Top pair production near threshold

M. Beneke (TU München)

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Outline

- · Introduction and theoretical framework
- Top production near threshold in e⁺e⁻ collisions QCD effects and parameter variations

MB, Kiyo, Schuller, 1312.4791 [hep-ph] and in preparation;
 MB, Piclum, Rauh, 1312.4792 [hep-ph];
 MB, ..., Steinhauser, in preparation; MB, Piclum, Rauh, in progress.

Beyond QCD – Non-resonant effects

MB, Jantzen, Ruiz-Femenia 1004.2188 [hep-ph]; Jantzen, Ruiz-Femenia 1307.4337 [hep-ph]





Pair production threshold - Coulomb force and Weak Decay

Ultra-precise mass measurement Unique QCD dyanmics



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Smallest <u>structure</u> in particle physics known to exist (10^{-17}m) . Direct "spectroscopic" mass and width measurement.

Additional smearing of the resonance due to beam luminosity spectrum (collider-specific) and ISR



Most recent study for ILC/CLIC [Seidel, Simon, Tesai, Poss, 2013] assumes 10 fb⁻¹ at 10 points.

 $[\delta m_t]_{\text{thr}} = 27 \,\text{MeV}$ [simultaneous fit of α_s]

Theory

Non-perturbative but weak coupling. Expansion in α_s and $v = \sqrt{\frac{E}{m}} = \sqrt{\frac{\sqrt{q^2} - 2m_t}{m_t}}$, while $\alpha_s/v = O(1)$

$$R \sim v \sum_{k} \left(\frac{\alpha_s}{v}\right)^k \cdot \left\{ 1 \text{ (LO)}; \alpha_s, v \text{ (NLO)}; \alpha_s^2, \alpha_s v, v^2 \text{ (NNLO)}; \dots \right\}$$
$$(q_{\mu}q_{\nu} - q^2 g_{\mu\nu}) \Pi(q^2) = i \int d^4x e^{iq \cdot x} \langle 0|T(j_{\mu}(x)j_{\nu}(0))|0\rangle, \qquad j^{\mu}(x) = [\bar{Q}\gamma^{\mu}(\gamma_5)Q](x)$$

Summation through Schrödinger equation.



Tools

Non-relativistic effective field theory and threshold expansion (defines the matching procedure!) See arXiv:1312.4791 [hep-ph]

Relevant scales: $m_t \approx 175 \text{ GeV}$ (hard), $m_t \alpha_s \approx 30 \text{ GeV}$ (soft, potential) and the ultrasoft scale (us) $m_t \alpha_s^2 \approx 2 \text{ GeV}$.

$$\mathcal{L}_{\text{QCD}} \left[Q(h, s, p), g(h, s, p, us) \right] \qquad \mu > m_t$$

$$\downarrow$$

$$\mathcal{L}_{\text{PNRQCD}} \left[Q(p), g(us) \right] \qquad \mu < m_t v$$

See mulit-loop fixed-order calculations to match PNRQCD, then perturbation theory in Coulomb background in PNRQCD. Can be extended systematically to any order. 3rd order is current technological limit.

$$\begin{split} \Pi^{(\nu)}(q^2) &= \frac{N_c}{2m^2} c_{\nu} \left[c_{\nu} - \frac{E}{m} \left(c_{\nu} + \frac{d_{\nu}}{3} \right) \right] G(E) + \dots \\ G(E) &= \frac{i}{2N_c(d-1)} \int d^d x \, e^{iEx^0} \left\langle 0 \right| T(\left[\chi^{\dagger} \sigma^i \psi\right](x) \left[\psi^{\dagger} \sigma^i \chi\right](0)) \left| 0 \right\rangle_{|\text{PNRQCD}} \,, \end{split}$$

- Bound state quantities (S-wave)
 - E_n Kniehl, Penin, Smirnov, Steinhauser (2002); MB, Kiyo, Schuller (2005); Penin, Sminrov, Steinhauser (2005)
 - $|\psi_n(0)|^2$ MB, Kiyo, Schuller (2007); MB, Kiyo, Penin (2007)
- Matching coefficients
 - a₃ Anzai, Kiyo, Sumino (2009); Smirnov, Sminrov, Steinhauser (2009)
 - C3 Marquard, Piclum, Seidel, Steinhauser (2013) [2009]
- Continuum (PNRQCD correlation function)
 - ultrasoft MB, Kiyo (2008)
 - potential MB, Kiyo, Schuller, in preparation (2014) [2007]
 - P-wave MB, Piclum, Rauh (2013)

Note: logarithmically enhanced 3rd order terms known before or resummed [Hoang et al. 2001-2013; Pineda et al. 2002-2007]

2nd order available since end of 1990s.

I. Top production near threshold in e^+e^- collisions – QCD effects and parameter variations

MB, Kiyo, Schuller, 1312.4791 [hep-ph] and in preparation; MB, ..., Steinhauser, in preparation (third-order QCD); MB, Rauh, Piclum, in preparation (Higgs, top-Yukawa interaction) • Mass definition: Potential-subtracted mass [MB, 1998]

$$m_{\rm PS}(\mu_f) \equiv m_{\rm pole} + \frac{1}{2} \int_{|\vec{q}| < \mu_f} \frac{d^3 \vec{q}}{(2\pi)^3} \, \tilde{V}_{\rm Coulomb}(\vec{q})$$

Cancellation of large perturbative contributions from the IR. In the following: $m_{t,PS}(20 \text{ GeV}) = 171.5 \text{ GeV}.$



QCD at 3rd order, P-wave (Z only) and Higgs at NLO \simeq NNNLO QCD

Cross section near threshold, from 2nd to 3rd order



 $m_{t,PS}(20 \text{ GeV}) = 171.5 \text{ GeV}, \Gamma_t = 1.33 \text{ GeV}$





No convergence for $\mu \lesssim 50 \, {\rm GeV}$.



Compare recent NNLO+(N)NLL analysis [Hoang, Stahlhofen, 2013] Coulomb corrections only (for $m_{t,PS} = 175$ GeV, $\Gamma_t = 1.5$ GeV). Scale depedence at third order and exact solution.

Parameter dependence





II. Beyond QCD - non-resonant effects

MB, Jantzen, Ruiz-Femenia 1004.2188 [hep-ph]; Jantzen, Ruiz-Femenia 1307.4337 [hep-ph]

Finte-width divergences and "electroweak effects"

The QCD-only result usually discussed is far from reality

• Finite-width divergences (overall log divergence, already at NNLO):

Since $E = \sqrt{s} - 2m_t + i\Gamma$, the divergence survives in the imaginary part:

$$\mathrm{Im} \left[\delta G(E)\right]_{\mathrm{overall}} \propto m_t \times \frac{\alpha_s \alpha_{ew}}{\epsilon}$$

• Electroweak effect. Must consider $e^+e^- \rightarrow W^+W^-b\bar{b}$.

$$\sigma_{e^+e^- \to W^+W^-b\bar{b}} = \underbrace{\sigma_{e^+e^- \to [\bar{t}\bar{t}]_{\text{res}}}(\mu_w)}_{\text{pure (PNR)QCD}} + \sigma_{e^+e^- \to W^+W^-b\bar{b}_{\text{nonres}}}(\mu_w)$$

Non-resonant starts at NLO (overall linear divergence) [MB, Jantzen, Ruiz-Femenia, 2010; Penin, Piclum, 2011]. Finite-width scale dep must cancel. Need consistent dim reg calculation.

From QCD $t\bar{t}$ threshold to a realistic prediction

At 3rd order QCD effects under control. Then

 σ<sub>e+e−→W+W−bb
nonres</sub> (μ_w) to (at least) NNLO in consistent scheme. Mostly inclusive.

NLO [MB, Jantzen, Ruiz-Femenia, 2010; Penin, Piclum, 2011] Partial results only at NNLO [Hoang, Reisser, Ruiz-Femenia, 2010; Jantzen, Ruiz-Femenia, 2013]

- QED effects (from NLO) [Pineda, Signer, 2006; MB, Jantzen, Ruiz-Femenia, 2010]
- Electroweak matching coefficients absorptive parts [Hoang, Reisser, 2004]
- Initial state radiation (formalism in MB, Falgari, Schwinn, Signer, Zanderighi, 2007).

Finally, need four-loop (formally even five-loop) relation between the pole and the MSbar mass to convert the precise value for the PS mass to an equally precise value of the MSbar mass.

[MB, Jantzen, Ruiz-Femenia, 2010]



QED and non-resonant corrections relative to the $t\bar{t}$ LO cross section in percent: $\sigma_{\text{QED}}^{(1)}/\sigma_{t\bar{t}}^{(0)}$ (upper solid blue), $\sigma_{\text{non-res}}^{(0)}/\sigma_{t\bar{t}}^{(0)}$ for the total cross section (lower solid red) and $\Delta M_t = 15$ GeV (lower dashed red). The relative size of the sum of the QED and non-resonant corrections is represented by the middle (black) lines, for $\Delta M_{t,\text{max}}$ (solid) and $\Delta M_t = 15$ GeV (dashed). $m_{t,\text{pole}} = 172$ GeV.

Large correction below threshold. Much larger than QCD scale-dependence at 3rd order ($\pm(2-3)\%$)

$e^+e^- \rightarrow W^+W^-b\bar{b}$ near $s = 4m_t^2$

NLO + NNLO singular terms [Jantzen, Ruiz-Femenia, 2013; see also Hoang, Reisser, Ruiz-Femenia, 2010] (Singular refers to expansion in Λ/m_t where Λ is an invariant mass cut such that $m_t\Gamma_t \ll \Lambda^2 \ll m_t^2$.)



NNLO non-resonant still -2% at threshold and larger below. Accurate description of region below peak is required for precise determination of m_t .

- I $e^+e^- \rightarrow t\bar{t}X$ cross section near threshold now computed at NNNLO in (PNR)QCD
 - Sizeable 3rd order corrections and reduction of theoretical uncertainty.
 - Parameter dependences $(m_t, \Gamma_t, y_t, \alpha_s)$ can be studied.
- II Realistic predictions for $e^+e^- \rightarrow W^+W^-b\bar{b}$ near top-pair threshold
 - NLO available, including cuts invariant mass cuts. Should (and can) be done to NNLO. Residual uncertainty should then be small.
 - ISR should be done by theorists. Known in principle, in practice accuracy may not yet be sufficient.
- III What is limiting the mass measurement: statistics? experimental systematics? theory? Still open.