Deformations of shell Λ hypernuclei with AMD

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Grand challenges of hypernuclear physics

Interaction: To understand baryon-baryon interaction

- 2 body interaction between baryons (nucleon, hyperon)
 - hyperon-nucleon (YN)
 - hyperon-hyperon (YY)

Structure: To understand many-body system of nucleons and hyperon

- Addition of hyperon(s) shows us new features of nuclear structure Ex.) Structure change by hyperon(s)
 - No Pauli exclusion between N and Y
 "Hyperon as an impurity in nuclei"
 - YN interaction is different from NN



Today's talk: deformation of Λ hypernuclei

Recent achievements in (hyper)nuclear physics

<u>Knowledge of ΛN effective interaction</u>

- Study of light (s, p-shell) Λ hypernuclei
 - Accurate solution of few-body problems ^[1]
 - ΛN G-matrix effective interactions ^[2]
 - Increases of experimental information ^[3]

Development of theoretical models

Through the study of unstable nuclei

Ex.: Antisymmetrized Molecular Dynamics (AMD)^[4]

- AMD can describe dynamical changes of various structure
- No assumption on clustering and deformation

Recent developments enable us to study structure of Λ hypernuclei

[1] E. Hiyama, NPA **805** (2008), 190c, [2] Y. Yamamoto, *et al.*, PTP Suppl. **117** (1994), 361., [3] O. Hashimoto and H. Tamura, PPNP **57** (2006), 564., [4] Y. Kanada-En'yo *et al.*, PTP **93** (1995), 115.



Toward heavier and exotic Λ hypernuclei

Experiments at J-PARC, JLab and Mainz etc.

• Hypernuclear chart will be extended to heavier regions "Structure of hypernuclei"



Superdeformation



Superdeformeation

- With 1:2 axis ratio of nuclear deformation
- Observed in several sd-pf shell nuclei

Example: Superdeformed (SD) states

 Ex.) ⁴⁰Ca
 J. R. MacDonald, et al., PRC3, 219(1971), E. Ideguchi, et al., PRL87, 222501(2011)

 W. Gerace and A. Green, NPA93, 110(1967); NPA123, 241 (1969)

• Observed in ⁴⁰Ca, while the ground state is spherical



If a Λ particle is added, what's happen?

AMD calculation for ⁴⁰Ca: basically follows Y. Taniguchi, et al., PRC 76, 044317 (2007)

B_{Λ} with different deformations(structures)

- Λ binging energy (B_{Λ}) in light Λ hypernuclei with cluster structure
- Λ coupled to the compact state is more deeply bound



How about difference of B_{Λ} in superdeformed states?

Superdeformation



Purposes: to reveal difference of B_{Λ} and deformation change by Λ Our method: antiysymmetrized molecular dynamics (AMD)

We extended the AMD to hypernuclei

HyperAMD (Antisymmetrized Molecular Dynamics for hypernuclei)

Hamiltonian

$$\hat{H} = \hat{T}_{N} + \hat{V}_{NN} + \hat{T}_{\Lambda} + \hat{V}_{\Lambda N} - \hat{T}_{g}$$

NN: Gogny D1S Λ N: YNG-ESC08c

Wave function

- Nucleon part: Slater determinant Spatial part of single particle w.f. is described as Gaussian packet
- Single-particle w.f. of Λ hyperon: Superposition of Gaussian packets
- Total w.f.:

$$\psi(\vec{r}) = \sum_{m} c_{m} \varphi_{m}(r_{\Lambda}) \otimes \frac{1}{\sqrt{A!}} \det[\varphi_{i}(\vec{r}_{j})]$$

$$\varphi_{N}(\vec{r}) = \frac{1}{\sqrt{A!}} \det[\varphi_{i}(\vec{r}_{j})]$$

$$\varphi_{i}(r) \propto \exp\left[-\sum_{\sigma=x,y,z} v_{\sigma}(r-Z_{i})_{\sigma}^{2}\right] \chi_{i}\eta_{i} \quad \chi_{i} = \alpha_{i}\chi_{\uparrow} + \beta_{i}\chi_{\downarrow}$$

$$\varphi_{\Lambda}(r) = \sum_{\sigma=x,y,z} c_{m}\varphi_{m}(r)$$

$$\varphi_{m}(r) \propto \exp\left[-\sum_{\sigma=x,y,z} \mu v_{\sigma}(r-z_{m})_{\sigma}^{2}\right] \chi_{m} \quad \chi_{m} = a_{m}\chi_{\uparrow} + b_{m}\chi_{\downarrow}$$

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Theoretical Framework: HyperAMD M. Isaka, et al., PRC83(2011) 044323 M. Isaka, et al., PRC83(2011) 054304

Procedure of the calculation

Variational Calculation $\frac{dX_i}{dt} = \frac{\kappa}{\hbar} \frac{\partial H^{\pm}}{\partial X_i^*}$ $\kappa < 0$ • Imaginary time development method $\frac{dX_i}{dt} = \frac{\kappa}{\hbar} \frac{\partial H^{\pm}}{\partial X_i^*}$ $\kappa < 0$ • Variational parameters: $X_i = Z_i, z_i, \alpha_i, \beta_i, a_i, b_i, v_i, c_i$



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Angular Momentum Projection

$$\left|\Phi_{K}^{s};JM\right\rangle = \int d\Omega D_{MK}^{J^{*}}(\Omega) R(\Omega) \Phi^{s+}$$

Generator Coordinate Method(GCM)

•Superposition of the w.f. with different configuration •Diagonalization of $H^{J\pm}_{sK,s'K'}$ and $N^{J\pm}_{sK,s'K'}$

$$H_{sK,s'K'}^{J\pm} = \left\langle \Phi_{K}^{s}; J^{\pm}M \left| \hat{H} \right| \Phi_{K'}^{s'}; J^{\pm}M \right\rangle \qquad \left| \Psi^{J\pm M} \right\rangle = \sum_{sK} g_{sK} \left| \Phi_{K}^{s}; J^{\pm}M \right\rangle$$
$$N_{sK,s'K'}^{J\pm} = \left\langle \Phi_{K}^{s}; J^{\pm}M \left| \Phi_{K'}^{s'}; J^{\pm}M \right\rangle \qquad \left| \Psi^{J\pm M} \right\rangle = \sum_{sK} g_{sK} \left| \Phi_{K}^{s}; J^{\pm}M \right\rangle$$

Results and Discussions



- 1. Energy curves and Λ single particle energy
- 2. B_{Λ} and excitation spectra

M.I., K. Fukukawa, M. Kimura, E. Hiyama, H. Sagawa, and Y. Yamamoto, PRC89, 024310 (2014)

ND and SD states of ⁴⁰Ca

• Ground, normal deformed and superdeformed states are obtained



Energy curves of ${}^{41}{}_{\Lambda}\text{Ca}$ as a function of β

- $^{41}_{\Lambda} Ca$ "GS⊗Λ", "ND⊗Λ" and "SD⊗Λ" curves are obtained → SD states will appear in $^{41}_{\Lambda} Ca$
 - Energy (local) minima are almost unchanged



Λ single particle energy

- Definition: $\epsilon_{\Lambda}(\beta) = E(^{46}_{\Lambda}Sc)(\beta) E(^{45}Sc)(\beta)$
- \bullet General trend: ϵ_{Λ} changes within 1 2 MeV as β increases



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- General trend: ϵ_{Λ} changes within 1 2 MeV as β increases

➤ Similar to the p shell Λ hypernuclei





Excitation spectra of ${}^{46}{}_{\Lambda}Sc$ and ${}^{48}{}_{\Lambda}Sc$

- \bullet Difference of \mathbf{B}_{Λ} leads to the energy shift up of the deformed states
- Similar phenomena in ${}^{48}{}_{\Lambda}$ Sc



We hope these states in Sc Λ hypernuclei are observed at JLab

Results and Discussions



M. I., M. Kimura, A. Dote and A. Ohnishi, PRC 85 (2012), 034303.

Triaxial deformation

Ex.) ²⁴Mg Candidate of triaxial deformed nuclei

- Low lying $K^{\pi}=2^+$ band: a sign of triaxial deformation
- $K^{\pi} = 0^+$ and $K^{\pi} = 2^+$ bands have almost the same (triaxial) deformation



Results: Excitation spectra of ²⁵ Mg

 $K^{\pi}=2^+\otimes \Lambda$ band is shifted up by about 200 keV due to the difference of B_{Λ} between the $K^{\pi}=0^+$ and $K^{\pi}=2^+$ bands 12 α threshold (Theor.) α threshold (Exp.) 10 10 6_{2} Excitation Energy (MeV) $9/2^+_3, 11/2^+_2$ 6⁺ $7/2^+_{3}, 9/2^+_{2}$ $13/2_1^+$ 6_{1}^{+} $5/2^+_{3}, 7/2^+_{2}$ $11/2_1^+$ $3/2^+_{2,5}/2^+_{2}$ 2^{+}_{2} 5 5 $\mathbf{K}^{\pi} = 2^{+} \otimes \Lambda \mathbf{s}$ 200keV $K^{\pi} = 2^{+}$ $7/2_1^+, 9/2_1^+$ $K^{\pi} = 2^{+}$ $3/2_1^+, 5/2_1^+ - 1/2_1^+$ 0 $K^{\pi} = 0^+ \otimes \Lambda s$ $K^{\pi} = 0^{+}$ $K^{\pi} = 0^{+}$ 24 Mg (Exp.) 24 Mg (AMD) $^{25}_{\Lambda}$ Mg (HyperAMD)

Two bands: Almost the same (triaxial) deformation \longrightarrow Why shifted up?

Results: Reasons for the shift up

- Difference of the Λ binding energy B_{Λ}
- B_{Λ} in the $K^{\pi}=0^+\otimes \Lambda$ band is larger than that in the $K^{\pi}=2^+\otimes \Lambda$ band
- Difference of B_{Λ} comes from the difference of deformation change between the two bands



We hope the shift up of the $K^{\pi}=2^+$ band is observed at J-PARC

Summary

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• AMD + GCM was used to study deformations of *sd-pf* shell Λ hypernuclei

Λ added to SD states: ${}^{\mathbf{41}}{}_{\Lambda}\mathbf{Ca}$, ${}_{\Lambda}\mathbf{Sc}$

- Prediction of SD states in Λ hypernuclei: ${}^{41}_{\Lambda}$ Ca and ${}^{46}_{\Lambda}$ Sc
 - Deformation is almost unchanged by Λ in Ca and Sc hypernuclei
 - B_{Λ} is different depending on deformations: smaller in SD states
- Λ added to triaxial deformation: $^{\rm 25}{}_{\Lambda}{\rm Mg}$
- B_{Λ} is different between the $K^{\pi}=0^+$ and $K^{\pi}=2^+$ bands, while they have almost the same (triaxial) deformation
 - This is due to the difference of deformation change

Future plan

- To predict the production cross sections
- Comparison of B_{Λ} with cluster states: ${}^{13}{}_{\Lambda}C$ (Hoyle + Λ)