

Sterile neutrinos at TLEP

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FCC-ee/TLEP physics workshop (TLEP7)

5 types of sterile neutrinos

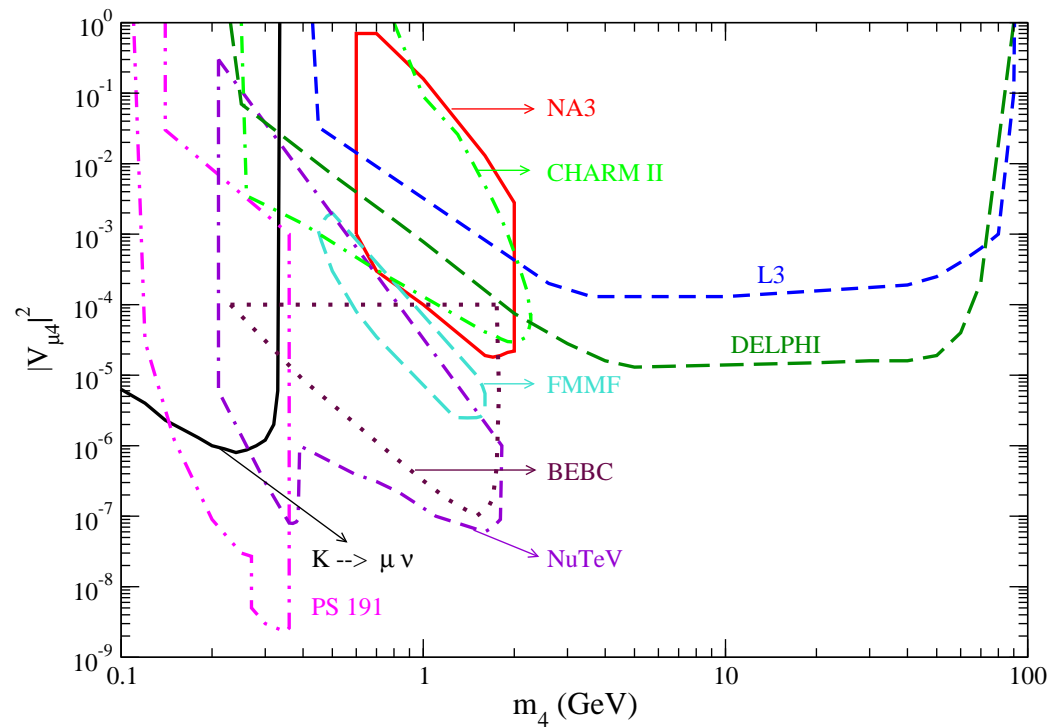
sterile neutrino for $M \gg 1 \text{ eV}$ = HNL - heavy neutral lepton

- just so HNL : if there are left-handed neutrinos why not to have right-handed neutrinos?
- eV scale sterile neutrinos, motivated by short base-line neutrino anomalies (LSND, etc).
- HNL, giving rise to active neutrino masses
- HNL, giving rise to active neutrino masses and leading to baryon asymmetry of the universe
- HNL, giving rise to active neutrino masses, leading to baryon asymmetry of the universe, and to dark matter

Interesting mass range: $M_N < M_Z$, to search for in Z -decays.

Interaction strength : mixing angle U^2

Survey of constraints on just so HNL



From arXiv:0901.3589, Atre et al

HNL and neutrino masses

Notations: M_D - Dirac mass, M - Majorana mass, F - Yukawa coupling, v - Higgs vev

One flavour see-saw formula:

$$m_\nu \simeq \frac{m_D^2}{M} \simeq \frac{F^2 v^2}{M} = U^2 M \quad \Rightarrow \quad U^2 = \frac{m_\nu}{M} \simeq 10^{-10} \frac{\text{GeV}}{M}$$

Naive expectation: If HNL gives rise to neutrino masses, it would not be possible to find them at TLEP, even though the number of Z -decays will be 5 orders of magnitude larger than at LEP (10^{12} versus 10^7).

LEP limit: $U^2 < \text{few} \times 10^{-5}$.

Naive expectation is wrong!

We need at least **2** HNL to explain atmospheric and solar mass differences.

Consequence: instead of equality we have only lower bound!

$$U^2 > \frac{m_\nu}{M} \simeq 10^{-10} \frac{\text{GeV}}{M}$$

Physics of large mixing angles

MS '06; Kersten and Smirnov '07:

Consider SM + one extra massive **Dirac** spinor Ψ , which is singlet with respect to SM.

$$L = L_{SM} + \bar{\Psi} i \partial_\mu \gamma^\mu \Psi - F_\alpha \bar{L}_\alpha \Psi H - M \bar{\Psi} \Psi + h.c.,$$

Symmetry: lepton number conservation. For **any** F_α and M all active neutrinos are massless.

Small symmetry breaking terms \implies small active neutrino masses :

$$\Delta L = f_\alpha \bar{L}_\alpha \Psi^c H - m \bar{\Psi} \Psi^c + h.c.,$$

Active neutrino masses:

$$m_\nu \simeq \frac{Fv^2 \times f}{M} + \frac{F^2v^2 \times m}{M^2}$$

Consequence: **any** increase of experimental sensitivity may lead to discovery of HNL responsible for active neutrino masses!

Potential TLEP limit: $U^2 < \text{few} \times 10^{-10}$ for a range of masses above **20 GeV**, see below

HNL, ν masses and BAU

Akhmedov, Rubakov, Smirnov '98; Asaka, MS '05

HNL oscillations as a source of baryon asymmetry. Qualitatively:

- HNL are created in the early universe and oscillate in a coherent way with CP-breaking.
- Lepton number from HNL can go to active neutrinos.
- The lepton number of active left-handed neutrinos is transferred to baryons due to equilibrium sphaleron processes.

Sakharov condition

Rate of HNL equilibration $\Gamma \simeq \kappa F^2 T$ must be smaller than the rate of the Universe expansion at the sphaleron freeze-out $T = T_{sph} \simeq 130$ GeV, $H \simeq T^2/M_0$, $M_0 \sim M_{Pl}$ ($\kappa \simeq 3 \times 10^{-6}$ - some number following from solution of kinetic equations in the early universe):

$$\kappa F^2 \left(1 - \frac{M^2}{M_W^2}\right)^2 T_{sph} < \frac{T_{sph}^2}{M_0}$$

Numerically, $F < 8 \times 10^{-6}$, and

$$U^2 < 2 \times 10^{-6} \left(\frac{GeV}{M}\right)^2 \left(1 - \frac{M^2}{M_W^2}\right)^2$$

Expectations for TLEP

Processes: $Z \rightarrow N\nu$, $N \rightarrow lq\bar{q}$ (lepton + meson, lepton + 2 quark jets),

$$BR(Z \rightarrow \nu N) \simeq BR(Z \rightarrow \nu\nu)U^2, \quad \Gamma_N \simeq \frac{G_F^2 M^5}{192\pi^3} U^2 A$$

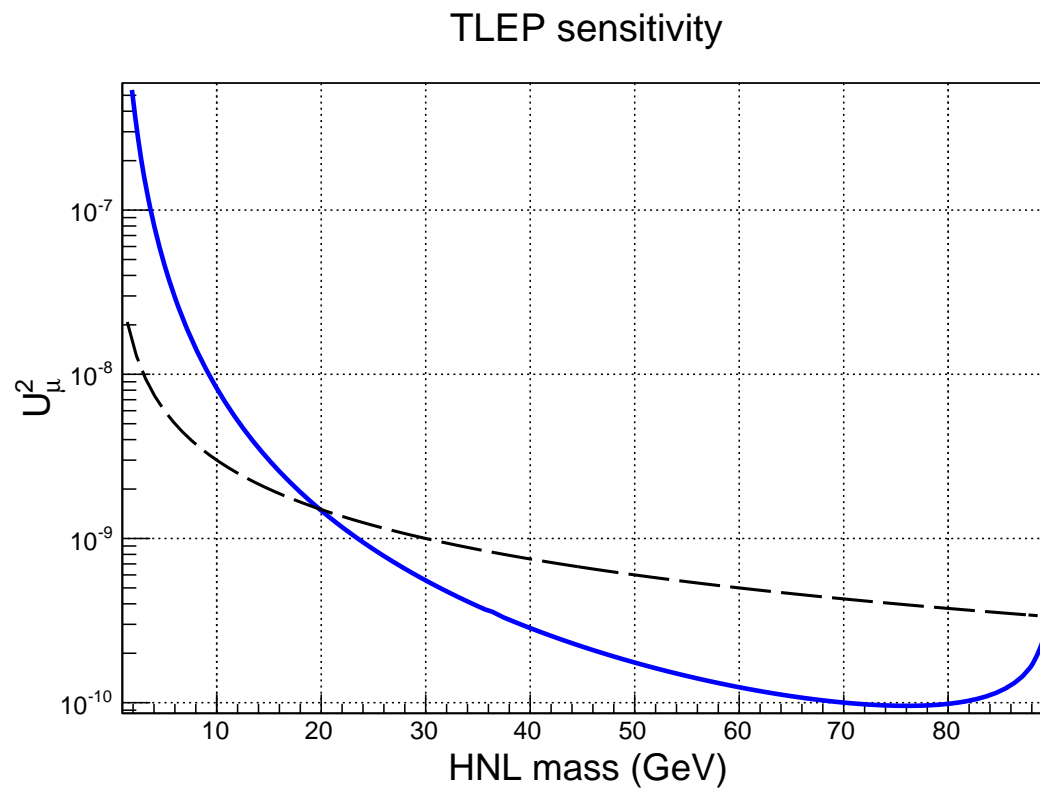
Coefficient A counts the number of open channels, $A \sim 10$ for $M > 10$ GeV

Assumptions: 10^{12} Z-decays in 3 m detector.

- “short lived” N: decay length < 3 m \implies constraint on U^2 may go down to $U^2 < 10^{-10}$ as the sensitivity will grow as the number of Z-decays! This works for $M \gtrsim 20$ GeV.
- “long lived” N: decay length exceeds the size of the detector \implies constraint on U^2 may go down to $U^2 < 4 \times 10^{-8}$ as the sensitivity will grow as the square root of the number of Z-decays. This works for lighter HNL.

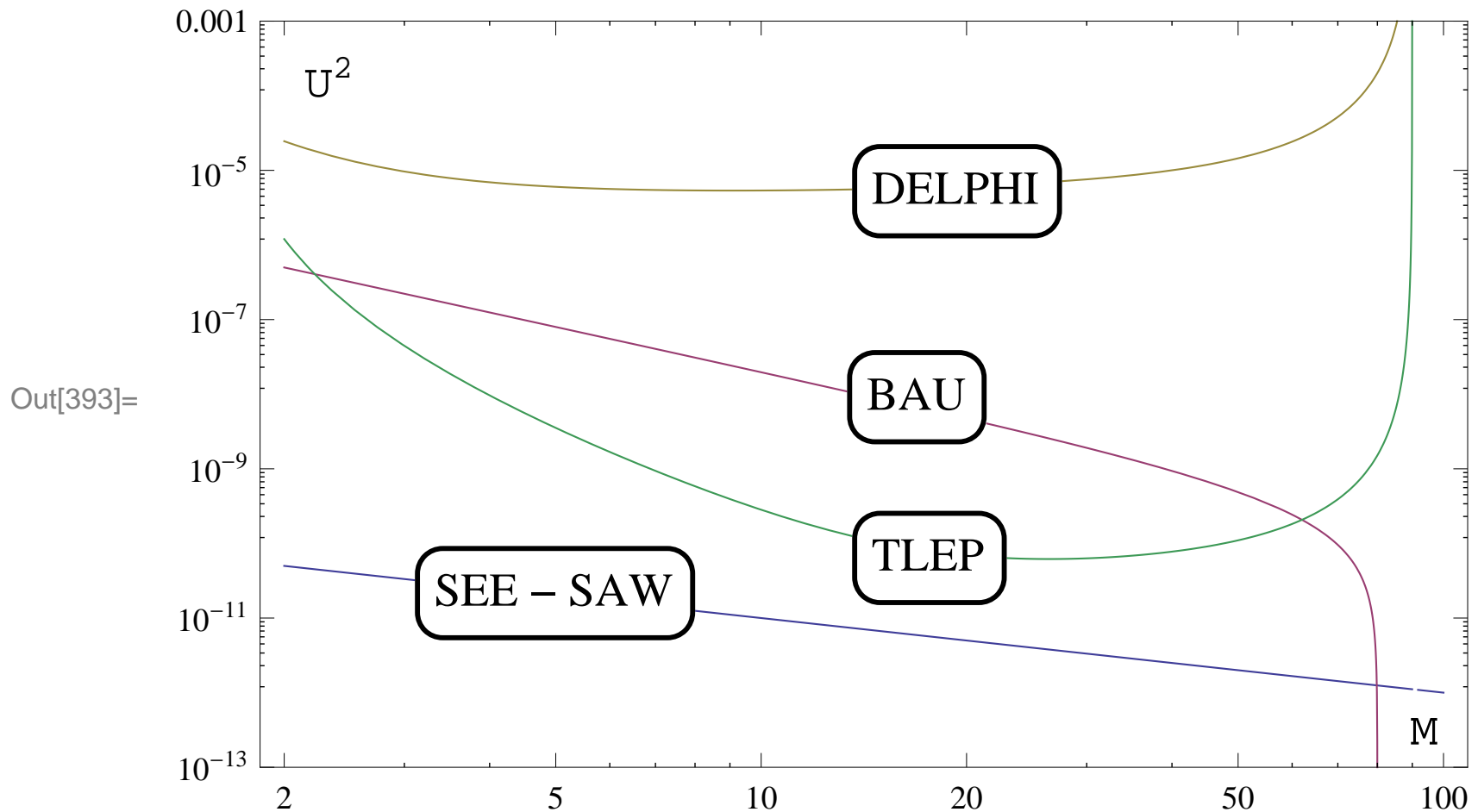
Sensitivity plot, very preliminary

Elena Graverini, Nico Serra, Walter Bonivento, 20.06.2014

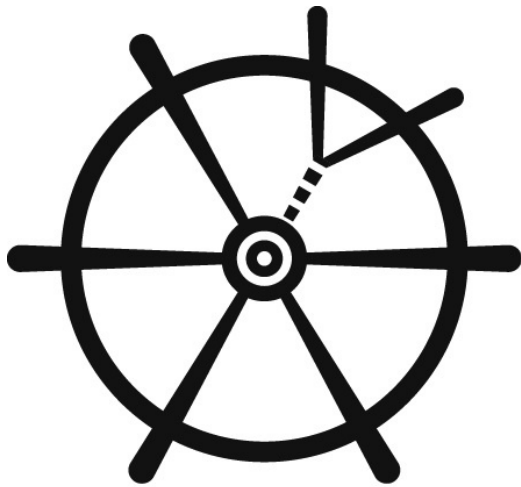


Sensitivity plot, very preliminary

MS, 20.06.2014



SHiP sensitivity

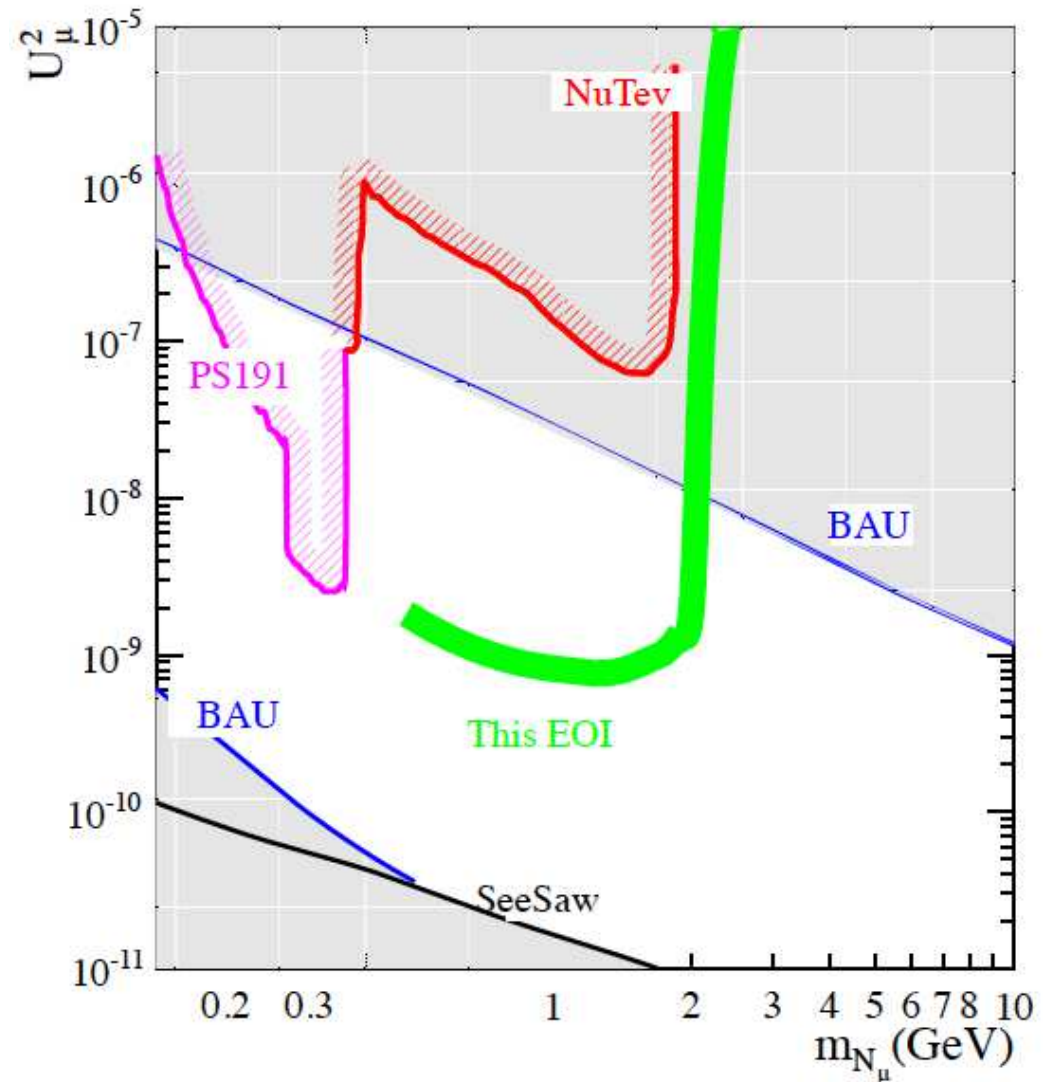


SHiP

Search for Hidden Particles



Imperial College
London



Conclusions

- TLEP can reveal the origin of neutrino masses if $M > 2 \text{ GeV}$ and the origin of baryon asymmetry of the Universe if $M > 3 \text{ GeV}$ (very indicative number!!!)
- Fix target experiment SHIP and collider experiments at TLEP can search for HNL in complementary regions: for SHIP $M < 2 \text{ GeV}$.