

FCC-ee/TLEP Physics Workshop (TLEP7)  
CERN, 19-21 June 2014

# ***Rare top decays at TLEP***

*( $\sim 10^6$  ttbar at threshold +  
 $\sim 10$  ab<sup>-1</sup> at 240 GeV)*

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with contributions by  
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19 June 2014

## Outline :

- **“inclusive” approach to rare top (exotic) decays:**
  - a) study of top recoil system in top pairs
  - b) excess in top total width
- **“exclusive” approach to rare top decays:**
  - a) “measurable” SM rare top decays
  - b) rare top decays measurable only in BSM
- **$e^+e^- \rightarrow t q$  (FCNC) at TLEP 240**

hard to conceive at hadronic colliders !

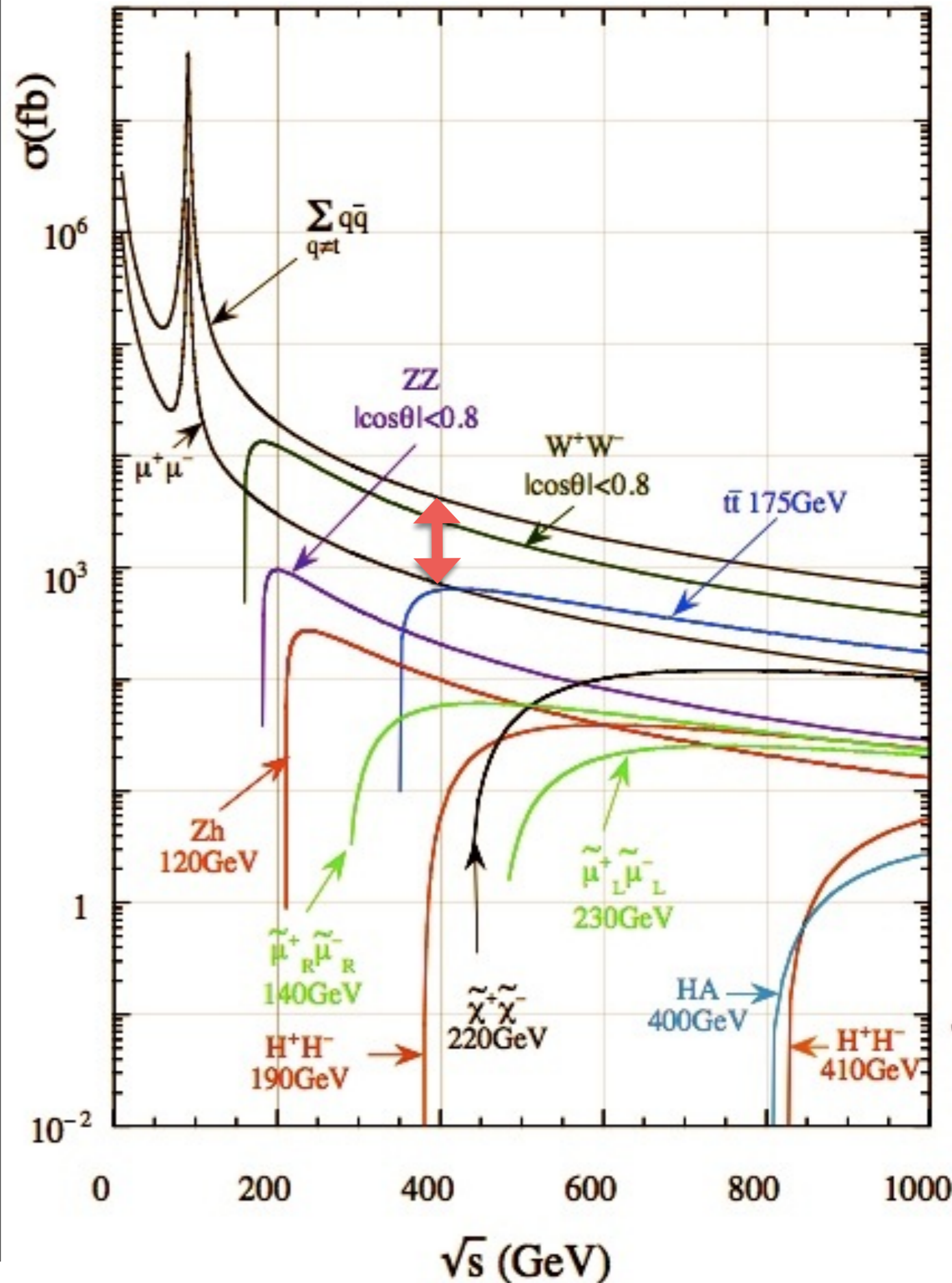
mainly a few thoughts  
not an exhaustive review !

## Main References

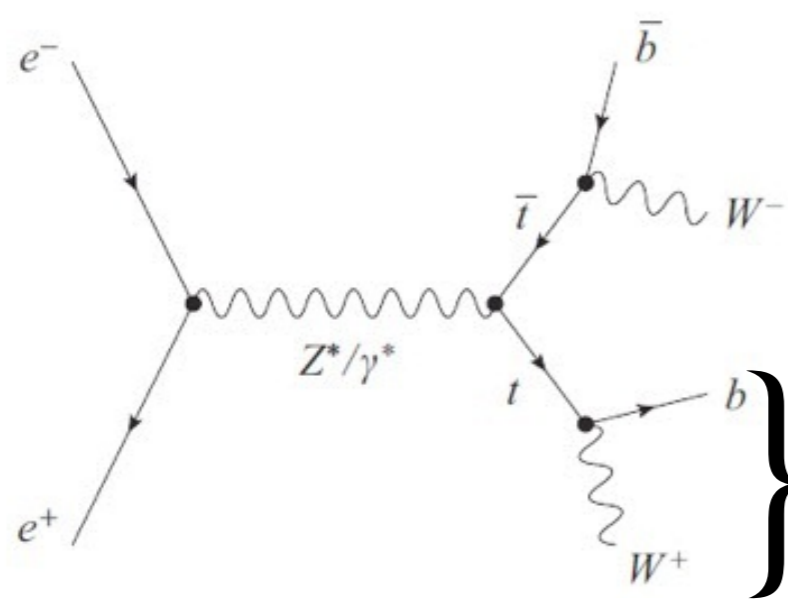
- TLEP Physics Case arxiv:1308.6176; TLEP<sup>n</sup><sup>th</sup> Workshop material, ILC TDR '13
- Snowmass Studies 2013 (arxiv:1307.8265,1311.2028)
- AWLC14, Fermilab, 12-16 May 2014 <http://www.linearcollider.org/awlc14/>

what makes unique  
 $e^+e^-$  environment wrt  
 (larger  $N_{ev}$ ) had collisions :

- democracy in  $\sigma$ 's  
 (all EW  $\sigma$ 's !)
- accurate TH predictions
- clean EXP environment
- untriggered operation
- can detect and  
 reconstruct "any"  
 hadronic final state
- can detect what is  
 invisible at LHC just  
 because we do not know  
 what to trigger on ...

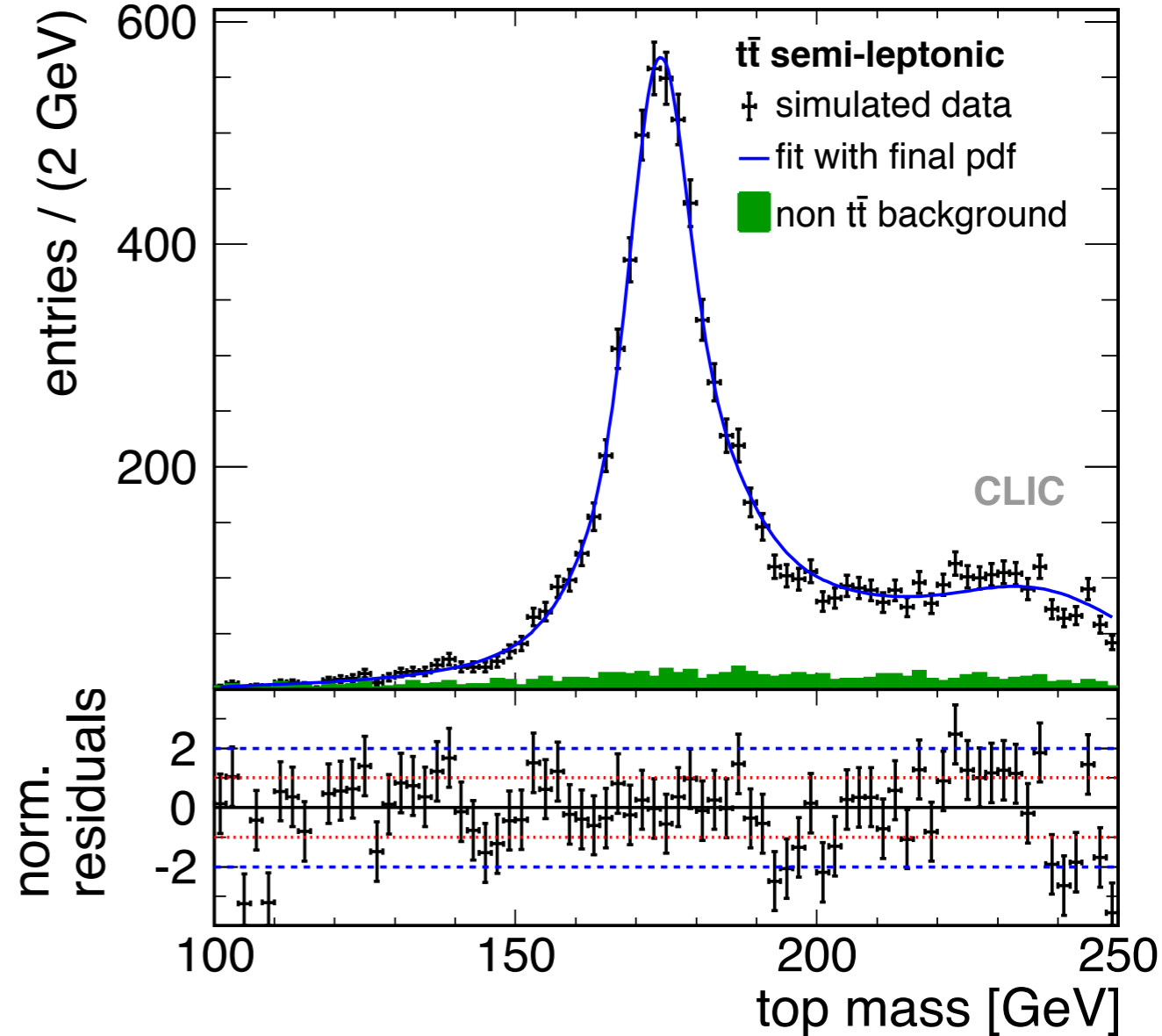
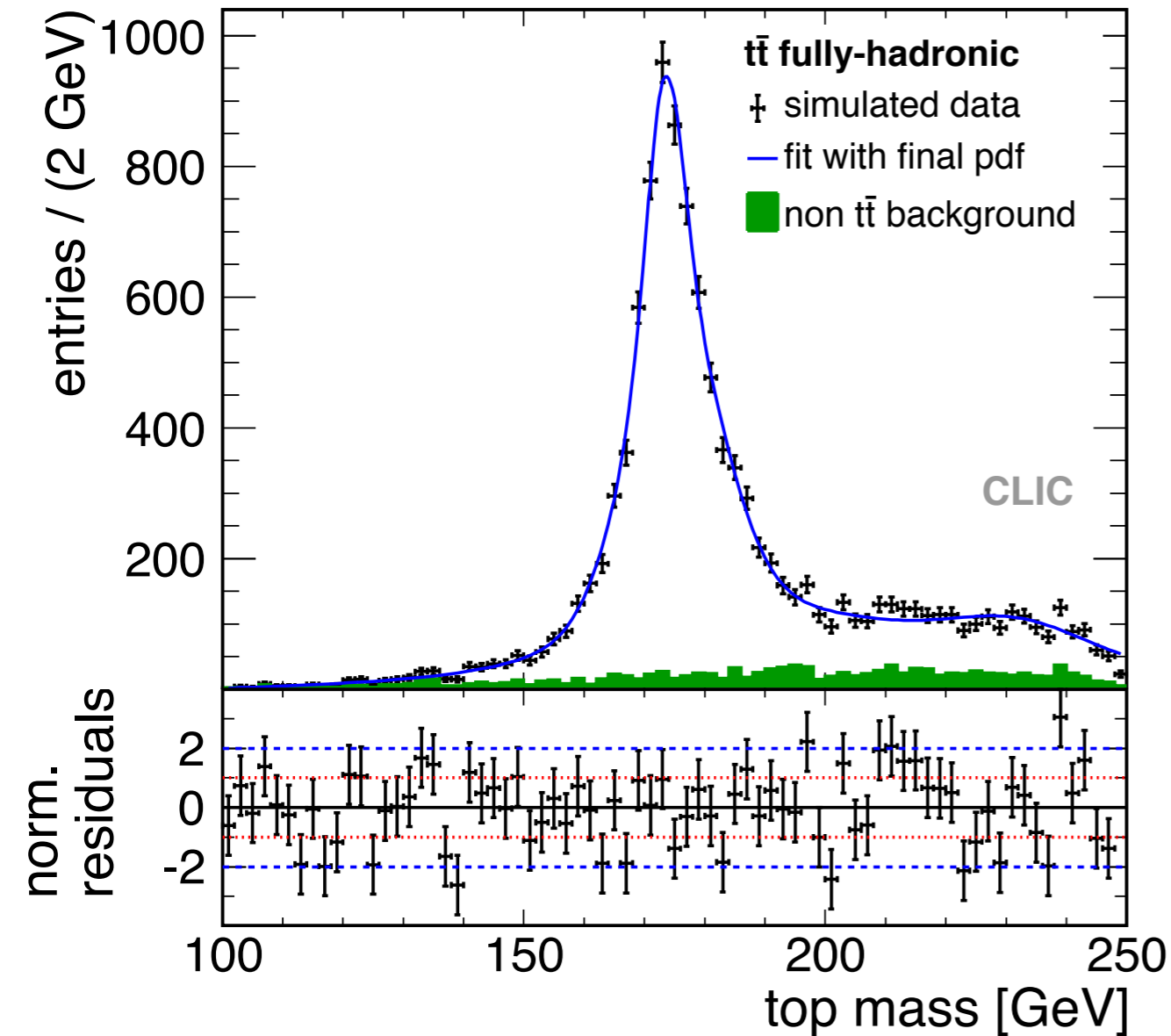


ttbar physics cleanness in  $e^+e^-$  collisions  
 well represented by plots below  
**(green is background !)**



**fully hadronic and semi-leptonic  
 top mass reconstruction**

Seidel et al. 1303.3758

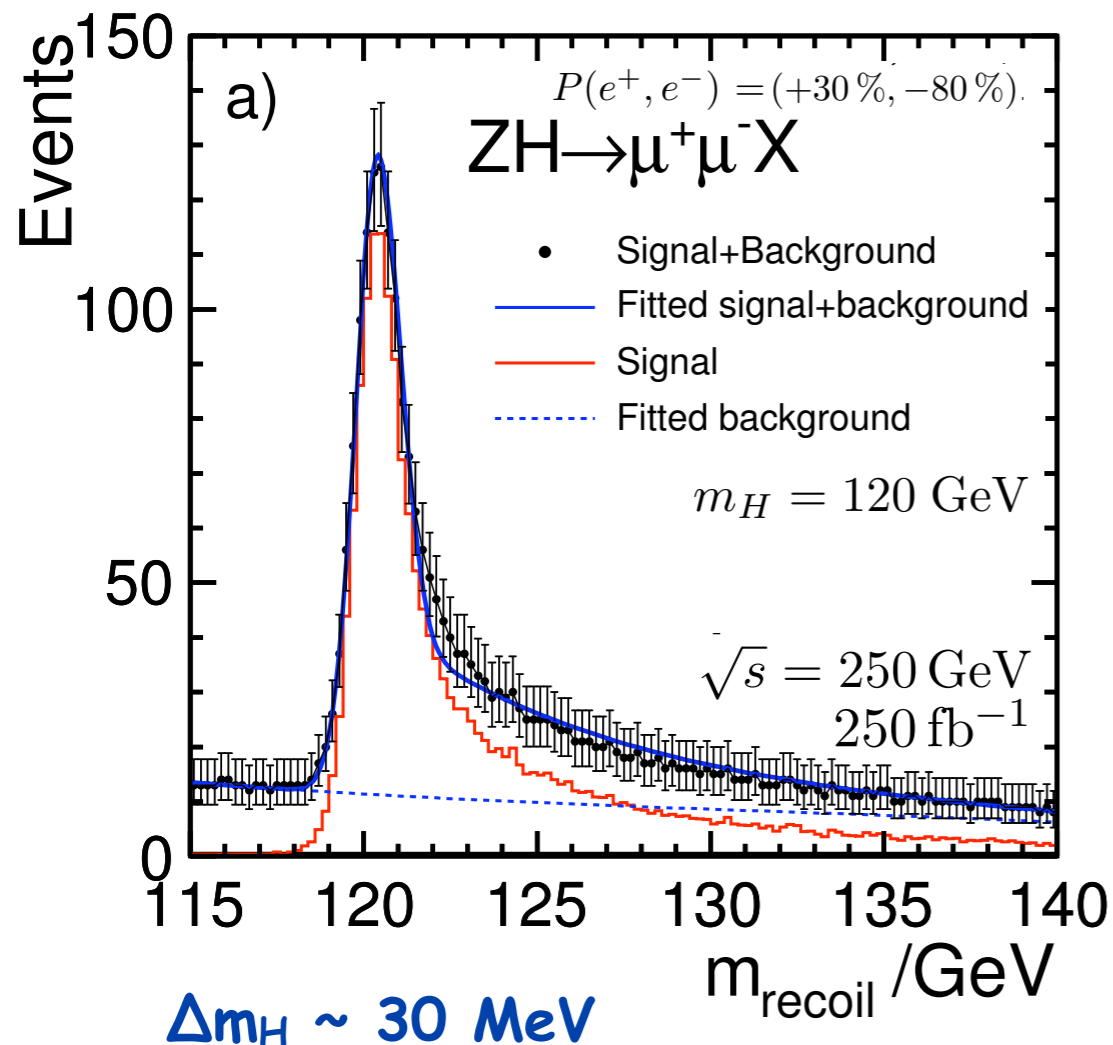
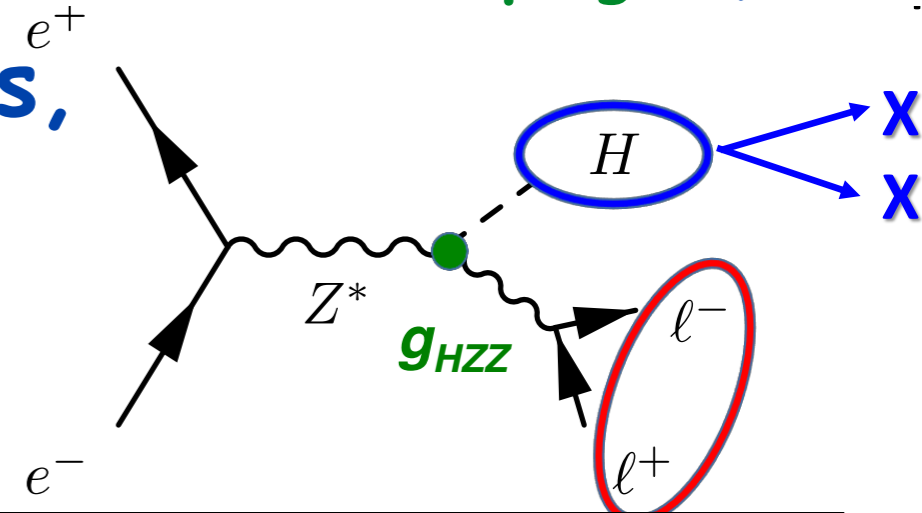


# inclusive Higgs studies through Z recoil system

**HZ** selected by just identifying Z decay products

(→ absolute  $\sigma_{\text{tot}} (\sim g_{\text{HZZ}}^2)$  measurement → model indep.  $g_{\text{HZZ}}$ )

→ direct access to **invisible H decays**,  
and **invisible-at-LHC decays**  
( $H \rightarrow cc$ ,  $H \rightarrow gg$ )



could we extend this technique to top pairs in  $e^+e^- \rightarrow t\bar{t}$  and make inclusive searches for exotic top final states by looking at recoil top system ?

# heavy top mass allows decays into new BSM states

examples :

$$t \rightarrow H^+ b \rightarrow \tau \nu b$$

$$t \rightarrow H^+ s \rightarrow c \bar{s} s$$

still allowed beyond 2HDM type II

$$t \rightarrow Z' c, Z' u \quad (\text{light neutral gauge bosons})$$

$$t \rightarrow \chi\chi c, \chi\chi u \quad (\text{dark matter})$$

$$t \rightarrow n \text{ jets} \neq bW \rightarrow bj\bar{j} \quad (???)$$

in principle one can have many different final states  
with unexpected kinematical features ...

can't find them at LHC unless you make assumptions on  
what you are looking for !

# inclusive searches for exotic $t$ decays through recoil system ( $e^+e^-$ )

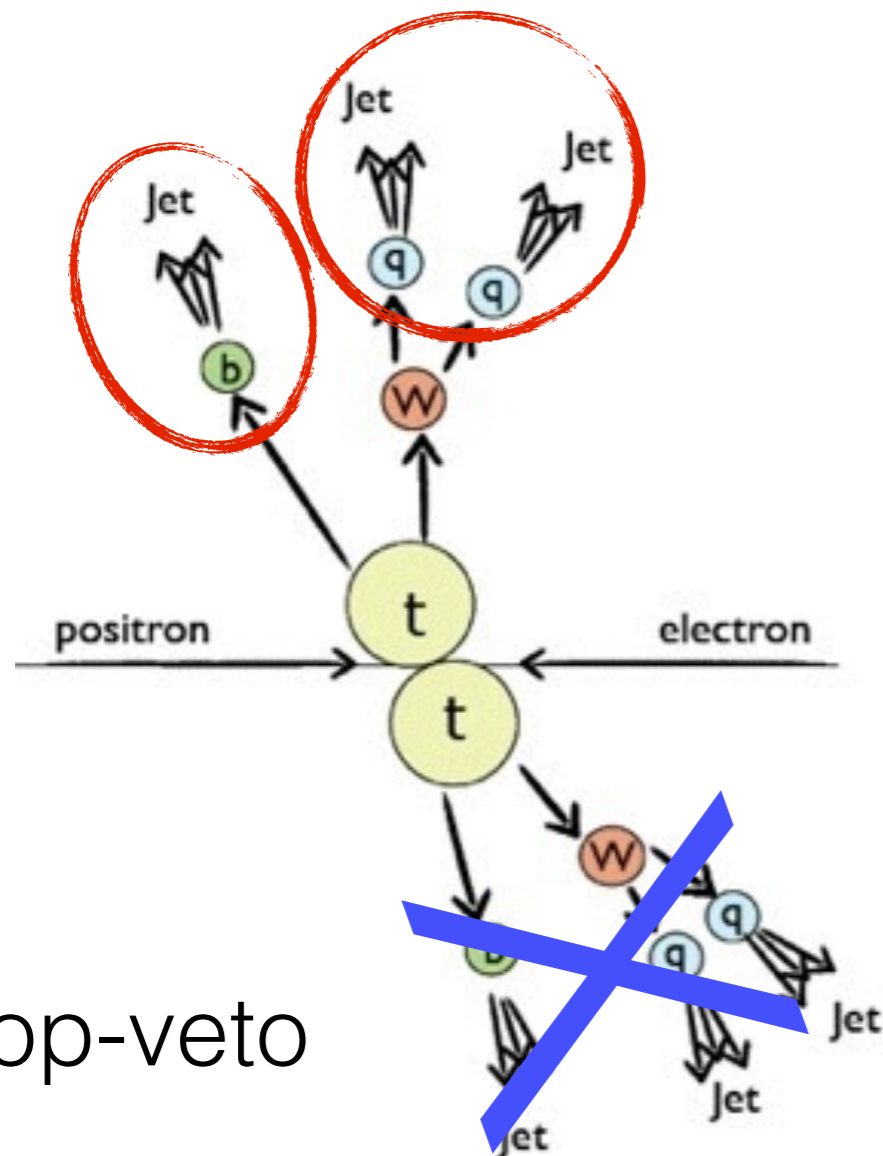
large variety  
of possible final states  
→ global analysis of the  
recoil system with a top-veto

a) define criteria to tag  
a  $Wb/Wj$  system  
as a (SM) top quark

b) look for events containing  
a top-system with  
a veto on a 2<sup>nd</sup> tag  
(i.e. recoil system does not  
pass the SM top-system  
criteria)

c) full simulation needed to  
assess sensitivity (  $< \% \sigma$  ? )

d) get model independent  
bounds on  $BR(\text{top})_{\text{exotica}}$  !



$E_{\text{cm}} > 350 \text{ GeV}$

# inclusive approach to exotic top decays (b)

bounds on  $\delta \Gamma_{top}$   (excess in top total width) can bound exotic decay widths

**SM:**

$$\Gamma_t = \frac{G_F m_t^3}{8\pi\sqrt{2}} |V_{tb}|^2 \left(1 - \frac{m_W^2}{m_t^2}\right)^2 \left(1 + 2\frac{m_W^2}{m_t^2}\right) \left[1 - \frac{2\alpha_s}{3\pi} \left(\frac{2\pi^2}{3} - \frac{5}{2}\right)\right]$$

**+ (b → s, d)**

$$\Gamma_t = 1.33 \text{ GeV} \quad (m_t = 172.5 \text{ GeV})$$

## top width measurement at hadron colliders :

combine  $\mathcal{R} = \mathcal{B}(t \rightarrow Wb) / \mathcal{B}(t \rightarrow Wq)$  with t-channel single top  $\sigma$  :

$$\Gamma_t = \frac{\sigma_{t\text{-ch.}}}{\mathcal{B}(t \rightarrow Wb)} \cdot \frac{\Gamma(t \rightarrow Wb)}{\sigma_{t\text{-ch.}}^{\text{theor.}}}$$

**assuming SM** (  $\sum_q \mathcal{B}(t \rightarrow Wq) = 1$  ) :

CMS 1404.2292

→  $\Gamma_t = 1.36 \pm 0.02 \text{ (stat.)}_{-0.11}^{+0.14} \text{ (syst.) GeV}$

$$\mathcal{R} = |V_{tb}|^2$$

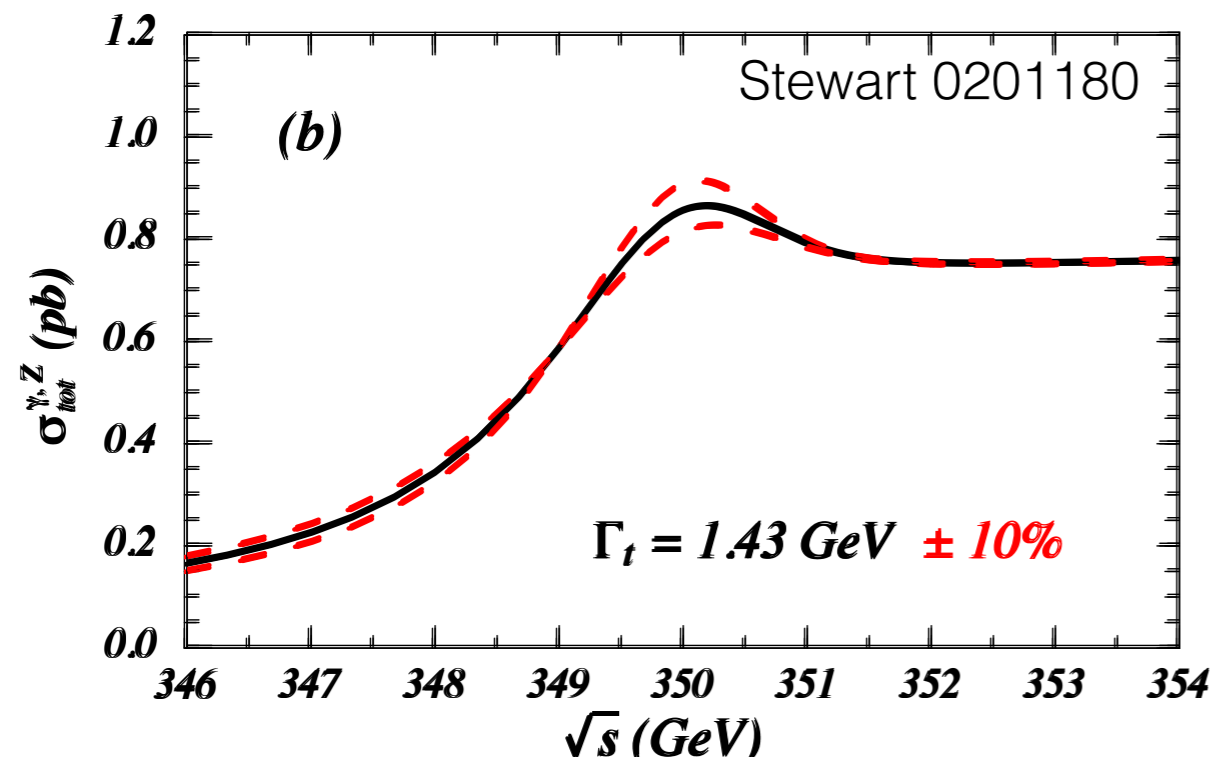
$$|V_{tb}| = 1.007 \pm 0.016 \text{ (stat.+syst.)}$$



# inclusive approach to exotic top decays (b)

(model independent)  $\Gamma_{\text{top}}$  measurement at 1% at TLEP

resonance cross section at threshold very sensitive to  $\alpha_s, m_t, \Gamma_t$  ;  
 peak at  $\sigma_{\text{res}} \sim \alpha_s^3 / (m_t \Gamma_t)$



	$m_{\text{top}}$	$\Gamma_{\text{top}}$	$\lambda_{\text{top}}$
TLEP	10 MeV	11 MeV	13%
ILC	31 MeV	34 MeV	40%

Martinez, Miguel 0207315

- TLEP Physics Case arxiv:1308.6176

bounds on  $\delta\Gamma_{\text{top}}$  can probe inclusively rare decays  
 for  $\text{BR}_{\text{exotic}} \geq \text{few } \%$  at TLEP

# exclusive approach :

a) measurable SM rare top decays

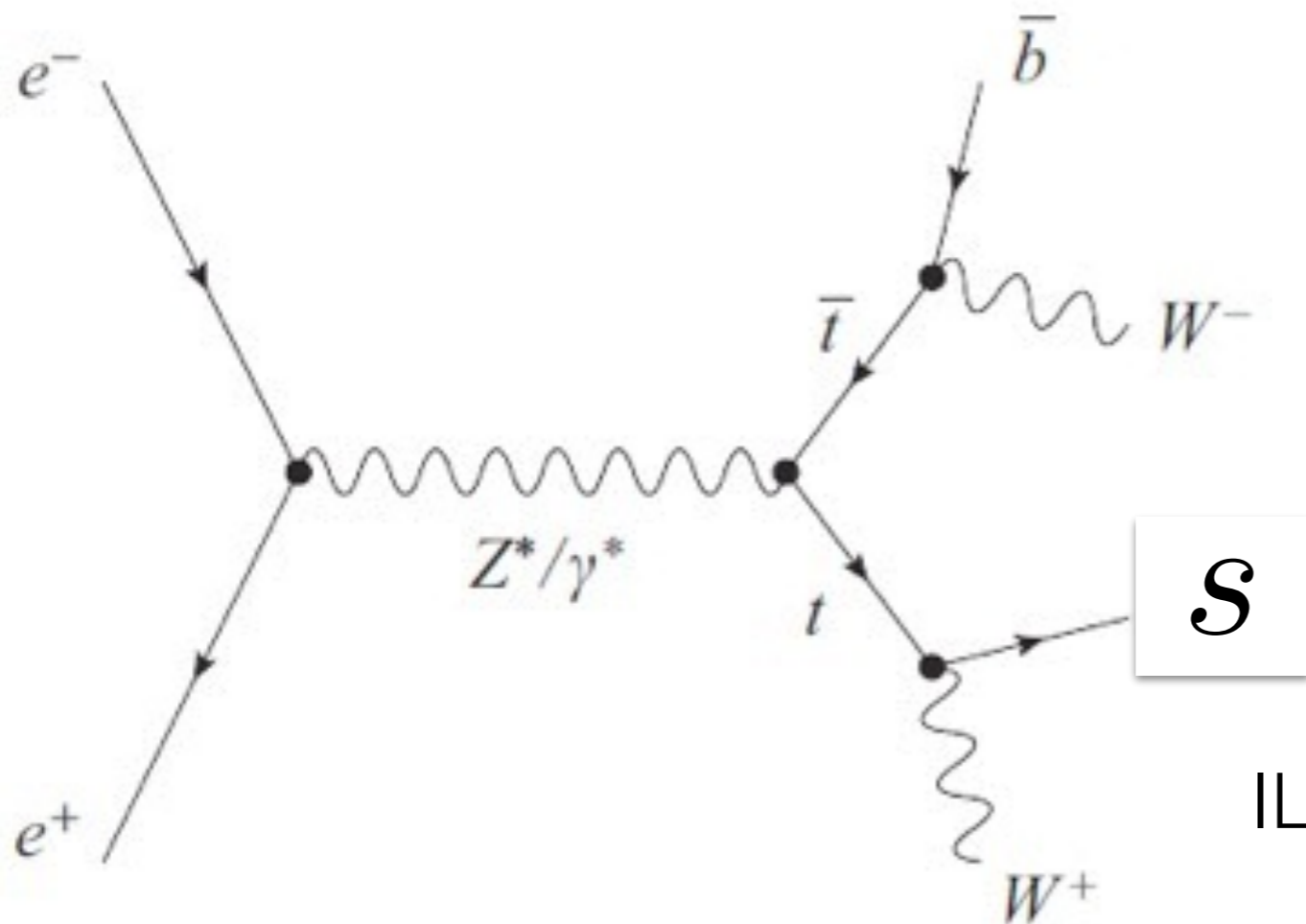
$$t \rightarrow s W$$

$BR(t \rightarrow s W) \sim 2 \times 10^{-3}$  direct observation of  $tsW$  coupl.

can also probe anomalous interac.s of type (b->s):

$$\mathcal{L} = -\frac{g}{\sqrt{2}} \bar{b} \gamma^\mu (V_L P_L + V_R P_R) t W_\mu^- - \frac{g}{\sqrt{2}} \bar{b} \frac{i\sigma^{\mu\nu} q_\nu}{M_W} (g_L P_L + g_R P_R) t W_\mu^- + h.c.$$

needs discriminants to suppress dominant  $t \rightarrow bW$



with  $10^6 \bar{t}t$  pairs

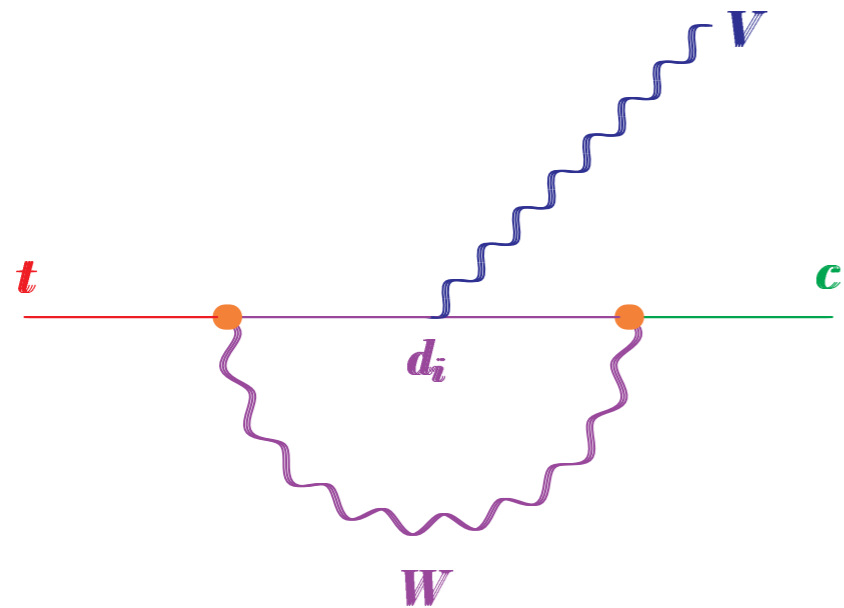
$$\Rightarrow 4 \cdot 10^3 \quad t \rightarrow sW$$

ILC sensitive to  $BR(t \rightarrow s W) \sim 10^{-3}$

(Diaz-Cruz et al 0712.3782)

# exclusive approach :

b) rare top decays measurable only in BSM frameworks



## FCNC top decays in the SM

GIM-suppressed by  $\left(\frac{m_b}{M_W}\right)^4$

$$\text{BR}(t \rightarrow c\gamma) \simeq 5 \times 10^{-14}$$

$$\text{BR}(t \rightarrow cg) \simeq 5 \times 10^{-12}$$

$$\text{BR}(t \rightarrow cZ) \simeq 1 \times 10^{-14}$$

$$\text{BR}(t \rightarrow ch) \simeq 3 \times 10^{-15}$$

$$\text{BR}(t \rightarrow u\gamma) \simeq 4 \times 10^{-16}$$

$$\text{BR}(t \rightarrow ug) \simeq 4 \times 10^{-14}$$

$$\text{BR}(t \rightarrow uZ) \simeq 8 \times 10^{-17}$$

$$\text{BR}(t \rightarrow uh) \simeq 2 \times 10^{-17}$$

$$(t \rightarrow ux)/(t \rightarrow cx) \simeq |V_{ub}/V_{cb}|^2 \simeq 0.008$$

(Aguilar-Saavedra hep-ph/0409342)

# New Physics models can hugely enhance predictions ! Different models predict different pattern of enhancements

	2HDM	MSSM	RS
$t \rightarrow cZ$	$\lesssim 10^{-6}$	$\lesssim 10^{-7}$	$\lesssim 10^{-5}$
$t \rightarrow c\gamma$	$\lesssim 10^{-7}$	$\lesssim 10^{-8}$	$\lesssim 10^{-9}$
$t \rightarrow cg$	$\lesssim 10^{-5}$	$\lesssim 10^{-7}$	$\lesssim 10^{-10}$
$t \rightarrow ch$	$\lesssim 10^{-2}$	$\lesssim 10^{-5}$	$\lesssim 10^{-4}$

Present LHC bounds from  $t\bar{t} \xrightarrow{\text{FCNC}} (Z/\gamma/g/h + j)(Wb)$   
 plus single-top anomalous production  $qg \rightarrow t(Z/\gamma/H)$

Process	Br Limit	Search	Dataset
$t \rightarrow Zq$	$5 \times 10^{-4}$	CMS $t\bar{t} \rightarrow Wb + Zq \rightarrow \ell\nu b + \ell\ell q$	19.7 fb <sup>-1</sup> , 8 TeV
$t \rightarrow Zq$	$7.3 \times 10^{-3}$	ATLAS $t\bar{t} \rightarrow Wb + Zq \rightarrow \ell\nu b + \ell\ell q$	2.1 fb <sup>-1</sup> , 7 TeV
$t \rightarrow gu$	$3.1 \times 10^{-5}$	ATLAS $qg \rightarrow t \rightarrow Wb$	14.2 fb <sup>-1</sup> , 8 TeV
$t \rightarrow gc$	$1.6 \times 10^{-4}$	ATLAS $qg \rightarrow t \rightarrow Wb$	14.2 fb <sup>-1</sup> , 8 TeV
$t \rightarrow \gamma u$	$1.6 \times 10^{-4}$	CMS $qg \rightarrow t\gamma \rightarrow Wb\gamma$	19.1 fb <sup>-1</sup> , 8 TeV
$t \rightarrow \gamma c$	$1.8 \times 10^{-3}$	CMS $qg \rightarrow t\gamma \rightarrow Wb\gamma$	19.1 fb <sup>-1</sup> , 8 TeV
$t \rightarrow hq$	$7.9 \times 10^{-3}$	ATLAS $t\bar{t} \rightarrow Wb + hq \rightarrow \ell\nu b + \gamma\gamma q$	20 fb <sup>-1</sup> , 8 TeV
$t \rightarrow hq$	$5.6 \times 10^{-3}$	CMS $t\bar{t} \rightarrow Wb + hq \rightarrow \ell\nu b + \ell\ell qX$	19.5 fb <sup>-1</sup> , 8 TeV

*single-top production more sensitive to u-type vertex*

# most general effective Lagrangian for FC tqV(H) interactions with terms up to dim 5

$$\begin{aligned}
 -\mathcal{L}^{\text{eff}} = & \frac{g}{2c_W} X_{qt} \bar{q} \gamma_\mu (x_{qt}^L P_L + x_{qt}^R P_R) t Z^\mu + \frac{g}{2c_W} \kappa_{qt} \bar{q} (\kappa_{qt}^v + \kappa_{qt}^a \gamma_5) \frac{i\sigma_{\mu\nu} q^\nu}{m_t} t Z^\mu \\
 & + e \lambda_{qt} \bar{q} (\lambda_{qt}^v + \lambda_{qt}^a \gamma_5) \frac{i\sigma_{\mu\nu} q^\nu}{m_t} t A^\mu + g_s \zeta_{qt} \bar{q} (\zeta_{qt}^v + \zeta_{qt}^a \gamma_5) \frac{i\sigma_{\mu\nu} q^\nu}{m_t} T^a q G^{a\mu} \\
 & + \frac{g}{2\sqrt{2}} g_{qt} \bar{q} (g_{qt}^v + g_{qt}^a \gamma_5) t H + \text{H.c.},
 \end{aligned}$$

$\sigma_{\mu\nu}$  terms grow with  $V^\mu$  momentum  $q^{\nu\mu}$

$$\text{Br}(t \rightarrow qZ)_\gamma = 0.472 X_{qt}^2, \quad \left. \begin{array}{l} > 5 \times 10^{-4} \\ > 2 \times 10^{-3} \end{array} \right\} \Rightarrow X_{qt} < 0.033$$

$$\text{Br}(t \rightarrow qZ)_\sigma = 0.367 \kappa_{qt}^2, \quad \Rightarrow \kappa_{qt} < 0.037$$

$$\text{Br}(t \rightarrow q\gamma) = 0.428 \lambda_{qt}^2, \quad > 2 \times 10^{-3} \Rightarrow \lambda_{qt} < 0.065$$

$$\text{Br}(t \rightarrow qq) = 7.93 \zeta_{qt}^2, \quad > 2 \times 10^{-4}$$

$$\text{Br}(t \rightarrow qH) = 3.88 \times 10^{-2} g_{qt}^2 > 6 \times 10^{-3}$$

LHC

# bounds on $tqZ$ and $tq\gamma$

## ILC versus full LHC

Process	Br Limit	Search	Dataset
$t \rightarrow Zq$	$2.2 \times 10^{-4}$	ATLAS $t\bar{t} \rightarrow Wb + Zq \rightarrow \ell\nu b + \ell\ell q$	300 fb <sup>-1</sup> , 14 TeV
$t \rightarrow Zq$	$7 \times 10^{-5}$	ATLAS $t\bar{t} \rightarrow Wb + Zq \rightarrow \ell\nu b + \ell\ell q$	3000 fb <sup>-1</sup> , 14 TeV
$t \rightarrow Zq$	$5 (2) \times 10^{-4} *$	ILC single top, $\gamma_\mu (\sigma_{\mu\nu})$	500 fb <sup>-1</sup> , 250 GeV
$t \rightarrow Zq$	$1.5 (1.1) \times 10^{-4} (-5)$	ILC single top, $\gamma_\mu (\sigma_{\mu\nu})$	500 fb <sup>-1</sup> , 500 GeV
$t \rightarrow Zq$	$1.6 (1.7) \times 10^{-3}$	ILC $t\bar{t}$ , $\gamma_\mu (\sigma_{\mu\nu})$	500 fb <sup>-1</sup> , 500 GeV
$t \rightarrow \gamma q$	$8 \times 10^{-5}$	ATLAS $t\bar{t} \rightarrow Wb + \gamma q$	300 fb <sup>-1</sup> , 14 TeV
$t \rightarrow \gamma q$	$2.5 \times 10^{-5}$	ATLAS $t\bar{t} \rightarrow Wb + \gamma q$	3000 fb <sup>-1</sup> , 14 TeV
$t \rightarrow \gamma q$	$6 \times 10^{-5} *$	ILC single top	500 fb <sup>-1</sup> , 250 GeV
$t \rightarrow \gamma q$	$6.4 \times 10^{-6}$	ILC single top	500 fb <sup>-1</sup> , 500 GeV
$t \rightarrow \gamma q$	$1.0 \times 10^{-4}$	ILC $t\bar{t}$	500 fb <sup>-1</sup> , 500 GeV

$\sigma_{\mu\nu}$  terms grow with  $V^\mu$  momentum  $q^{\mu}$  \* extrapolated  
 $\Rightarrow e^+e^- \rightarrow \gamma, Z(q^\mu) \rightarrow tq$  at ILC, most sensitive one!

**TLEP studies for FCNC single top in progress** →

$$e^+ e^- \rightarrow \gamma, Z(q^\mu) \rightarrow tq$$

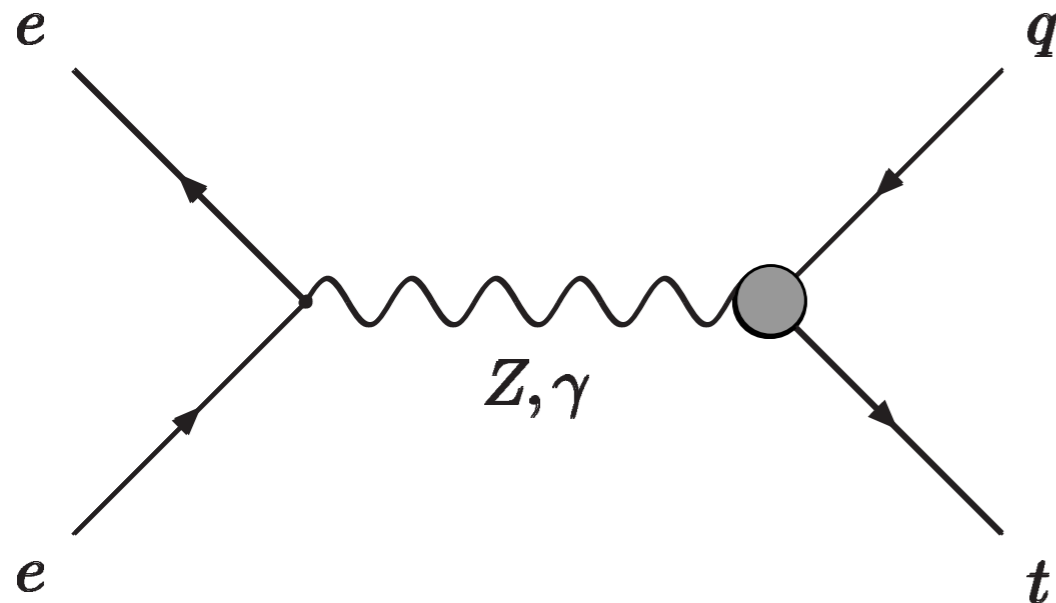
Han and Hewett 9811237

Bar-Shalom, Wudka 9905407

Aguilar-Saavedra, Riemann 0102197

(at LEP2 and ILC)

**main background from  $Wjj$**



**E= 240 GeV (max cross section and large lumi at TLEP)**

versus

**E= 500 GeV (lower bckgd and more sensitive to  $\sigma_{\mu\nu}$  terms)**

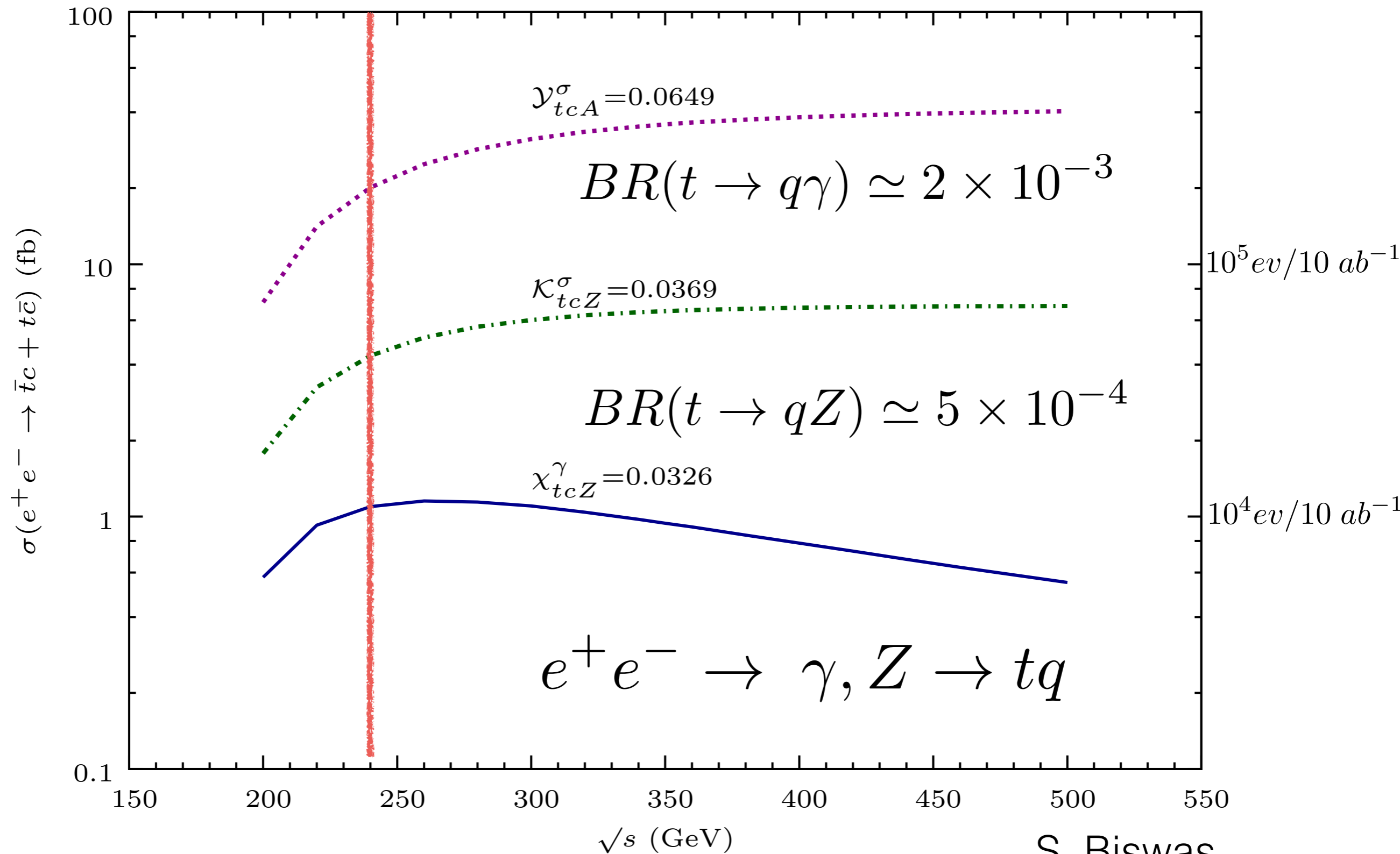
$$\begin{aligned}
 -\mathcal{L} = & \frac{g_W}{2c_W} X_{tq} \bar{t} \gamma_\mu (x_{tq}^L P_L + x_{tq}^R P_R) q Z^\mu + \frac{g_W}{2c_W} \kappa_{tq} \bar{t} (\kappa_{tq}^v - \kappa_{tq}^a \gamma_5) \frac{i\sigma_{\mu\nu} q^\nu}{m_t} q Z^\mu \\
 & + e \lambda_{tq} \bar{t} (\lambda_{tq}^v - \lambda_{tq}^a \gamma_5) \frac{i\sigma_{\mu\nu} q^\nu}{m_t} q A^\mu,
 \end{aligned}$$

**Updated analysis at TLEP 240 in progress**

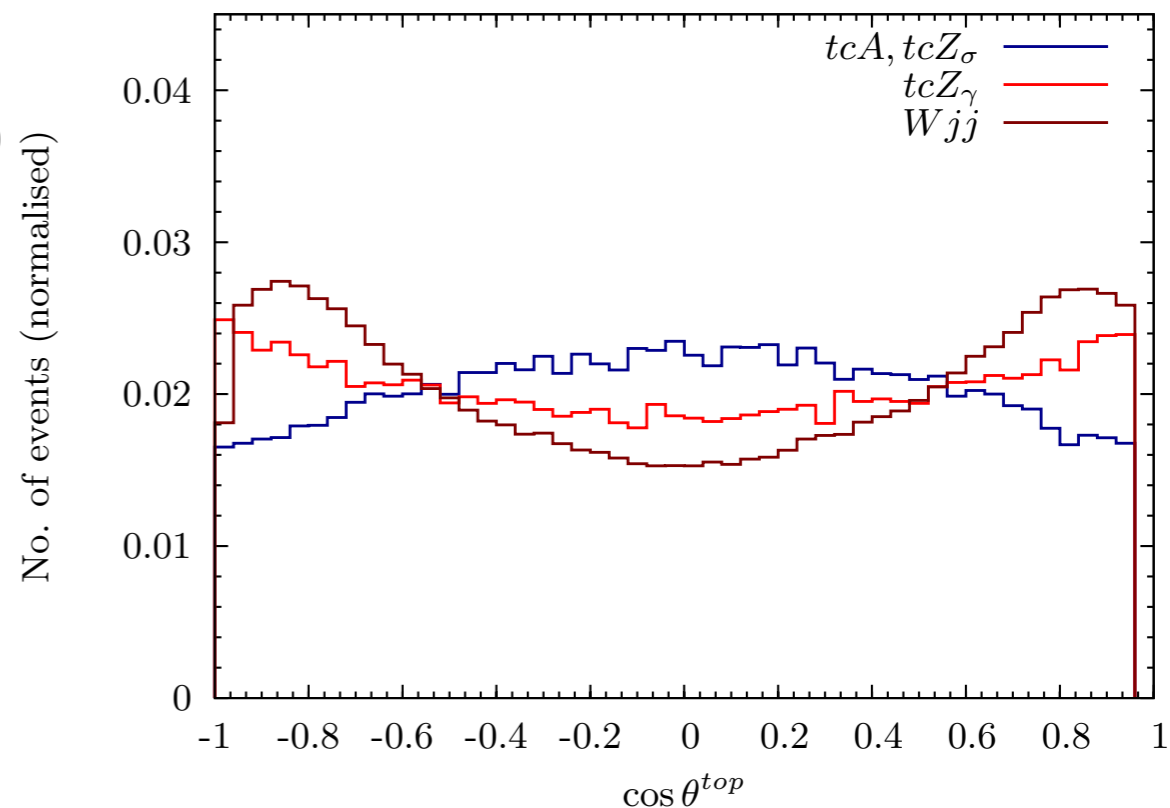
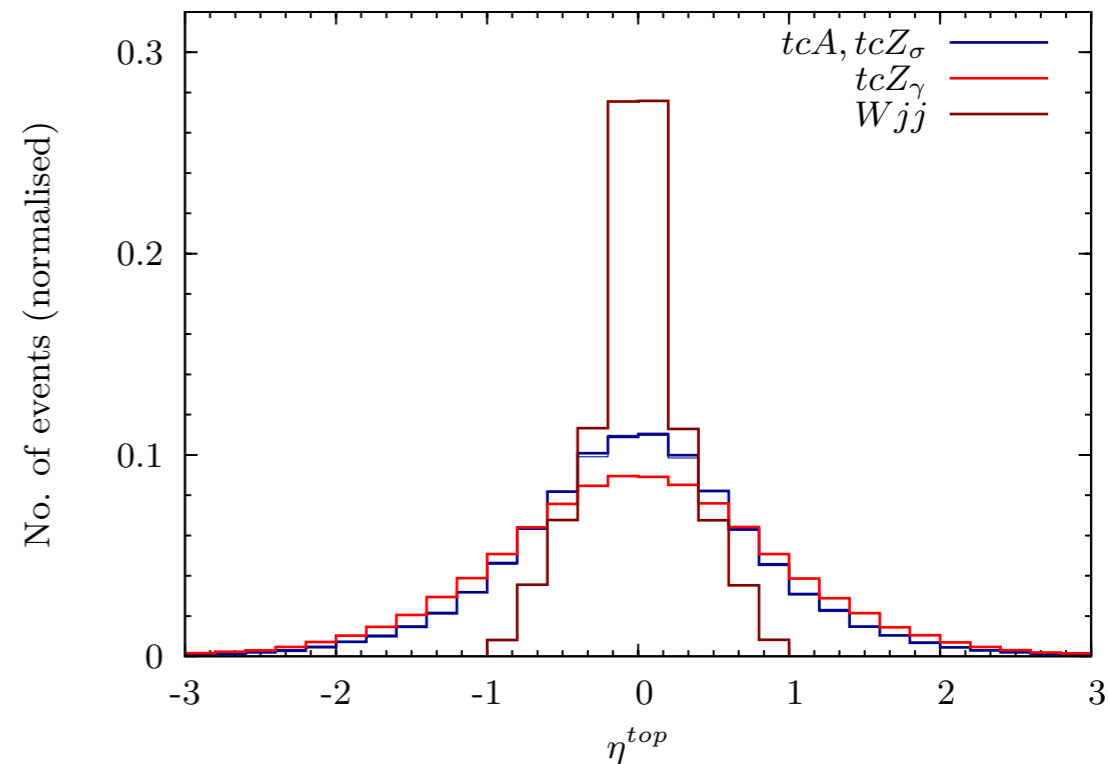
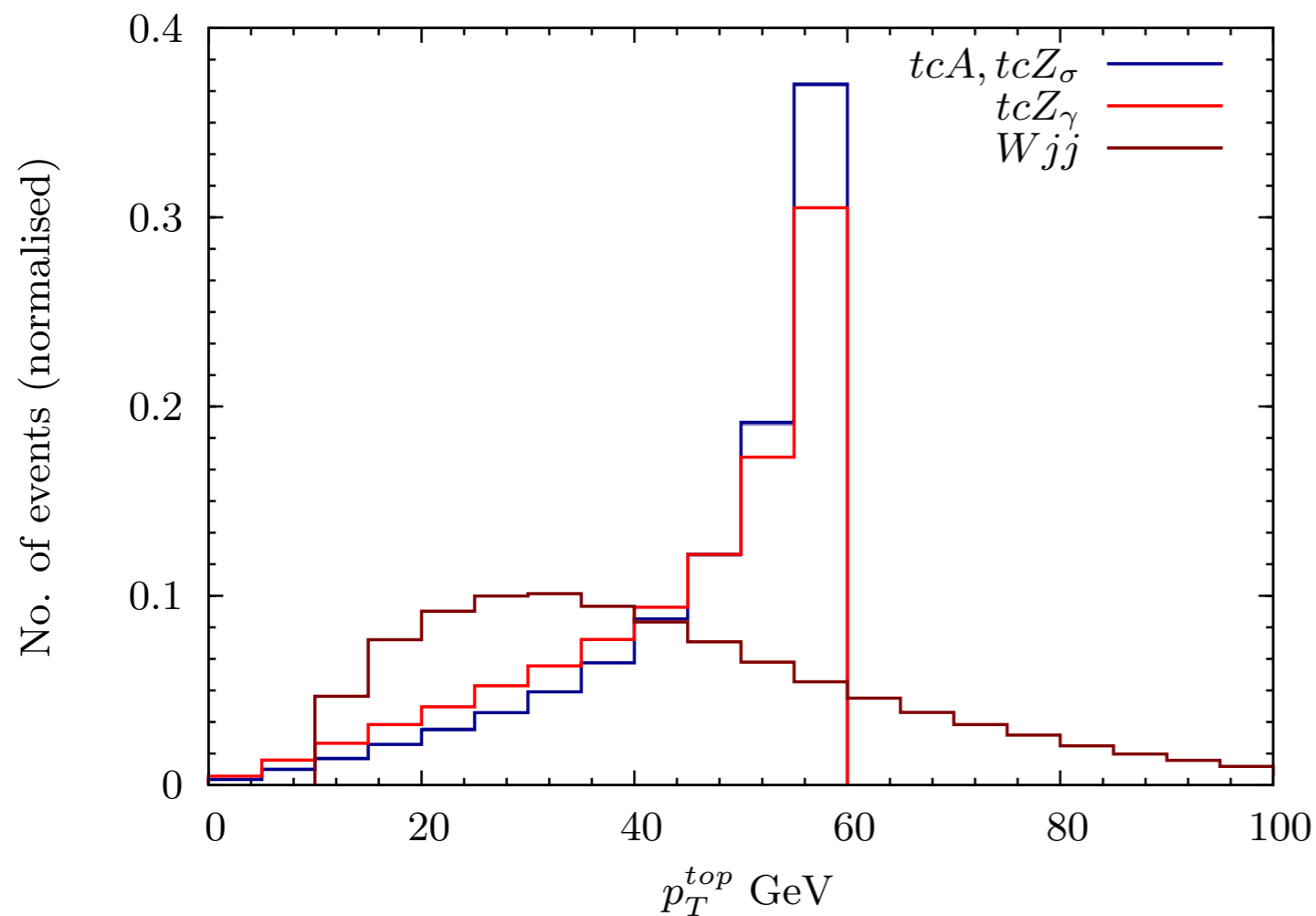
**(Azzi, Biswas, Margaroli, BM)**



# maximum allowed $\sigma$ 's by presents bounds on $BR(\text{top})^{\text{FCNC}}$



# signal vs Wjj bckgr : normalized distributions

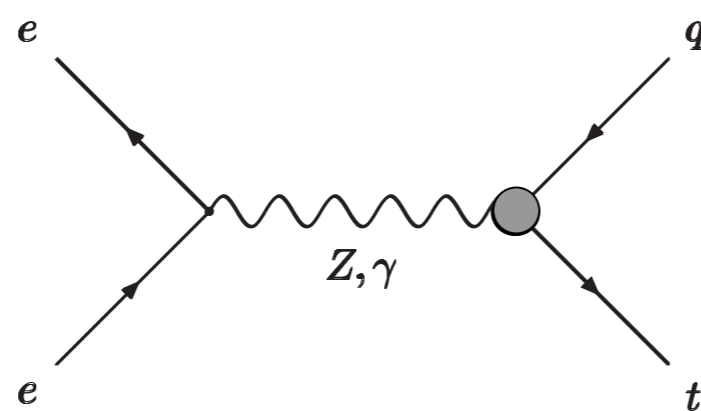


$$\sqrt{S} = 240 GeV$$

Madgraph5

$$p_T^j > 10 GeV$$

$$\Delta R(jj) > 0.4$$



disentangling photon from Z vertex can require looking at top FCNC decays at tt threshold

S. Biswas

# Outlook

- ever since its discovery, the top quark has never been produced and studied in such a clean environment as the one expected in  $e^+e^-$  collisions
- $e^+e^-$  collisions will almost allow to trace back top-quark final states on an event-by-event basis
- this will open the opportunity to look at details of top production and kinematics that is unthinkable in hadron collisions  
(relevant strategies mostly still to be developed ...)
- rare top decays is one of the (many) top physics chapters that would widely benefit from such spectacularly clean environment