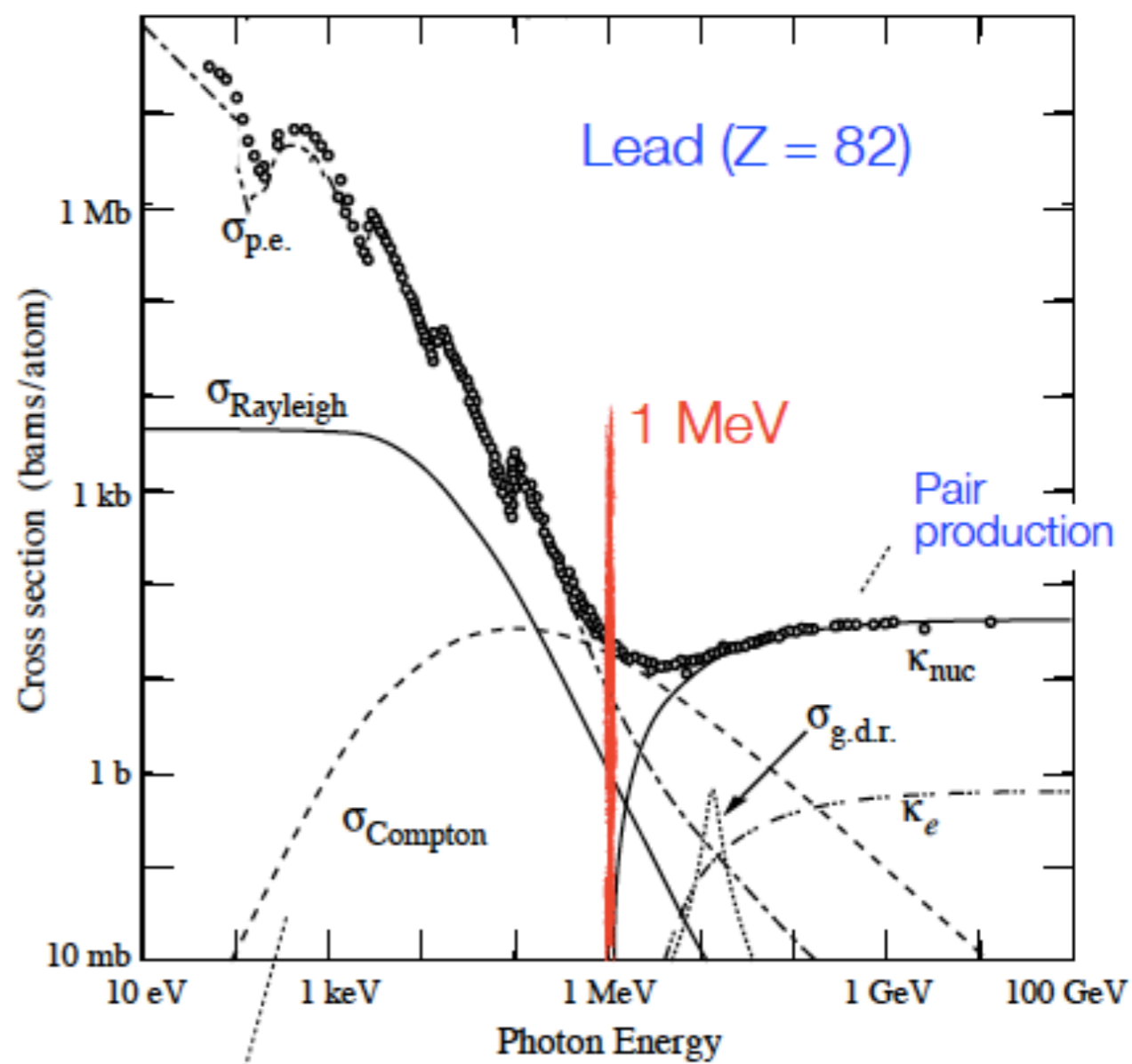
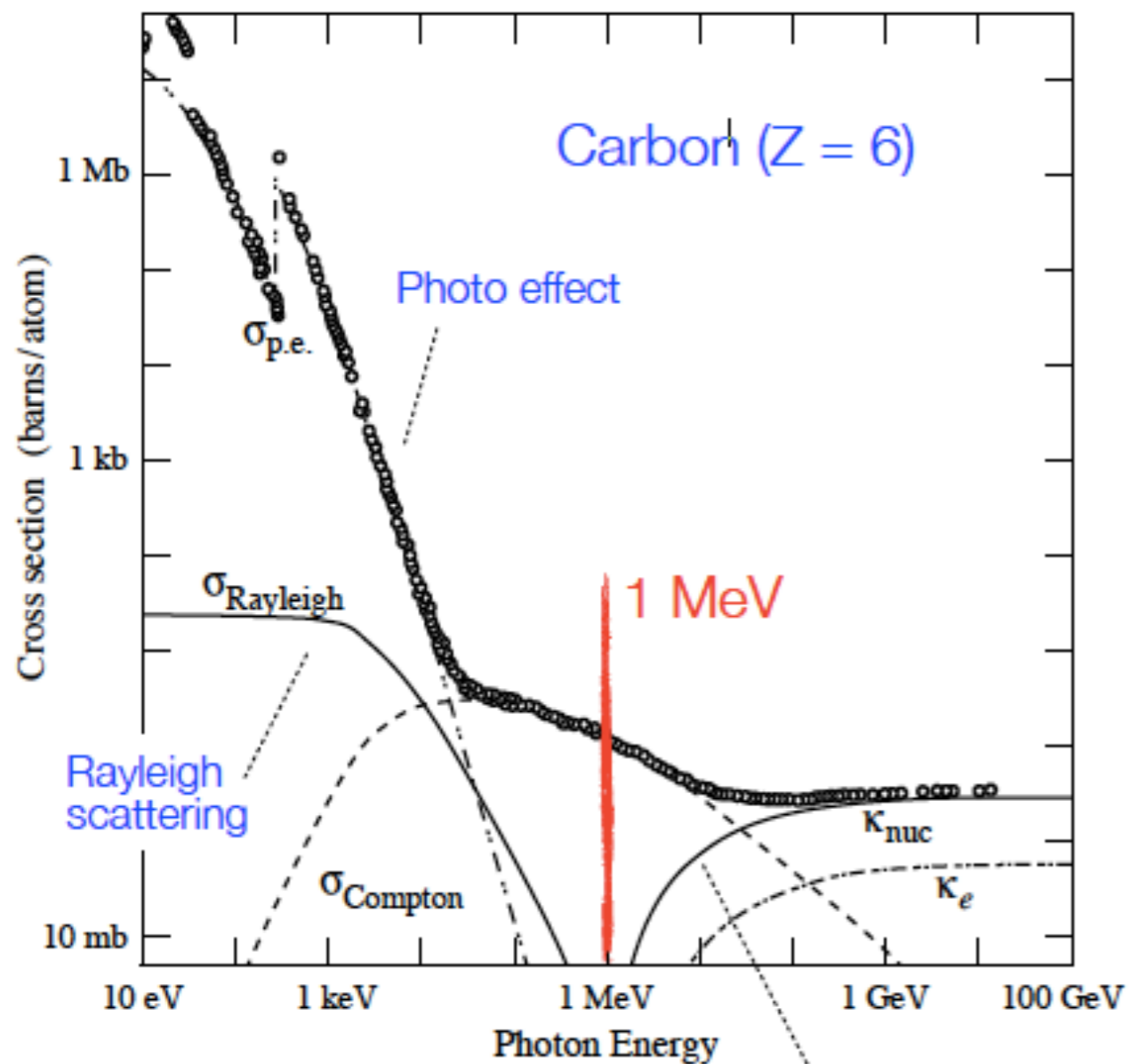
Neutrinos: react very weakly with matter

Cross section for $\nu_e + n \rightarrow e^- + p$ is around 10^{-43} cm^{-2}
in 1m Iron interaction probability 10^{-17}



I) Photons

Photon Total Cross Sections

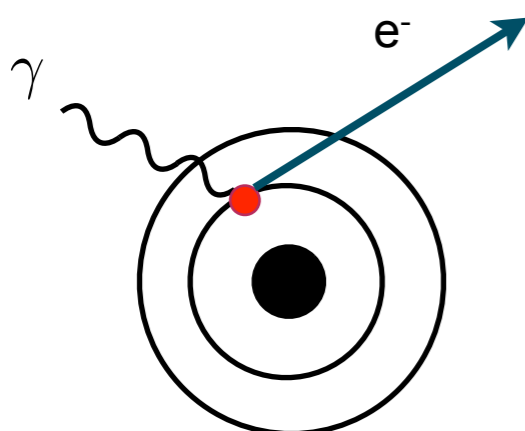




Interaction of photons with matter

- Most dominating effects - dependence on E_γ and material (Z)

Photo-Effect



Cross-section largest for $E_\gamma \approx K$ -shell energy
 Strongest E dependence
 for $I_0 < E_\gamma < m_e c^2$

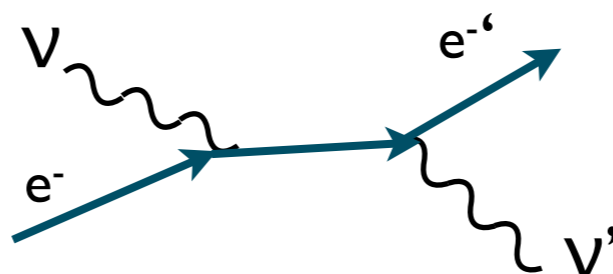
$$\sigma_{ph} = \alpha \pi a_B^2 Z^5 \left(\frac{E_\gamma}{I_0} \right)^{7/2}$$

$a_B = 0.53 \text{ \AA}$
 $I_0 = 13.6 \text{ eV}$

Softer for $E_\gamma > m_e c^2$

$$\sigma_{ph} = 2\pi r_e^2 \alpha^4 Z^5 (mc)^2 / E_\gamma$$

Compton-Scattering

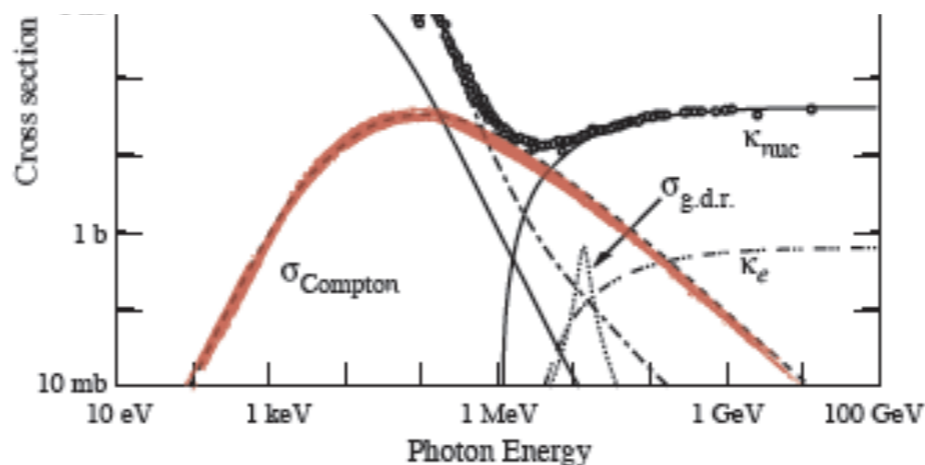


for $E_\gamma \ll m_e c^2$ $\sigma_c \propto \sigma_{Th} (1 - 2\varepsilon + \frac{26}{5}\varepsilon^2)$

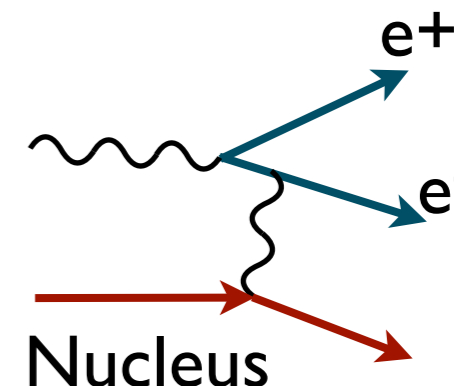
Thompson cross-section:
 $\sigma_{Th} = 8\pi/3 r_e^2 = 0.66 \text{ barn}$

$$\varepsilon = \frac{E_\gamma}{m_e c^2}$$

for $E_\gamma \gg m_e c^2$ $\sigma_c \propto \frac{\ln \varepsilon}{\varepsilon} Z$



Pair creation



Only possible if
 $E_\gamma \geq 2m_e c^2 \sim 1 \text{ MeV}$

$$\sigma_{pair} \sim \ln(E_\gamma)$$

$$\sigma_{pair} \sim Z^2$$

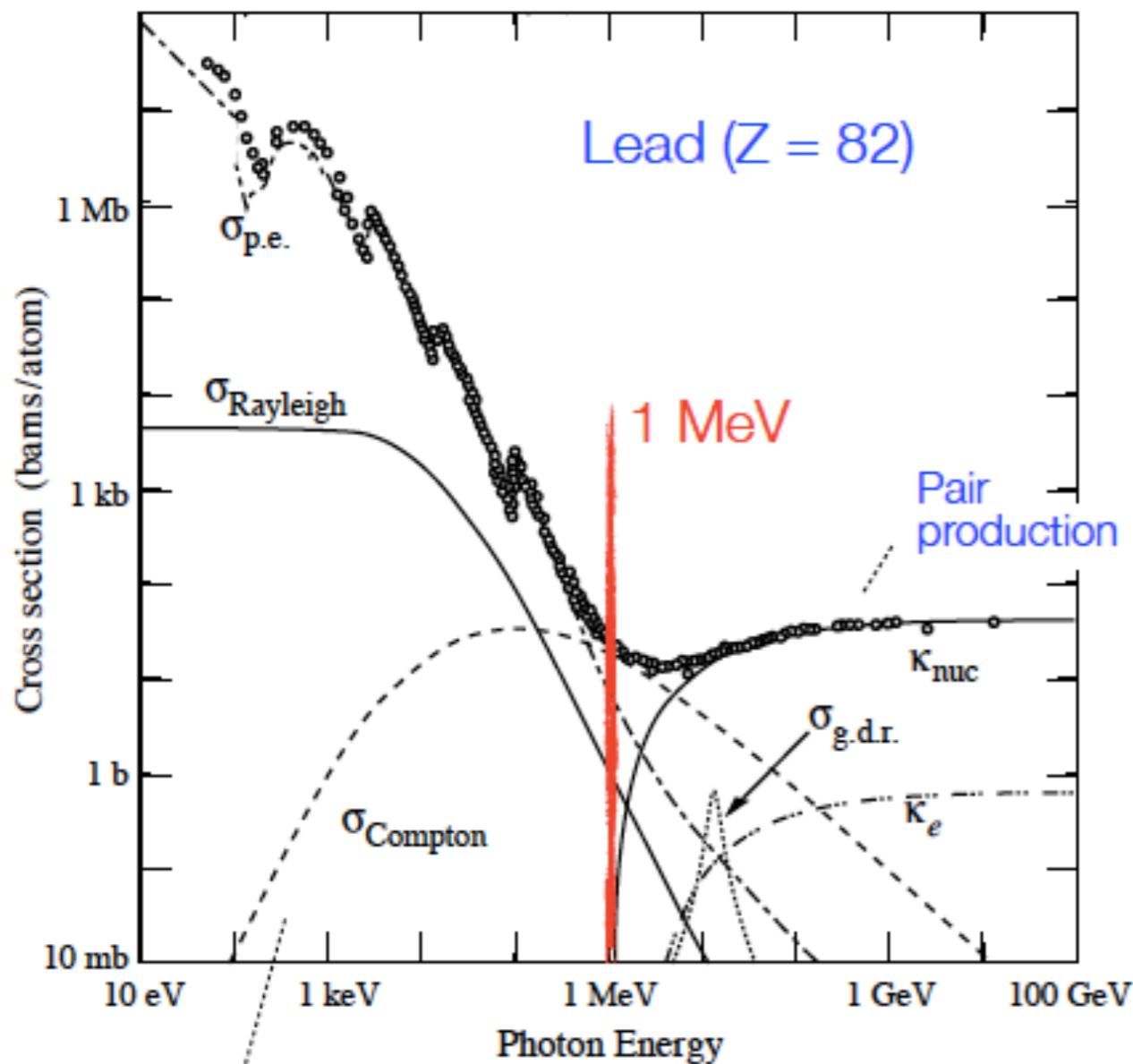
$$\sigma_{pair} = \frac{7}{9} \frac{N_A}{A} \cdot \frac{1}{X_0}$$

X_0 : Radiation length
 Radiation length

Interaction of photons with matter

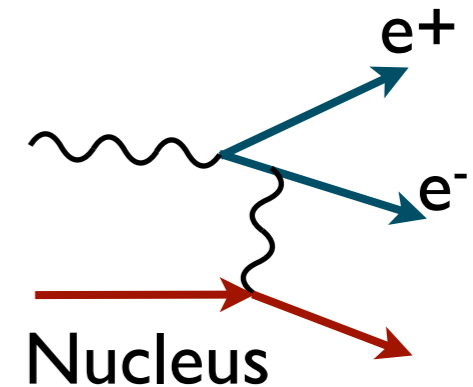


- Most dominating effects - dependence on E_γ and material (Z)



Compton scattering

Pair creation



Only possible if
 $E_\gamma \geq 2m_e c^2 \sim 1 \text{ MeV}$

$$\sigma_{\text{pair}} \sim \ln(E_\gamma)$$

$$\sigma_{\text{pair}} \sim Z^2$$

$$\sigma_{\text{pair}} = \frac{7}{9} \frac{N_A}{A} \cdot \frac{1}{X_0}$$

X_0 : Radiation length
 Radiation length

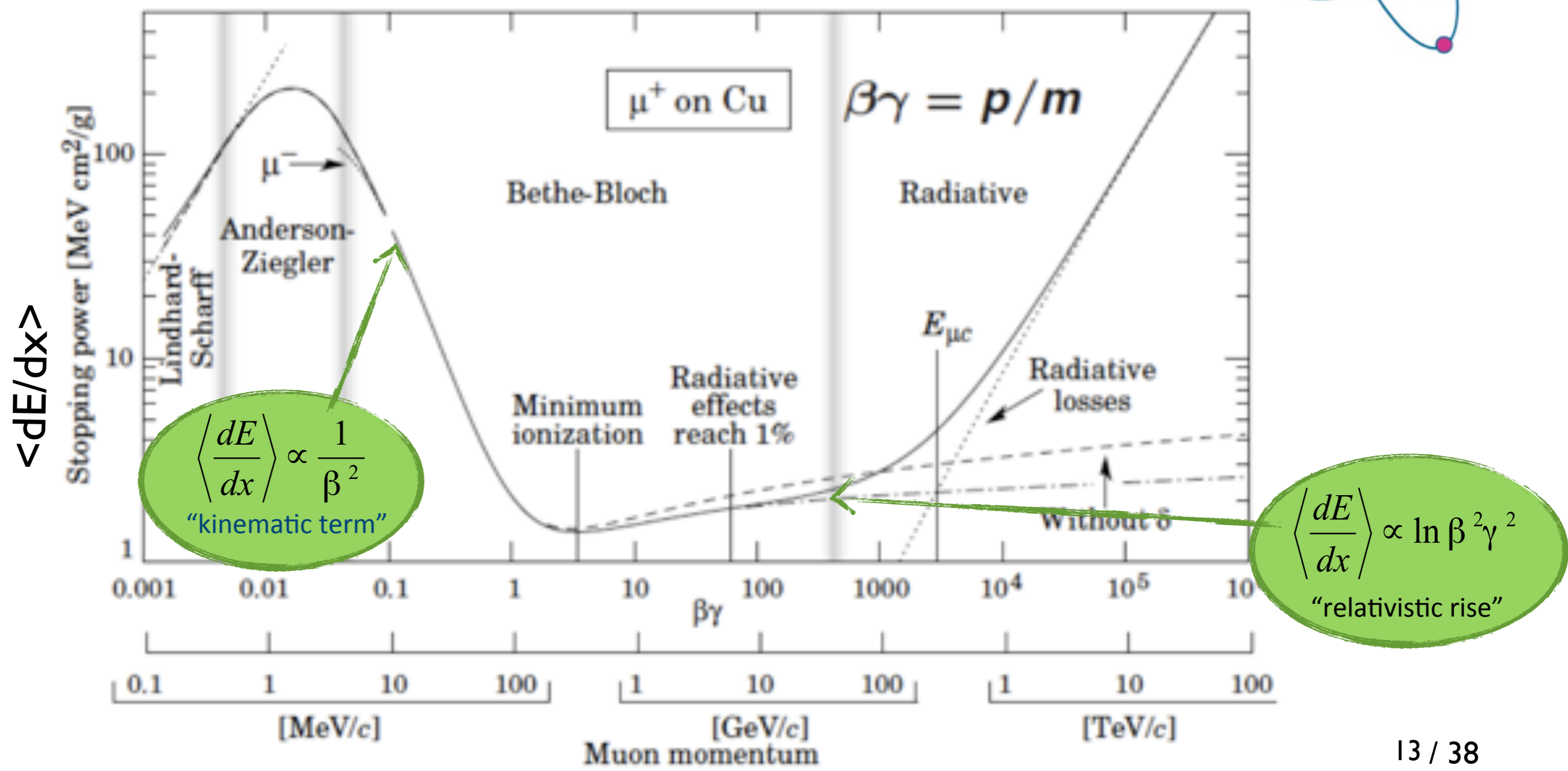
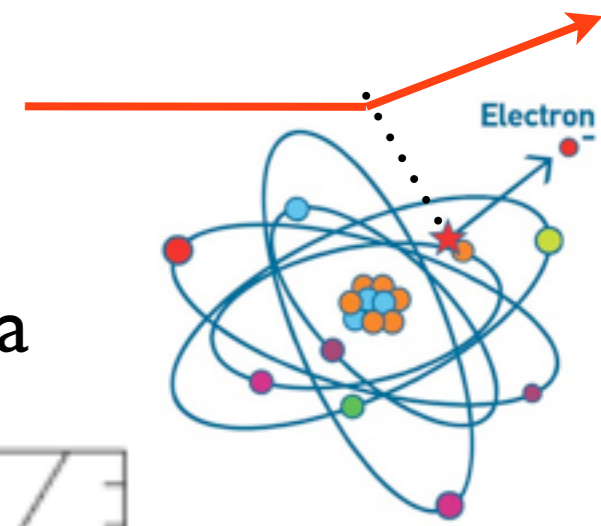


2) Heavy charged particles

$\mu^\pm, \pi^\pm, K^\pm, p^\pm$

Mostly energy transfer to the atomic electrons causing **ionization** and excitation

Mean energy loss is described by the **Bethe-Bloch** formula



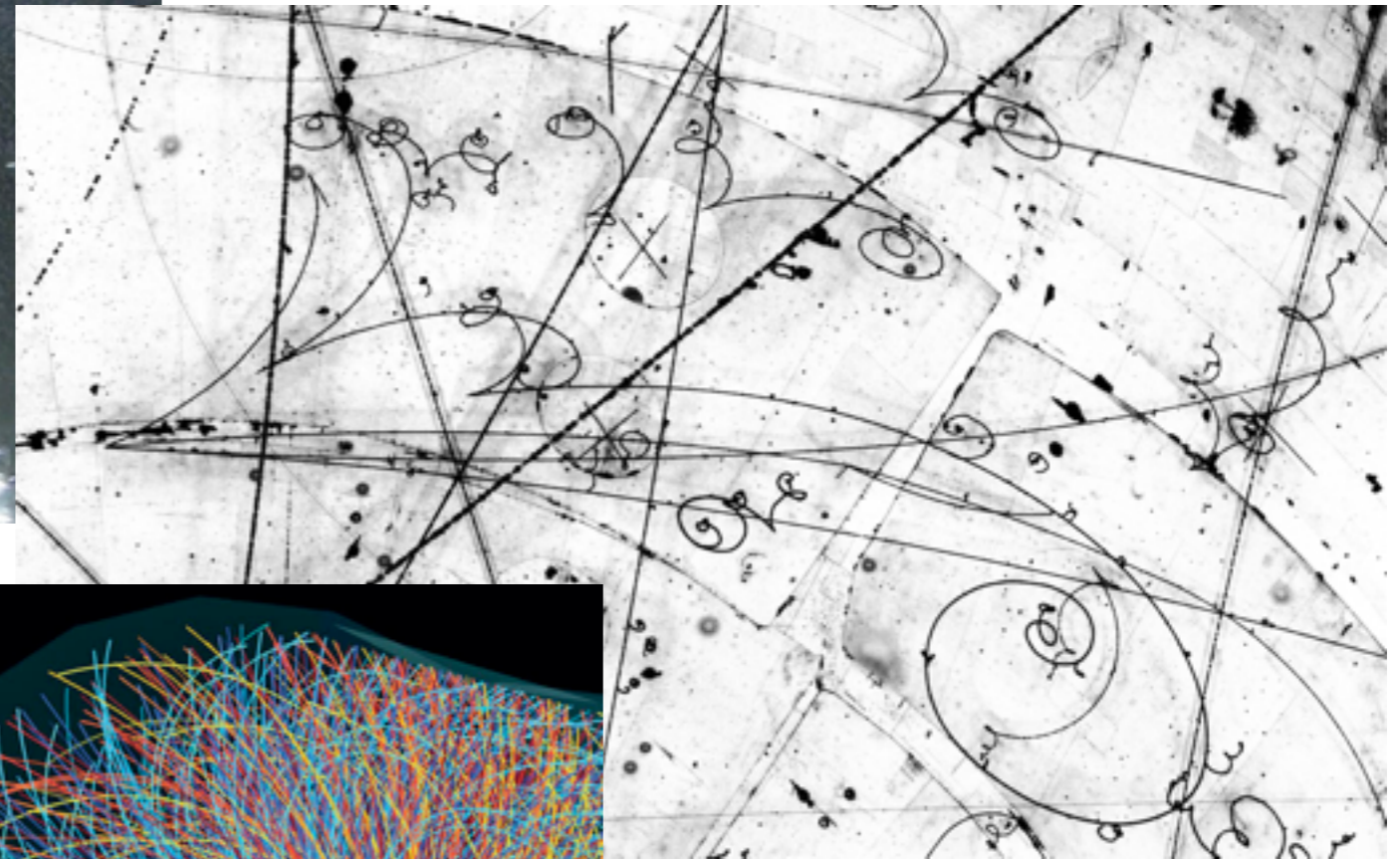


Detection via ionization

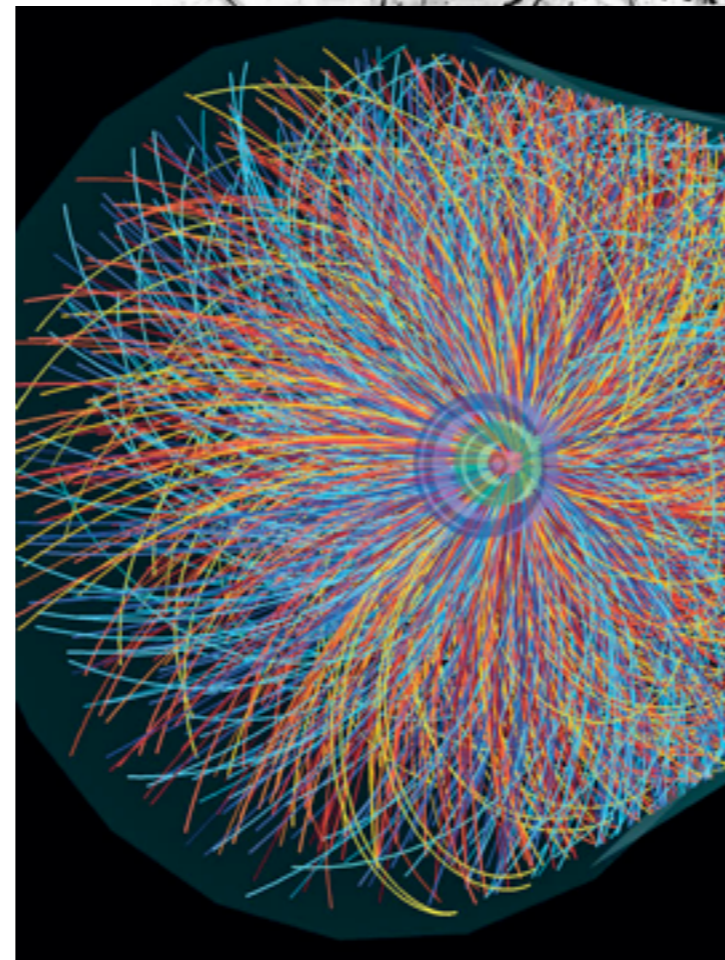


Radon220 decay in a cloud chamber

bubble chamber



spark chamber



ALICE
time projection chamber



3) Electron / positron

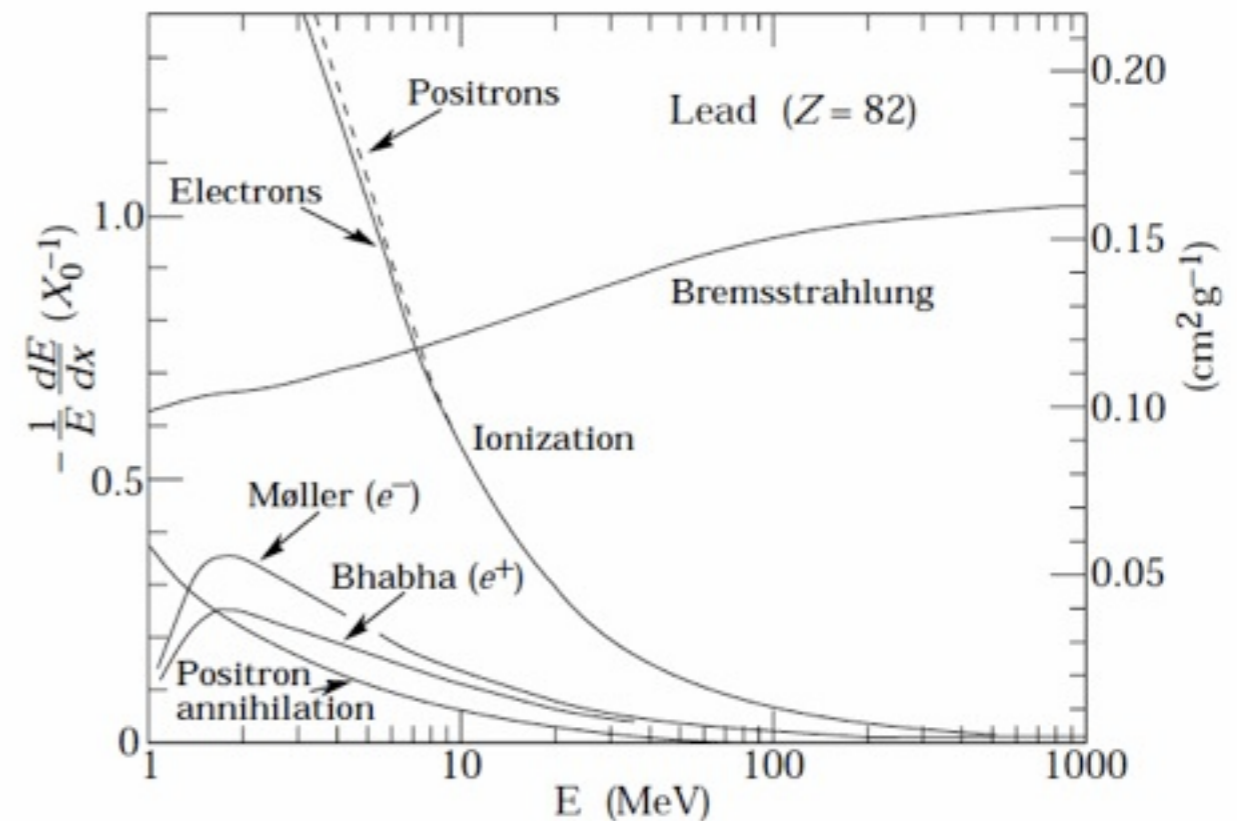
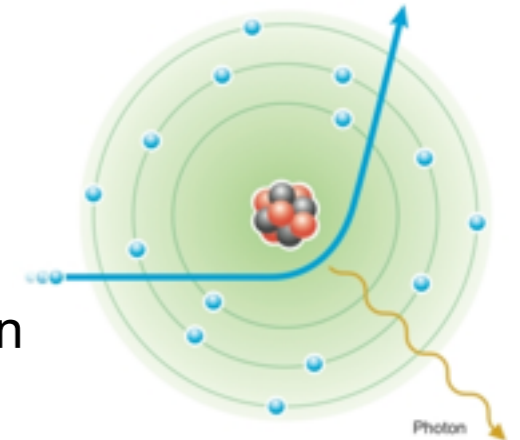


Deflection of a charge in a strong nuclear E-field → emission of a photon

$$-\frac{dE}{dx} \propto \frac{E}{m^2}$$

Effect plays a role only for e± and ultra-relativistic μ (> 1000 GeV)

Incident electron and Bremsstrahlung photon



Bremsstrahlung dominates for energies > O(10 MeV)



3) Electron / positron

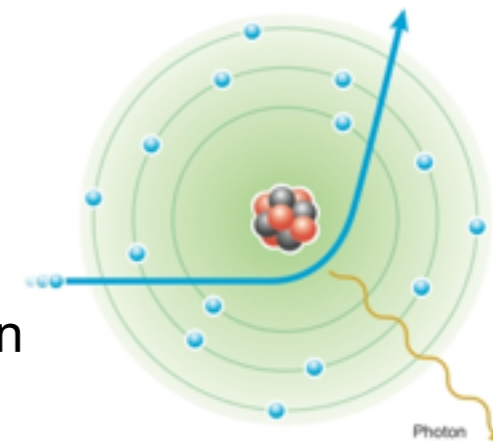


Deflection of a charge in a strong nuclear E-field → emission of a photon

$$-\frac{dE}{dx} \propto \frac{E}{m^2}$$

Effect plays a role only for e± and ultra-relativistic μ (> 1000 GeV)

Incident electron and Bremsstrahlung photon



$$-\frac{dE}{dx} = \frac{E}{X_0}$$

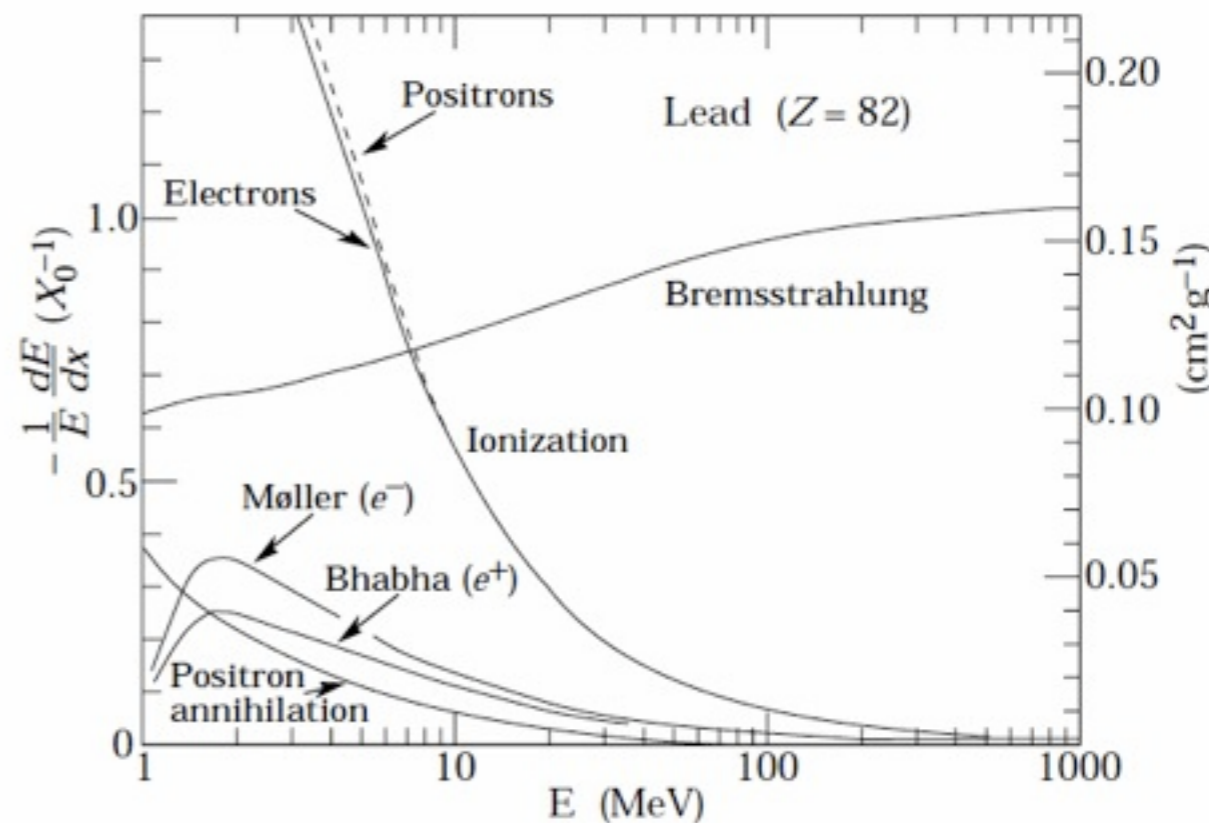


$$E = E_0 e^{-x/X_0}$$

$$X_0 = \frac{A}{4\alpha N_A Z^2 r_e^2 \ln \frac{183}{Z^{1/3}}}$$

Radiation length

- crucial parameter for detector design
- thickness of material an electron traverses till its energy is reduced by 1/e (~37%)
- depends only on material



Bremsstrahlung dominates for energies > O(10 MeV)



3) Electron / positron

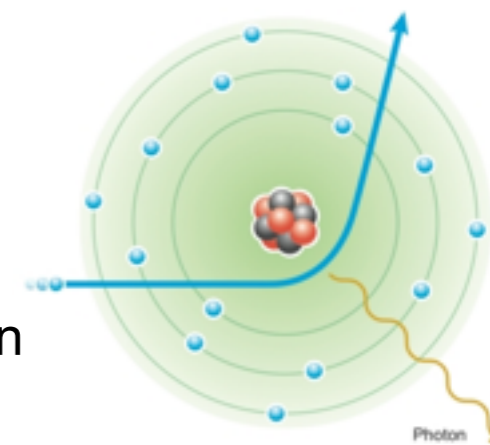


Deflection of a charge in a strong nuclear E-field → emission of a photon

$$-\frac{dE}{dx} \propto \frac{E}{m^2}$$

Effect plays a role only for e± and ultra-relativistic μ (> 1000 GeV)

Incident electron and Bremsstrahlung photon



$$-\frac{dE}{dx} = \frac{E}{X_0}$$



$$E = E_0 e^{-x/X_0}$$

$$X_0 = \frac{A}{4\alpha N_A Z^2 r_e^2 \ln \frac{183}{Z^{1/3}}}$$

Radiation length

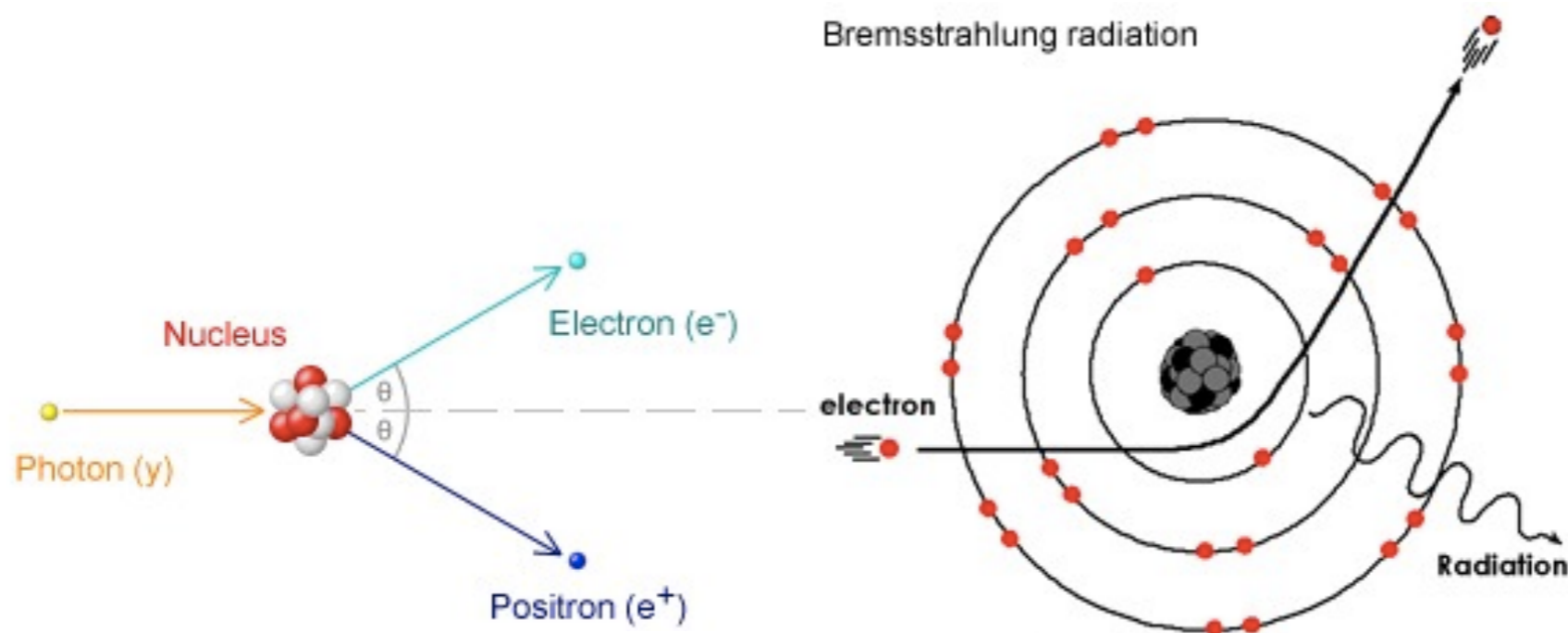
- crucial parameter for detector design
- thickness of material an electron traverses till its energy is reduced by 1/e (~37%)
- depends only on material

• Usually quoted in [g/cm²], typical values are:

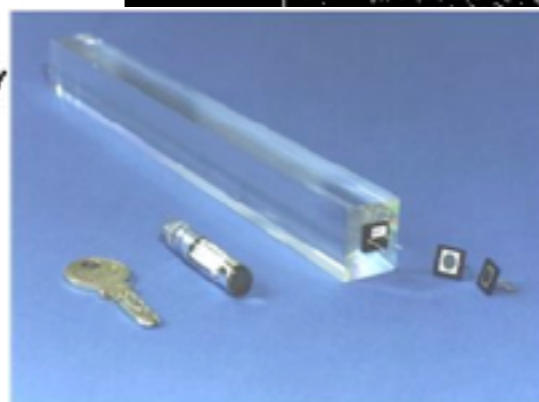
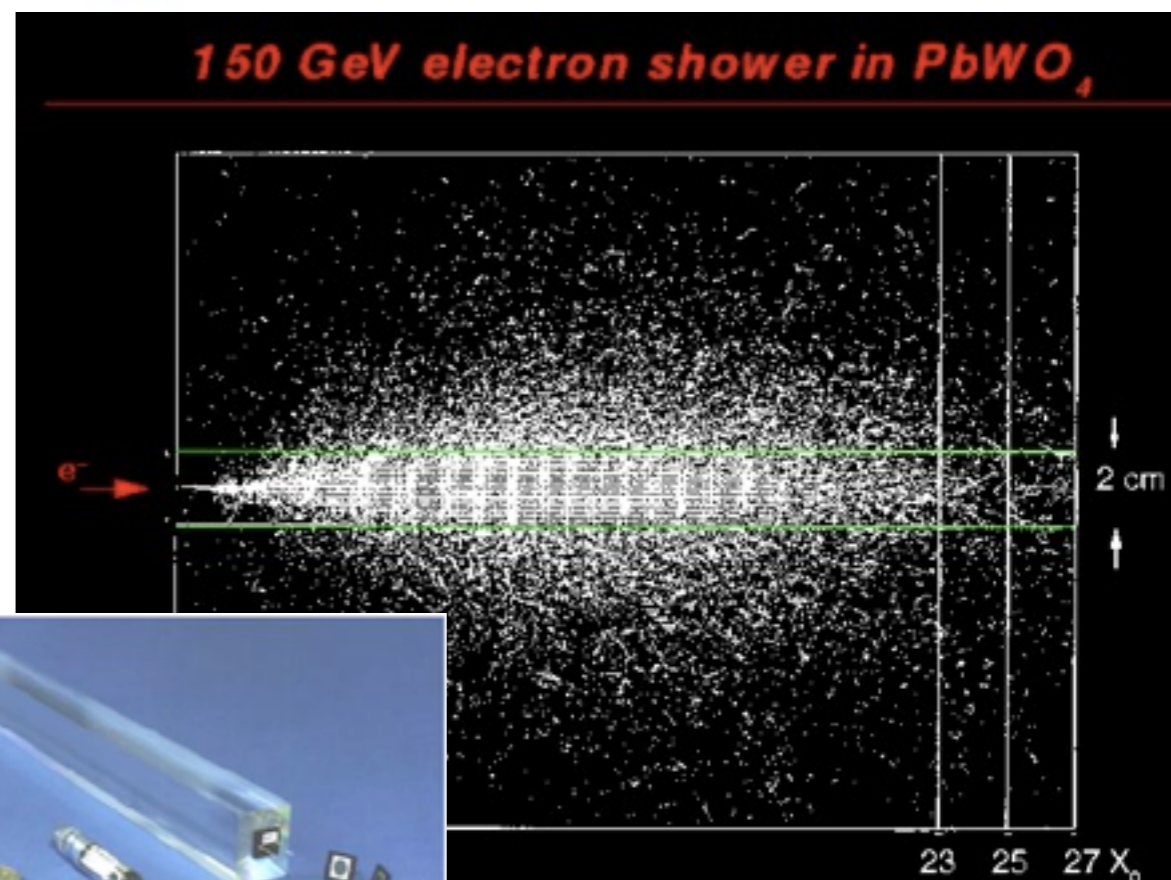
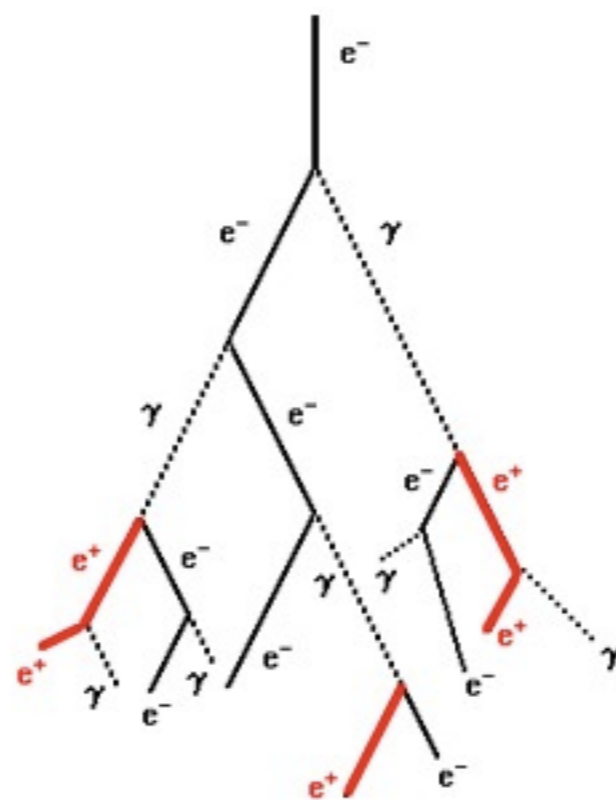
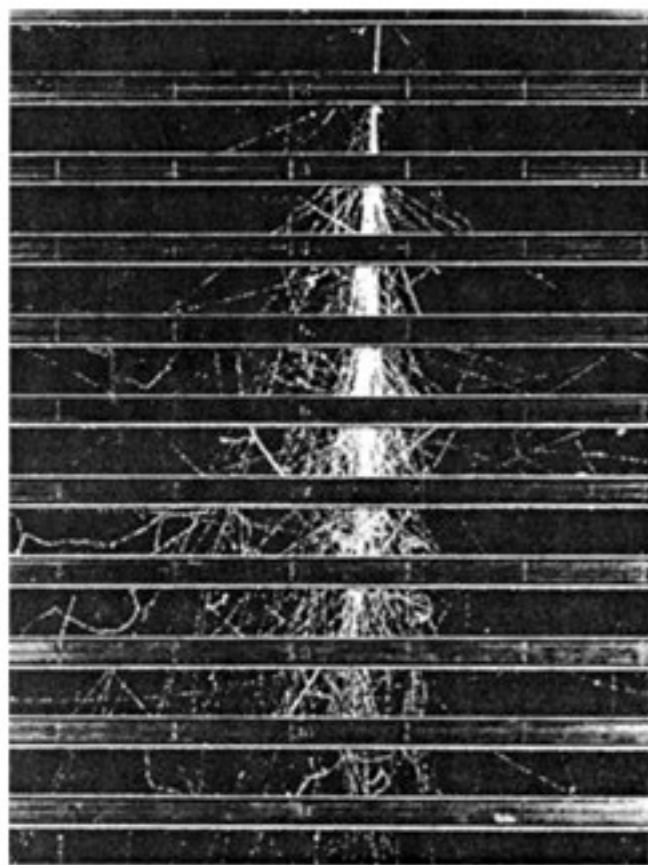
- Air: 36.66 g/cm² -> ~ 300 m
- Water: 36.08 g/cm² -> ~ 36 cm
- Silicon: 21.82 g/cm² -> 9.4 cm
- Aluminium: 24.01 g/cm² -> 8.9 cm
- Tungsten: 6.76 g/cm² -> 0.35 cm



Electromagnetic shower



$\sim \ln(E) !!!$
 $\sim X_0$



cloud chamber photograph (MIT cosmic ray group) of a particle traversing a series of brass plates

Simulated Shower in CMS- ECAL

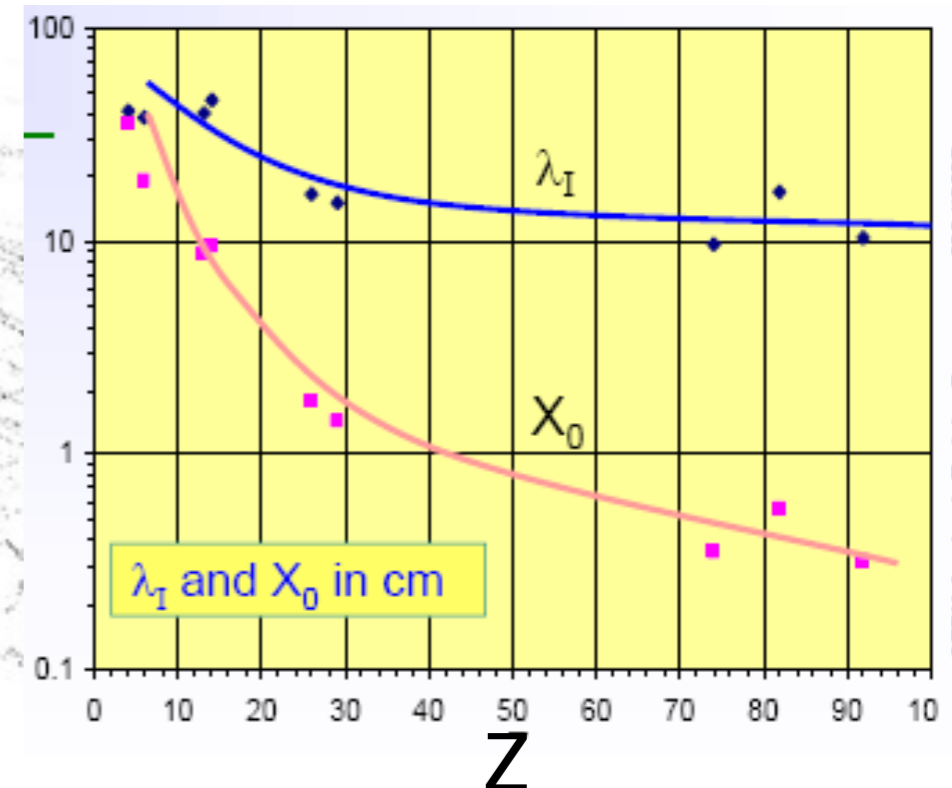
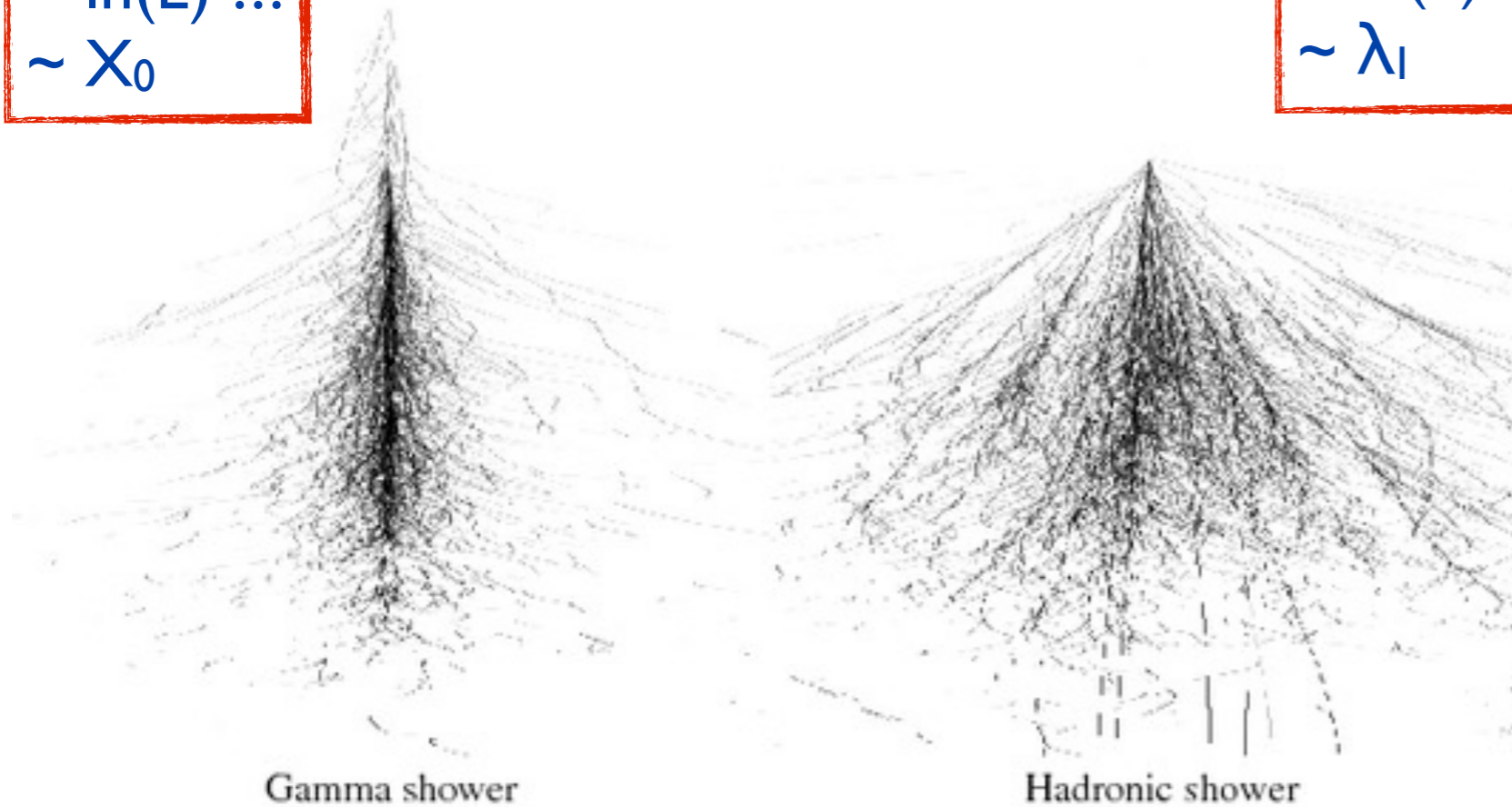
4) Neutrons (and charged hadrons)

- Neutron interaction is based only on strong (and weak) nuclear force
- To detect neutrons, one has to create charged particles
- In high energy physics **fast neutrons ($E_n > 100 \text{ MeV}$)** mainly **elastic / inelastic interactions** \rightarrow hadronic cascades

$\sim \ln(E) !!!$
 $\sim X_0$

$\sim \ln(E) !!!$
 $\sim \lambda_I$

Interaction length:
 $\lambda_I (\text{g cm}^{-2}) \propto A^{1/3}$

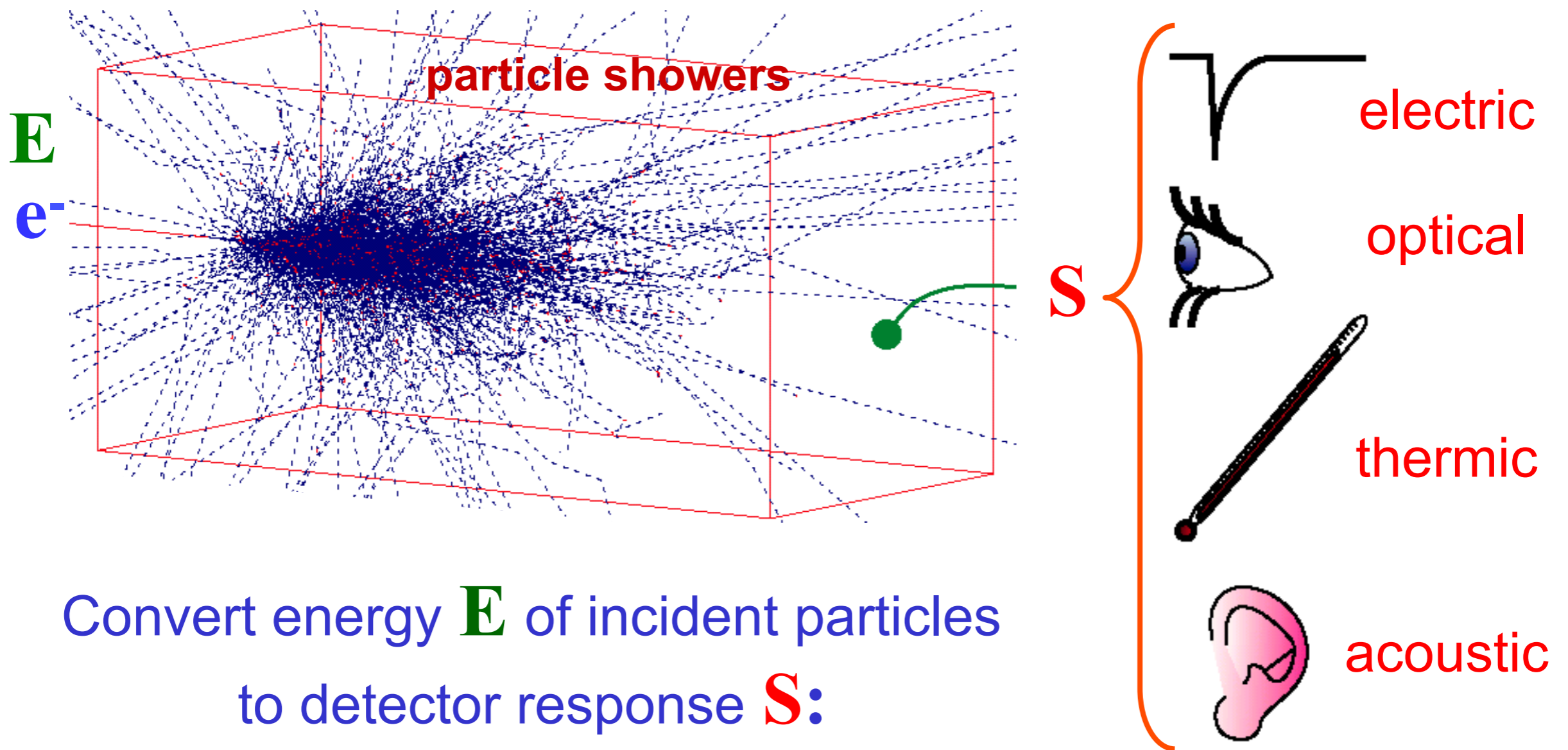




How do we measure a particle?



How do we measure a particle?



Convert energy **E** of incident particles
to detector response **S**:

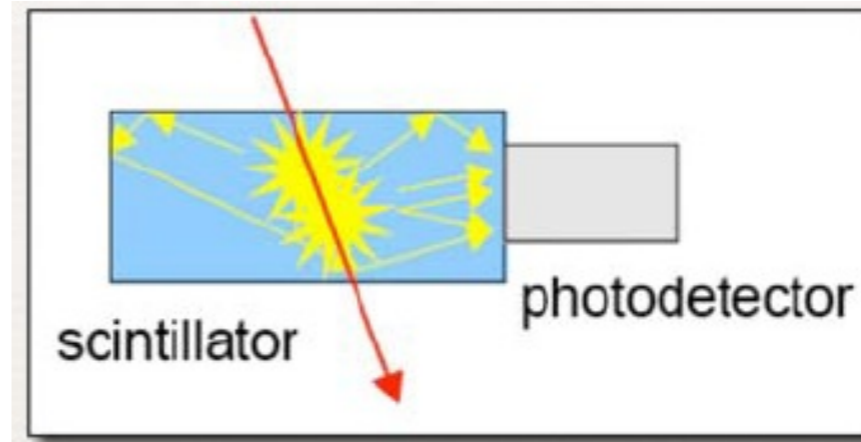
$$\mathbf{S} \propto \mathbf{E}$$



How do we measure a particle?

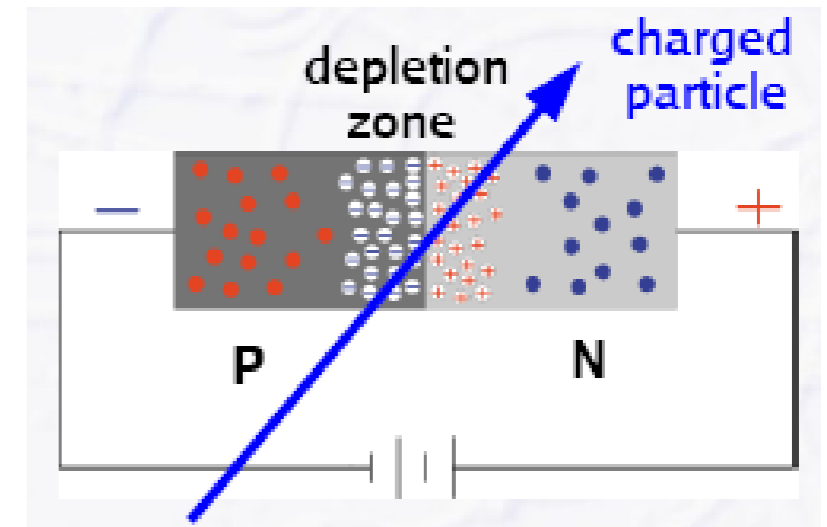
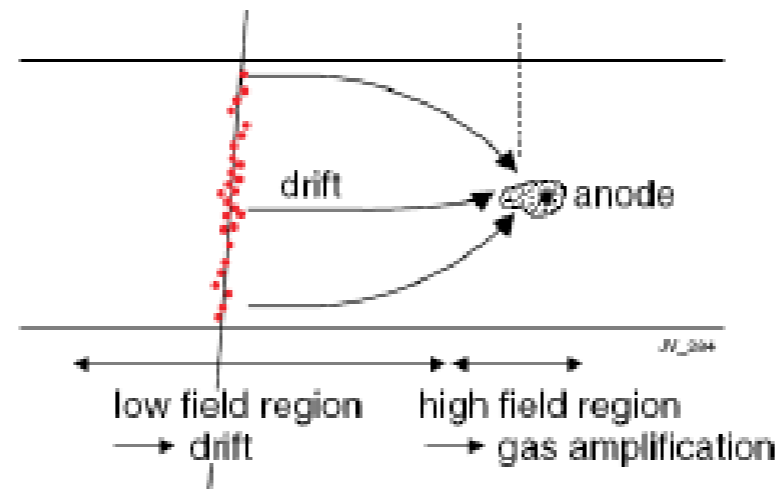
1) Convert particle energy to **light**:
scintillator (org. / in-org.)

& measure light:
PMT / APD / HPD / SiPM ...



2) Measure ionization E:
gas
noble liquids
semiconductors

& measure charge signal

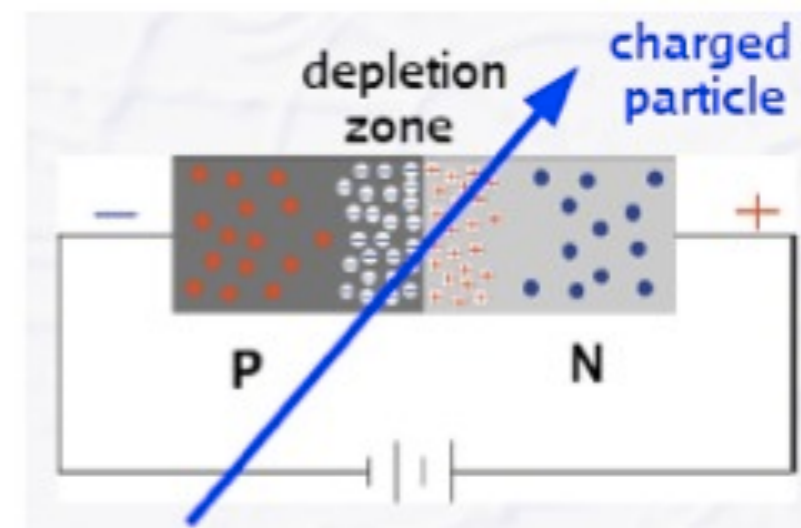
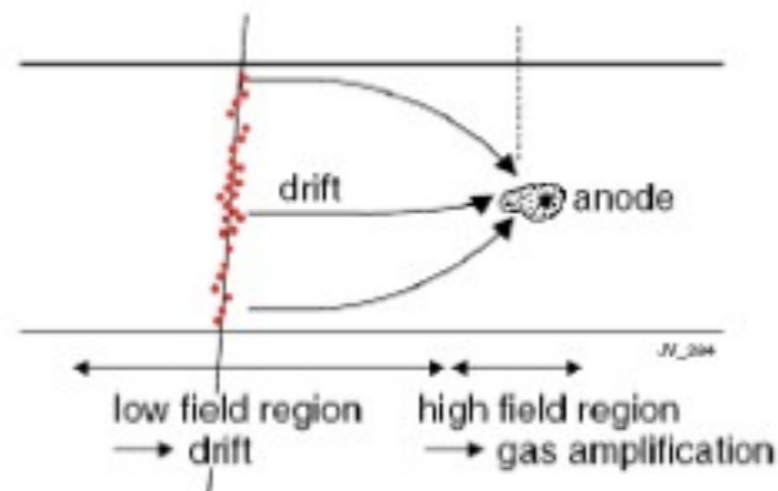
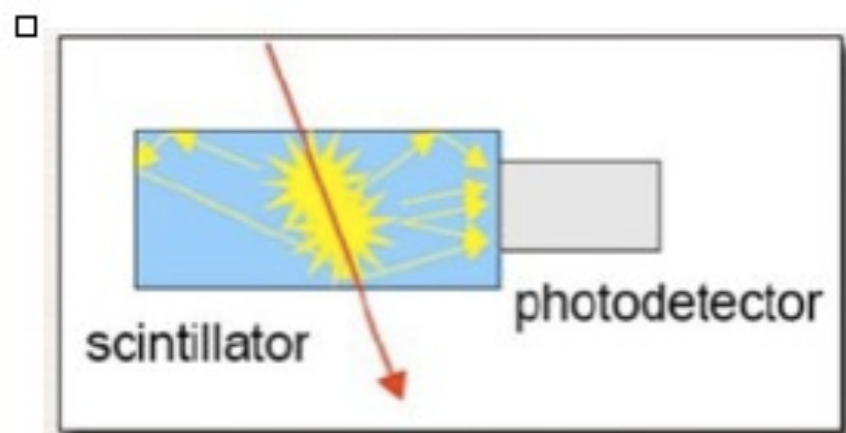


3) Measure temperature:

specialized detectors for: DM, solar ν s, magnetic monopoles, double β -decay
very precise measurements of small energy deposits
phenomena that play a role in the 1 Kelvin to few milli-Kelvin range



Measurement of ionization charge



Relevant parameter:

ionization energy (I_0) = energy needed to create a detectable quantum

semiconductors $I_0 = 1-5$ eV per e-hole pair

gas detectors $I_0 = 20-40$ eV per e-ion pair

scintillators $I_0 = 400-1000$ eV to create a photon (need to convert photon!)

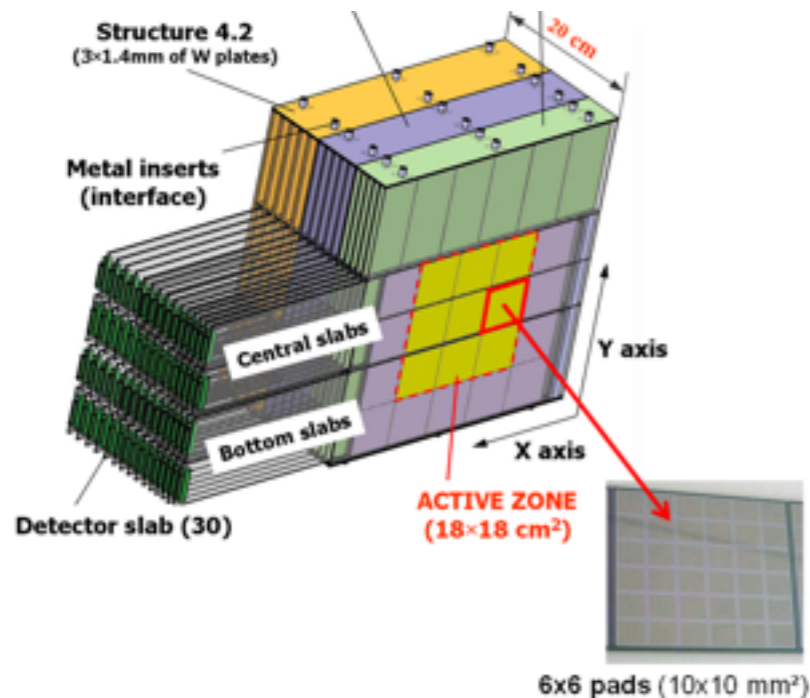


Historically

- **semiconductors** & **gas** mainly used in tracker detectors
→ p measurement (+ dE/dx)
- **scintillators** (organic/inorganic) mainly used in calorimeters
→ E measurement
- ... but exceptions exist

as detector developer be always open minded and daring !

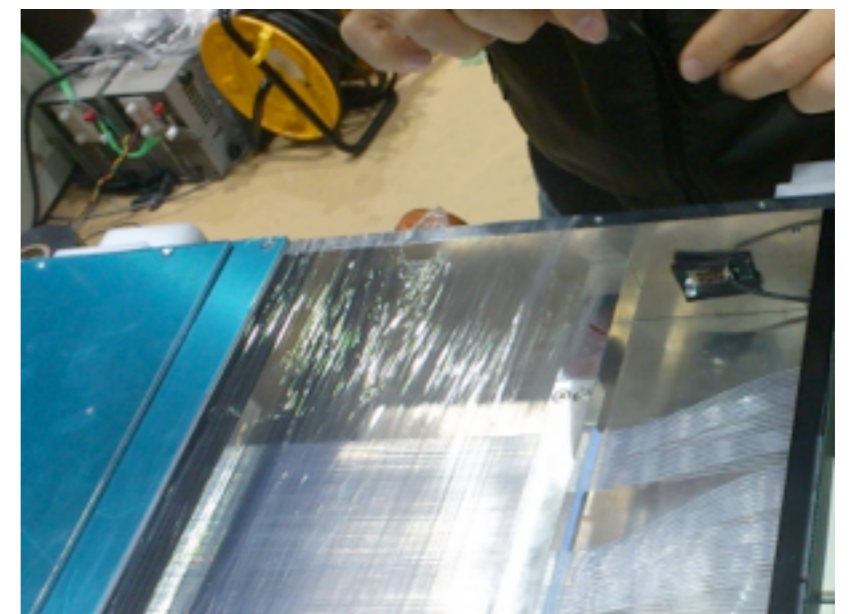
Silicon - ECAL



Gas readout for HCAL



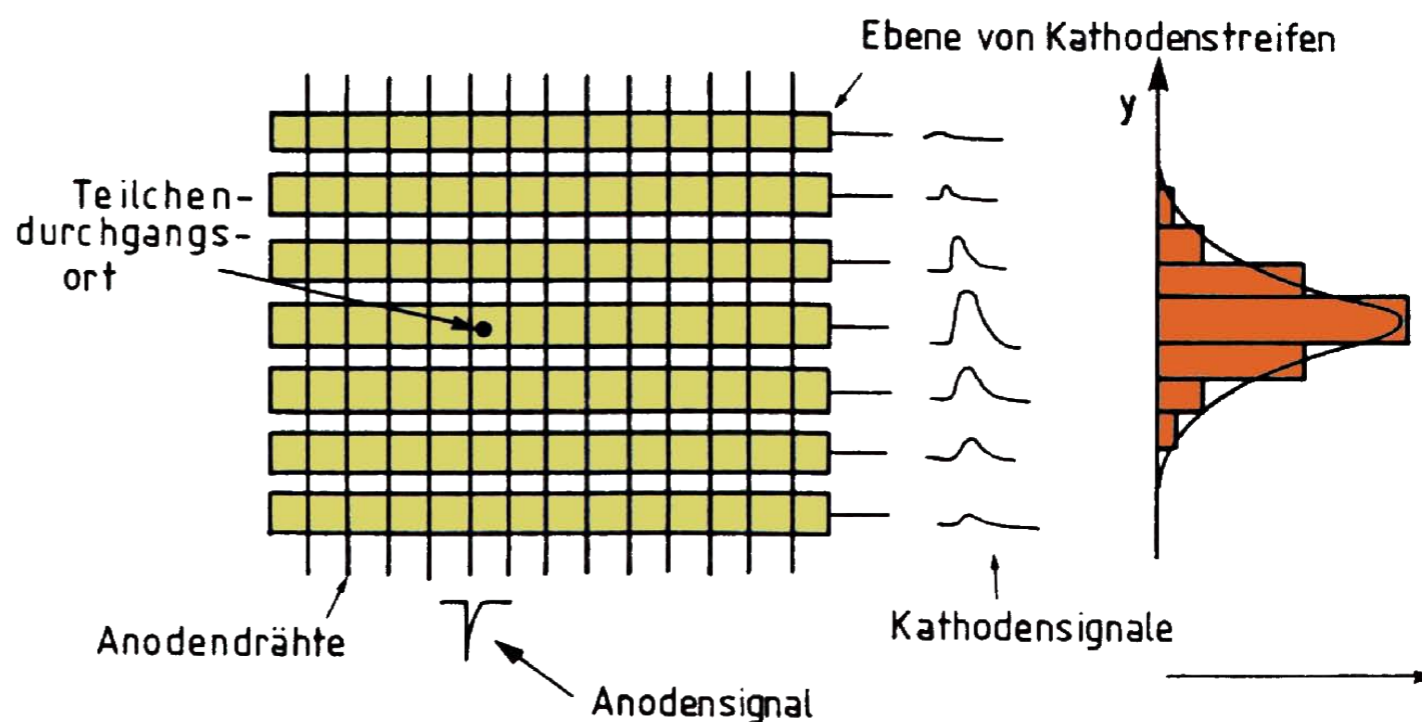
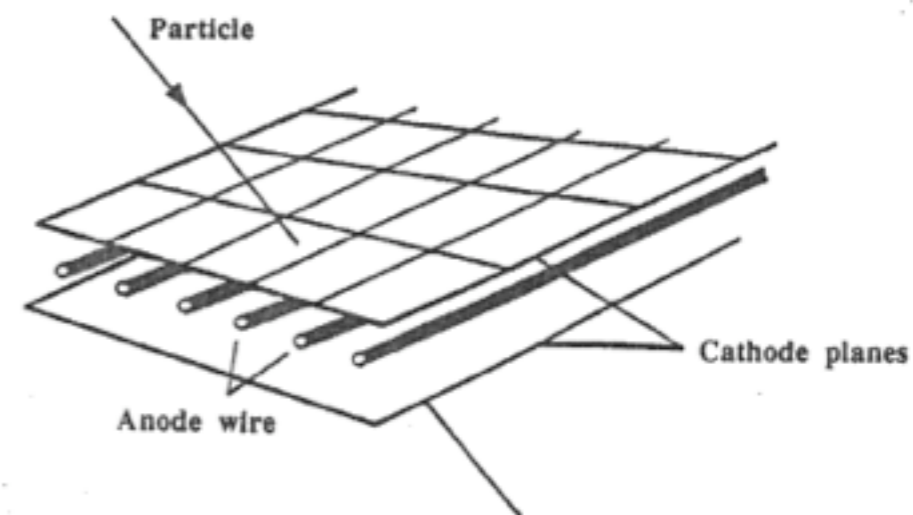
Fiber tracker





Track position: ionization chambers

- Extreme successful approach to provide good spatial resolution with gas detectors
- Multi wire proportional chamber (MWPC)
- Gas-filled box with a large number of parallel detectors wires, each connected to individual amplifiers
- G. Charpak 1968 (Nobel-Preis 1992)





Adding time: drift chamber

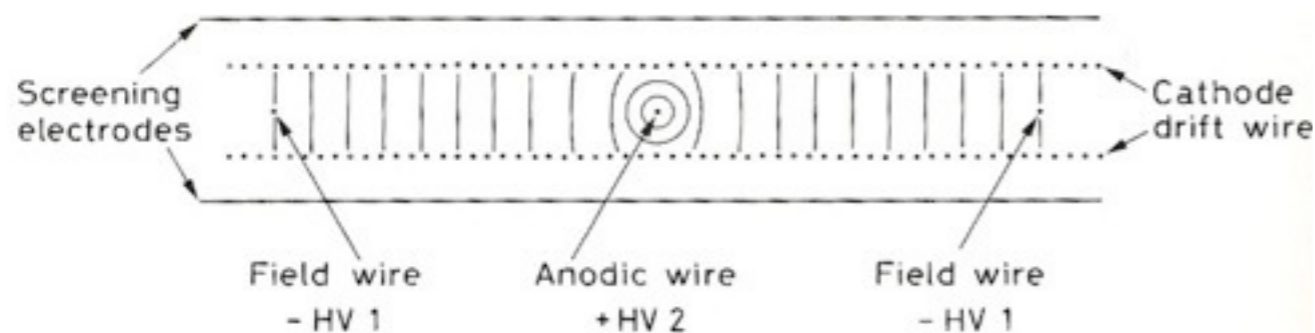
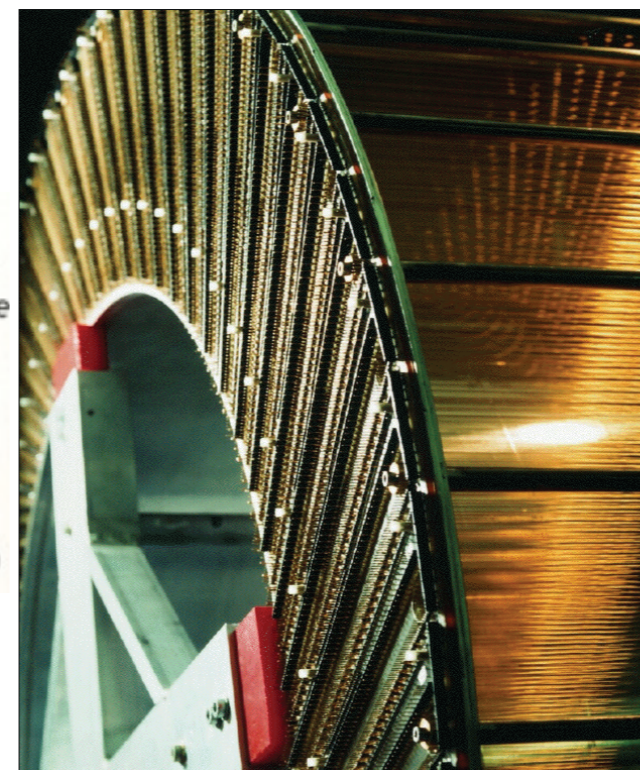


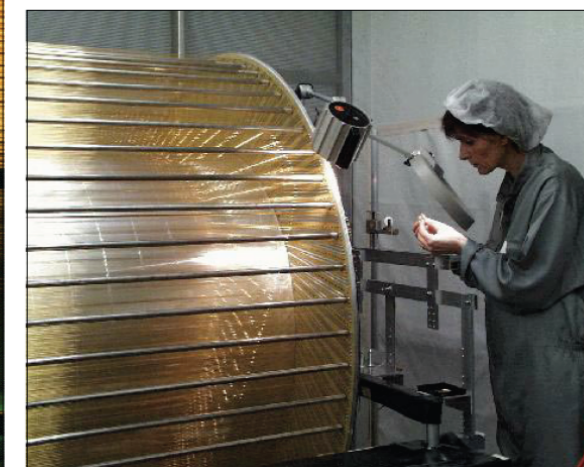
Fig. 6.16. Drift chamber design using interanode field wires (from *Breskin et al.* [6.22])



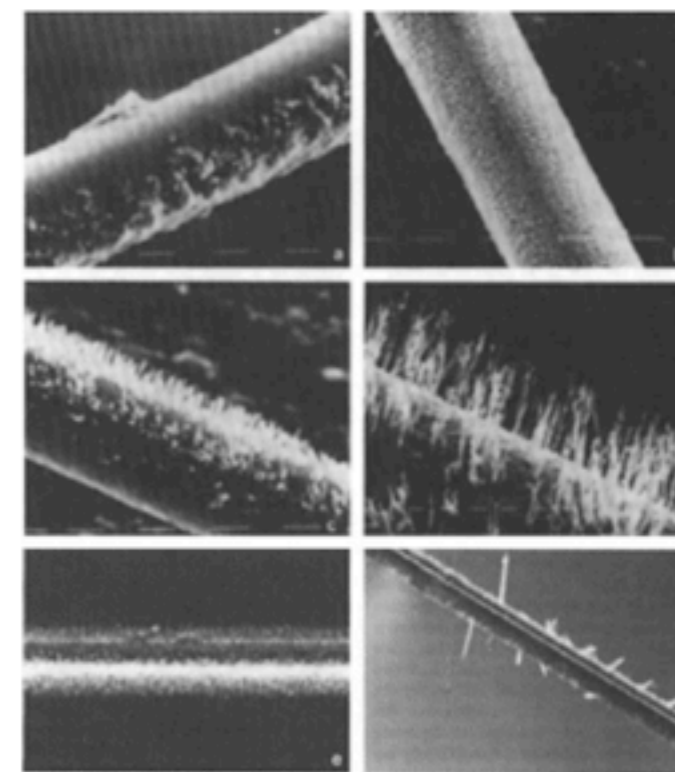
Cylindrical
Drift Chamber

[H1 Experiment]

Number of wires: ~ 15000
Total force from wire tension: ~ 6 t



- Electric field is designed in a way that electrons drift with a constant velocity and only amplify very close to the wire
- If time of arrival of a particle is known (trigger), one can derive from the signal arrival time at the anode the position of the track
- Condition: the HV field distribution and therefore the drift velocity within the gas is well known

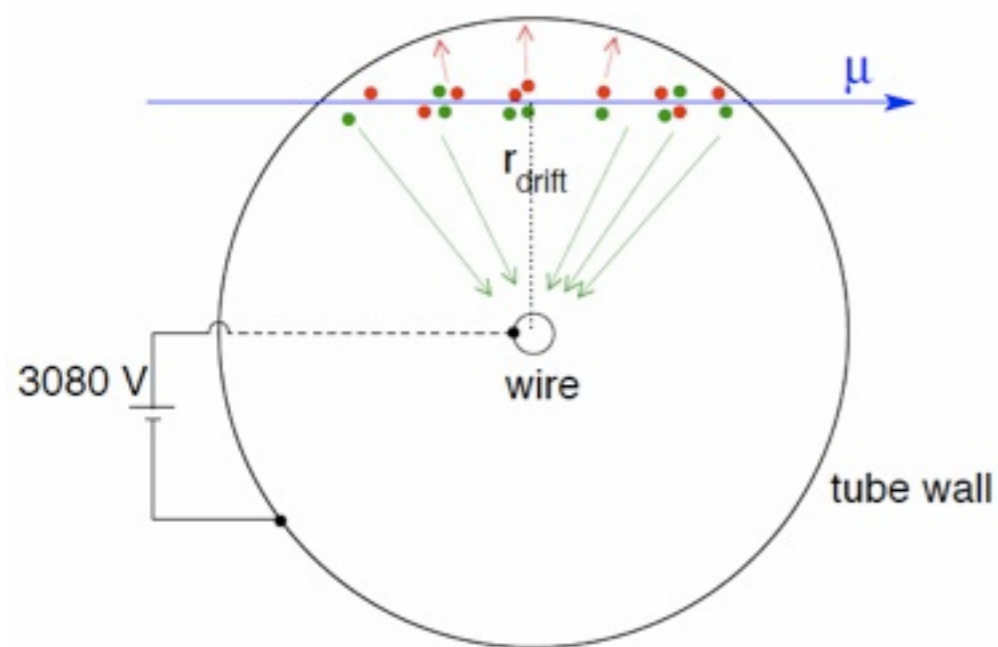


Wire aging due to gas contamination
23 / 38



Commonly used: Drift Tube

- Example: ATLAS Muon-System



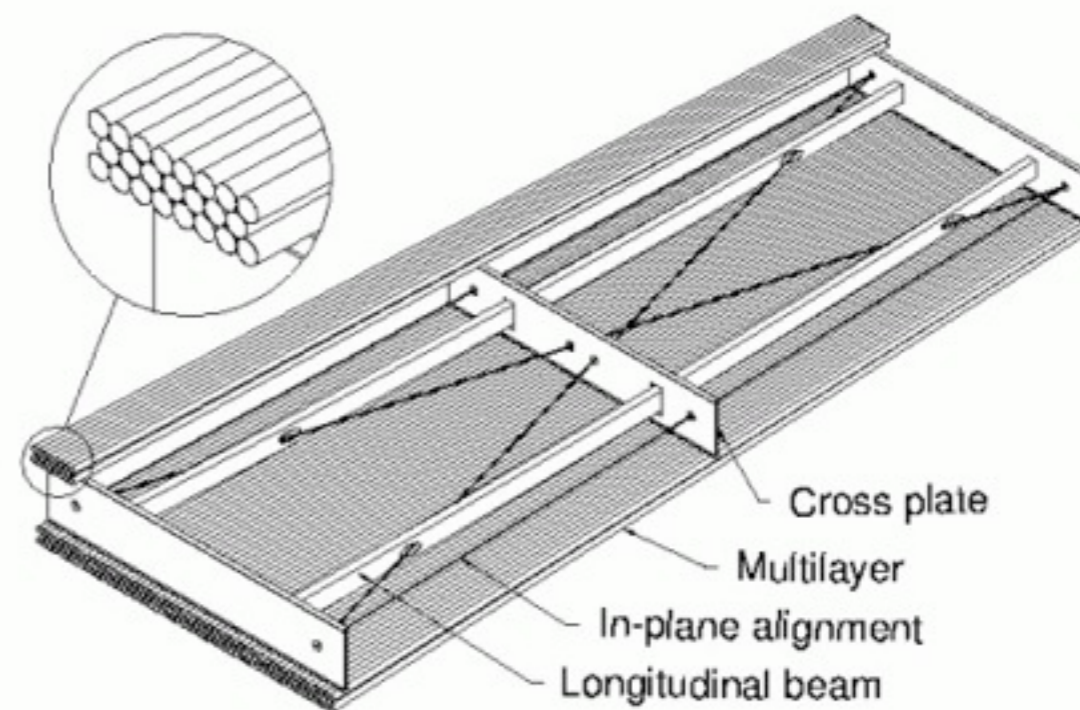
Measurement of the drift time: defines the smallest distance of the track to the wire

⇒ right/left ambiguity: multiple layers shifted to each other necessary

⇒ spacial resolution typically $\sim 100 \mu\text{m}$



Foto: CERN

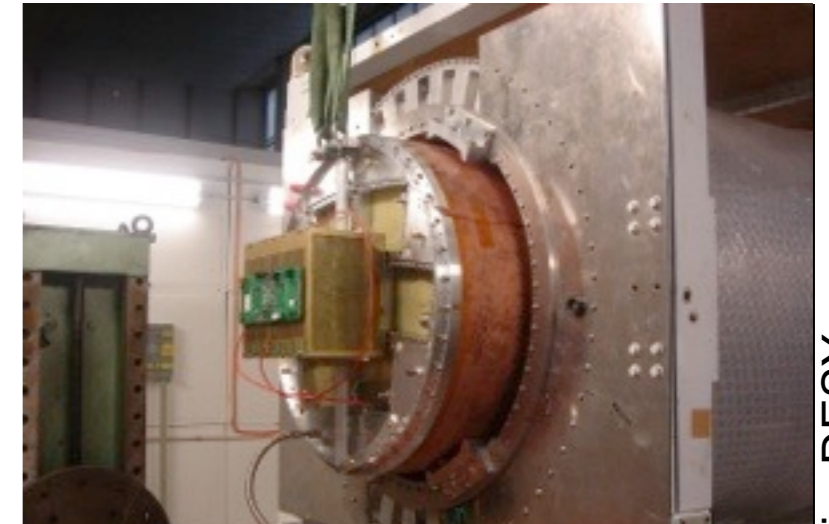
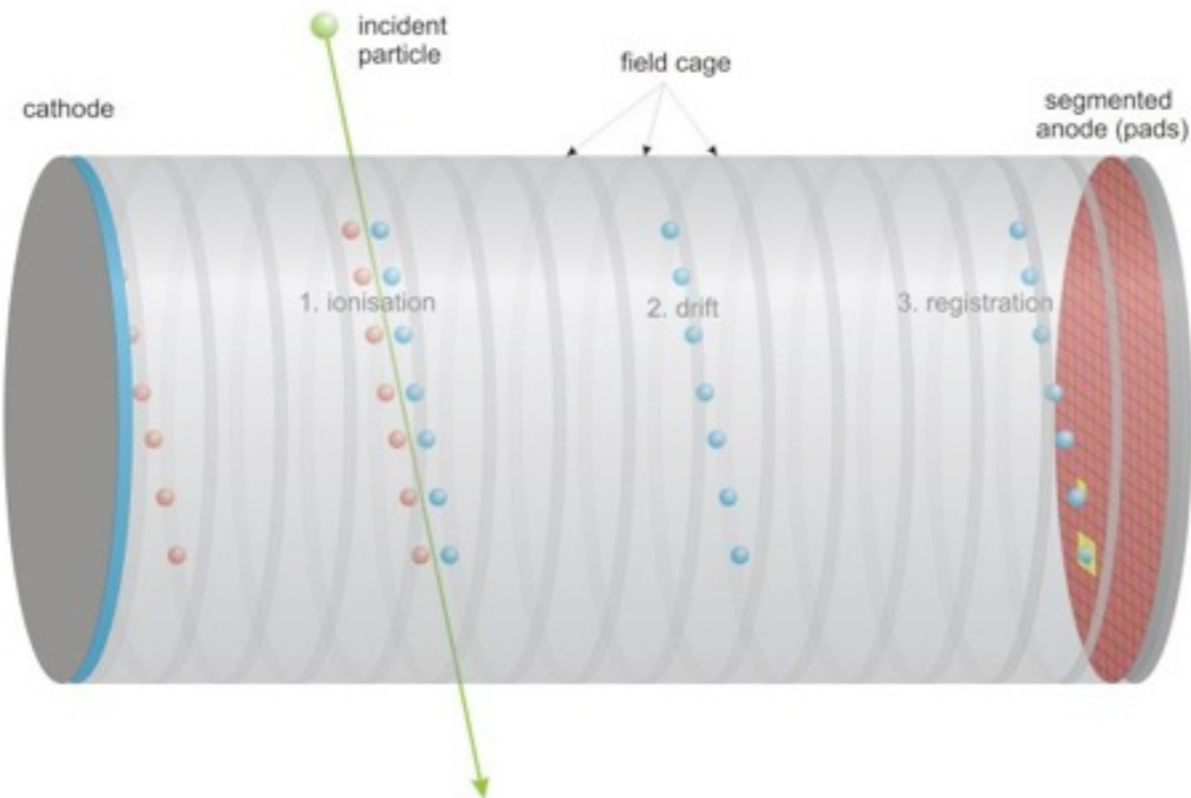




TPC- Time Projection Chamber: 3D

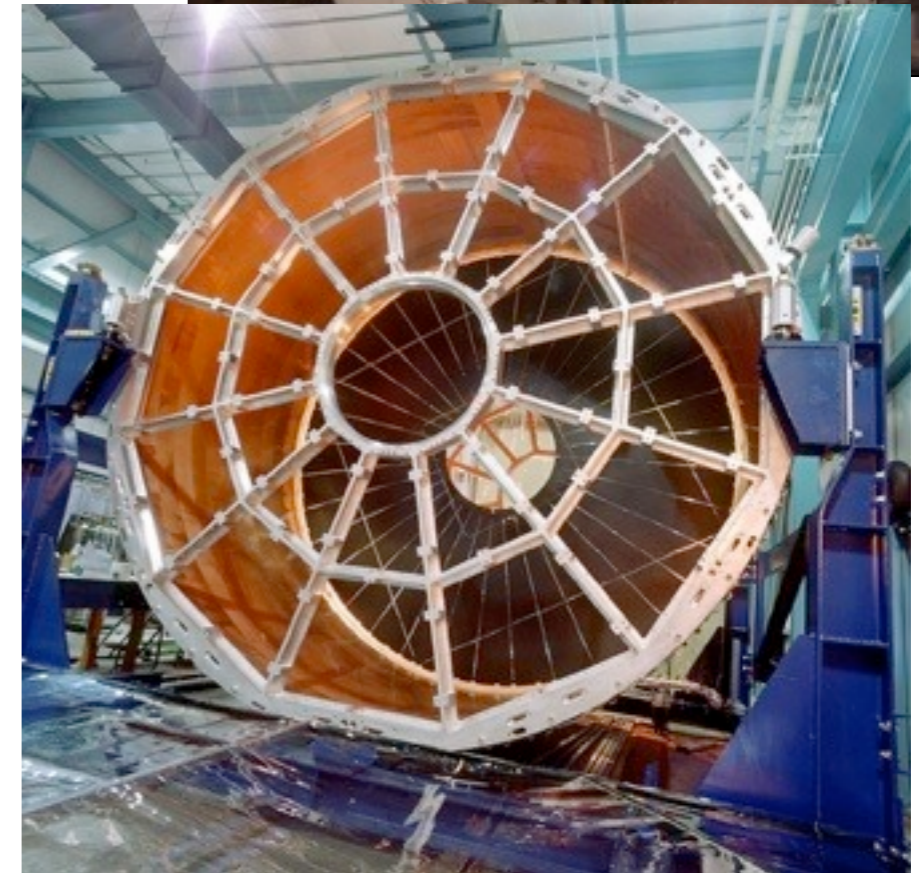
- Combination of the the 2D track information and the time results in a real 3D point

Pic: O. Schäfer



Pic: DESY

- Readout of the anode usually with multi-wire projection chambers
- Nowadays new developments for the readout: **Micro-Pattern Gas Detectors** i.e. gas electron multipliers (GEMs), micromegas, ...



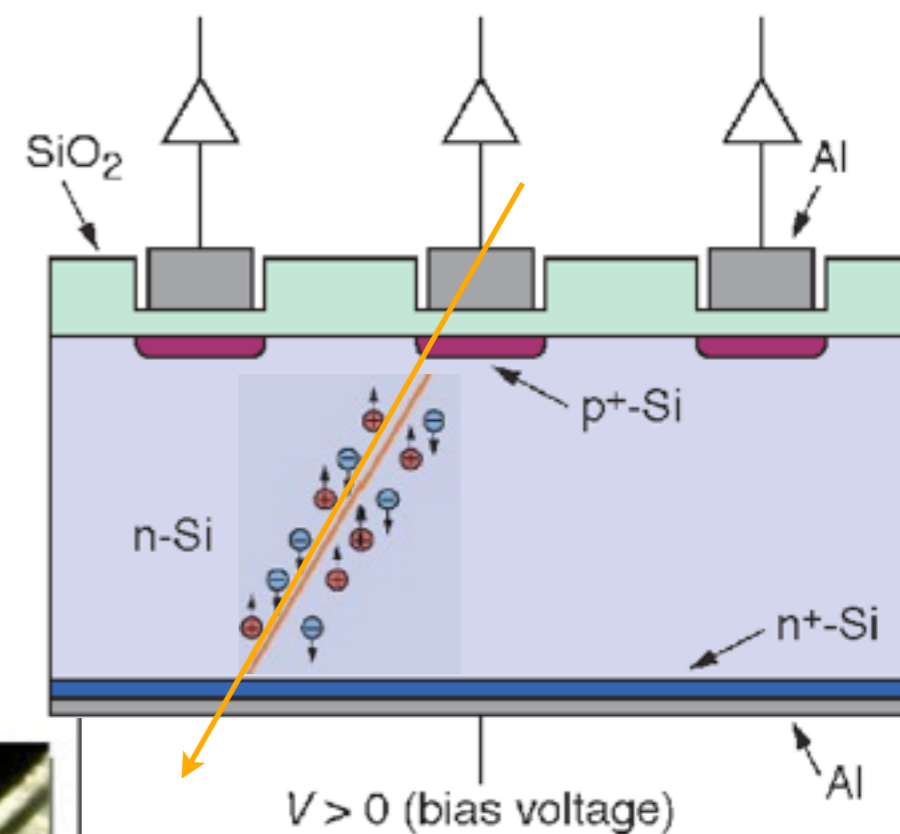
Pic: ALICE Collaboration



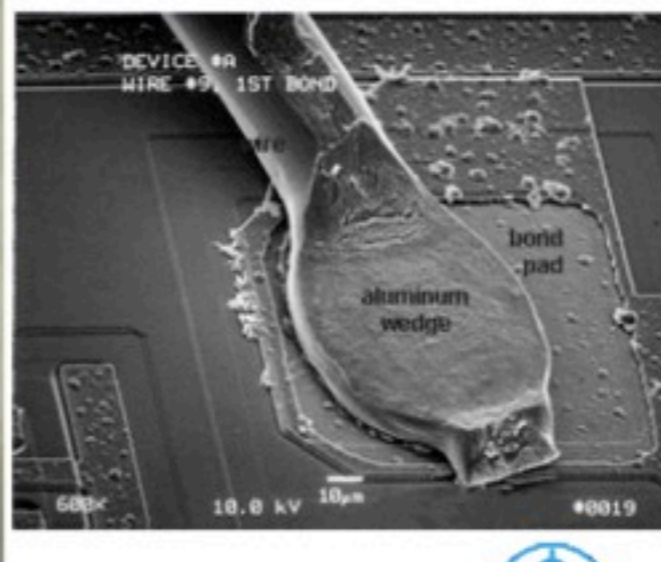
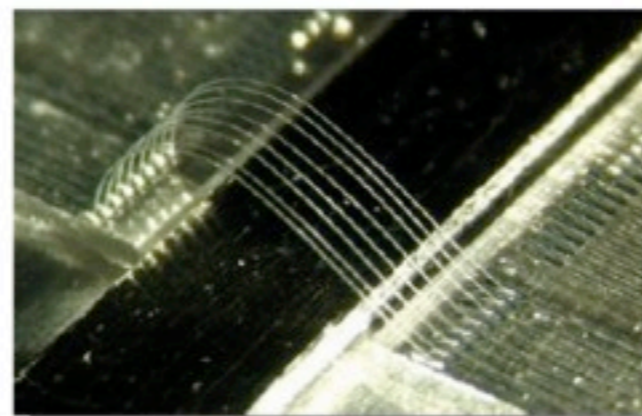
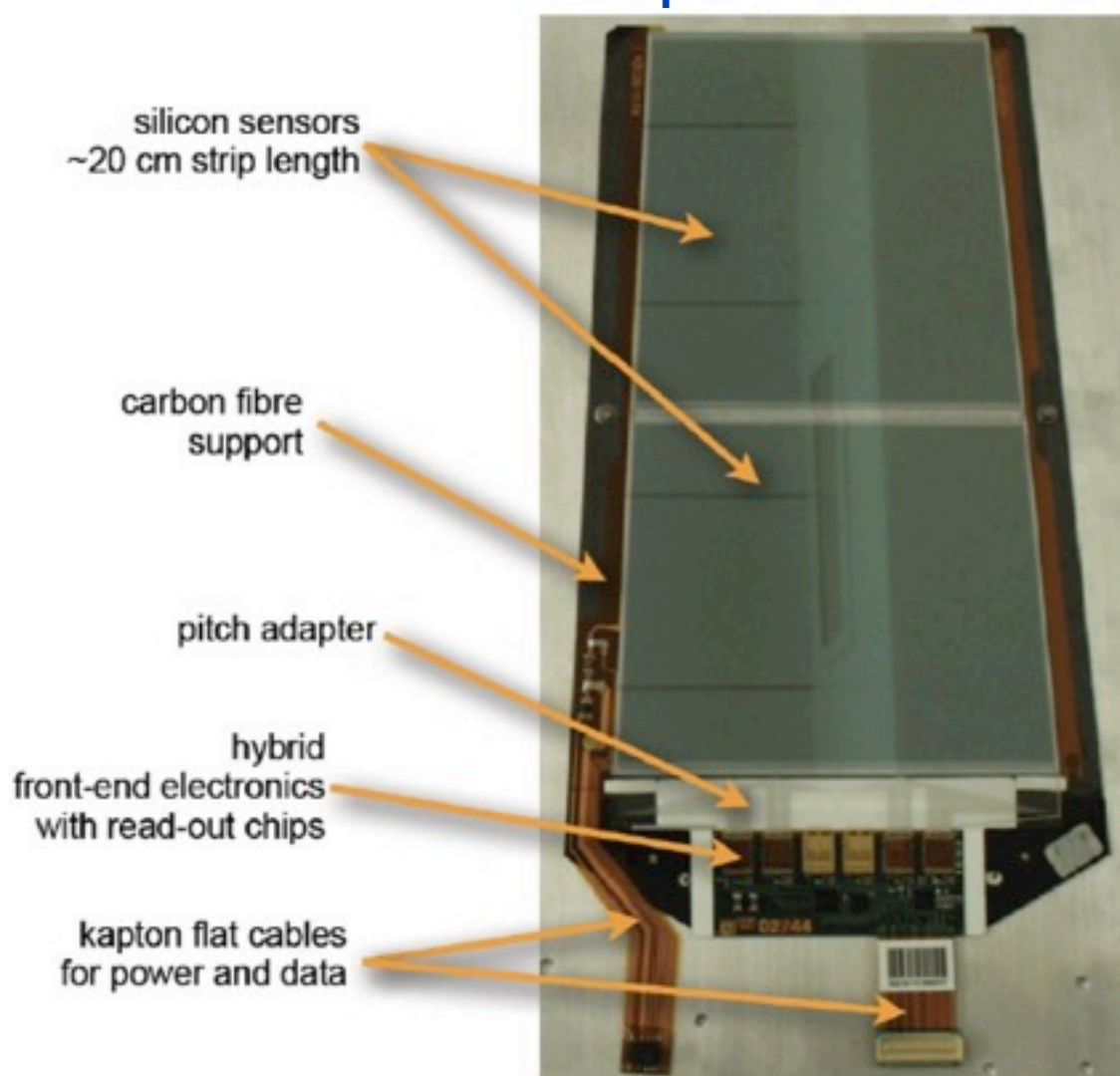
Track position: 1D silicon detectors

Charge carrier drift to electrodes and induce signal

By segmenting the implant we can reconstruct the position of the traversing particle in one dimension



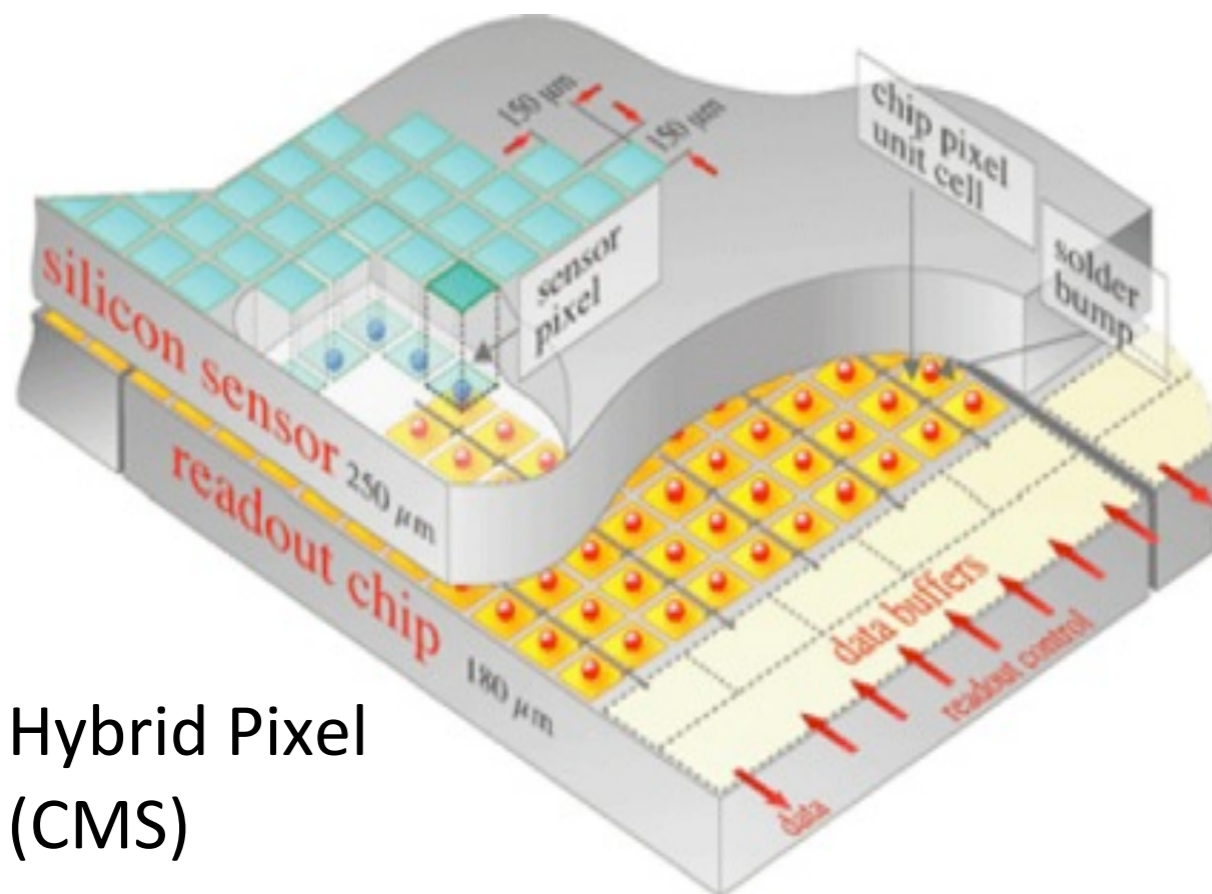
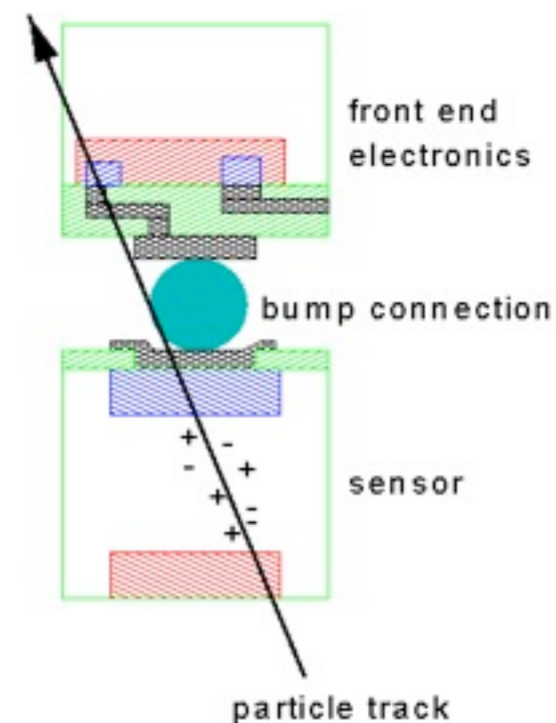
Strip module CMS





Track position: 2D silicon detectors

- 2D segmented implants = pixels
- Read-out chip mounted directly on top of the pixels (bump-bonding)
- Each pixel has its own read-out amplifier
- Fast read-out and radiation-tolerant
- **... but:**
- **Pixel area defined** by the size of the read-out chip
- **High material budget** and high power dissipation



Hybrid Pixel
(CMS)

- CMS Pixels: ~65 M channels
150 μm x 150 μm
- ATLAS Pixels: ~80 M channels
50 μm x 400 μm (long in z or r)
- Alice: 50 μm x 425 μm
- LHCb
- Phenix
-



Alternative designs

Planar Sensor

- current running horse for HEP experiments
- radiation hardness proven up to LHC requirements
- problem: HV might need to exceed 1000V

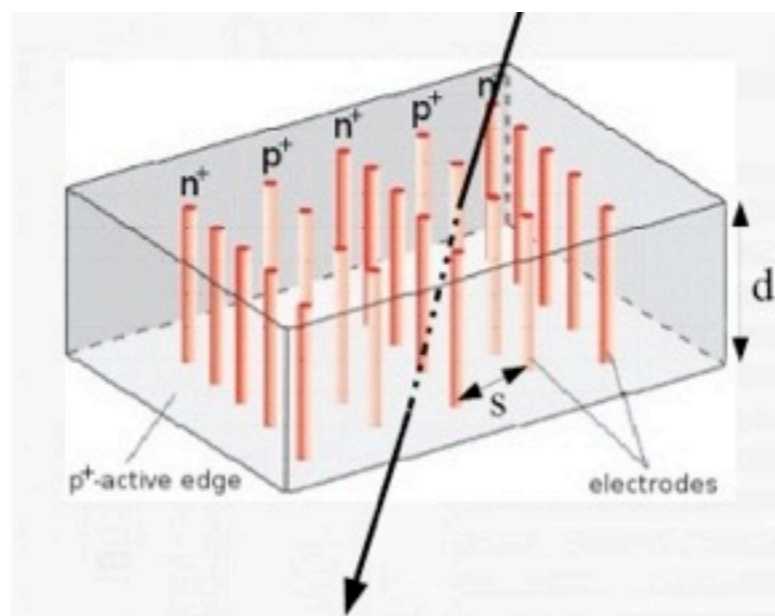
3D Silicon

- Both electrode types are processed inside the detector bulk instead of being implanted on the wafer's surface.
- Max. drift and depletion distance set by electrode spacing
- Reduced collection time and depletion voltage

CVD (Diamond)

- Large band gap 5.5 eV (no depletion zone required)
- Operation at room temperature possible
- Radiation hard material
- Drawback: 50% signal compared to silicon for same X_0 , but better S/N ratio (no dark current)

Very strong R&D efforts to develop sensors for future LHC/ILC applications!





Alternative semiconductors

- **Compound semiconductors**

two or more elements of the periodic table

- **important IV-IV compounds:**

- ***SiGe, SiC***

- **important III-V compounds:**

- ***GaAs***: Faster and probably more radiation resistant than Si. Drawback is less experience in industry and higher costs.

- GaP, GaSb, InP, InAs, InSb, InAlP

- **important II-VI compounds:**

- ***CdTe***: High atomic numbers (48+52) hence very efficient to detect photons

- ZnS, ZnSe, ZnTe, CdS, CdSe, Cd_{1-x}Zn_xTe, Cd_{1-x}Zn_xSe

	I	II	III	IV	V	VI	VII	VIII
1	1 H							2 He
2	3 Li	4 Be	5 B	6 C	7 N	8 O	9 F	10 Ne
3	11 Na	12 Mg	13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
4	19 K	20 Ca	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
5	37 Rb	38 Sr	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
6	55 Cs	56 Ba	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn
7	87 Fr	88 Ra	113 Uut	114 Uuq	114 Uup	115 Uuh	117 Uus	118 Uuo



Calorimetry: Overview

In nuclear and particle physics calorimetry refers to the detection of particles, and measurements of their properties, through **total absorption** in a block of matter: the **calorimeter**

Common feature of all calorimeters:
the measurement process is **destructive**
i.e. the particles are no longer available for
inspection after the calorimeter

The only exception concerns **muons**
→ mean for muon identification

Calorimetry works both for charged (**e^\pm** and **hadrons**) and neutral particles (**γ** , **n**) !

Principle:

- formation of **electromagnetic**
- or **hadronic** showers



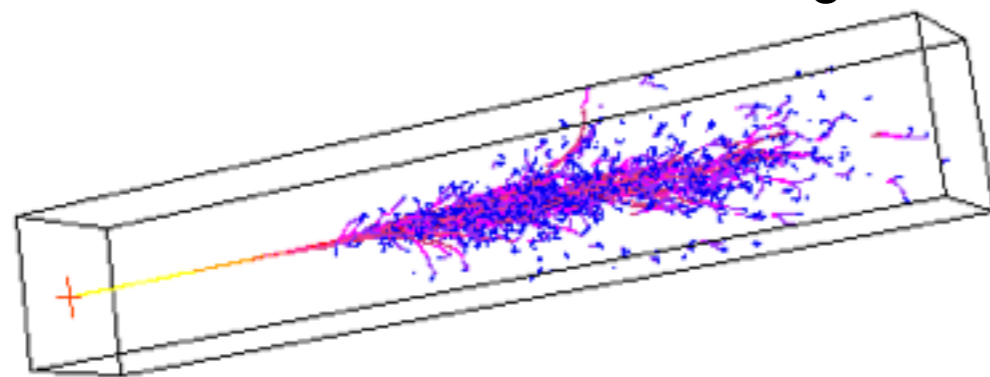
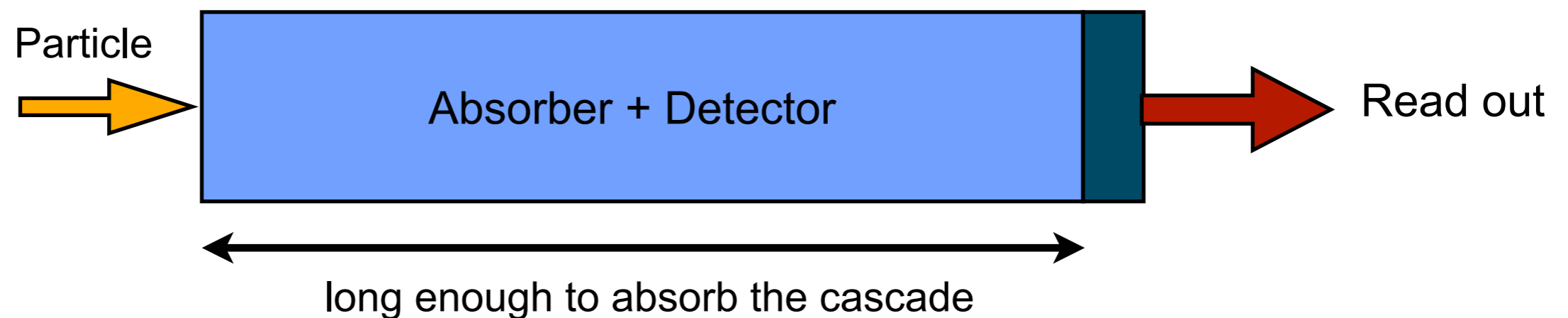


Calorimeter Types

Most commonly used: Homogeneous and Sampling Calorimeter

Homogeneous Calorimeter

- The absorber material is active; all deposited energy is converted into signal
- **Pro:** very good energy resolution
- **Contra:** segmentation difficult, selection of material is limited, expensive



Pic: Cornell

Example: CMS electromagnetic calorimeter

design not suitable for hadronic calorimeters

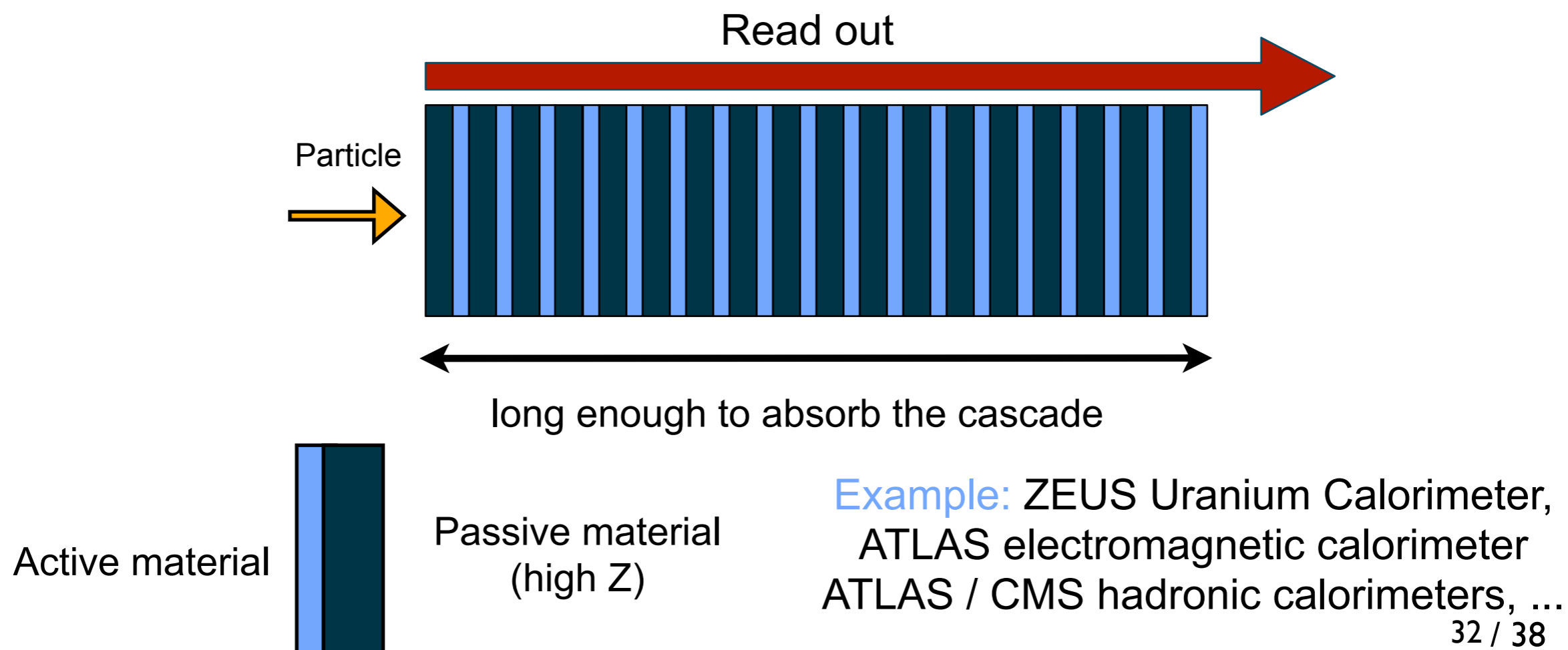


Calorimeter Types

Most commonly used: Homogeneous and Sampling Calorimeter

Sampling Calorimeter

- A structure of passive and active material; a fraction (**Sampling Fraction, f_s**) of the deposited energy is detected (1-5%)
- **Pro:** Segmentation, compact detectors by the usage of dense materials (W, U)
- **Contra:** Energy resolution is limited by fluctuations

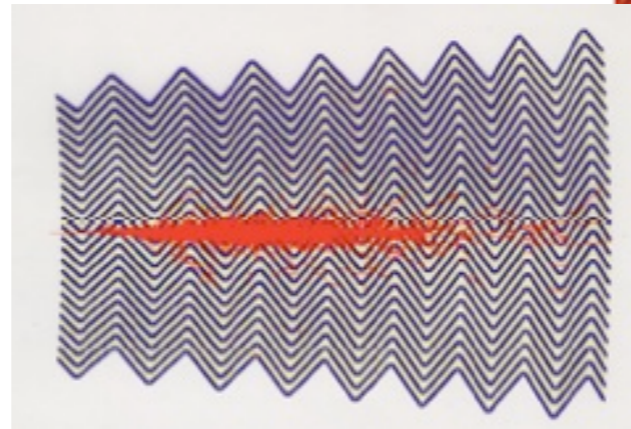
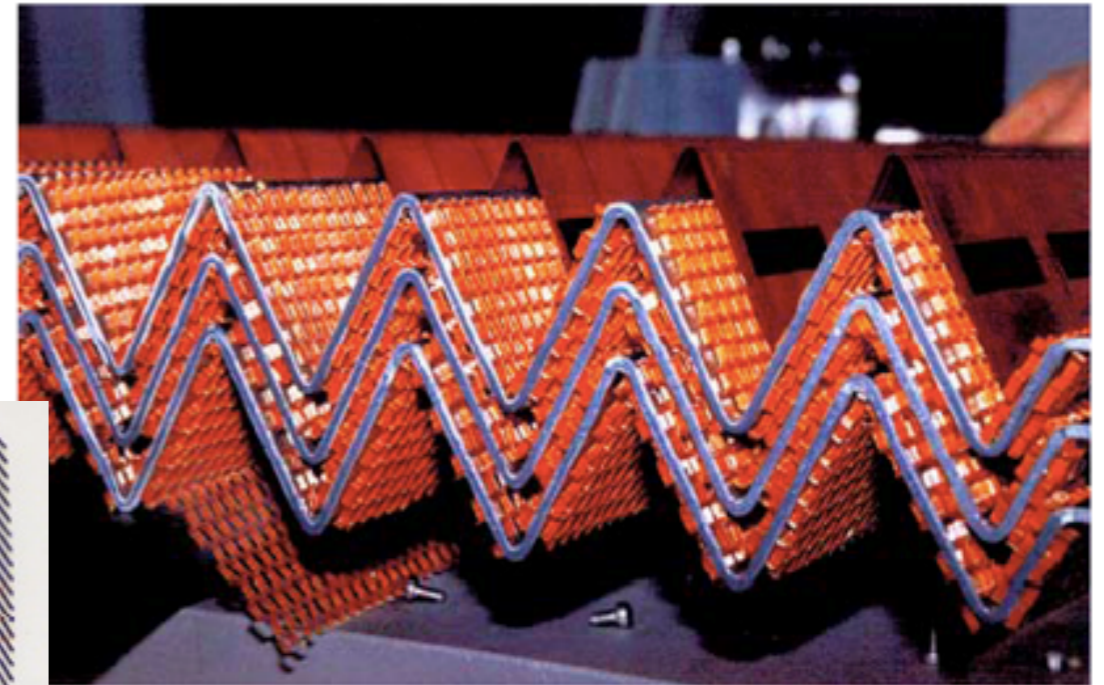




Examples of electromagnetic calorimeters

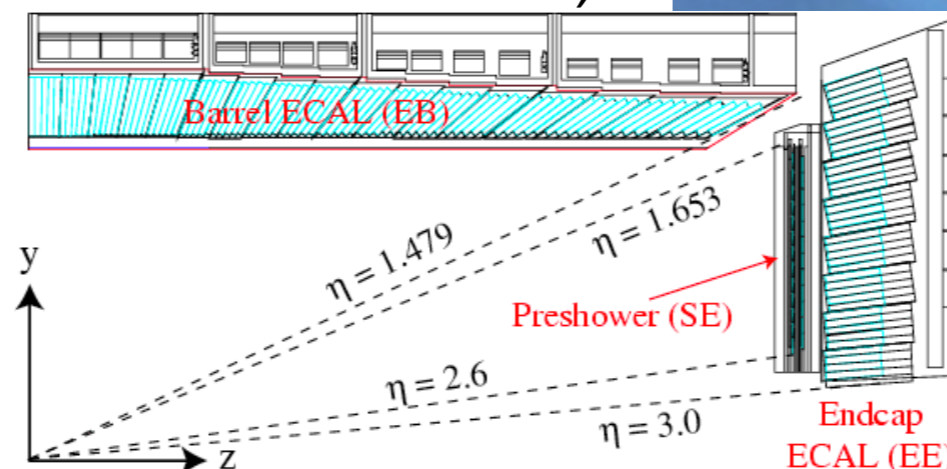
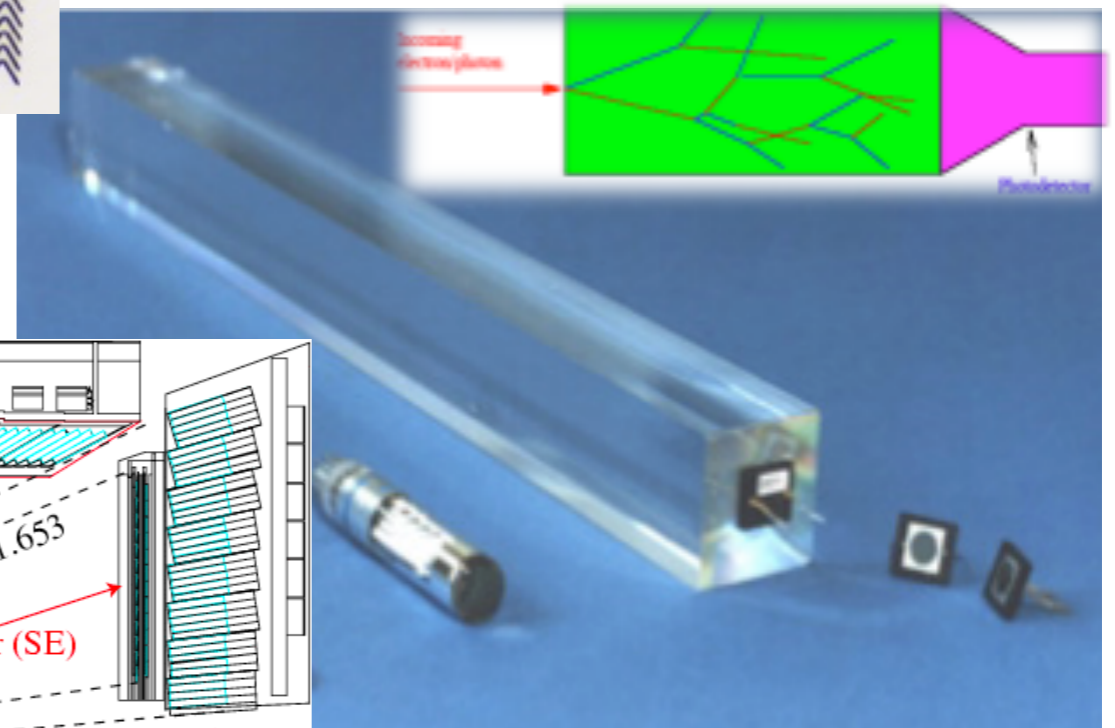
ATLAS EM barrel calorimeter

- Honeycomb spacers position the electrodes between the lead absorber plates
- Liquid Argon at 90°K flows through.
- Radiation resistant, no cracks in η
- Accordion structure with Pb-LAr sampling



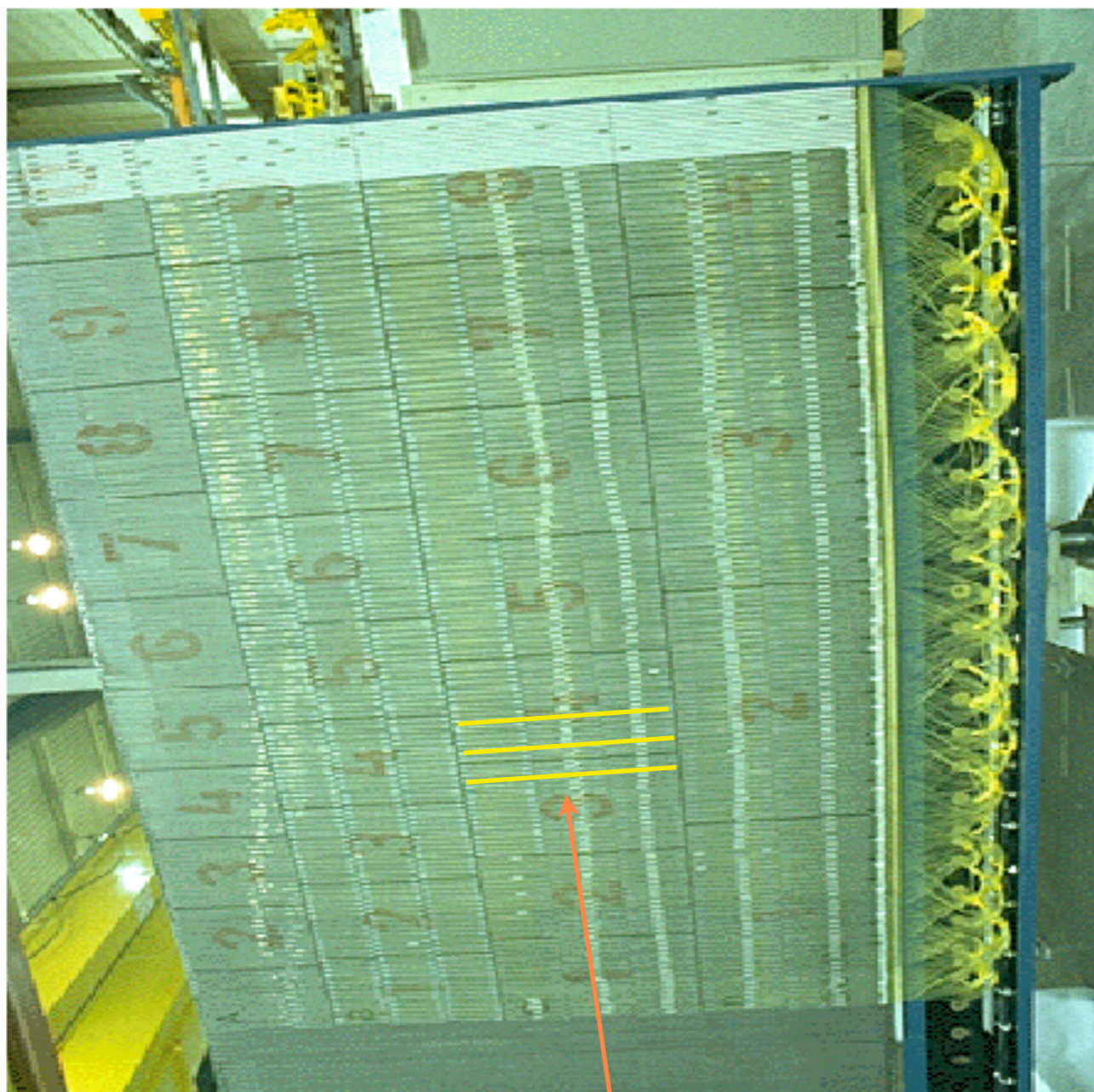
CMS EM barrel calorimeter

- PbWO₄ crystals (230x22x22 mm³)
- Read out by APD (Avalanche PhotoDiodes)
- Homogeneous

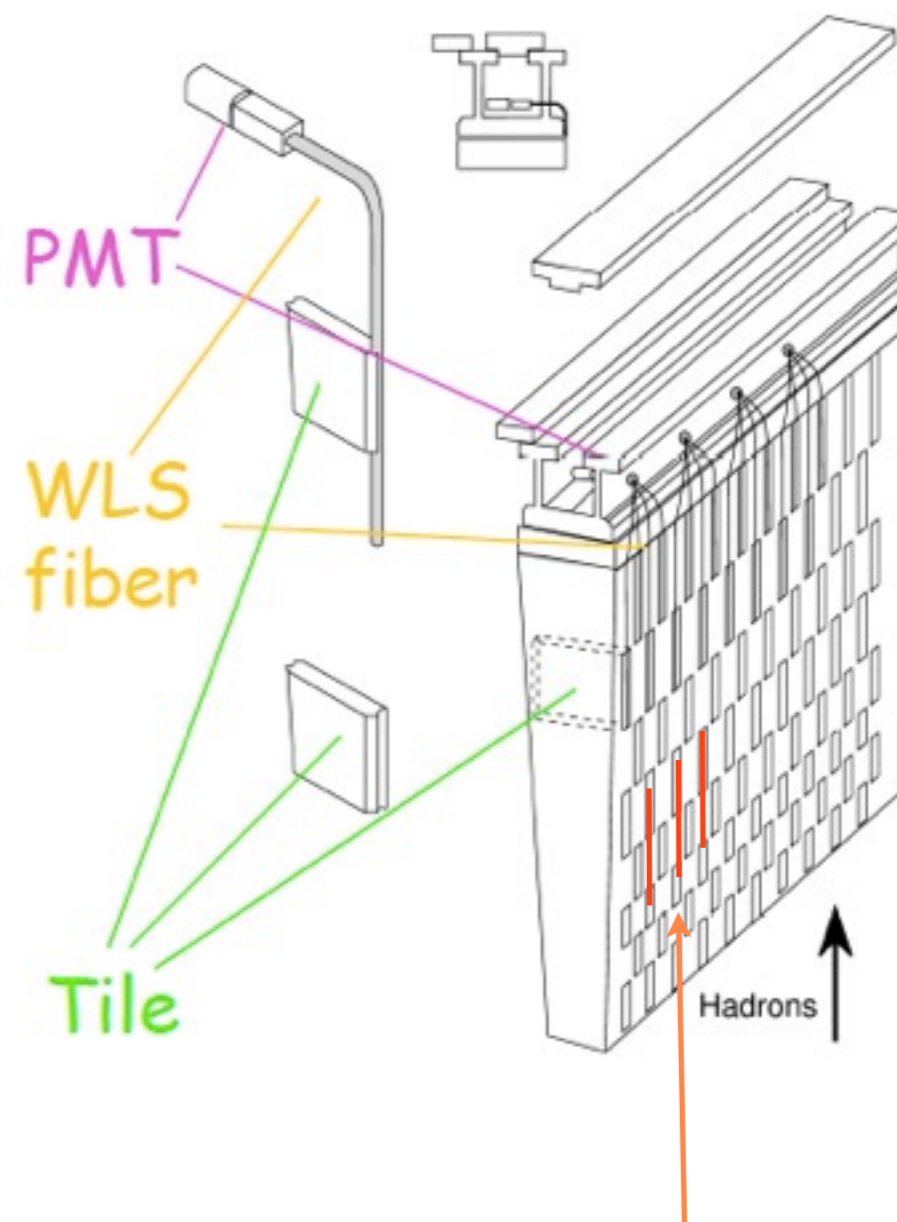




Examples of hadronic calorimeters



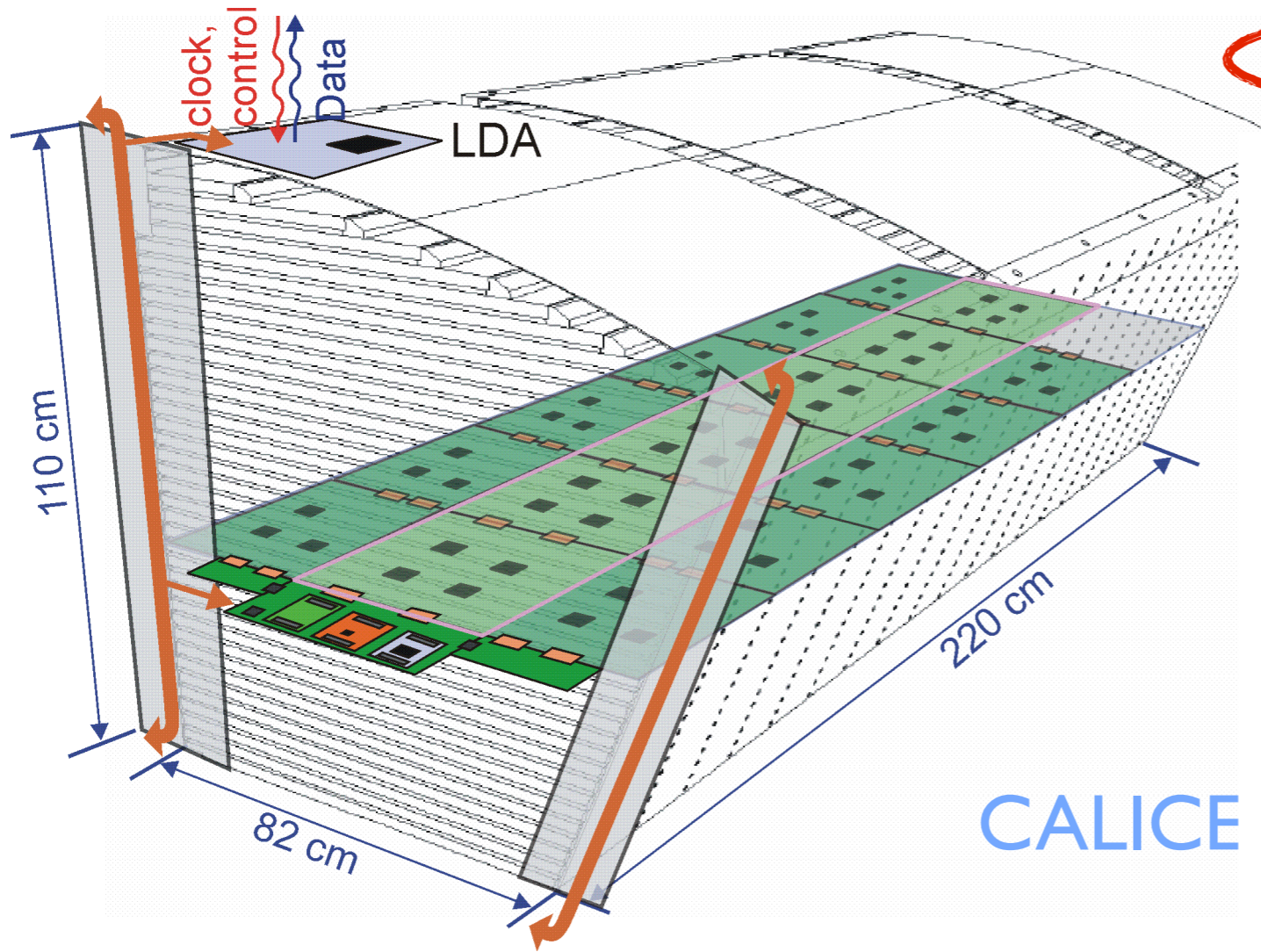
CMS: Brass/scintillator longitudinal orientation



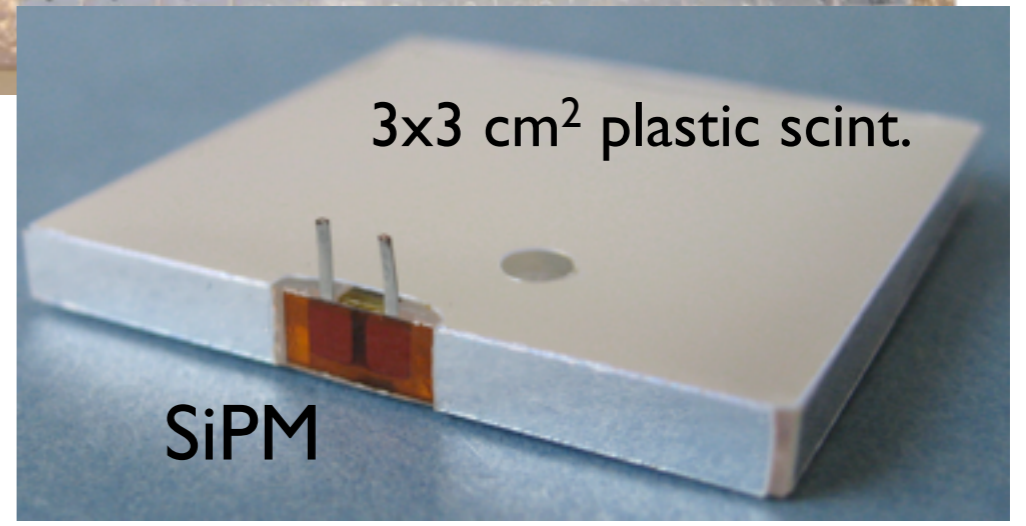
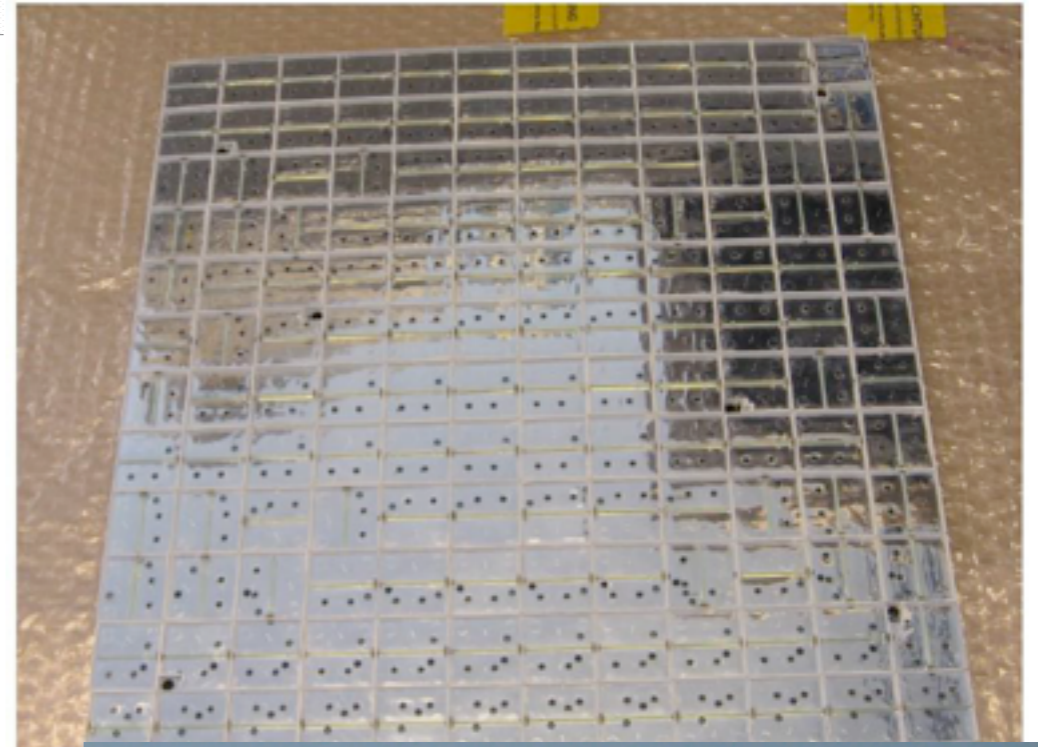
ATLAS: Fe/scintillator vertical orientation



New Concepts: highly-integrated granular calorimeters



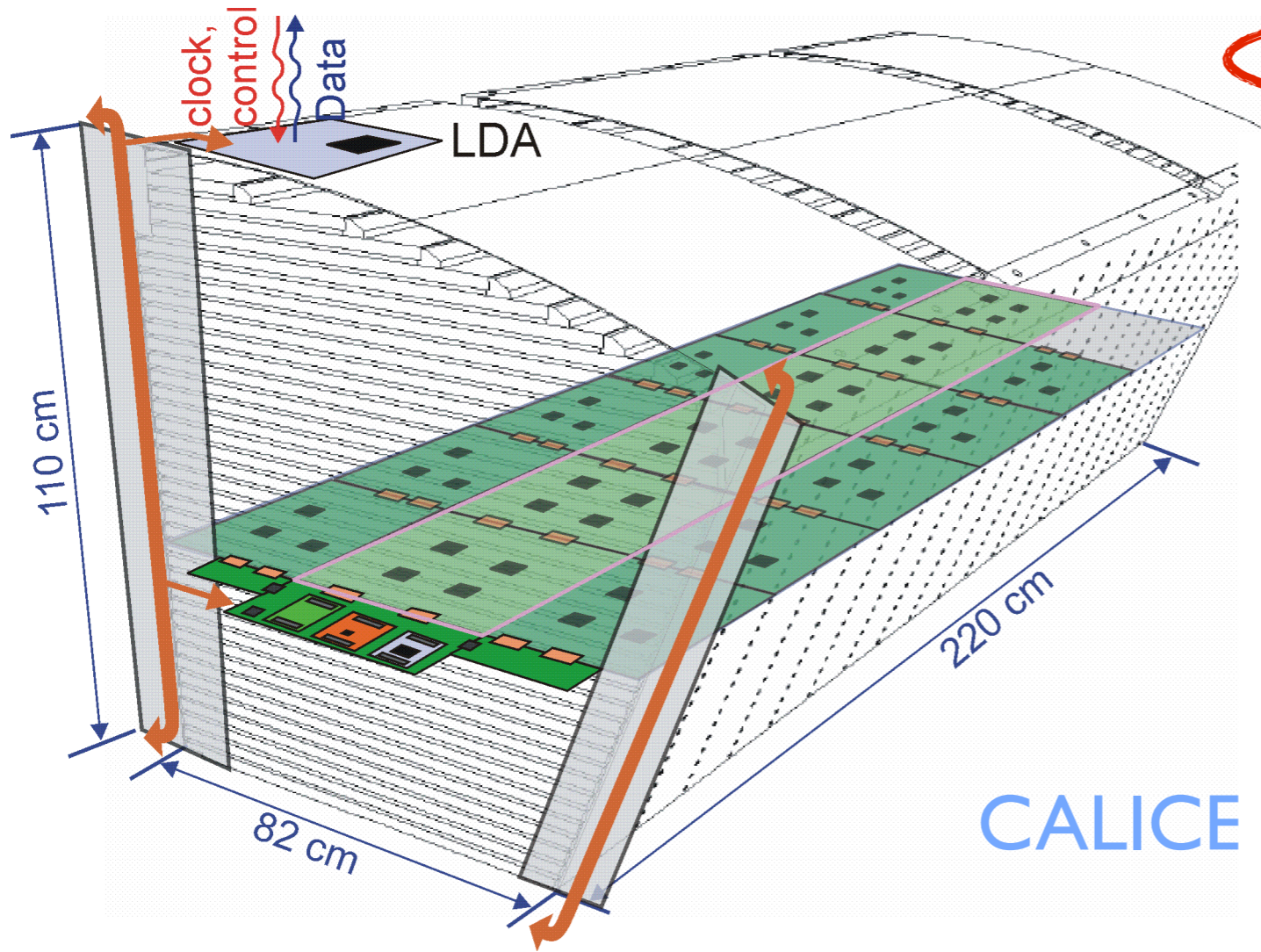
~ 1,000 calorimeter cells / m²



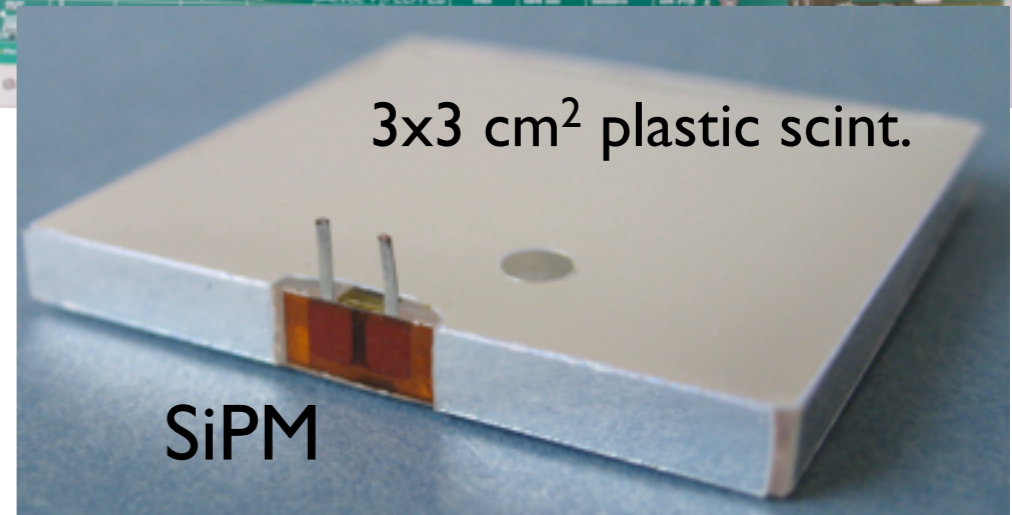
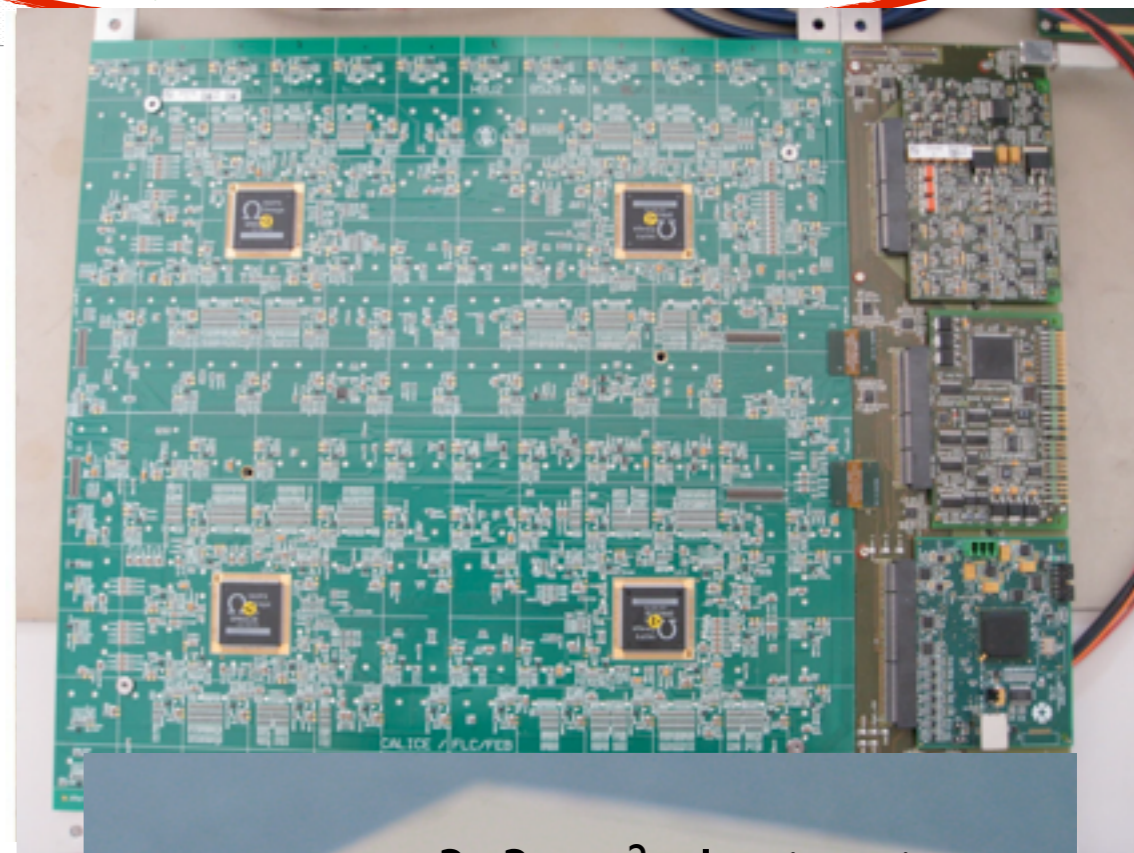
- sampling calorimeter Fe / plastic scintillator
- ~ 48 layers ~ 6 λ
- Front end electronics integrated in active layer



New Concepts: highly-integrated granular calorimeters



~ 1,000 calorimeter cells / m²

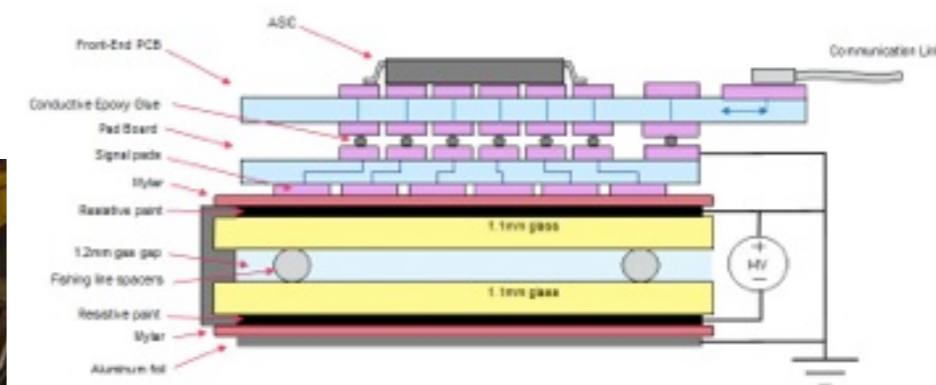
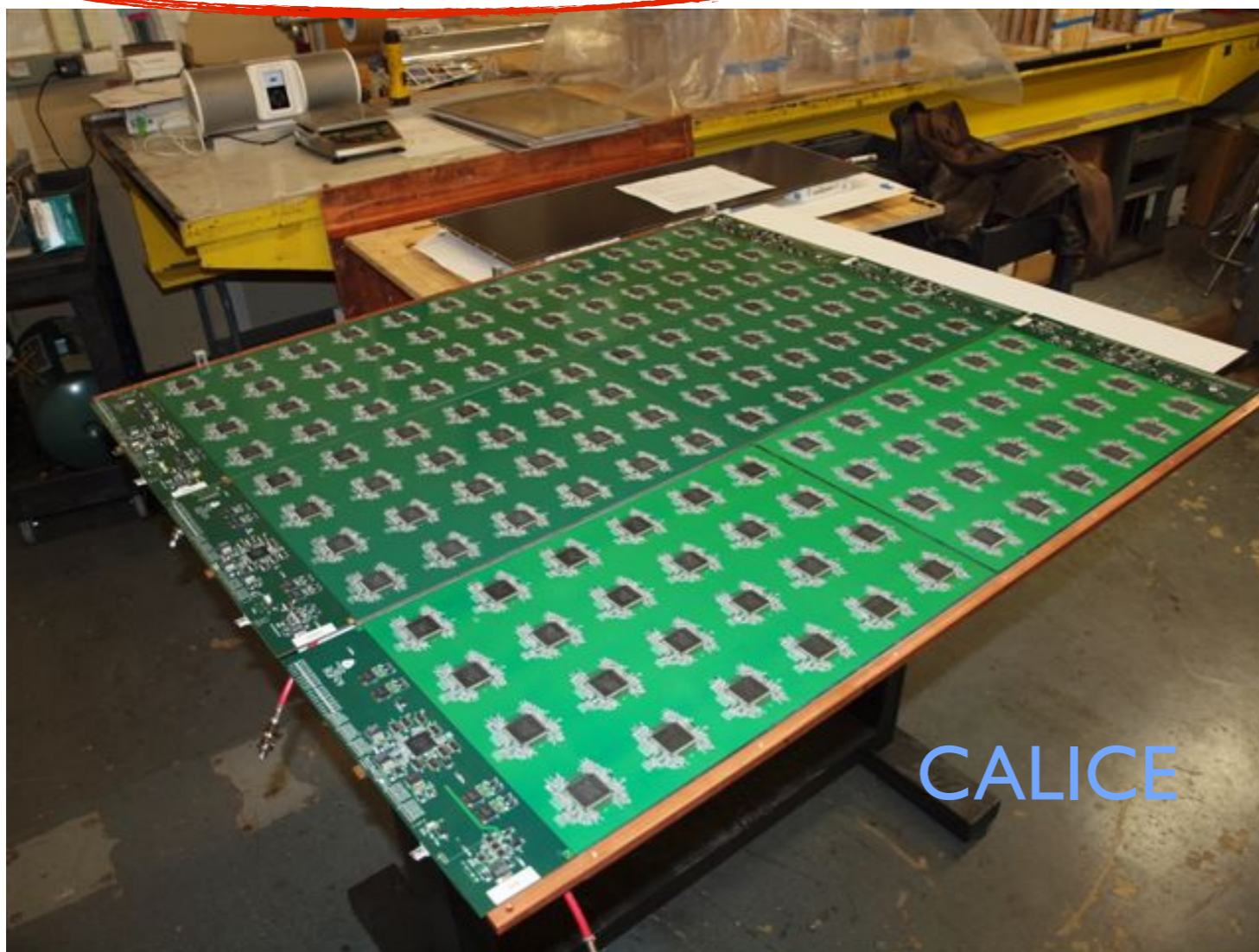


- sampling calorimeter Fe / plastic scintillator
- ~ 48 layers ~ 6 λ
- Front end electronics integrated in active layer



New Concepts: highly-integrated granular calorimeters

~ 10,000 calorimeter cells / m²



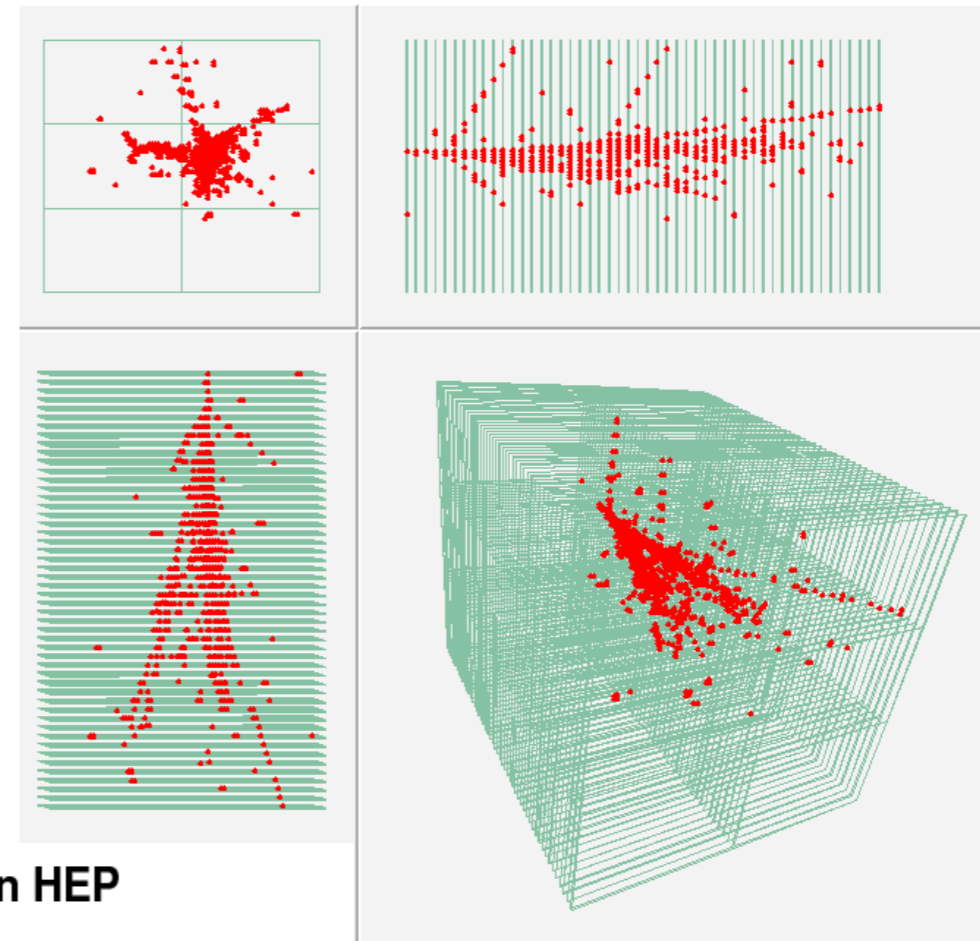
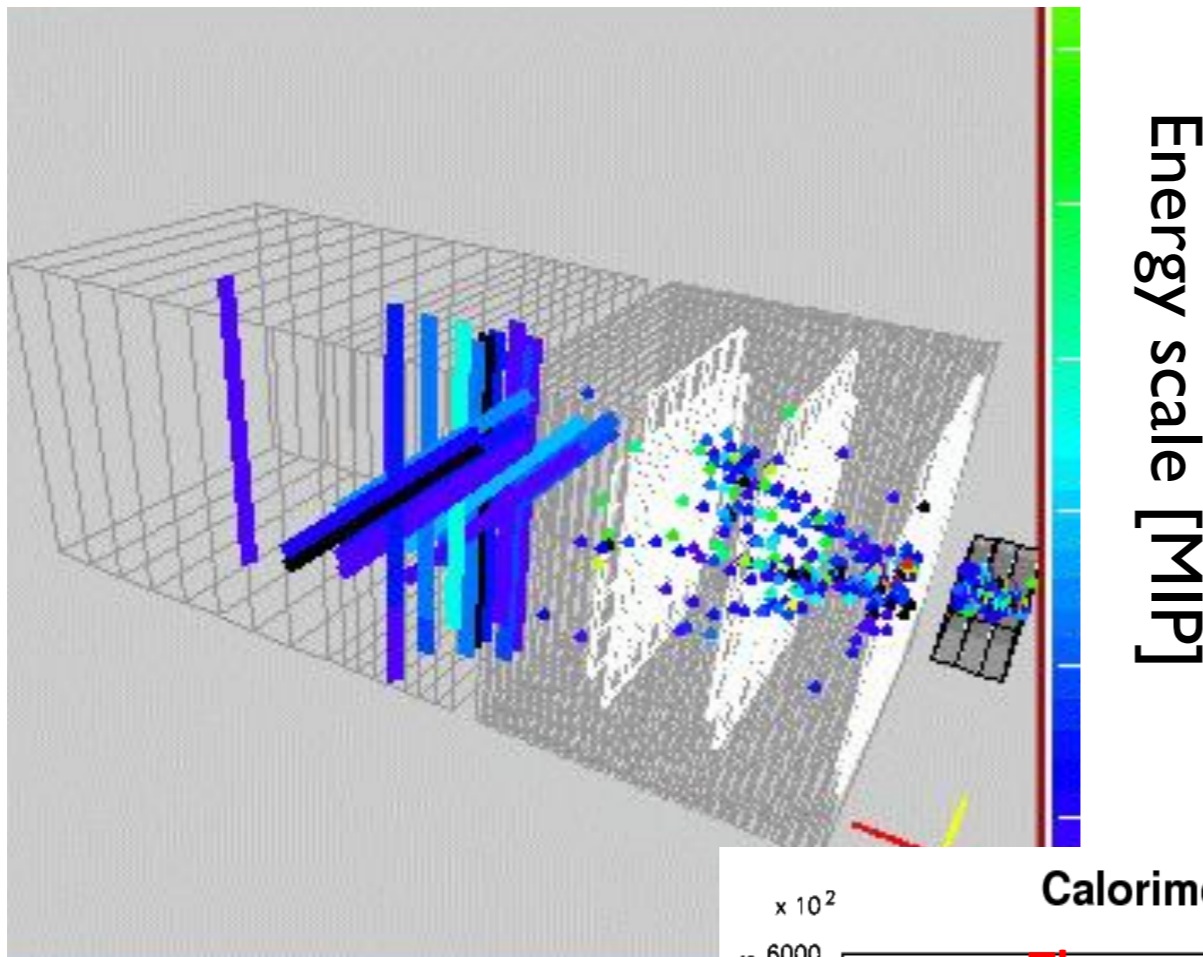
- sampling calorimeter Fe / gas layer (RPC)
~ 48 layers ~ 6 λ
- Front end electronics integrated in active layer



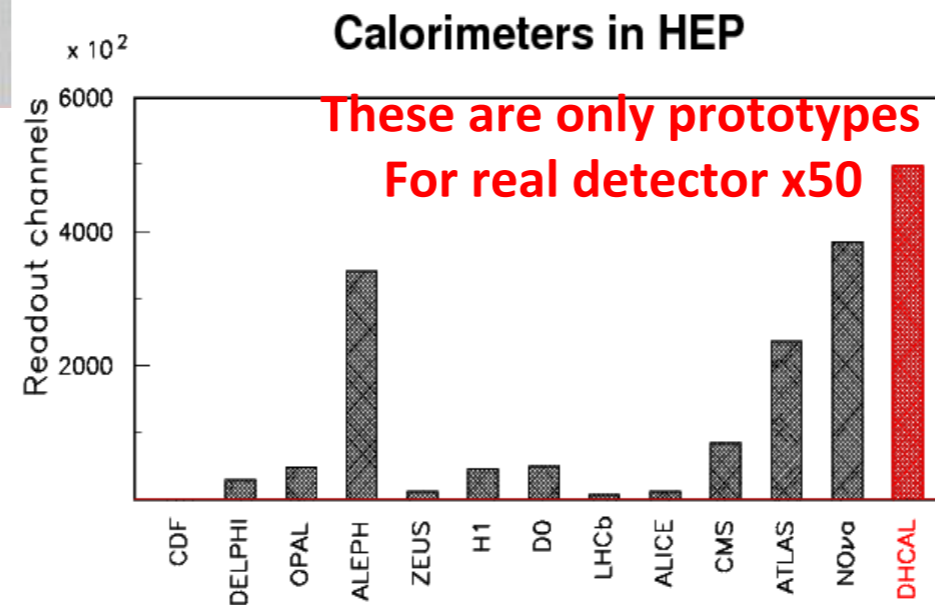
Single pion reconstruction with high granularity

Fe/scintillator with analog readout

Fe/gas with digital readout



CALICE





Summary & conclusions

- Stable particle detection = **measurement of** particle **E and p**
- Main interaction of high energetic particles with matter
 - heavy charged particles: **ionization** (Bethe-Block)
 - electrons (> 10 MeV): **Bremsstrahlung**
 - photons (> 1 MeV): **pair-production**
 - neutrons (> 20 MeV): **elastic and inelastic n-p scattering**
- p meas. : segmented gaseous or semiconductor detectors
- E meas. : calorimeters (homogeneous / inhomogeneous)