The 15th International Conference on B-Physics at Frontier Machines at the University of Edinburgh.

14th - 18th July 2014

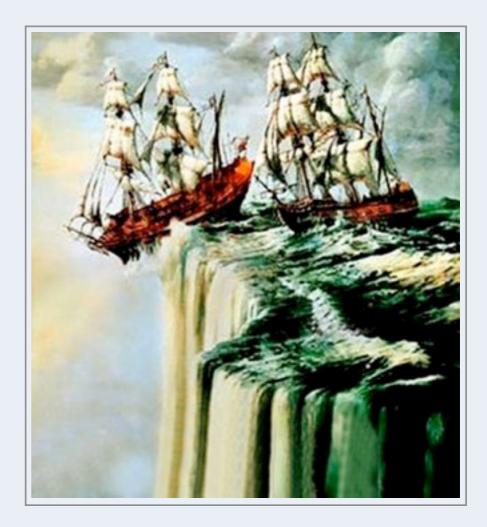
Beyond the Standard Model

Jernej F. Kamenik



Univerza v Ljubljani



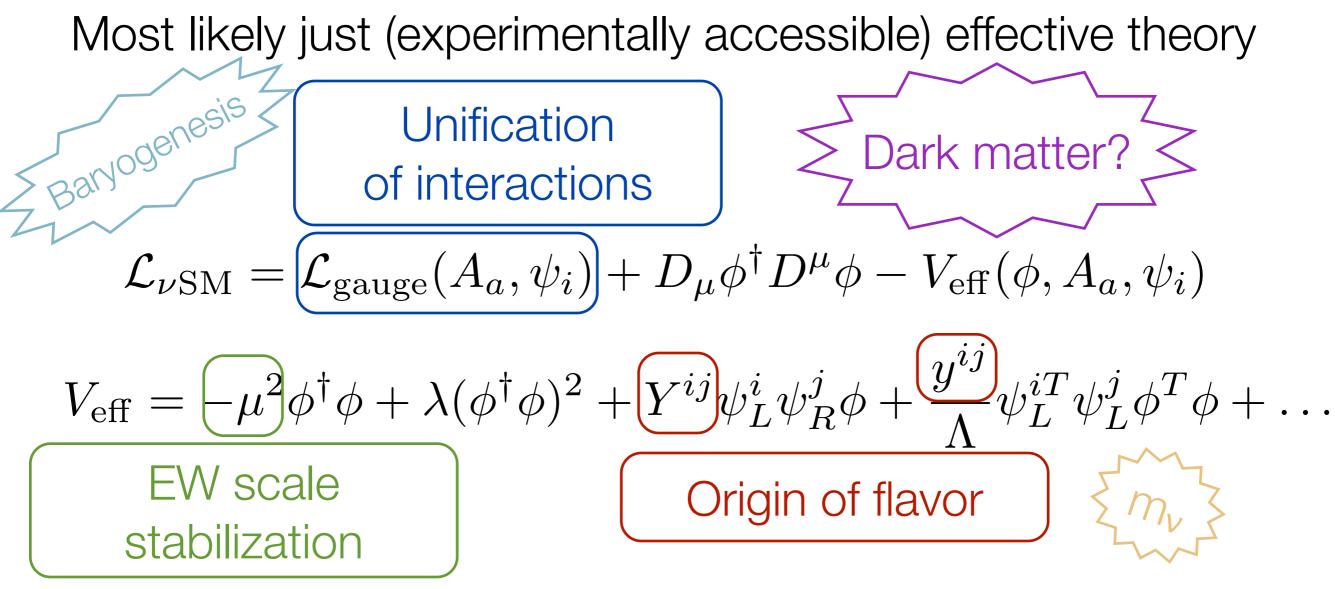


18/07/2014, Edinburgh

SM phenomenologically very successful

Most likely just effective theory

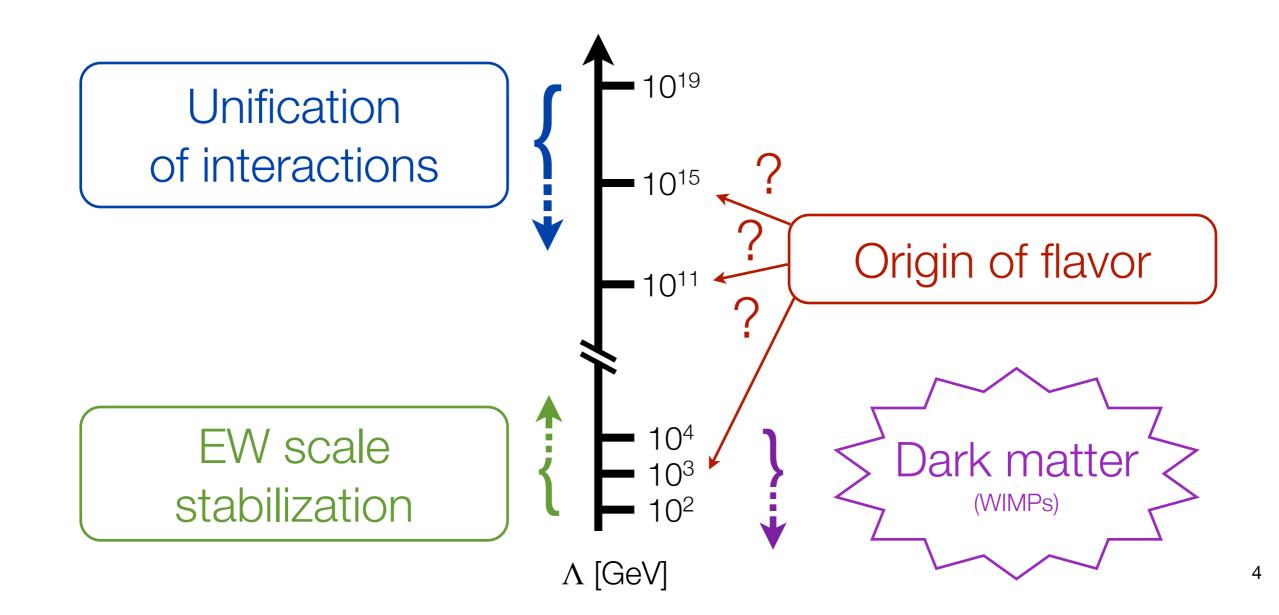
SM phenomenologically very successful



Need to understand/constrain size of <u>additional terms in</u> <u>series</u>

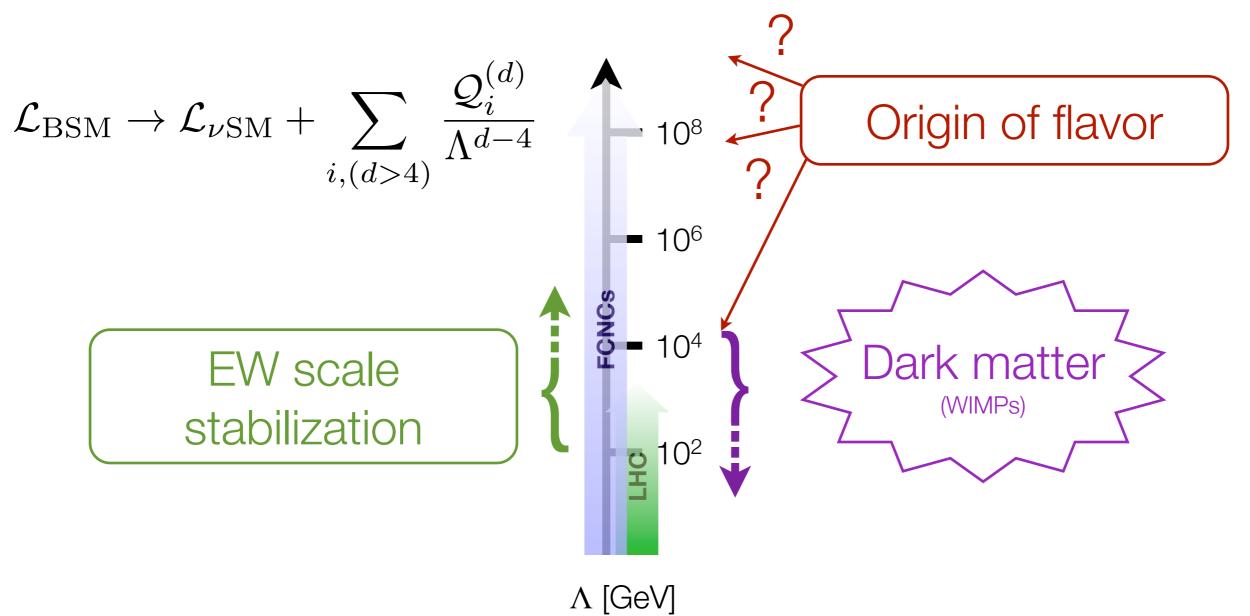
SM phenomenologically very successful

Most likely just (experimentally accessible) effective theory



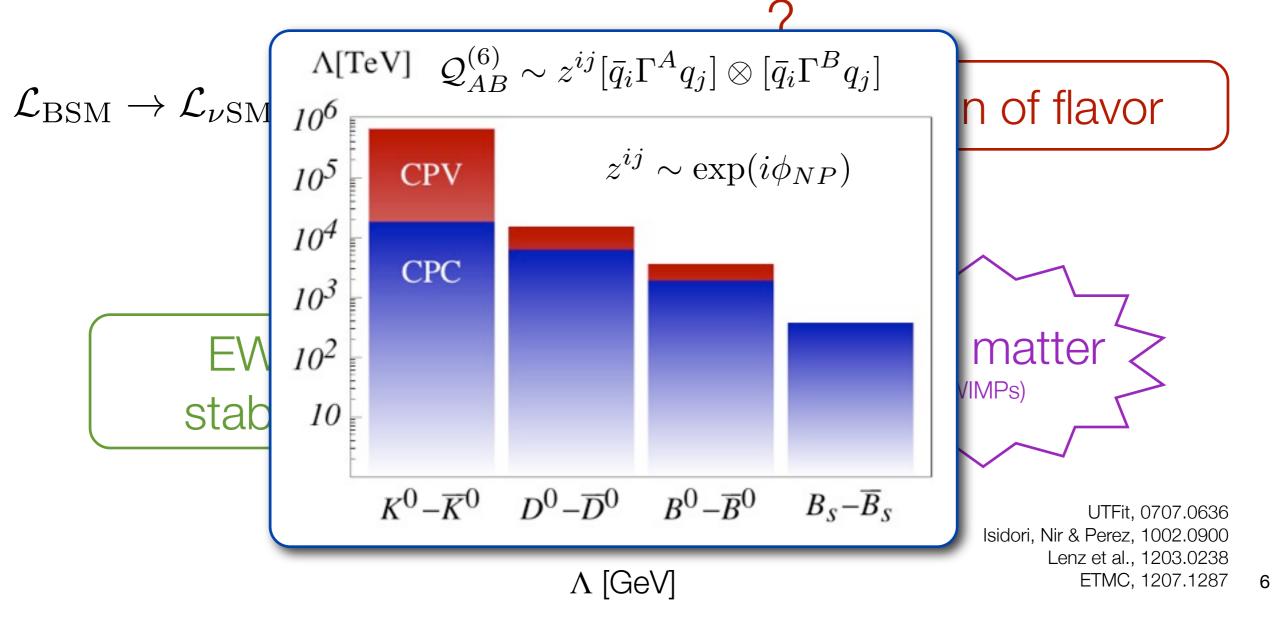
Twofold role of flavor (B) physics

(1) Indirect probe of BSM physics beyond direct reach



Twofold role of flavor (B) physics

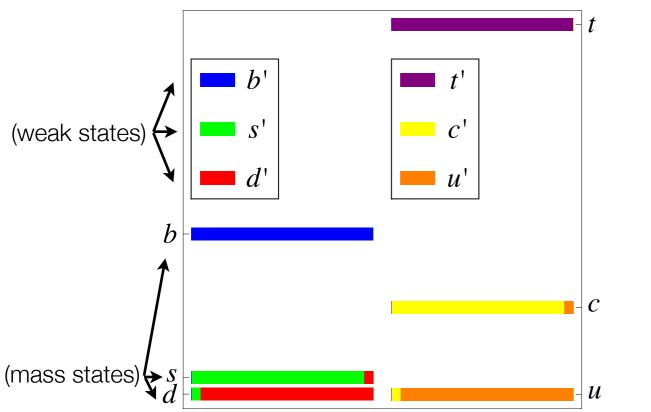
(1) Indirect probe of BSM physics beyond direct reach



Twofold role of flavor (B) physics

(2) Test sources of flavor symmetries & their violation

Suggestive pattern of masses and mixings



Accidental?

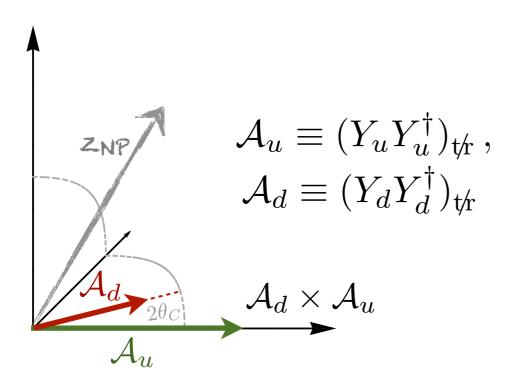
Dynamics?

Symmetries?

Twofold role of flavor (B) physics

(2) Test sources of flavor symmetries & their violation Global flavor symmetry of SM broken by Yukawas:

$$G_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$$

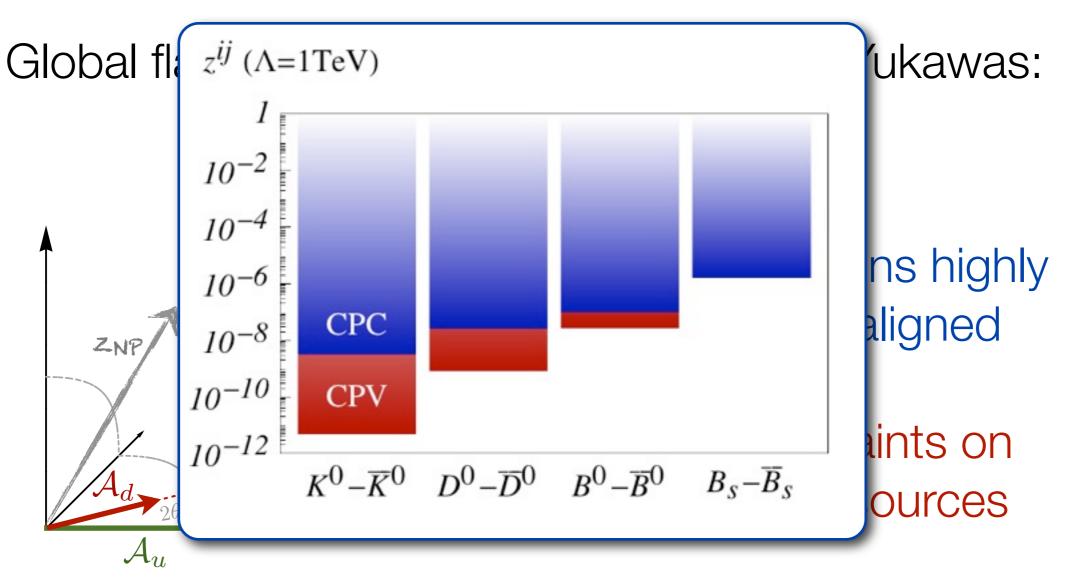


SM contributions highly hierarchical & aligned

Severe constraints on generic BSM sources

Twofold role of flavor (B) physics

(2) Test sources of flavor symmetries & their violation

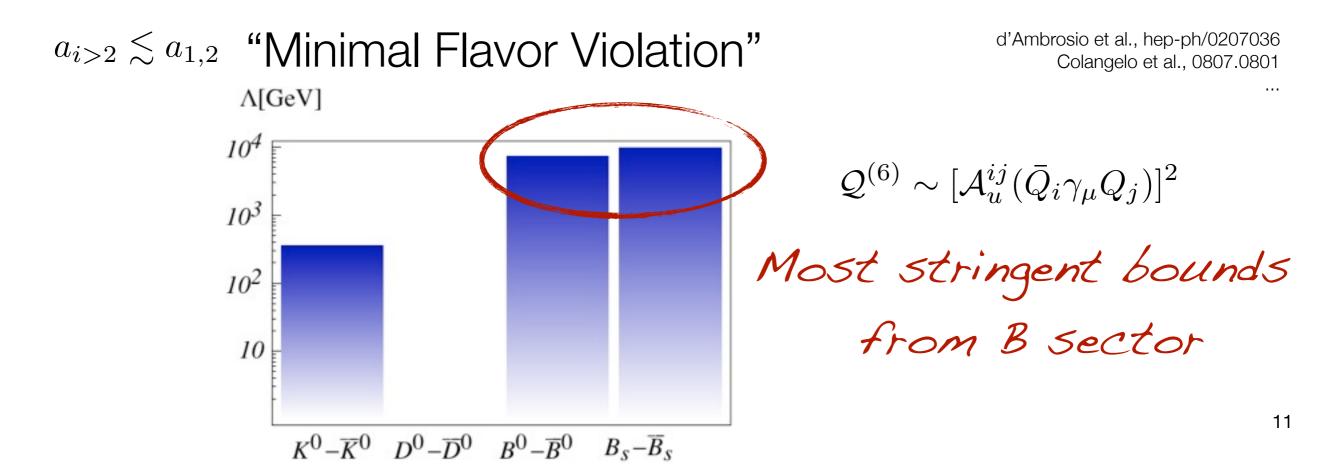


Global flavor symmetry of SM broken by Yukawas:

$$G_F = SU(3)_Q \times SU(3)_U \times SU(3)_D$$

Formally, NP flavor cannot be completely trivial $\int d^4x T \{Q_{\rm NP} \mathcal{H}_{\rm SM}\}$

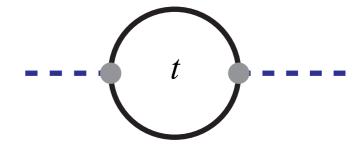
 $\mathbf{z} = \mathbf{1} + a_1 \mathcal{A}_u + a_2 \mathcal{A}_d + \dots$



Flavor triviality imposes degeneracy in NP spectra - problematic for naturalness@LHC

In SM, top Yukawa imposes largest fine-tuning in Higgs potential \Rightarrow

$$\delta m_h^2 \sim \frac{m_t^2}{v^2} \Lambda^2$$

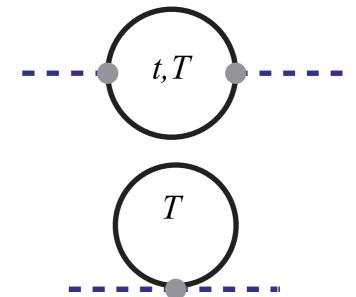


Flavor triviality imposes degeneracy in NP spectra - problematic for naturalness@LHC

In SM, top Yukawa imposes largest fine-tuning in Higgs potential \Rightarrow

$$\delta m_h^2 \sim \frac{m_T^2}{2} \Lambda^2 + \frac{m_t^2}{v^2} m_T^2 \log \frac{\Lambda^2}{m_T^2} + \dots$$

prefer light top partners ($m_T < 1 \text{TeV}$)



Flavor triviality imposes degeneracy in NP spectra - problematic for naturalness@LHC

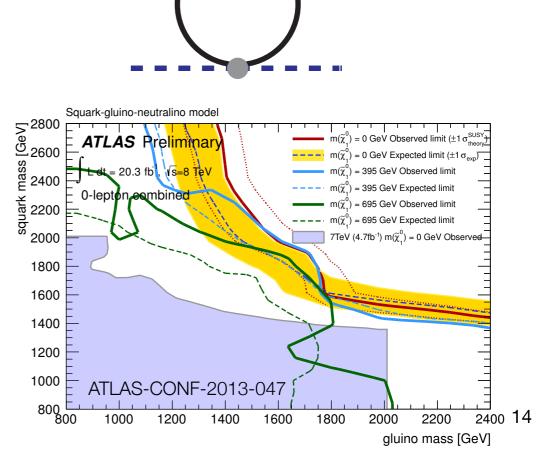
In SM, top Yukawa imposes largest fine-tuning in Higgs potential \Rightarrow

$$\delta m_h^2 \sim \frac{m_t^2}{2} \Lambda^2 + \frac{m_t^2}{v^2} m_T^2 \log \frac{\Lambda^2}{m_T^2} + \dots$$

prefer light top partners ($m_T < 1 \text{TeV}$)

avoiding flavor bounds though triviality \Rightarrow presence of u,d,... partners ($m_U \sim m_T$)

Strong LHC direct search constraints (MSSM example)



t, T

EW hierarchy stabilization only requires light 3rd generation partners \Rightarrow LHC bounds then imply flavor nontrivial spectra

Possible in flavor models mimicking the SM SU(3)/SU(2) flavor Kagan et al., 0903.1794 Buras & Girrbach, 1206.3878 Barbieri et al., 1105.2296

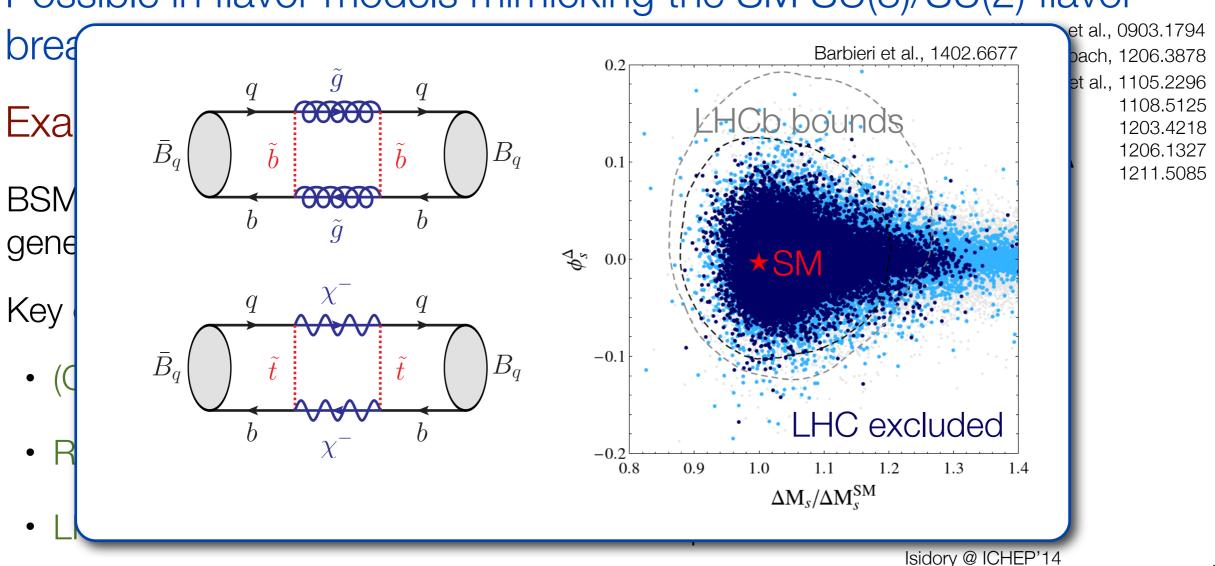
Example: natural SUSY

1211.5085 other BSM flavor effects mediated by 3rd t \widetilde{g} squark generation squarks (& sleptons) δī 1.5 TeV · Key observables: 1.0 TeV -• (CPV) in K(ε_K B mixing ($\Delta m_q, \phi_q$) W \widetilde{t}_{R} 0.5 TeV μ • Rare B decays $(B \rightarrow (X)l^+l^-, \nu\nu)$ $\widetilde{\mathbf{B}}$ LFV & EDMs

1108.5125

1203.4218 1206.1327

EW hierarchy stabilization only requires light 3rd generation partners \Rightarrow LHC bounds then imply flavor nontrivial spectra



Possible in flavor models mimicking the SM SU(3)/SU(2) flavor

EW hierarchy stabilization only requires light 3rd generation partners \Rightarrow LHC bounds then imply flavor nontrivial spectra

Possible in flavor models mimicking the SM SU(3)/SU(2) flavor Kagan et al., 0903.1794 Buras & Girrbach, 1206.3878

Example: natural SUSY

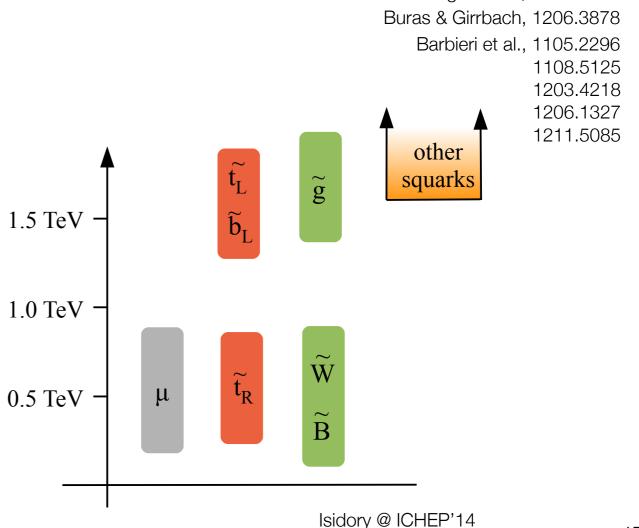
BSM flavor effects mediated by 3rd generation squarks (& sleptons)

Key observables:

• (CPV) in K(ε_K), B mixing ($\Delta m_q, \phi_q$)

• Rare B decays $(B \rightarrow (X)l^+l^-, \nu\nu)$

• LFV & EDMs



 $B_{s,d} o \mu^+ \mu^-$

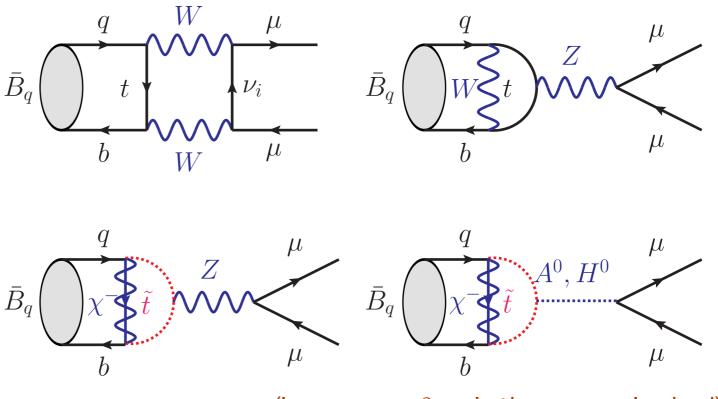
Particularly sensitive to FCNC scalar currents and FCNC Z penguins

Clean probe of the Yukawa interaction $(\Rightarrow$ Higgs sector)

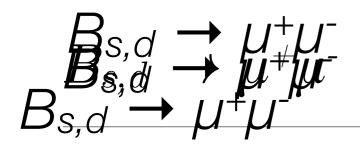
beyond tree level (virtually no long-distance contributions)

Leading SM contributions:

Examples in natural SUSY:



(large $tan\beta$ solutions excluded)



Particularly sensitive to FCNC scalar currents and FCNC Z penguins

Clean probe of the Yukawa interaction $(\Rightarrow$ Higgs sector)

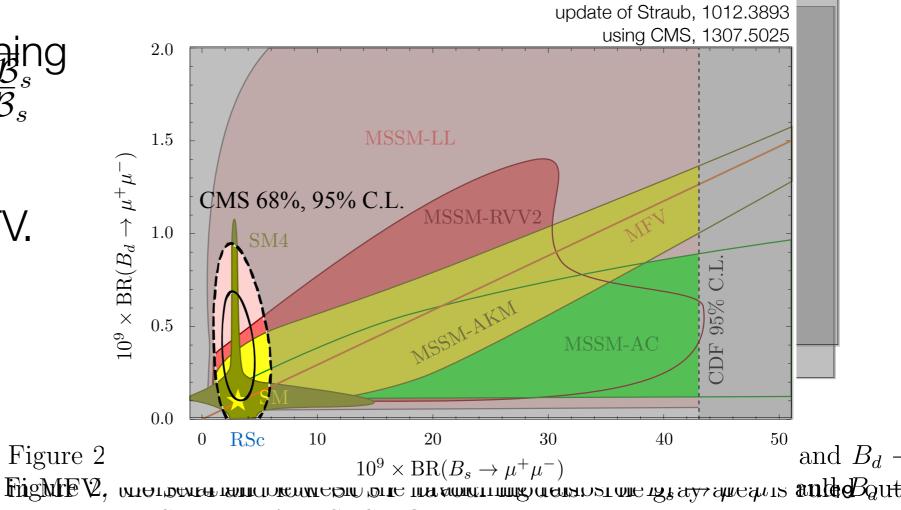
beyond tree level (virtually no long-distance contributions)

Latest results beginning to test possible $\mathcal{B}_d/\overline{\mathcal{B}}_s^s$ enhancement $\mathcal{B}_d/\overline{\mathcal{B}}_s$

Nontrivial test of MFV. Hurth et al., 0807.5039

Example:

SUSY flavor models



im MEMy. the SMM appointous SM Sted by our standels. The gray area is ruled out

 $B \rightarrow K^* \ell^+ \ell^-$ anomaly $b \rightarrow s \ \mu^+ \mu^-$

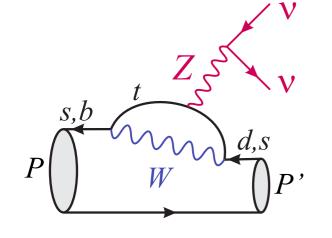
Similar to $B_s \to \mu^+ \mu^-$ in probing semileptonic $b \to s$ FCNCs and $\overline{\mathcal{Q}}_{\bar{9}}$ penguins $\mathcal{Q}_{L+R} \sim z_{ij} \bar{Q}^i \gamma_\mu Q^j (\bar{L} \gamma^\mu L + e_R \gamma^\mu e_R)$

<u>Caveat:</u> LD (charming penguin) contributions need to be controlled!

Complementary observables: $b \rightarrow s \nu \nu$

No non-local QCD contributions (but need precise form factor estimates)

c.f. Altmannshofer et al., 0902.0160



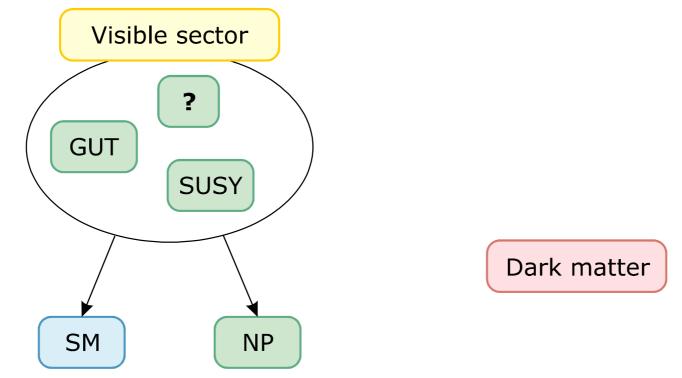
Key measurements for BelleII

Flavor portals to dark sector

Are there only SM particles at low-energy?

• Experimentally:

- Even very light states could be missed if very weakly interacting,
- There is dark matter in the Universe; it could be relatively light.
- Theoretically: Plenty of models predict new light particles
 - Pseudo-Goldstone scalars (axion, familon,...),
 - U(1) vectors (string, ED,...),
 - Hidden sectors & messengers (SUSY, mirror worlds,...)
 - Many others: millicharged fermions, dilaton, majoron, neutralino, sterile neutrino, gravitino,...

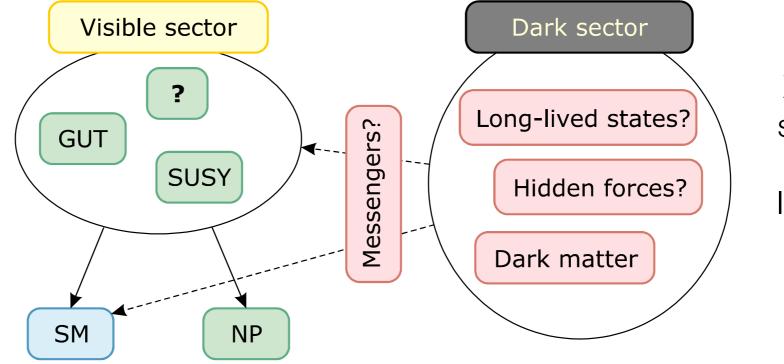


taken from C. Smith @ LPC - Clermont-Ferrand, 4/2012

 Heavy NP can be projected on to effective gauge-invariant operators built in terms of SM fields.

$$\mathcal{L}_{SM} + \frac{c_{\nu}}{\Lambda} (HL)^2 + \frac{c_i}{\Lambda^2} Q_i + \dots$$

Buchmuller & Wyler, Nucl.Phys. B268 (1986) 621 Grzadkowski et al., arXiv:1008.4884



X = dark sector state connected to the SM, or a light messenger.

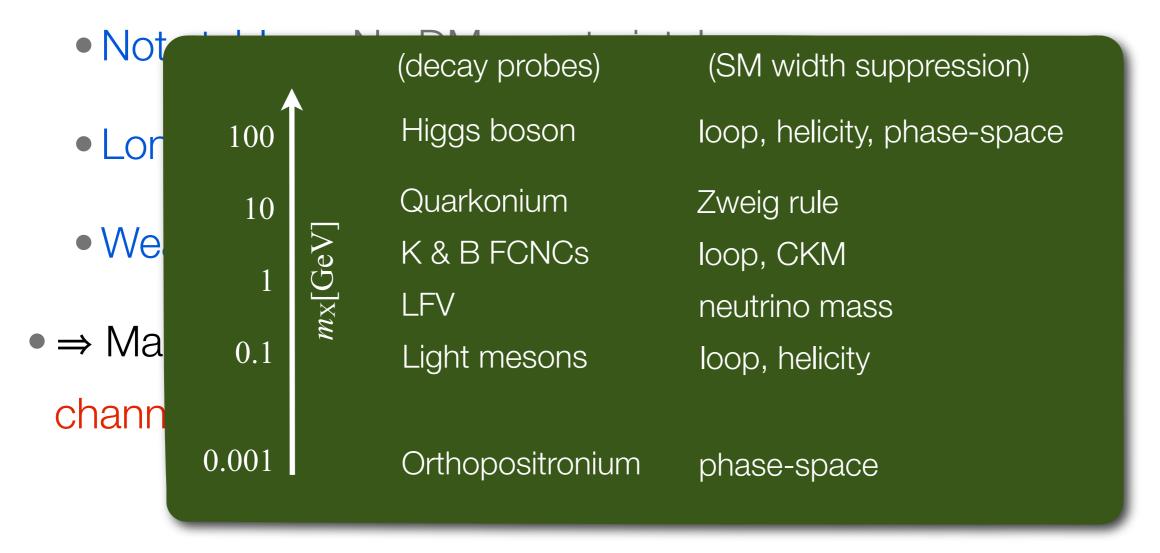
taken from C. Smith @ LPC - Clermont-Ferrand, 4/2012

• Take X as neutral, but include all possible interactions as SM gaugeinvariant effective operators.

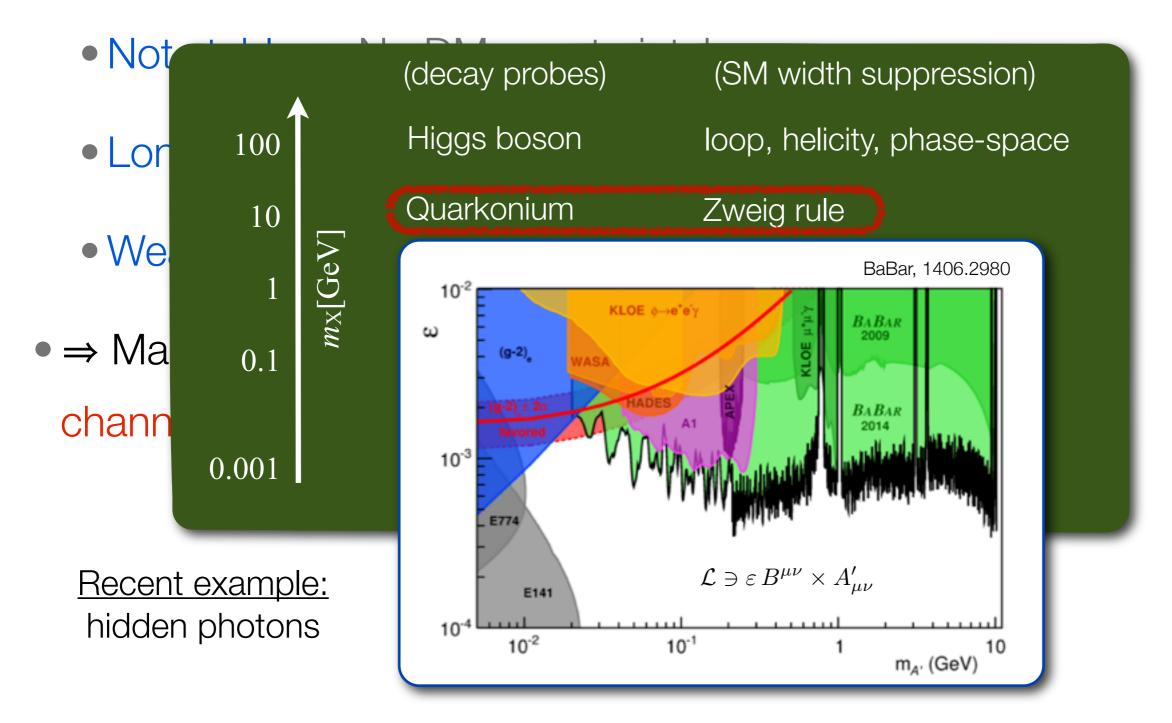
$$\mathcal{L}_{SM} + \frac{c_v}{\Lambda} (HL)^2 + \frac{c_i}{\Lambda^2} Q_i + \dots + \sum_{d \ge 3} \frac{c_i}{\Lambda^{d-4}} Q'_i + \dots$$

- Assumptions about the dark state X :
 - Not stable \Rightarrow No DM constraints!
 - Long-lived \Rightarrow Escapes as missing energy.
 - Weakly coupled \Rightarrow Does not affect SM processes.
- → Main impact is then to open new decay and production channels.

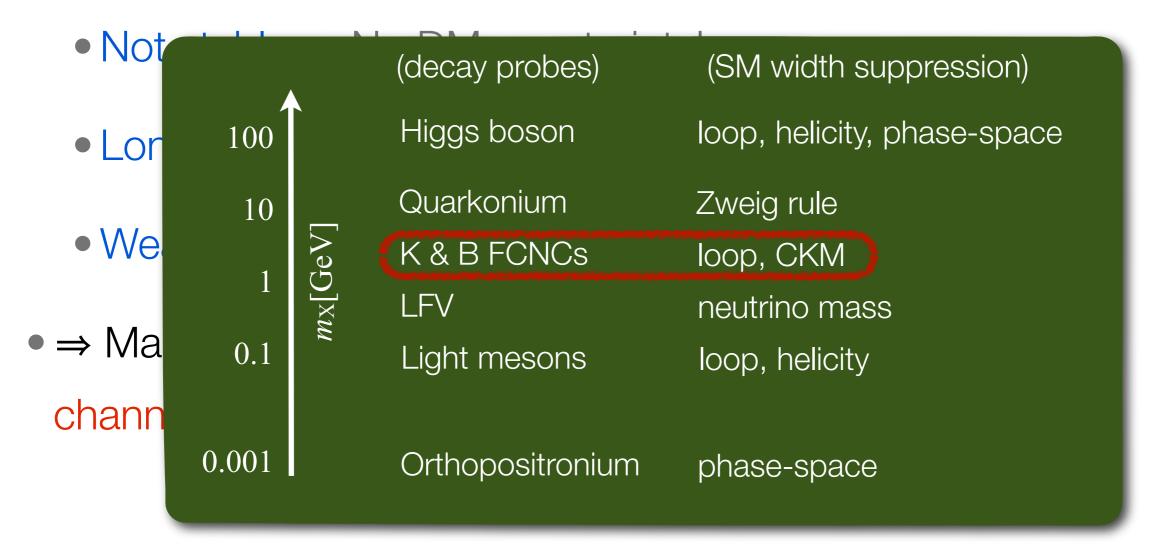
Assumptions about the dark state X :



Assumptions about the dark state X :



Assumptions about the dark state X :



Flavor probes of the invisible

• FCNC meson decays with E_{miss} CKM suppressed in SM

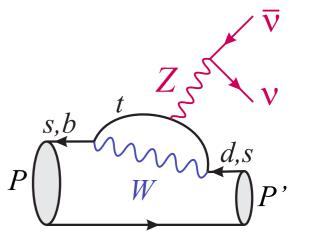
$$\begin{split} d^{I} \rightarrow d^{J}X : & \frac{g^{2}}{M_{W}^{2}} \frac{g^{2}}{16\pi^{2}} |V_{tI}^{*}V_{tJ}| , \\ \mathcal{B}(K \rightarrow \pi E_{miss}) \sim 10^{-11} & P \underbrace{Z \swarrow \mathbf{v}}_{W} \\ \mathcal{B}(B \rightarrow K^{(*)}E_{miss}) \sim 10^{-6} & P \underbrace{W }_{W} P, \end{split}$$

Flavor probes of the invisible

• FCNC meson decays with E_{miss} CKM suppressed in SM

$$d^{I} \rightarrow d^{J}X : c^{IJ} \frac{m_{I}^{n-6}}{\Lambda^{n-4}} \approx \frac{g^{2}}{M_{W}^{2}} \frac{g^{2}}{16\pi^{2}} |V_{tI}^{*}V_{tJ}| ,$$

(n-dim X-NP \approx SM)



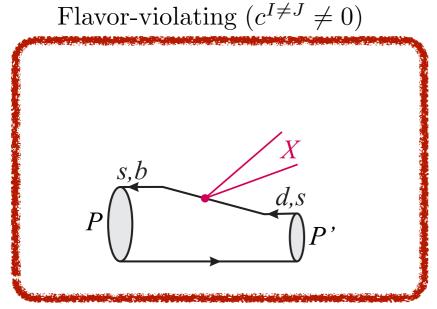
• Potentially very high X-operator scales probed:

 $c^{I \neq J} \sim \mathcal{O}(1)$

	n = 5	n = 6	n = 7
$s \to d$	$3.3 \cdot 10^7 \text{ TeV}$	$130 { m TeV}$	$2.0 { m TeV}$
$b \to d$	$1.3 \cdot 10^5 { m TeV}$	$26 { m TeV}$	$1.5 { m TeV}$
$b \rightarrow s$	$2.7 \cdot 10^4 \text{ TeV}$	$12 { m TeV}$	$0.9~{\rm TeV}$

$$\mathcal{H}_{eff}(q^I \to q^J X) = \frac{c^{IJ}}{\Lambda^n} \bar{q}^I q^J \times X$$

Flavor - based classification of dark operators

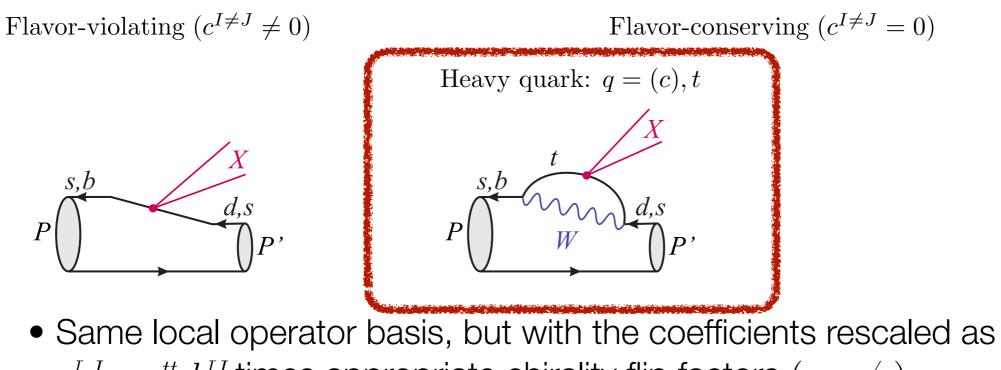


• Bounds directly derived from $d_I \rightarrow d_J X$ processes.

• When MFV holds, $c^{IJ} \sim \lambda^{IJ}$ times appropriate chirality flip factors $(m_{I,J}/v)$.

$$\lambda^{IJ} = \mathbf{Y}_{u}^{\dagger} \mathbf{Y}_{u} \approx V_{tI}^{*} V_{tJ} \rightarrow \begin{cases} \lambda^{sd} \approx (-3.1 + i1.3) \times 10^{-4} ,\\ \lambda^{bd} \approx (7.8 - i3.1) \times 10^{-3} ,\\ \lambda^{bs} \approx (-4.0 - i0.07) \times 10^{-2} \end{cases}$$

Flavor - based classification of dark operators



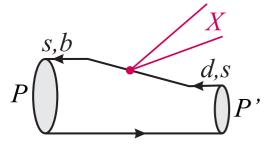
 $c^{IJ} \sim c^{tt} k^{IJ}$ times appropriate chirality flip factors $(m_{I,J}/v)$.

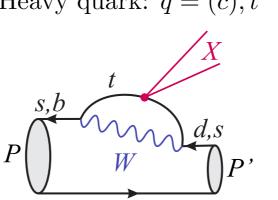
$$k^{IJ} = \frac{g^2}{16\pi^2} \lambda^{IJ} \rightarrow \begin{cases} k^{sd} \approx (-0.8 + i0.4) \times 10^{-6} \\ k^{bd} \approx (2.1 - i0.8) \times 10^{-5} \\ k^{bs} \approx (-1.1 - i0.02) \times 10^{-4} \end{cases}$$

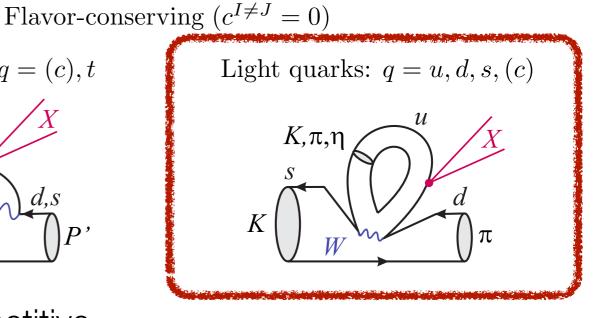
Flavor - based classification of dark operators

Flavor-violating $(c^{I\neq J}\neq 0)$

Heavy quark: q = (c), t







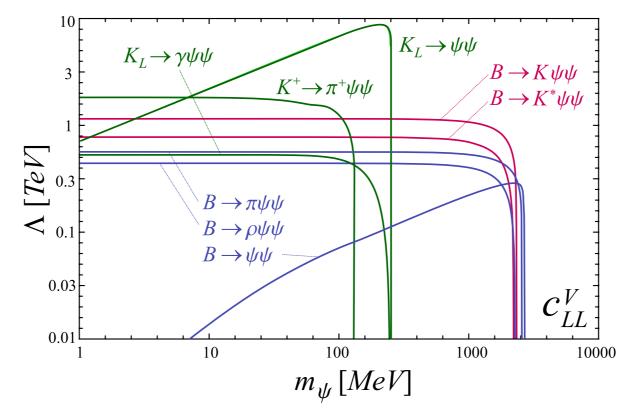
- Due to small V_{ub} , B decays not competitive.
- For *K* decays, *q* = *u* contributions are dominant but non local, and require controlling long-distance hadronic effects.

Beyond the scaling argument: Kinematics

Importance of combining several available modes

MFV example:

$$\frac{V_{ti}^* V_{tj}}{\Lambda^2} \bar{Q}^i \gamma_\mu Q^j \times \bar{\psi}_L \gamma^\mu \psi_L$$



Current bounds assuming flat exp. acceptances

Observational hints



Several recent intriguing flavor anomalies

$$\begin{array}{ccc} \Delta A_{CP} & P_5' & R_K \\ & A_{sl}^b & & R_{\tau/\ell} \end{array} & \begin{array}{c} R_K \\ \end{array} \end{array}$$

 P'_5

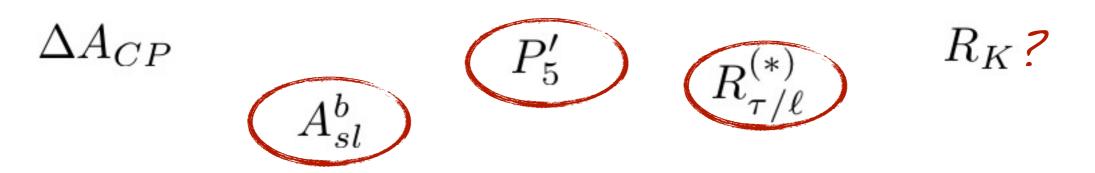


 A^b_{sl}

Theorists often taken by surprise ⇒ Don't listen to us too much







Theorists often taken by surprise

Proposed BSM explanations not always very elegant





 A^b_{sl}



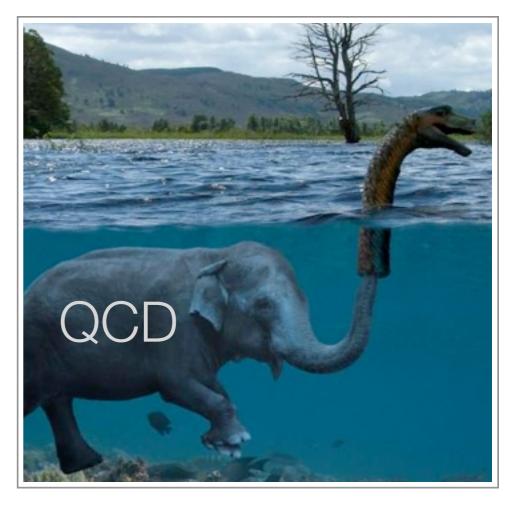
 $R_{\tau/\ell}^{(*)}$

 R_K

Theorists often taken by surprise

Proposed BSM explanations not always very elegant

Prompt reevaluation of possible SM effects



Theorists often taken by surprise

Proposed BSM explanations not always very elegant

Prompt reevaluation of possible SM effects

Unfortunately, some deviations of limited duration



If any such hint would be confirmed, tremendous implications for high prsearches Flavor as guide to high p_{T}

In SM weak charged current interactions are lepton flavor universal e^{-v}

• Tested directly at colliders via *W* decays ~1%

Additional charged (scalar) interactions could induce LFU violation in processes at low energies

Can be predicted accurately even in hadronic processes, since most QCD uncertainties cancel in ratios

 Pion, kaon, D processes well consistent with LFU expectations ~(0.1-1)%

c.f. HFAG, 1010.1589

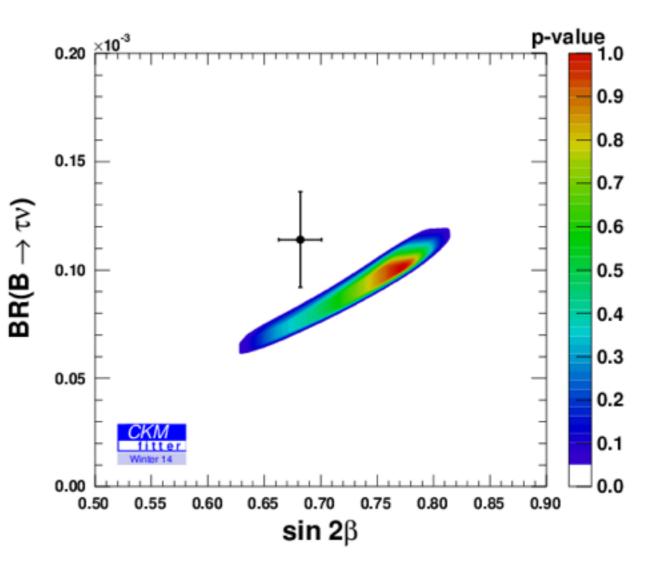
W-

Apparent tension in global CKM unitarity fits

Discrepancy between $|V_{ub}|$ determinations see talks by Kwon, Bharucha

Most pronounced for taunic *B* decay

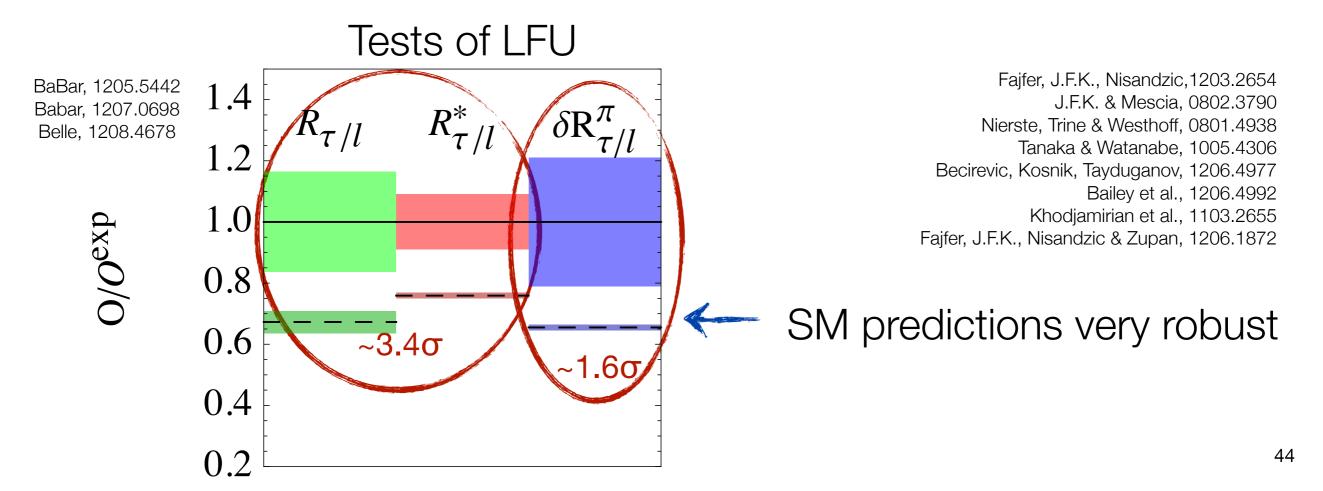
Somewhat reduced with updated Belle result Belle, 1208.4678



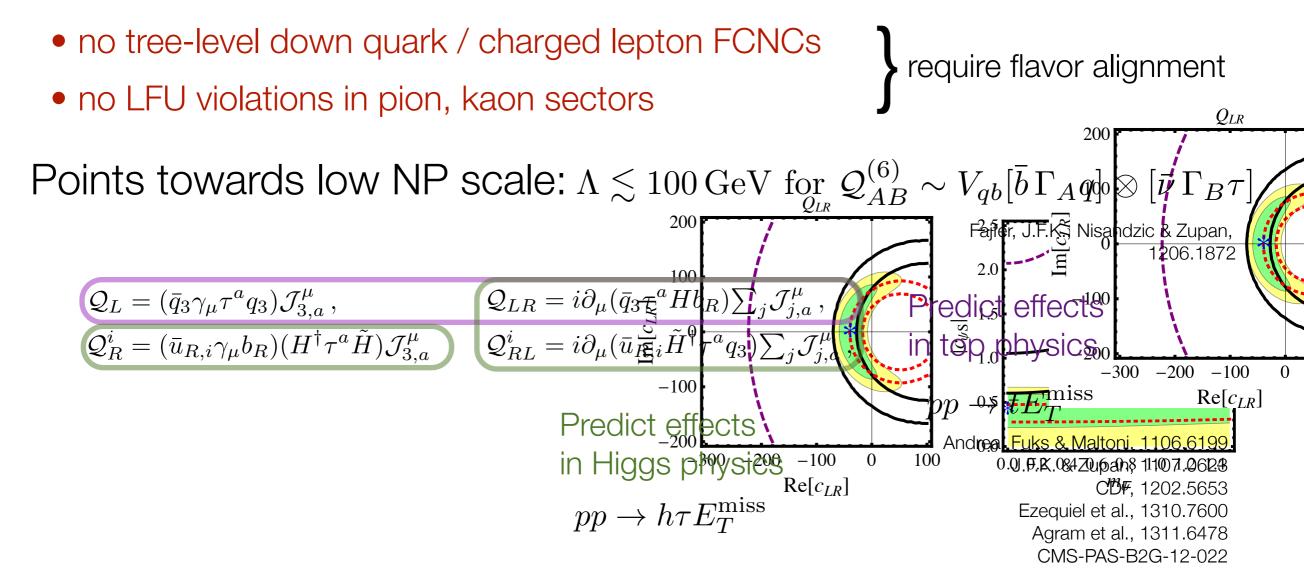
However, maybe not CKM issue at all

Can eliminate $|V_{ub}|$ in ratio $\Delta \mathcal{R}^{\pi}_{\tau/\ell} \equiv \frac{\tau(B^0)}{\tau(B^-)} \frac{\mathcal{B}(B^- \to \tau^- \bar{\nu})}{\Delta \mathcal{B}(\bar{B}^0 \to \pi^+ \ell^- \bar{\nu})}$

Similarly in semitauonic decays $\mathcal{R}_{\tau/\ell} \equiv \frac{\mathcal{B}(B \to D \tau \nu)}{\mathcal{B}(B \to D \ell \nu)} \quad \mathcal{R}^*_{\tau/\ell} \equiv \frac{\mathcal{B}(B \to D^* \tau \nu)}{\mathcal{B}(B \to D^* \ell \nu)}$



Can it be NP? Need to satisfy severe constraints:



Some of these implications remain largely unexplored Interesting future research directions

Conclusions

Success of SM in describing flavor-changing processes implies that large new sources of flavor symmetry breaking at TeV scale are mostly excluded.

However, NP at TeV scale need not be flavor trivial!

If (properly aligned) new sources of flavor breaking present

- Precision B observables may hide NP signals @10% level in well motivated NP models (natural SUSY)
- can significantly affect & guide NP searches high p_T
- have implications for EW fine-tuning

Conclusions

Success of SM in describing flavor-changing processes implies that large new sources of flavor symmetry breaking at TeV scale are mostly excluded.

However, NP at TeV scale need not be flavor trivial!

There are sectors of the theory that are just starting to be tested

- Measurements of $B_{s,d} \rightarrow \mu^+ \mu^-$ probe the Higgs Yukawa sector at loop level
- Several recent anomalies in rare *B* decays need to be clarified.

Worth to improve searches of exotic flavor-violating effects!

• Example: $b \rightarrow s E_{miss}$ modes as portals to dark BSM sectors



Thank you!