



# Measurements of $\gamma$ from LHCb

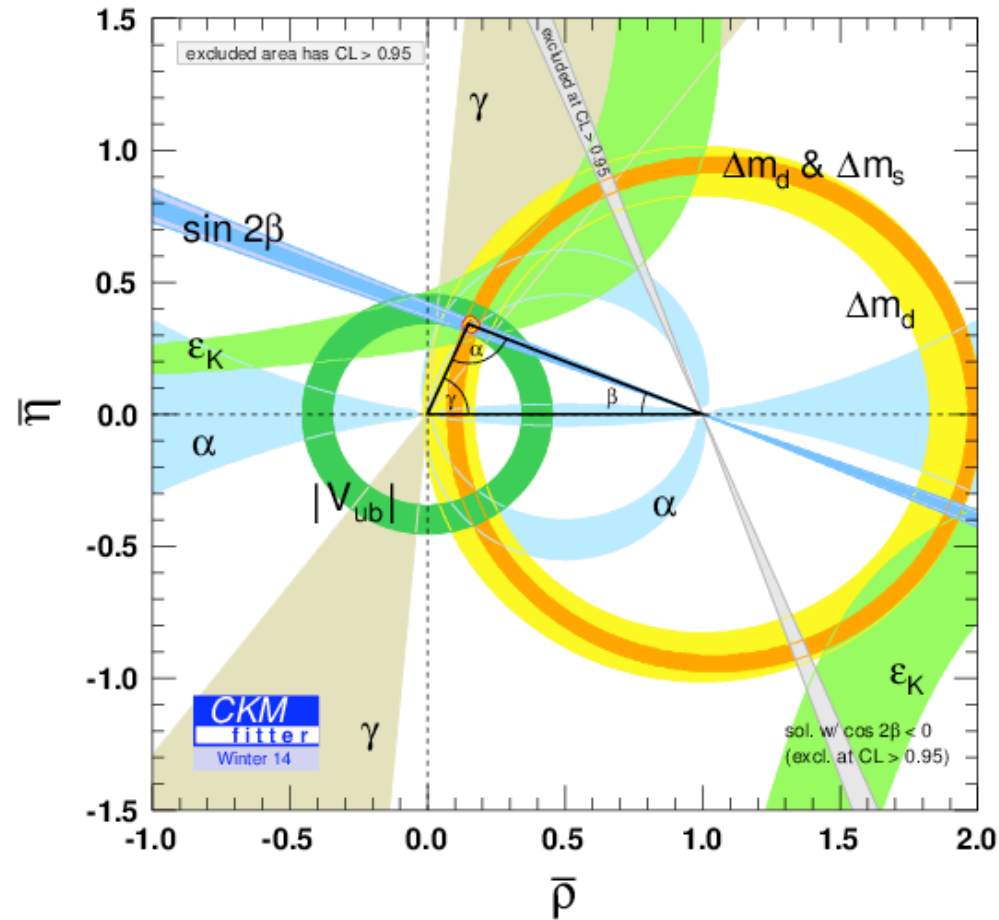
Sneha Malde

University of Oxford

On behalf of the LHCb Collaboration

Beauty 2014: 14th -18th July

# The CKM angle $\gamma$



CKM matrix parameterises quark couplings

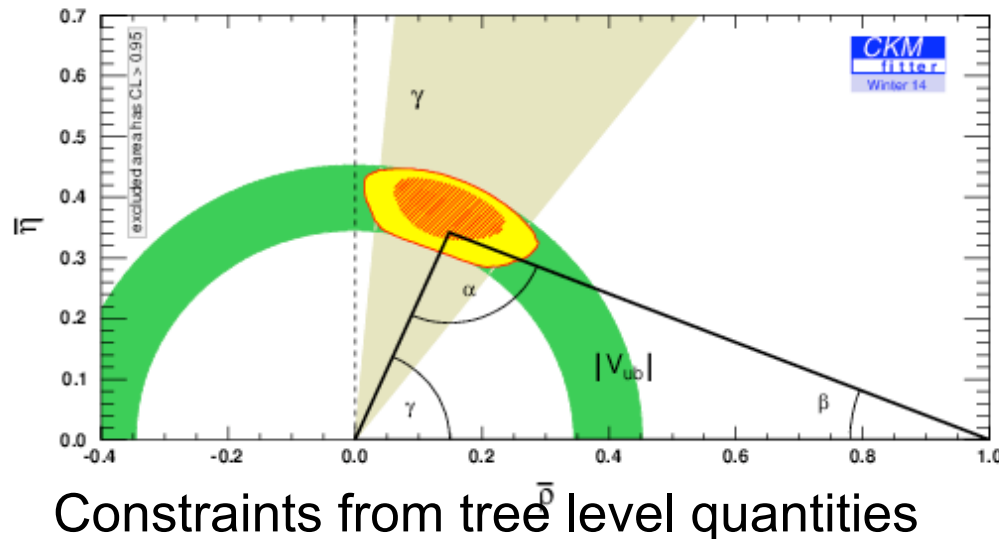
The matrix has one complex phase that results in CPV

Unitarity triangle is a representation of this CPV

$$\gamma = -\arg\left(\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*}\right)$$

$\gamma$  is the least well known angle

# Precision measurement → New Physics test bed



$\gamma$  is the only angle directly accessible in tree decays

“Standard Model” measurement possible

Direct measurements (all results combined):

$$70.0^{+7.7^\circ}_{-9.9}$$

Indirect precision from global CKM fit:

$$66.5^{+1.3^\circ}_{-2.5}$$

Despite recent progress in the direct measurements better precision still required to test for New Physics.

Goal : Improve the direct precision

# This talk

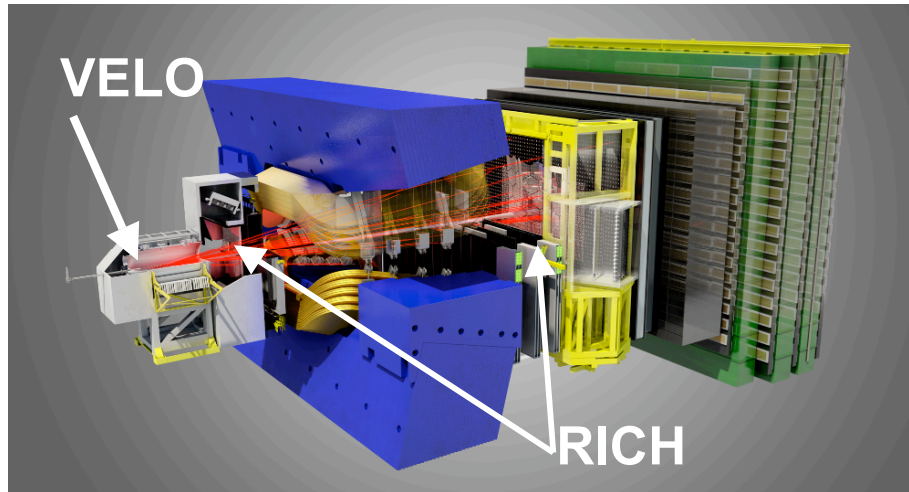
4 new exciting recent results from LHCb

- Model Independent  $B \rightarrow DK$ ,  $D \rightarrow K_S hh$  3 fb<sup>-1</sup>
- Model dependent  $B \rightarrow DK$ ,  $D \rightarrow K_S \pi\pi$  1 fb<sup>-1</sup>
- $B^0 \rightarrow DK^{0*}$ ,  $D \rightarrow hh$
- $B_s \rightarrow D_s K$

NEW

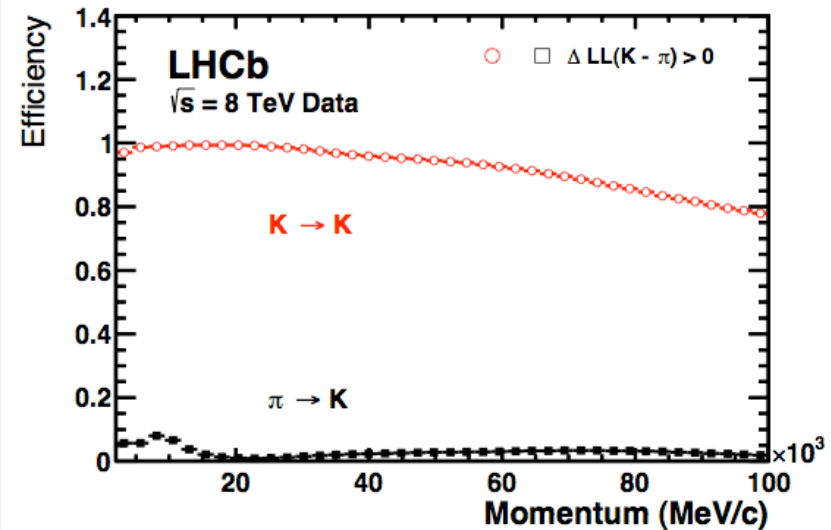
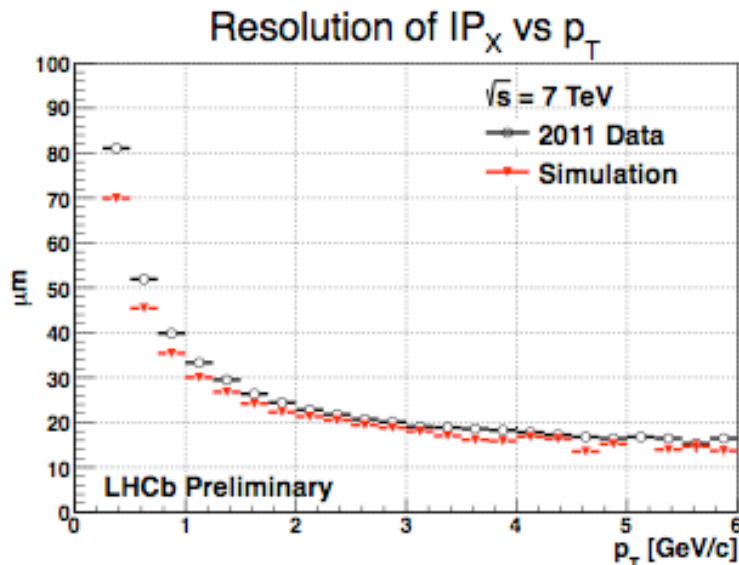
All papers to be submitted to journals soon

# The LHCb detector

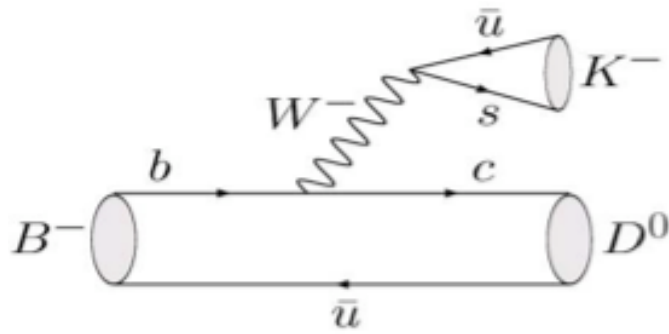


All measurements profit from the VELO for displaced vertices and the RICH for particle identification

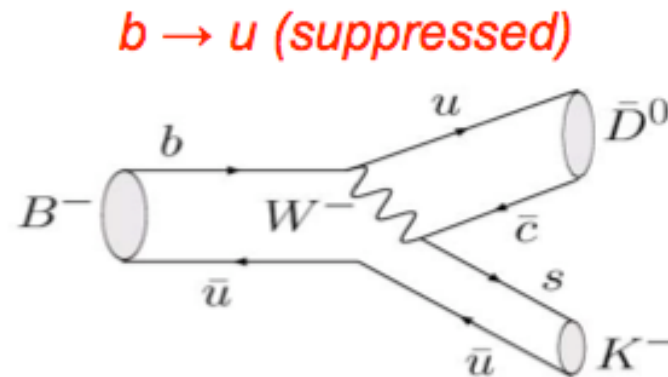
$B_s \rightarrow D_s K$  benefits in particular from the excellent time resolution.



# B → DK, D → K<sub>s</sub>hh - the golden mode



$b \rightarrow c$  (favoured)

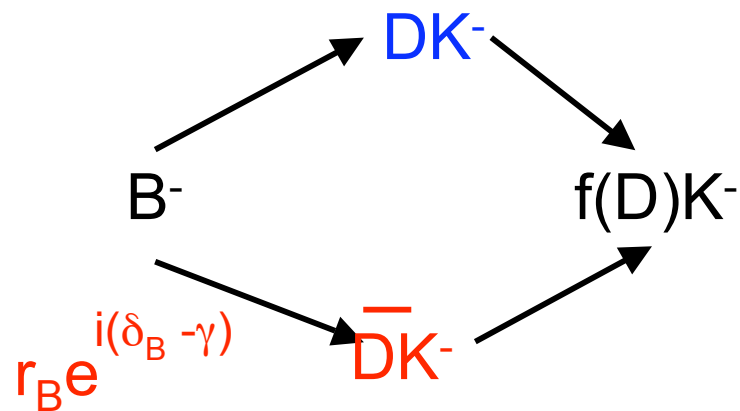


$b \rightarrow u$  (suppressed)

Sensitivity to  $\gamma$  from  $b \rightarrow c$  and  $b \rightarrow u$  interference

Require  $D^0$  and  $\bar{D}^0$  to decay to same final state

$$\gamma = -\arg\left(\frac{V_{ud} V_{ub}^*}{V_{cd} V_{cb}^*}\right)$$



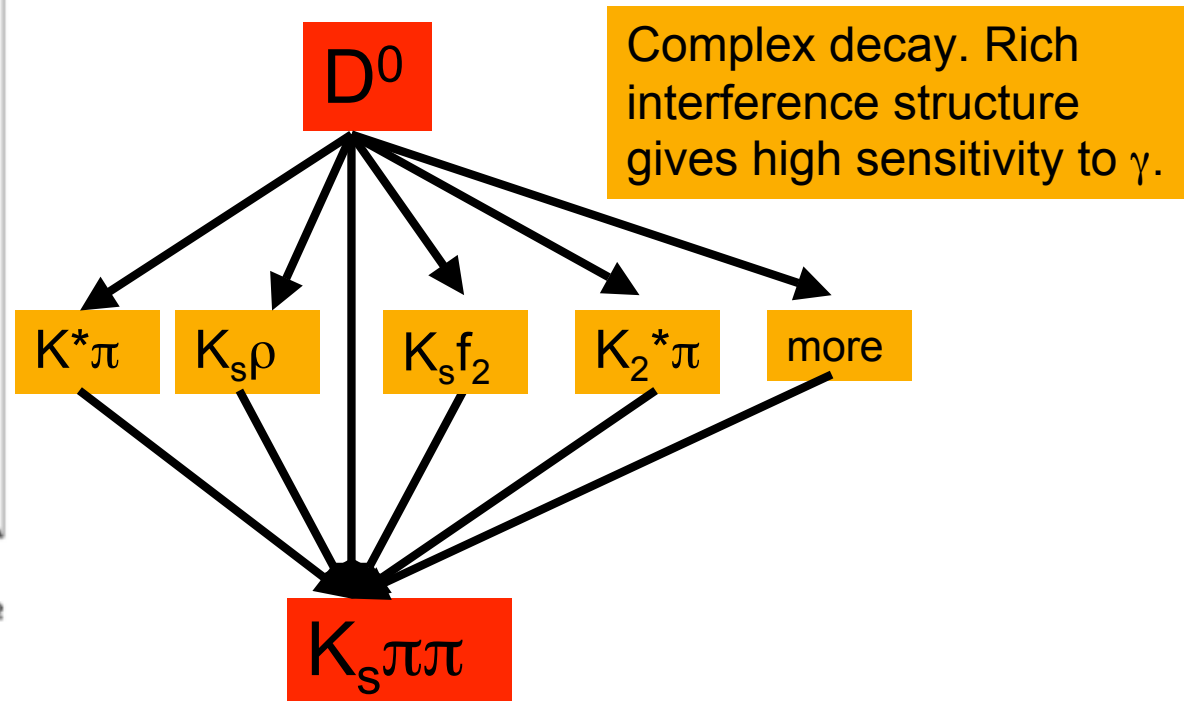
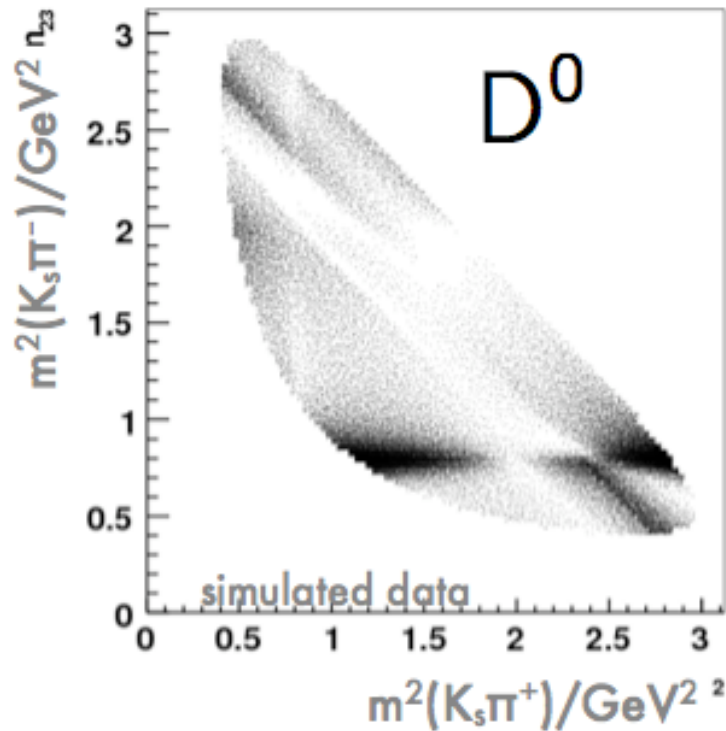
Self-conjugate three body final states are particularly sensitive “GGSZ”

$D \rightarrow K_s \pi \pi$

$D \rightarrow K_s K K$

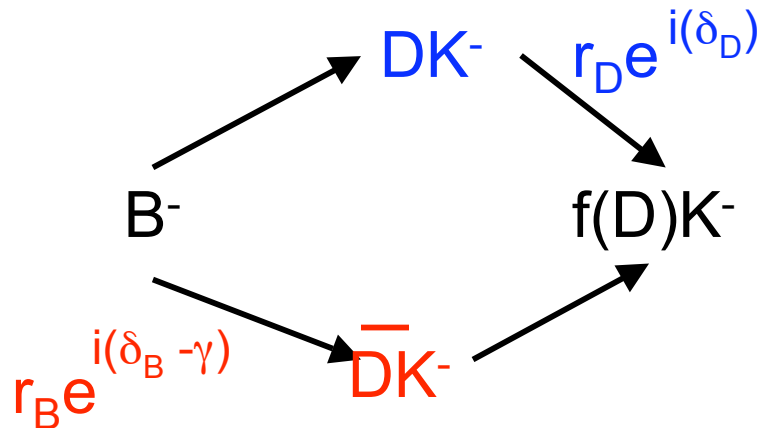
$r_B \sim 0.1$

# $D \rightarrow K_s \pi \pi$ : accessible from both D flavours



Compare the Dalitz plots of  $B^+$  and  $B^-$  decays. Difference driven by  $r_B$ ,  $\delta_B$ ,  $\gamma$

# Accounting for the D strong phase difference



Complication:

D strong phase variation on Dalitz plot

2 methods of dealing with that

both pursued at LHCb

## Model-independent: Count and compare signal yields in Dalitz plot regions

Counting is (relatively) easy

Necessary strong phase information available from dedicated analysis of CLEO-c data.

Well understood systematic uncertainties

Loss of statistical precision (ok if minimal)

## Model-dependent: Fit the Dalitz plot

Maximises statistical power

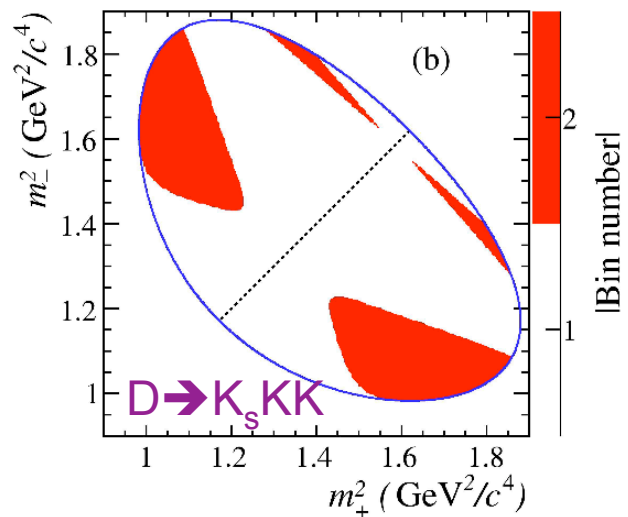
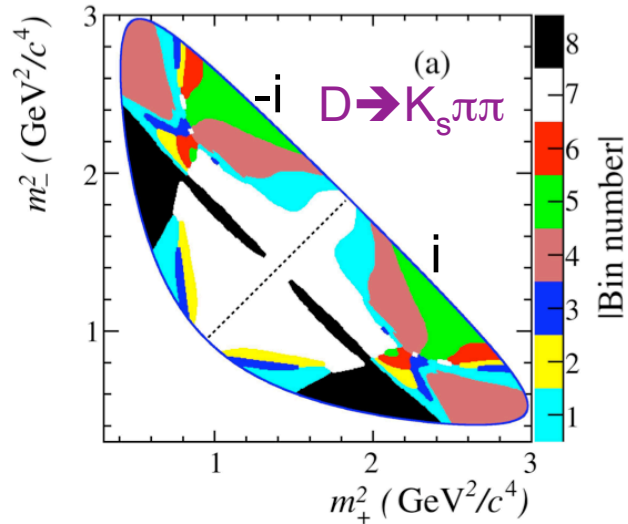
D strong phase variation given by a model

Systematic uncertainties due to this are hard to quantify.

Problematic in the precision era.



# Model-independent method

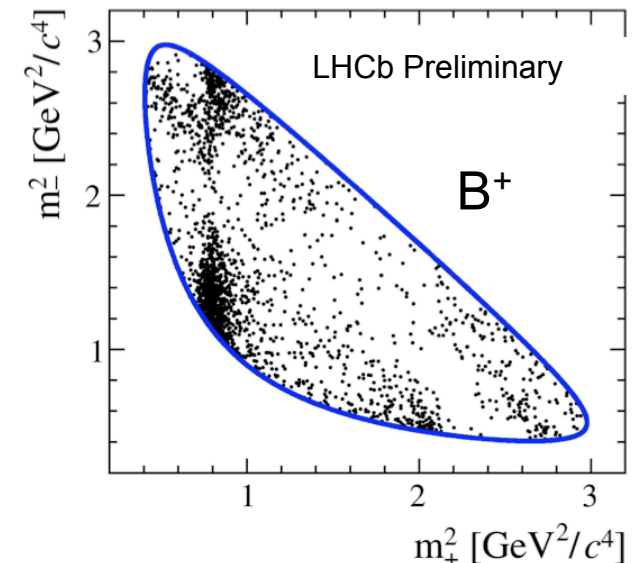
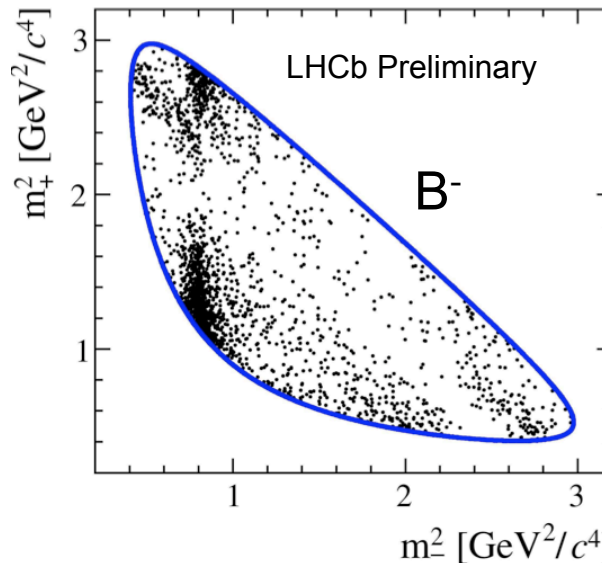


Binning shapes optimised for statistical sensitivity.

Symmetry of Dalitz plot defines positive and negative bins.

Reduces the analysis to a counting experiment in bins of the Dalitz plot

Data from  $D \rightarrow K_S K K$  easily added as two additional bins.



# Overall strategy

To determine  $\gamma$  : Count the number of observed events in a region of the Dalitz plot.

$$N_i^\pm = h \left( K_{\pm i} + r_B^2 K_{\mp i} + 2\sqrt{K_i K_{-i}} \left[ x_\pm c_i \pm y_\pm s_i \right] \right)$$

D from  $B^\pm$  events in bin  $i$  of Dalitz plot

Fraction of events in bin for pure  $D^0$  sample with the efficiency profile of signal

$x_\pm = r_B \cos(\delta_B \pm \gamma)$   
 $y_\pm = r_B \sin(\delta_B \pm \gamma)$

$c_i$  and  $s_i$  are inputs from CLEO - they are measurements of the cosine and sign of the average strong phase difference

$K_i$  are inputs from other LHCb decays - use  $B^0 \rightarrow D^{*+} \mu^- \nu$ ,  $D^{*+} \rightarrow D^0 \pi^+$

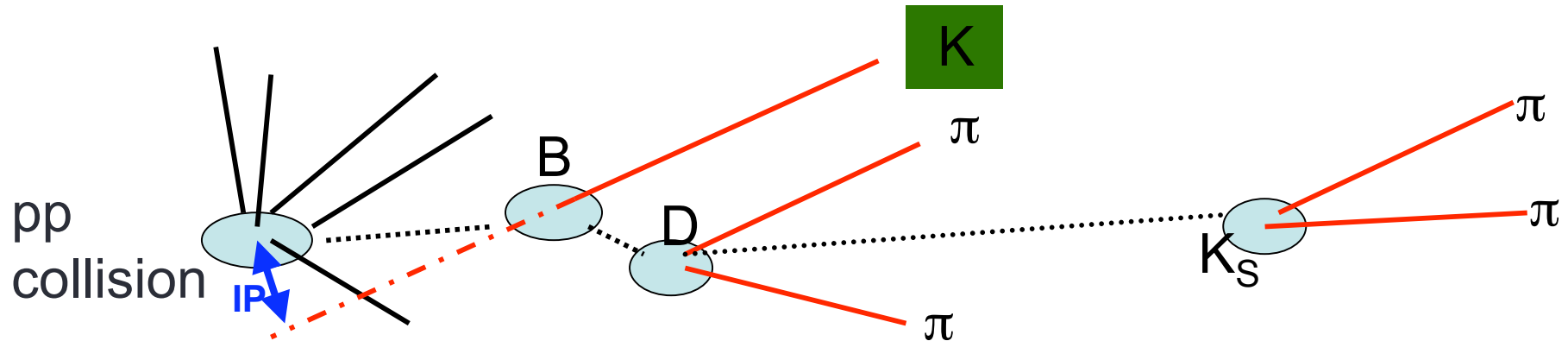
Charge of the  $\pi$  tags the  $D^0$  flavour.

Simultaneous mass fit to candidates in all bins to extract best  $x, y$

Combined measurement of  $3 \text{ fb}^{-1}$  data

Supersedes the previous  $1 \text{ fb}^{-1}$  measurement or the  $(1+2) \text{ fb}^{-1}$  preliminary result

# Selecting events based on topology



Separate the topology of interest from random combinations

Useful variables include:

Impact parameters

Flight distances. (B travels a  $\sim$ cm,  $K_S$  many cm)

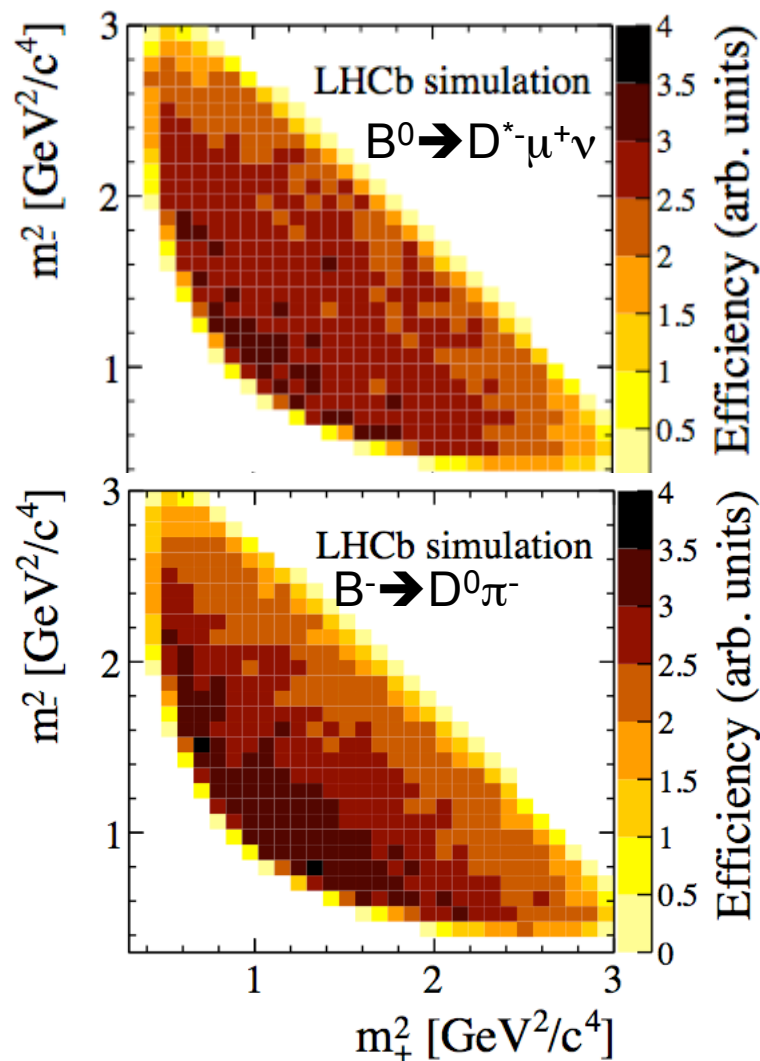
Vertex quality

Particle ID

Specific vetos against particular backgrounds

All analyses shown here employ similar strategies

# Accounting for the efficiency



$B \rightarrow D^* \mu \nu$  best choice for control mode to determine  $K_i$

D decay is flavour tagged. Similar efficiency profile to the signal channel. High purity

Small differences between  $B \rightarrow Dh$  and  $B \rightarrow D^* \mu \nu$  efficiencies observed in MC

Previously used  $B \rightarrow D\pi$ , but CPV exists in this decay → large systematic uncertainties

To determine correct  $K_i$

Determine bin yields of  $B \rightarrow D^* \mu \nu$

Modify them with correction factor derived from MC

# Invariant mass distribution

First fit mass distribution combining all Dalitz Plot bins

Split data by D decay channel and  $K_s$  decay position

In total across all categories ~2600 signal candidates

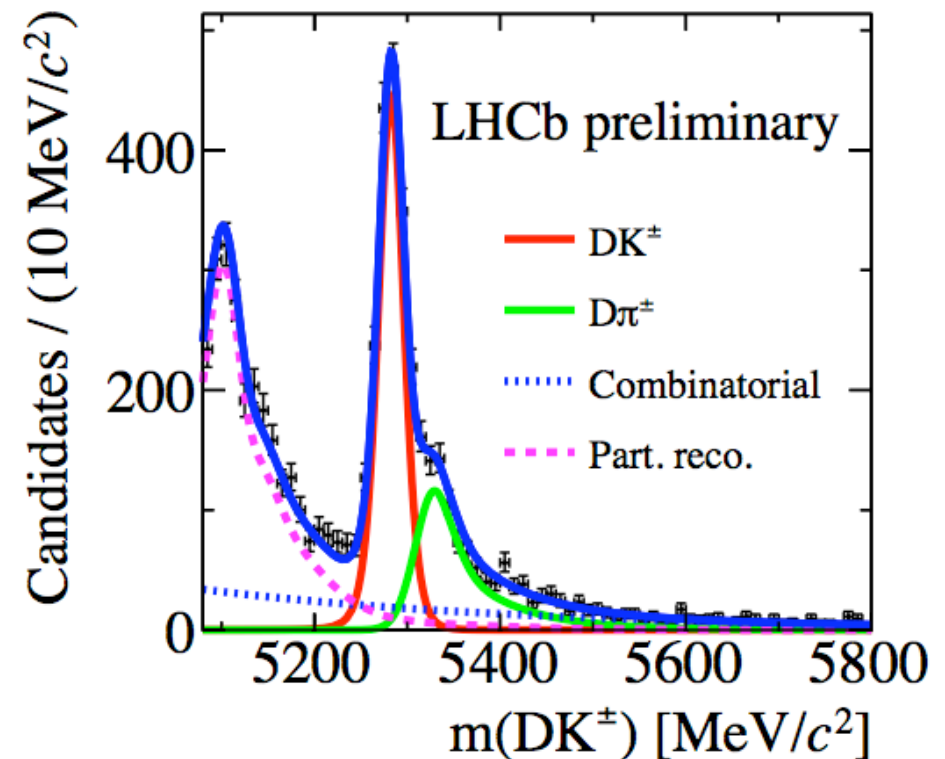
Purity ~ 75% in signal region

Second fit uses fixed mass model PDFs to determine yields in each bin simultaneously and determine best  $x$ ,  $y$  values

$$N_{+i}^+ = n_{B^+} [K_{-i} + (x_+^2 + y_+^2)K_{+i} + 2\sqrt{K_{+i}K_{-i}}(x_+c_{+i} - y_+s_{+i})]$$

$$x_{\pm} = r_B \cos(\delta_B \pm \gamma), y_{\pm} = r_B \sin(\delta_B \pm \gamma)$$

$D \rightarrow K_s \pi \pi$ , where the  $K_s$  decays outside of the VELO



# Results on CP parameters

Preliminary

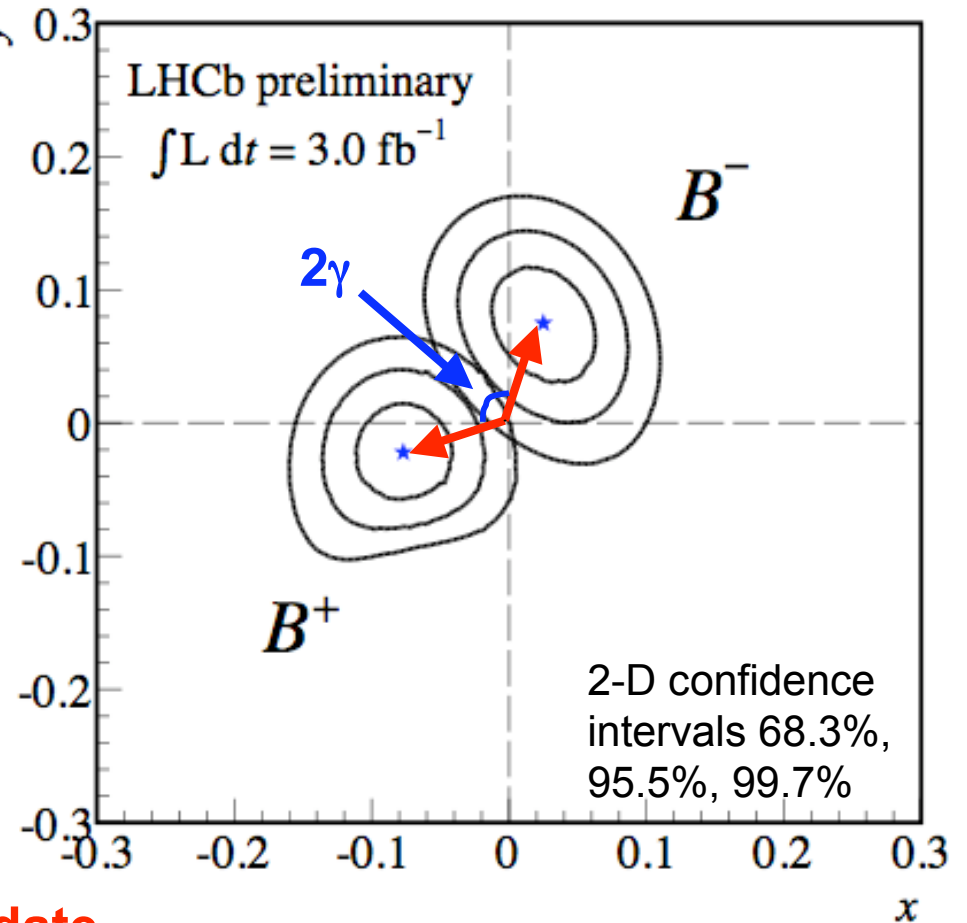
$$x_+ = (-7.7 \pm 2.4 \pm 1.0 \pm 0.4) \times 10^{-2},$$

$$x_- = (2.5 \pm 2.5 \pm 1.0 \pm 0.5) \times 10^{-2},$$

$$y_+ = (-2.2 \pm 2.5 \pm 0.4 \pm 1.0) \times 10^{-2},$$

$$y_- = (7.5 \pm 2.9 \pm 0.5 \pm 1.4) \times 10^{-2},$$

Corrections for D mixing,  $K_s$   
CPV ignored - negligible effect



**Most precise measurement of x,y to date**

# Comparing to 1 fb<sup>-1</sup> result

Statistical uncertainties reduce due to increased data sample

	x <sub>+</sub>	x <sub>-</sub>	y <sub>+</sub>	y <sub>-</sub>
2011	0.045	0.043	0.037	0.052
2013	0.024	0.025	0.025	0.029

Strong phase systematic reduces due to increased sample size

	x <sub>+</sub>	x <sub>-</sub>	y <sub>+</sub>	y <sub>-</sub>
2011	0.014	0.006	0.030	0.023
2013	0.004	0.005	0.010	0.014

Experimental systematic uncertainty reduced due to change in control mode. MC efficiency correction dominates.

	x <sub>+</sub>	x <sub>-</sub>	y <sub>+</sub>	y <sub>-</sub>
2011	0.018	0.015	0.008	0.008
2013	0.010	0.010	0.004	0.005

$$\begin{aligned}
 x_+ &= (-7.7 \pm 2.4 \pm 1.0 \pm 0.4) \times 10^{-2}, \\
 x_- &= (2.5 \pm 2.5 \pm 1.0 \pm 0.5) \times 10^{-2}, \\
 y_+ &= (-2.2 \pm 2.5 \pm 0.4 \pm 1.0) \times 10^{-2}, \\
 y_- &= (7.5 \pm 2.9 \pm 0.5 \pm 1.4) \times 10^{-2},
 \end{aligned}$$

**Most precise measurement of x,y to date**

# Physics parameters of interest

Determine (stat+syst) confidence intervals for  $\gamma$ ,  $r_B$ ,  $\delta_B$

Two-fold ambiguity

1fb<sup>-1</sup> result:  $\gamma = 44^{+43}_{-38}$

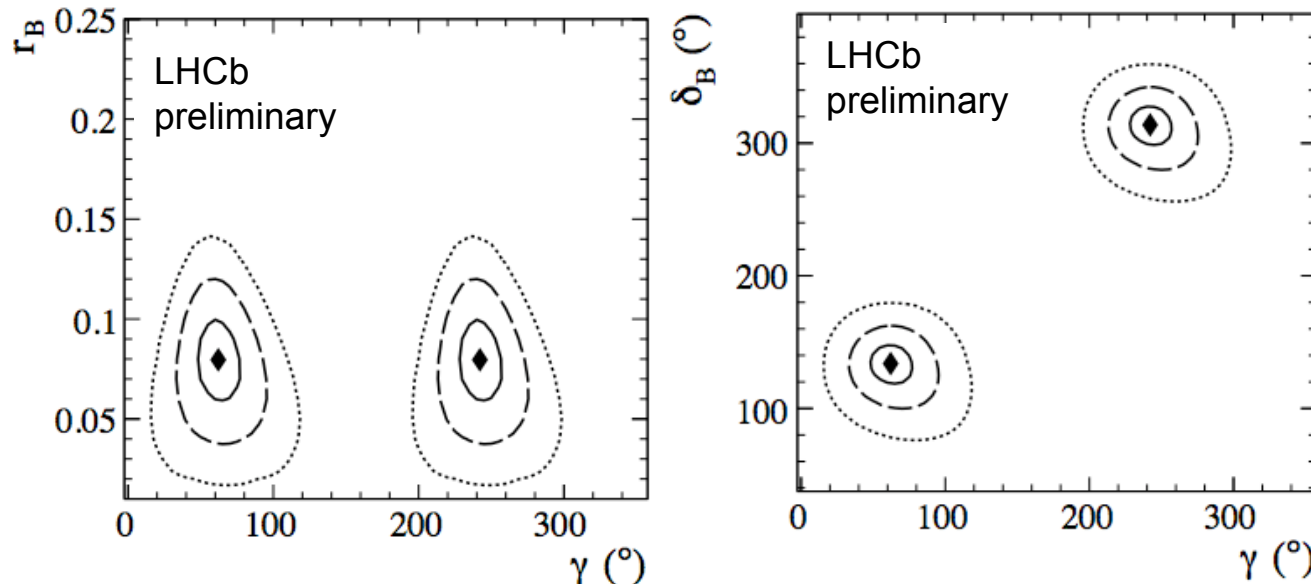
Improvement threefold - more data, lower systematic uncertainties and higher central  $r_B$  value

$$\gamma = (62^{+15}_{-14})^\circ$$

$$r_B = (8.0^{+1.9}_{-2.1}) \times 10^{-2}$$

$$\delta_B = (134^{+14}_{-15})^\circ$$

**Precision matches that of either B factory  $\gamma$  combinations**



Projection of contours onto 1-D gives 68.3, 95.5, 99.7 % CL

Preliminary



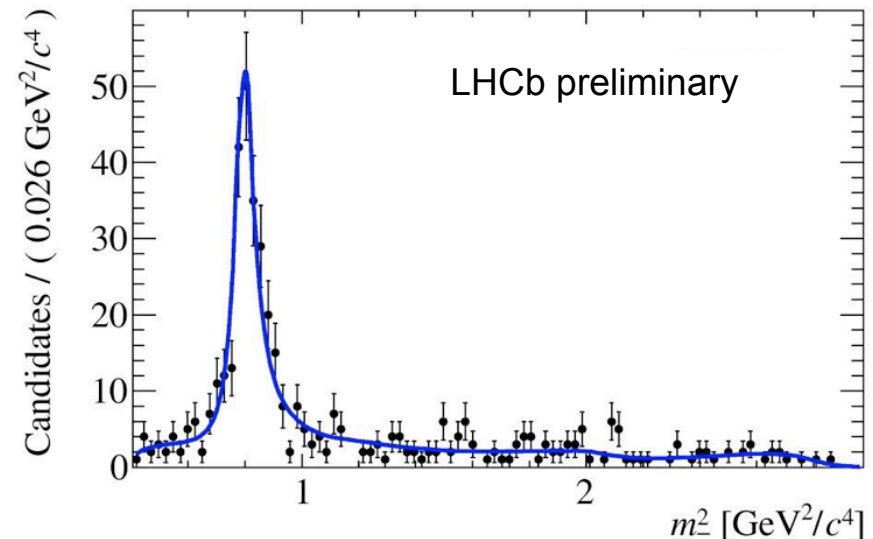
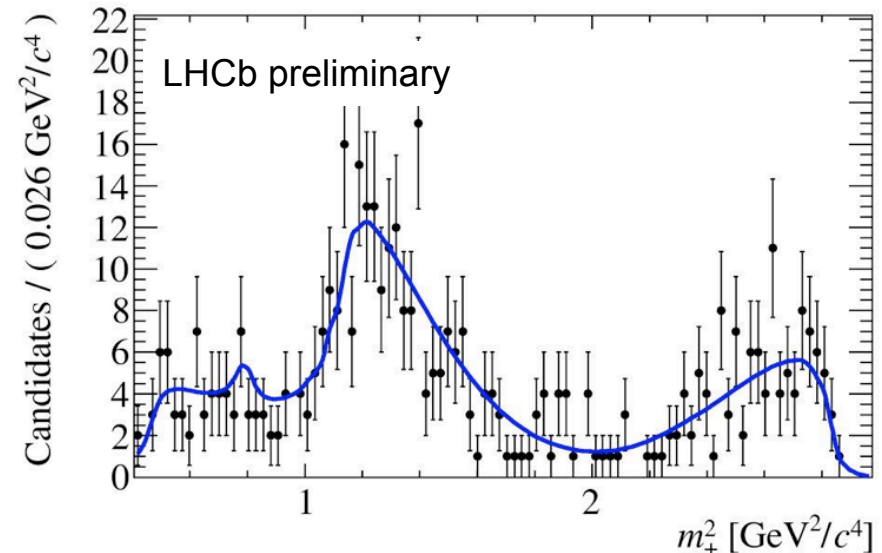
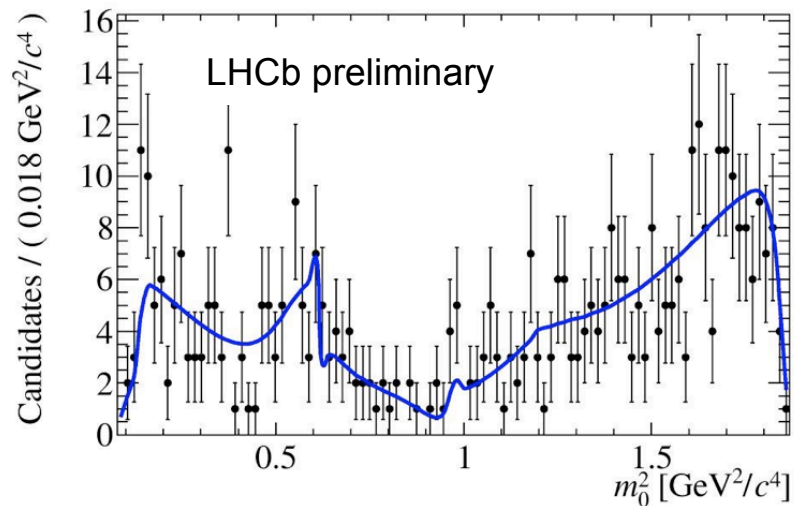
# Model-dependent method

D strong phase differences accounted for by using an amplitude model

The latest BaBar amplitude model is chosen.

Only consider the  $K_S\pi\pi$  decay ( $1 \text{ fb}^{-1}$ )

Efficiency on the Dalitz Plot determined from  $B \rightarrow D\pi$  assuming no CPV



# Model-dependent results

Preliminary

$$x_- = +0.027 \pm 0.044_{-0.008}^{+0.010} \pm 0.001,$$

$$y_- = +0.013 \pm 0.048_{-0.007}^{+0.009} \pm 0.003,$$

$$x_+ = -0.084 \pm 0.045 \pm 0.009 \pm 0.005,$$

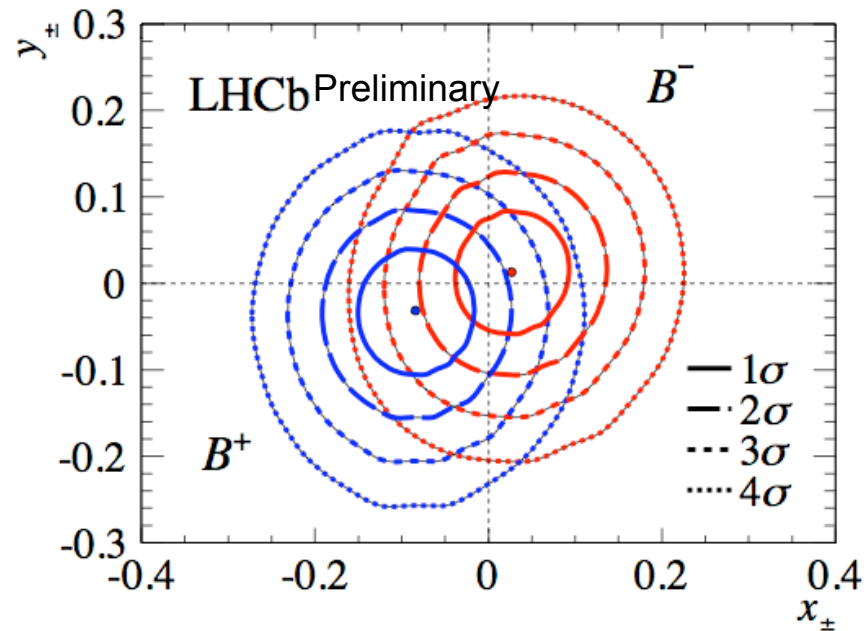
$$y_+ = -0.032 \pm 0.048_{-0.009}^{+0.010} \pm 0.008,$$

Final uncertainty is due to model. Determined by using alternate models.

Leading experimental contributions are efficiency and background distribution uncertainties.

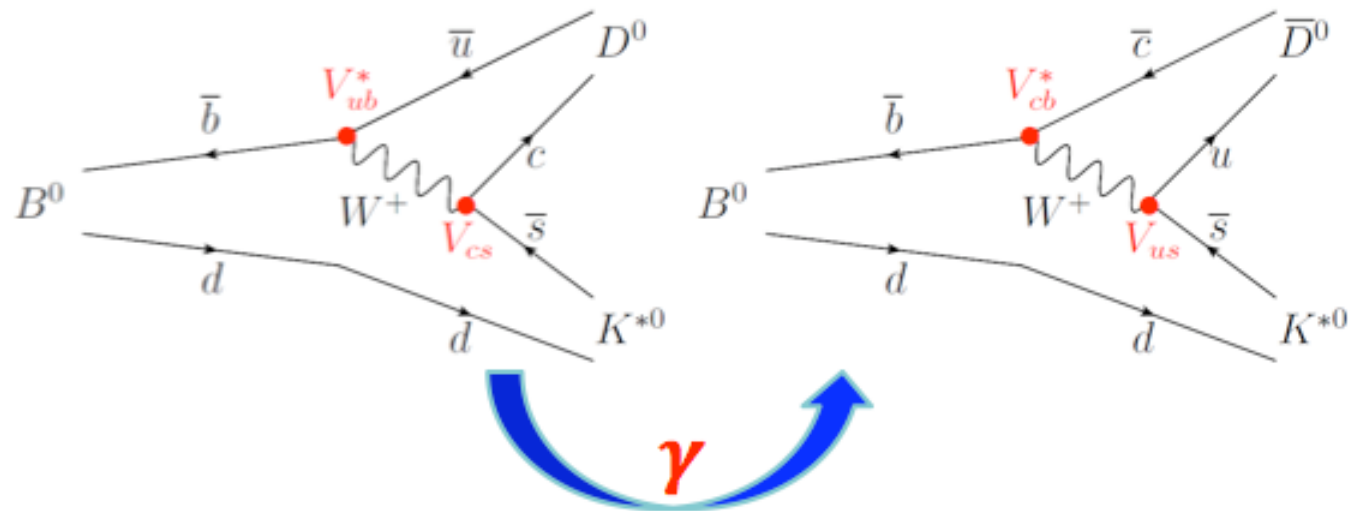
x, y measurement consistent with the  $1\text{fb}^{-1}$  MI result

$$\gamma = 84_{-42}^{+49}$$



2-D confidence intervals 68.3%, 95.5%, 99.7%

# $B^0 \rightarrow DK^{0*}$



Both diagrams colour suppressed  $\rightarrow$  large  $r_B$  (large interference)

Initial  $B^0$  flavour tagged by the charge of the kaon in the  $K^{0*} \rightarrow K\pi$

(no need for time-dependent analysis)

Interference requires same D final state

So far two-body D decay modes considered with  $3 \text{ fb}^{-1}$

$\gamma$  common to all analyses; here different B decay means different  $r_B$  and  $\delta_B$

# $B^0 \rightarrow DK^{0*}$ observables

Can construct a number of observables that have sensitivity to  $\gamma$

Asymmetries in  $D \rightarrow$  CP eigenstates (KK,  $\pi\pi$  easily accessible at LHCb)

$$A_d^{hh} \equiv \frac{\Gamma(\bar{B}^0 \rightarrow D(h^+h^-)\bar{K}^{*0}) - \Gamma(B^0 \rightarrow D(h^+h^-)K^{*0})}{\Gamma(\bar{B}^0 \rightarrow D(h^+h^-)\bar{K}^{*0}) + \Gamma(B^0 \rightarrow D(h^+h^-)K^{*0})} = \frac{2r_D \kappa \sin \delta_B \sin \gamma}{1 + r_B^2 + 2r_D \kappa \cos \delta_B \cos \gamma}$$

Can also consider the CF decay  $D^0 \rightarrow K\pi^+$  and the DCS decay  $D^0 \rightarrow K^+\pi^-$

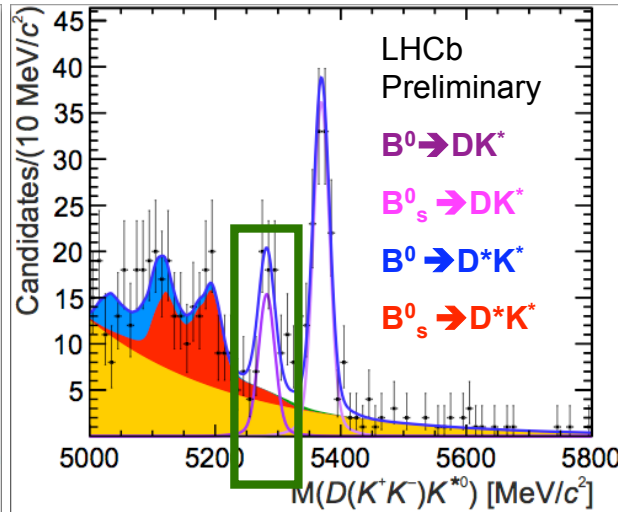
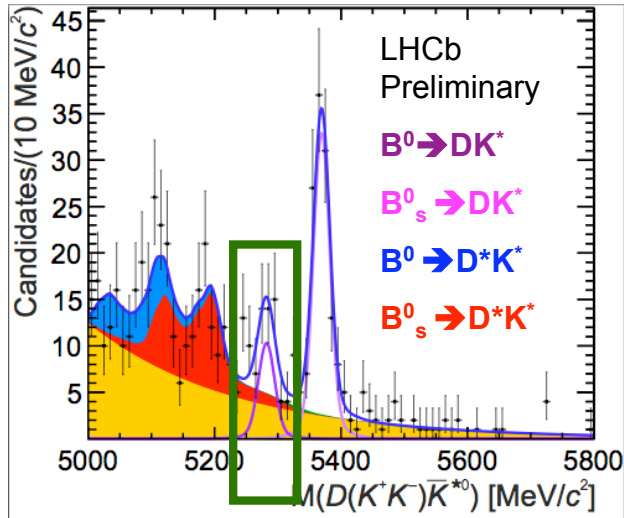
$$\mathcal{R}_d^+ \equiv \frac{\Gamma(B^0 \rightarrow D(\pi^+K^-)K^{*0})}{\Gamma(B^0 \rightarrow D(K^+\pi^-)K^{*0})} = \frac{r_B^2 + r_D^2 + 2r_B r_D \kappa \cos(\delta_B + \delta_D + \gamma)}{1 + r_B^2 r_D^2 + 2r_B r_D \kappa \cos(\delta_B - \delta_D + \gamma)}$$

$$\mathcal{R}_d^- \equiv \frac{\Gamma(\bar{B}^0 \rightarrow D(\pi^-K^+)\bar{K}^{*0})}{\Gamma(\bar{B}^0 \rightarrow D(K^-\pi^+)\bar{K}^{*0})} = \frac{r_B^2 + r_D^2 + 2r_B r_D \kappa \cos(\delta_B + \delta_D - \gamma)}{1 + r_B^2 r_D^2 + 2r_B r_D \kappa \cos(\delta_B - \delta_D - \gamma)}$$

Other observables include asymmetry in  $\bar{B}^0$  and  $B^0$  in the CF decay, and ratio of CP even to the CF decay.

Relations involve  $\kappa$ , the coherence factor to take into account non-resonant yield in the  $K^{0*}$  region

# $B^0 \rightarrow DK^{0*}$ $D \rightarrow KK, \pi\pi$

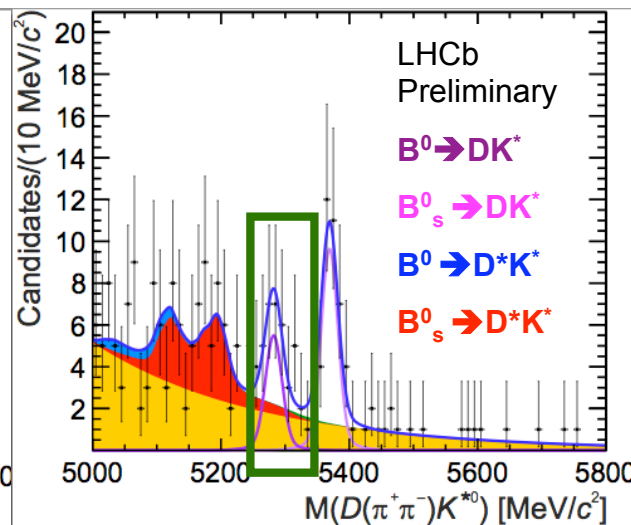
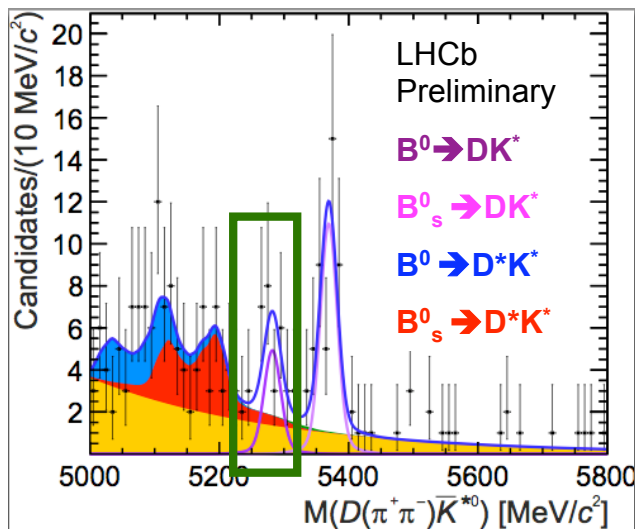


**KK - significance of signal:**  
 $8.6\sigma$

**$\pi\pi$  - significance of signal:**  
 $5.8\sigma$

Corrections applied for production and selection efficiencies.

Leading systematic uncertainties from charge detection asymmetry in the K from  $K^{0*}$

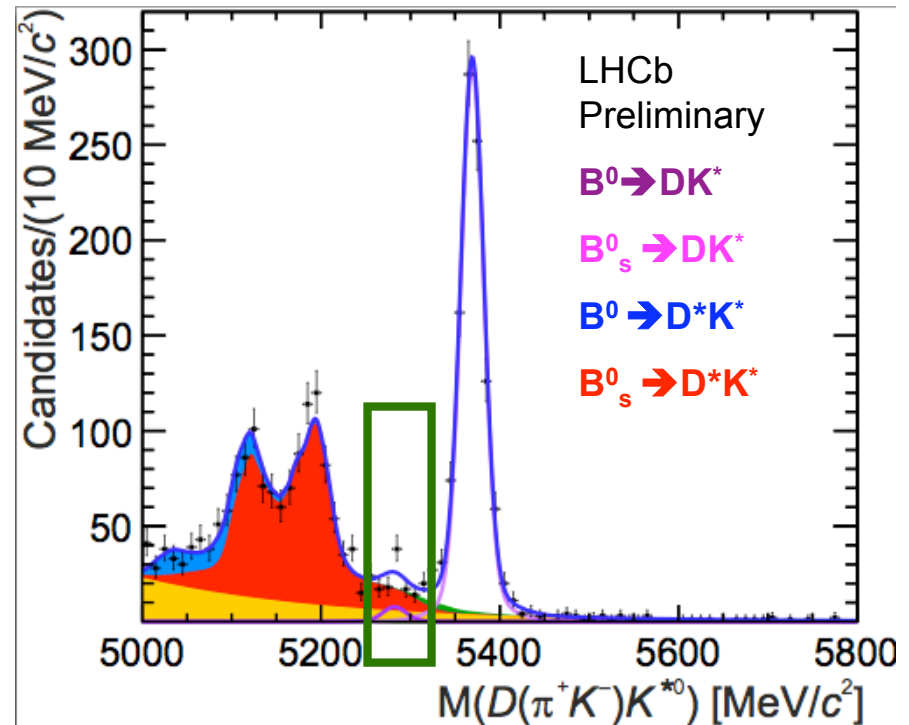
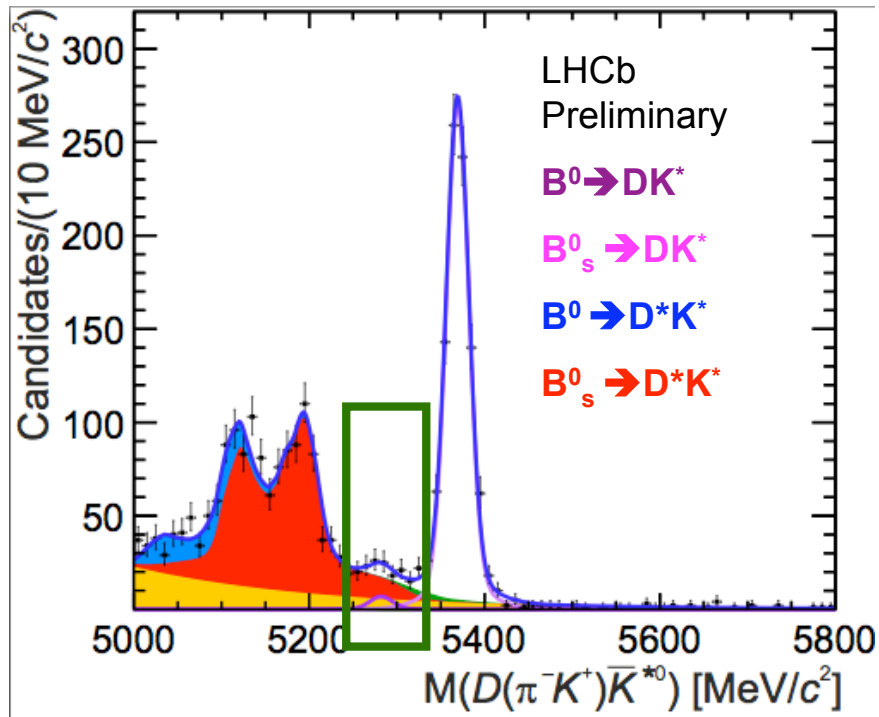


$$A_d^{KK} = -0.198^{+0.144+0.019}_{-0.145-0.020}$$

$$A_d^{\pi\pi} = -0.092^{+0.217+0.019}_{-0.217-0.019}$$

Preliminary

# $B^0 \rightarrow DK^{0*}, D \rightarrow \pi K$



Combined signal significance is  $2.9\sigma$ .

$$R_d^+ = 0.057^{+0.029+0.009}_{-0.027-0.012}$$

$$R_d^- = 0.056^{+0.032+0.009}_{-0.030-0.012}$$

Preliminary

# Constraints on physics parameters

Coherence factor  $\kappa$  determined from simulation of a realistic model of the resonance content of  $B^0 \rightarrow DK\pi$

$$\kappa = 0.95 \pm 0.03$$

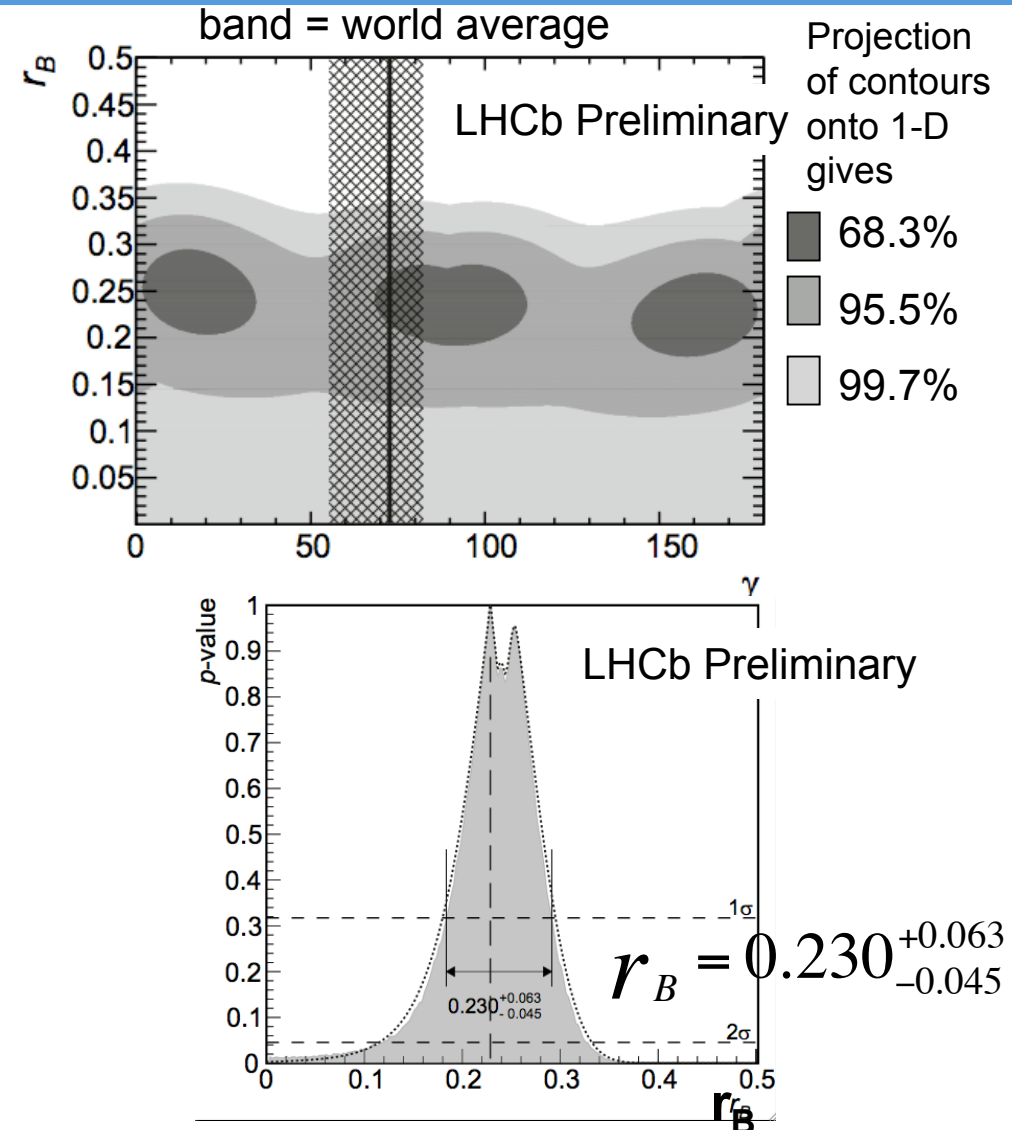
Measurements of all observables combined to determine  $r_B$ ,  $\delta_B$ ,  $\gamma$

Ambiguities from the trigonometric relations

Some constraints can be set at the 68.3% CL

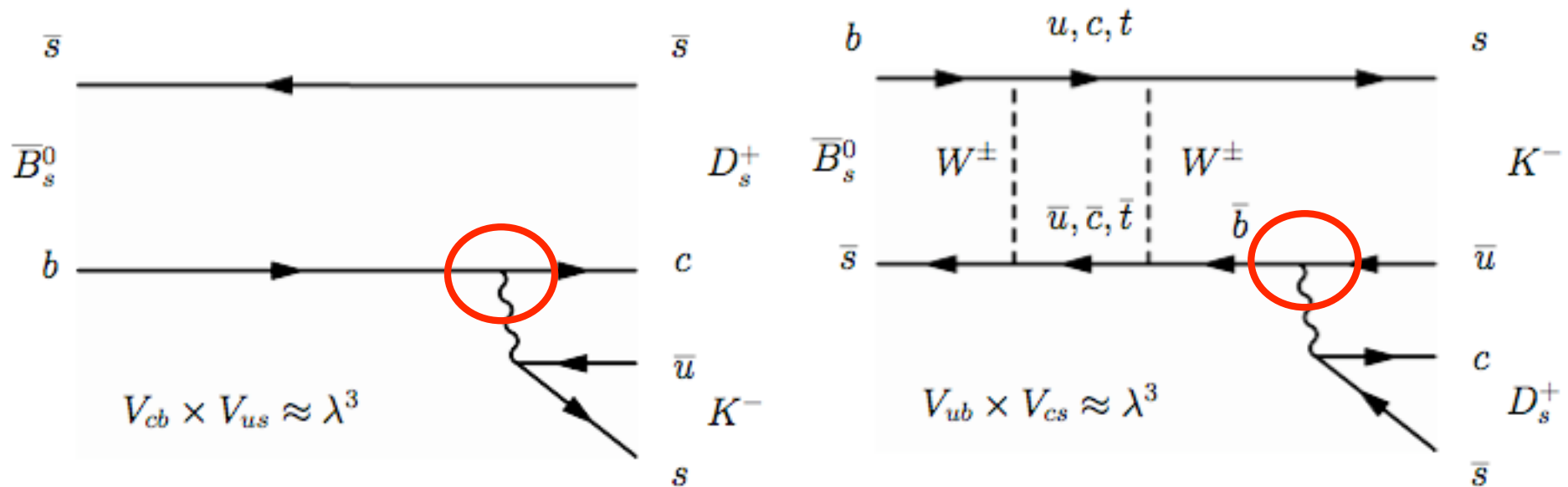
$r_B$  value larger than that for  $B \rightarrow DK$

Promising decay to study further



# $B_s \rightarrow D_s K$

Measure CP violation in the interference of mixing and decay



Both decay amplitudes  $\sim \lambda^3 \rightarrow$  Large interference

Tree level process like other analyses shown

Time-dependence increases the complexity of the analysis



# CP Observables

$$\frac{d\Gamma_{B_s^0 \rightarrow f}(t)}{dt} = \frac{1}{2}|A_f|^2(1 + |\lambda_f|^2)e^{-\Gamma_s t} \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + A_f^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) + C_f \cos(\Delta m_s t) - S_f \sin(\Delta m_s t) \right],$$

$$\lambda_f \equiv \frac{q}{p} \left( \frac{\bar{A}_f}{A_f} \right)$$

$$\frac{d\Gamma_{\bar{B}_s^0 \rightarrow f}(t)}{dt} = \frac{1}{2}|A_f|^2 \left| \frac{p}{q} \right|^2 (1 + |\lambda_f|^2)e^{-\Gamma_s t} \left[ \cosh\left(\frac{\Delta\Gamma_s t}{2}\right) + A_f^{\Delta\Gamma} \sinh\left(\frac{\Delta\Gamma_s t}{2}\right) - C_f \cos(\Delta m_s t) + S_f \sin(\Delta m_s t) \right],$$

$A_f$  is the decay amplitude for  $B_s$  to decay to final state  $f$

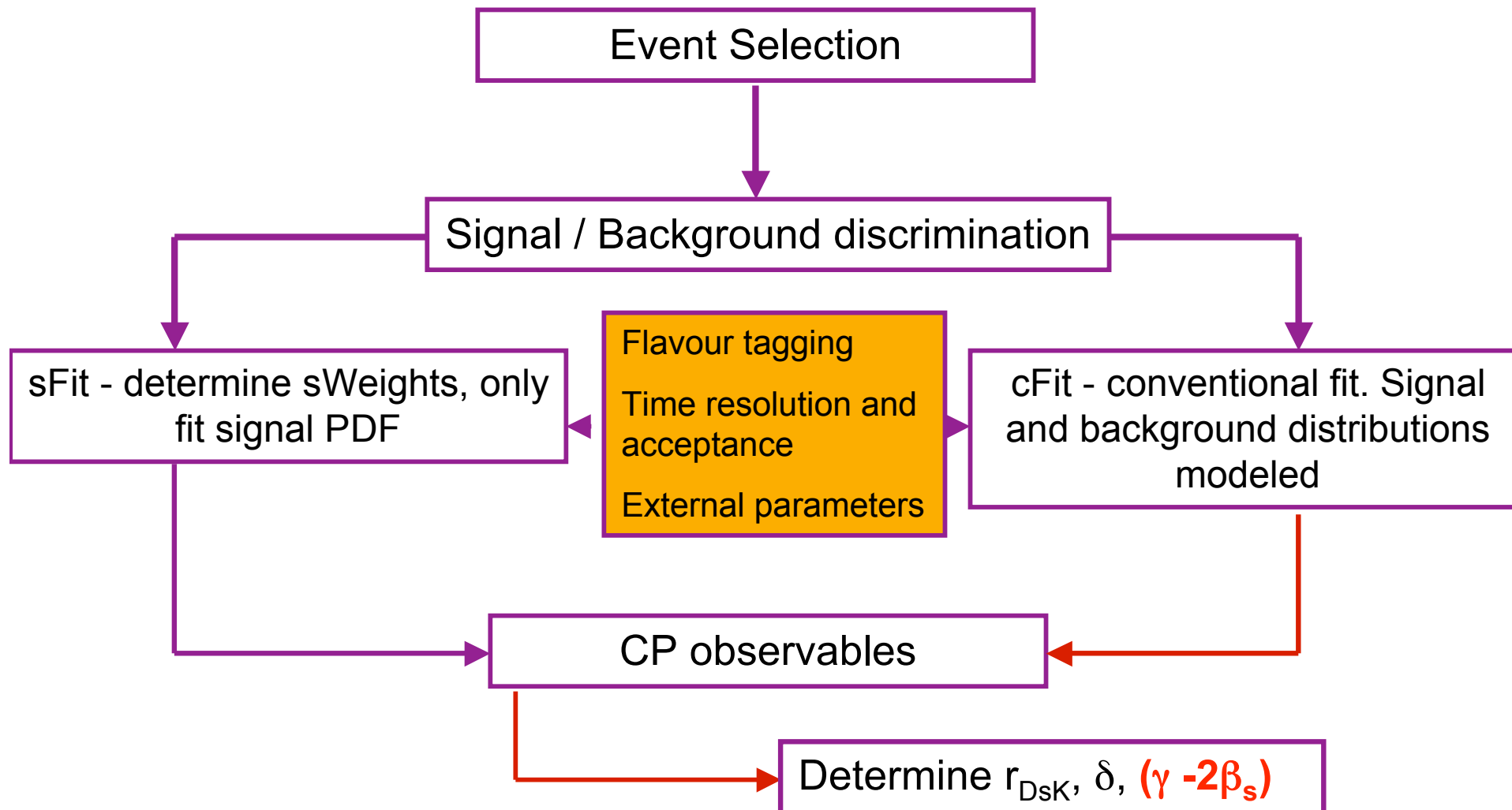
$$C_f = \frac{1 - r_{D_s K}^2}{1 + r_{D_s K}^2},$$

$$A_f^{\Delta\Gamma} = \frac{-2r_{D_s K} \cos(\delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad \bar{A}_f^{\Delta\Gamma} = \frac{-2r_{D_s K} \cos(\delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2},$$

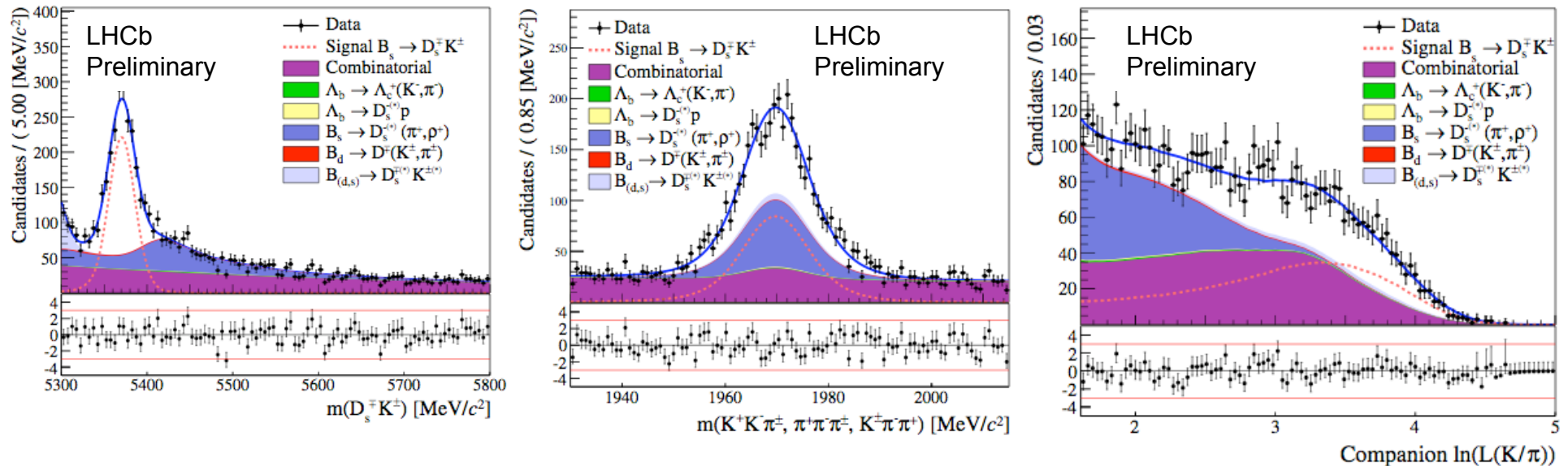
$$S_f = \frac{2r_{D_s K} \sin(\delta - (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}, \quad \bar{S}_f = \frac{-2r_{D_s K} \sin(\delta + (\gamma - 2\beta_s))}{1 + r_{D_s K}^2}.$$

$\beta_s$  - mixing phase

# Analysis Strategy



# Signal/Background discrimination



Three  $D_s^-$  decays considered:  $K^- K^+ \pi^-$ ,  $\pi^- \pi^+ \pi^-$ ,  $K^- \pi^+ \pi^-$ : Plots show all  $D_s^-$  states combined  
 Simultaneous fit in 3 variables:  $M(B_s)$ ,  $M(D_s)$  and PID variable on the Kaon from the B  
 Allows for signal/background discrimination, and for determination of signal weights

# Other Inputs

## Flavour tagging:

Combination of SS and OS taggers

Efficiency of tagging an event = 67.5%

Effective tagging power = 5.07%

## Decay time acceptance:

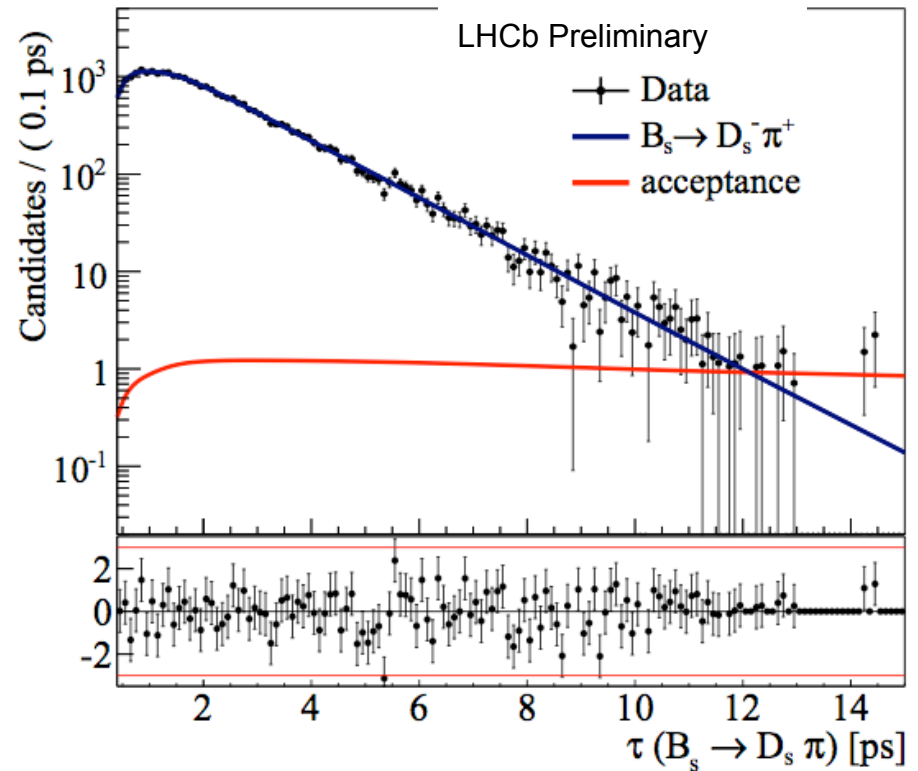
$B_s \rightarrow D_s \pi$  with additional corrections from simulation

## Decay time resolution:

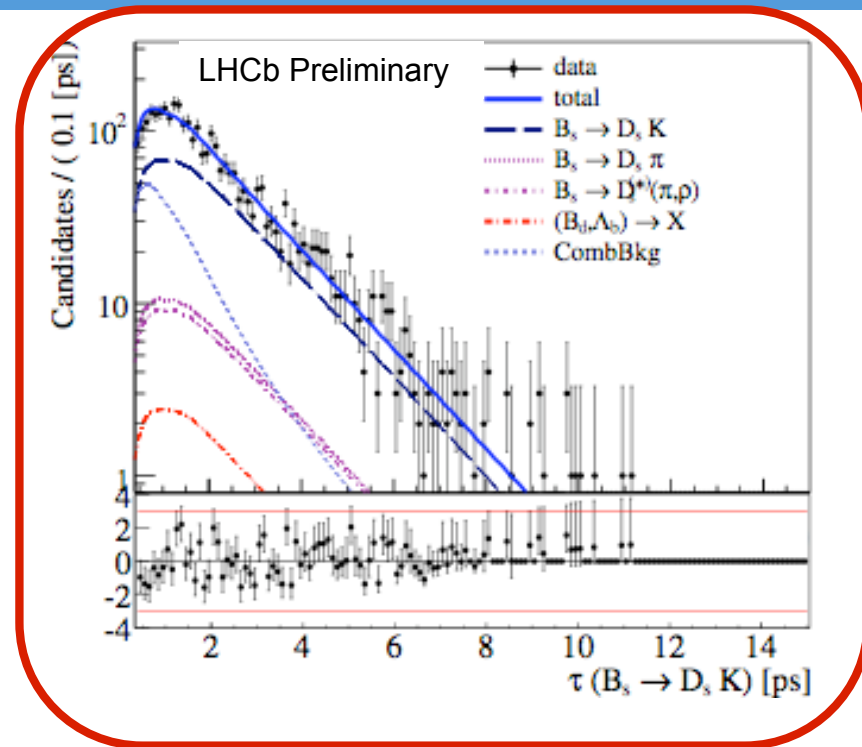
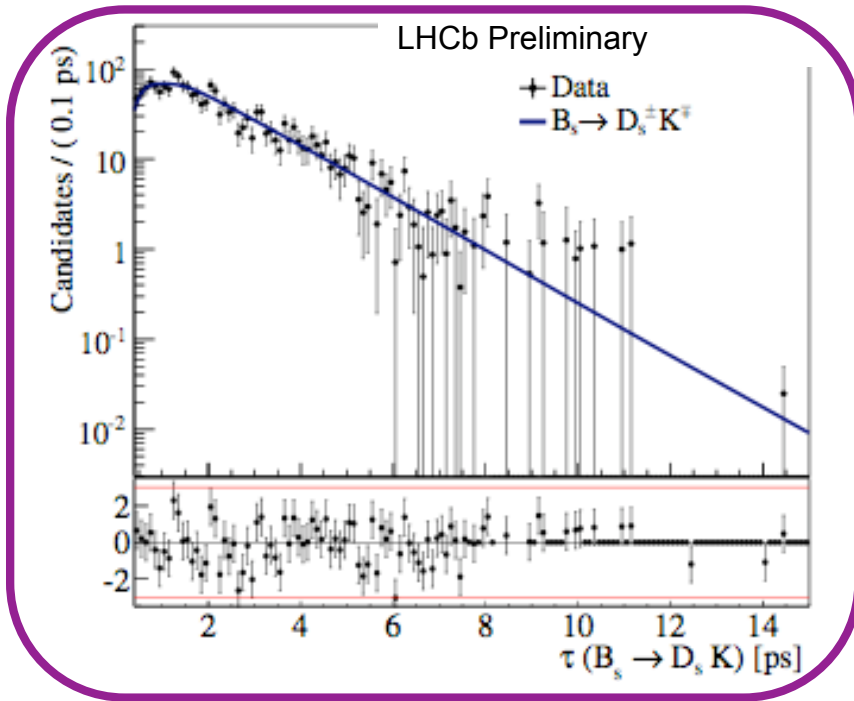
Use the per-event error. Average resolution is 47 fs

## External Inputs:

$\Gamma_s$ ,  $\Delta\Gamma_s$ ,  $\Gamma_d$ ,  $\Gamma_{\Lambda b}$ ,  $\Delta m_s$  all fixed from other measurements.



# Decay time fit



Parameter	<i>sFit</i> fitted value	<i>cFit</i> fitted value
$C_f$	$0.52 \pm 0.25 \pm 0.04$	$0.53 \pm 0.25 \pm 0.04$
$A_f^{\Delta\Gamma}$	$0.29 \pm 0.42 \pm 0.17$	$0.37 \pm 0.42 \pm 0.20$
$A_{\bar{f}}^{\Delta\Gamma}$	$0.14 \pm 0.41 \pm 0.18$	$0.20 \pm 0.41 \pm 0.20$
$S_f$	$-0.90 \pm 0.31 \pm 0.06$	$-1.09 \pm 0.33 \pm 0.08$
$S_{\bar{f}}$	$-0.36 \pm 0.34 \pm 0.06$	$-0.36 \pm 0.34 \pm 0.08$

Good agreement

*sFit* is a nice cross check to *cFit*

# Measurement of angle $\gamma$

$\gamma$  determination based on cFit results

Statistical and systematic uncertainties  
+ correlations taken into account

First measurement from  $B_s \rightarrow D_s K$

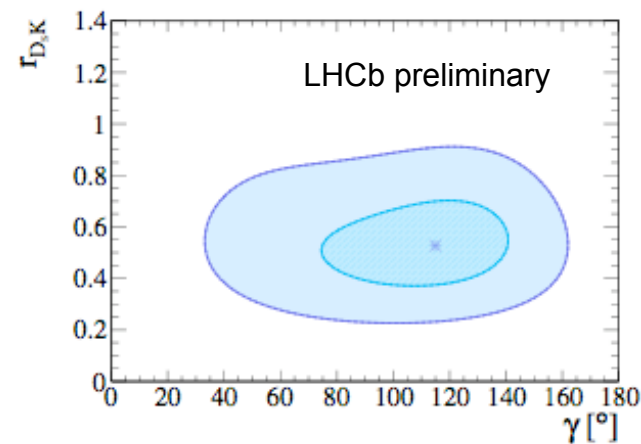
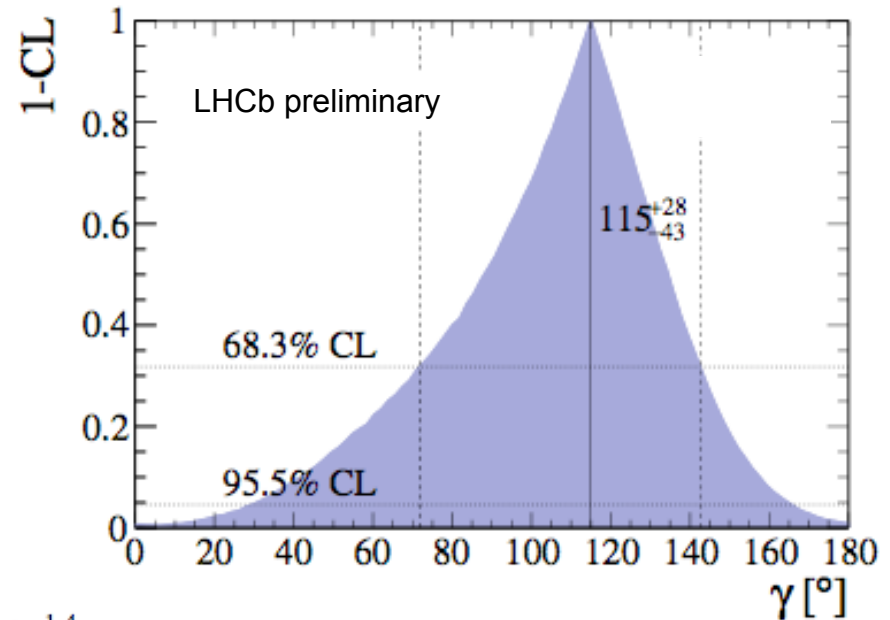
Only  $1\text{fb}^{-1}$  - more available

$$\gamma = (115^{+28}_{-43})^\circ$$

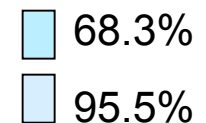
$$r_{D_s K} = (0.53^{+0.17}_{-0.16})$$

$$\delta_{D_s K} = (3^{+19}_{-20})^\circ$$

Preliminary



Projection  
of contours  
onto 1-D  
gives



# Summary and outlook

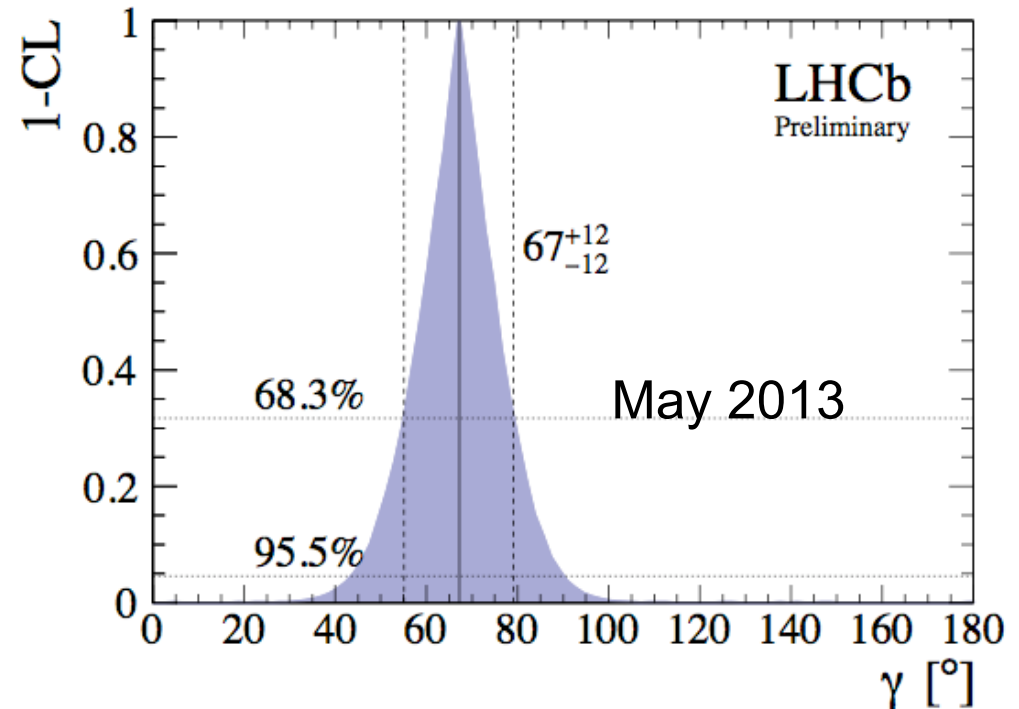
Many tree level  $\gamma$  measurements being pursued

Demonstrated that the complex analyses like  $B_s \rightarrow D_s K$  and GGSZ can be done at LHCb

$B^0 \rightarrow DK^{0*}$ , and  $B_s \rightarrow D_s K$  will add to the combination and continue to drive the direct uncertainty down

Updated  $3\text{fb}^{-1}$  GGSZ measurement has better overall precision than  $(1+2)\text{fb}^{-1}$  already included in the combination and will also continue to improve precision.

Other  $B \rightarrow DK$ ,  $B \rightarrow DX$  analyses also pursued, either updating  $1\text{fb}^{-1} \rightarrow 3\text{fb}^{-1}$  or new D decays.



**Precision on direct  $\gamma$  measurements from LHCb will continue to reduce**