$pp \rightarrow t\overline{t}\gamma\gamma$ 23 years after

Zoltán Trócsányi

University of Debrecen and MTA-DE Particle Physics Research Group in collaboration with

A. Kardos



Zoltan 70, Zürich May 23, 2014





Outline

- Motivation
- Method
- Predictions
- Conclusions

Motivation

- \circ Zoltan's excitement about a clear signal of observing a Higgs-boson of intermediate mass (70 < m_H c²/GeV < 140)
- irreducible background ttyy final-state was unknown

- \circ Zoltan's excitement about a clear signal of observing a Higgs-boson of intermediate mass (70 < m_H c²/GeV < 140)
- irreducible background t̄τγγ final-state was unknown
- my first phenomenology paper signal/background study of this channel
- message: works with excellent resolution of photons

- Zoltan's excitement about a clear signal of observing a Higgs-boson of intermediate mass $(70 < m_H c^2/GeV < 140)$
- irreducible background t̄τγγ final-state was unknown
- my first phenomenology paper signal/background study of this channel
- message: works with excellent resolution of photons
- ... actually achieved by ATLAS and CMS

Physics Letters B 271 (1991) 247-255 North-Holland

PHYSICS LETTERS B

Clear signal of intermediate mass Higgs boson production at LHC and SSC

Z. Kunszt, Z. Trócsányi 1

Theoretical Physics, ETH, CH-8093 Zurich, Switzerland

and

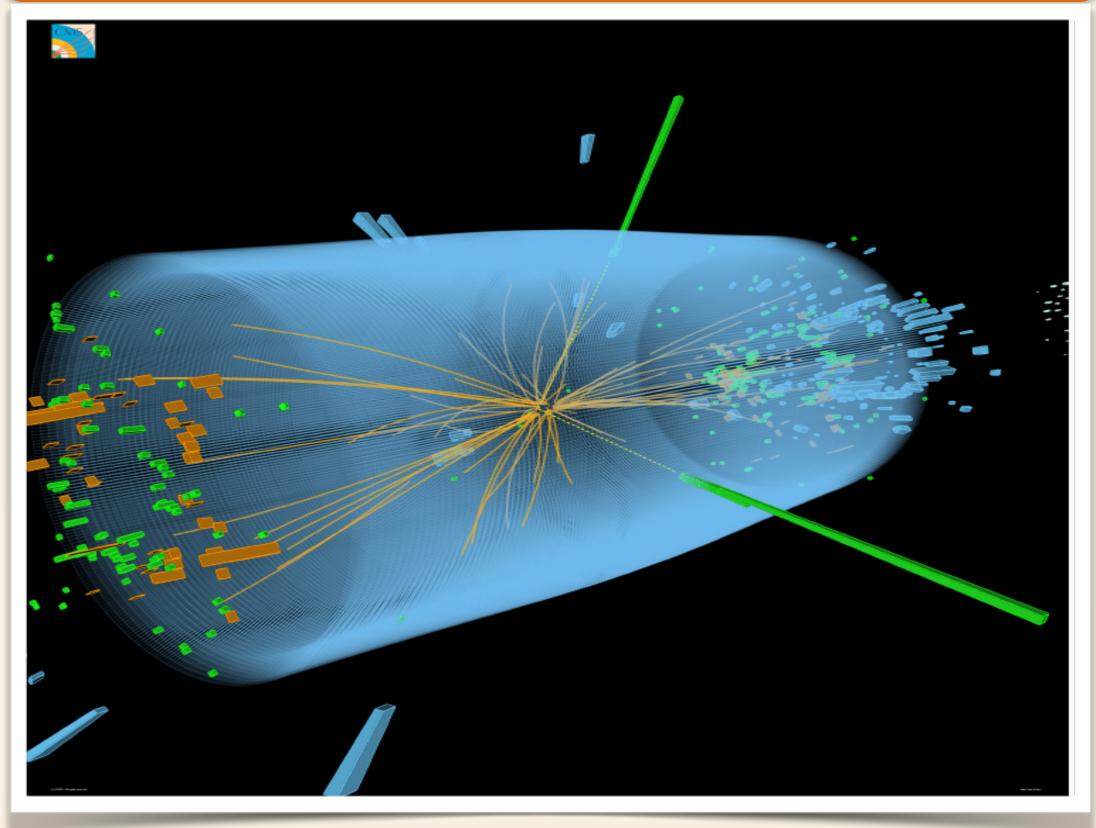
W.J. Stirling

Departments of Physics and Mathematical Sciences, University of Durham, Durham DH1 3LE, UK

Received 11 July 1991

We compute the event rates for two potentially important background processes $pp \rightarrow t\bar{t}\gamma\gamma X$ and $pp \rightarrow b\bar{b}\gamma\gamma X$ to the recently proposed signature for Higgs production of one isolated lepton and two isolated photons in the intermediate-mass range. We find that the background can be suppressed assuming good ($\approx (2-3)\%$) mass resolution in the invariant mass of the two photons, and assuming that the isolation criteria for the lepton and photons can be efficiently implemented experimentally. We reanalyse the signal to background ratios using realistic experimental cuts and find that by measuring the inclusive production of one isolated lepton and two isolated photons at the LHC or the SSC we can obtain a clear experimental signal for the production of the Higgs boson in the mass region $70 < M_H < 140$ GeV.

Higgs boson has been discovered



Higgs boson has been discovered

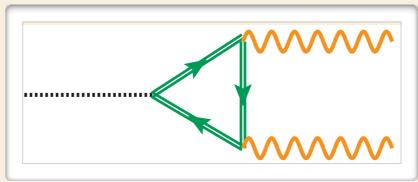
- $m_H [GeV]=125.5\pm0.2_{stat}\pm0.6_{syst} (ATLAS 2013)$ $125.7\pm0.3_{stat}\pm0.3_{syst} (CMS 2013)$
- All measured properties are consistent with SM expectations within experimental uncertainties
 - branching ratios as predicted
 - spin zero
 - parity +
 - \circ couples to masses of W and Z (with $c_v=1$ within experimental uncertainty)

t-quark: potential tool for discovery

- The t-quark is heavy, Yukawa coupling ~1 m_t [GeV]=173.34±0.64 (LHC+TeVatron, 2014) $(\Rightarrow y_t=0.997\pm0.003)$
 - \Rightarrow plays important role in Higgs physics (more tantalizing: $m_t m_Z = (125.7 \pm 0.3)^2 \text{ GeV}^2$)
- \circ y_t cannot be measured in H \rightarrow tT decay (m_t > m_H)

How to measure yt?

 \bullet H \to $\gamma\gamma$ is sensitive to y_t through t-quark loop, but rates are small and W loop also contributes



• $gg \rightarrow H$ is sensitive to y_t through t-quark loop if only SM model particles contribute (so far xsec is consistent with SM)

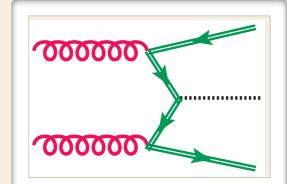
ullet gg o H is sensitive to BSM physics if y_t is measured separately

000000

tTH hadroproduction

 \circ y_t can be measured in pp \rightarrow tTH through many decay

channels (all very difficult):

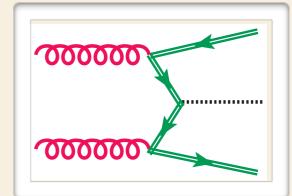


- ullet hadrons with single lepton: ${
 m t}
 ightarrow {
 m b} \ell
 u, {
 m ar t}
 ightarrow {
 m ar b} j j, {
 m H}
 ightarrow {
 m b} {
 m ar b}$
- ullet hadrons with dilepton: $t o b\ell
 u, \, ar{t} o ar{b}\ell
 u, \, H o bar{b}$
- ullet hadrons with hadronic tau: ${
 m t}
 ightarrow {
 m b}\ell
 u, {
 m ar t}
 ightarrow {
 m ar b}jj, {
 m H}
 ightarrow au_h^+ au_h^-$
- ullet diphoton with lepton: ${
 m t}
 ightarrow {
 m b}\ell
 u, \, {
 m ar t}
 ightarrow {
 m ar b} jj, \, {
 m H}
 ightarrow \gamma \gamma$
- ullet diphoton with hadrons: ${
 m t}
 ightarrow {
 m b} \, jj, \, {
 m ar t}
 ightarrow {
 m ar b} jj, \, {
 m H}
 ightarrow \gamma \gamma$
- ullet same sign dilepton: $\mathrm{t} o \mathrm{b}\, jj,\, \mathrm{ar{t}} o \mathrm{ar{b}} jj,\, \mathrm{H} o \ell
 u \ell[
 u]$
- ullet 3 leptons with di, trilepton: ${
 m t}
 ightarrow {
 m b}\ell
 u, {
 m ar t}
 ightarrow {
 m ar b}jj, {
 m H}
 ightarrow \ell[
 u]\ell[
 u]$
- ullet 4 lepton with di, trilepton: $t o b\ell
 u, ar t o ar b\ell[
 u], H o \ell[
 u]\ell[
 u]$

tTH hadroproduction

y $_{t}$ can be measured in pp \rightarrow tTH through many decay

channels (all very difficult):



- ullet hadrons with single lepton: ${
 m t}
 ightarrow {
 m b} \ell
 u, {
 m ar t}
 ightarrow {
 m ar b} j j, {
 m H}
 ightarrow {
 m b} {
 m ar b}$
- ullet hadrons with dilepton: $t o b\ell
 u, \, ar{t} o ar{b}\ell
 u, \, H o bar{b}$
- ullet hadrons with hadronic tau: ${
 m t}
 ightarrow {
 m b}\ell
 u, {
 m ar t}
 ightarrow {
 m ar b}jj, {
 m H}
 ightarrow au_h^+ au_h^-$
- \circ diphoton with lepton: $\mathrm{t} o \mathrm{b}\ell
 u, \, \mathrm{ar{t}} o \mathrm{ar{b}} jj, \, \mathrm{H} o \gamma \gamma$
- ullet diphoton with hadrons: $\mathrm{t} o \mathrm{b} \ jj, \ \mathrm{ar{t}} o \mathrm{ar{b}} jj, \ \mathrm{H} o \gamma\gamma$
- ullet same sign dilepton: $\mathrm{t} o \mathrm{b} \, jj, \, \mathrm{ar{t}} o \mathrm{ar{b}} jj, \, \mathrm{H} o \ell
 u \ell[
 u]$
- ullet 3 leptons with di, trilepton: ${
 m t}
 ightarrow {
 m b}\ell
 u,\, {
 m ar t}
 ightarrow {
 m ar b}jj,\, {
 m H}
 ightarrow \ell[
 u]\ell[
 u]$
- 4 lepton with di, trilepton: $t\to b\ell \nu,\, ar t \to ar b\ell[\nu],\, H\to \ell[\nu]\ell[\nu]$

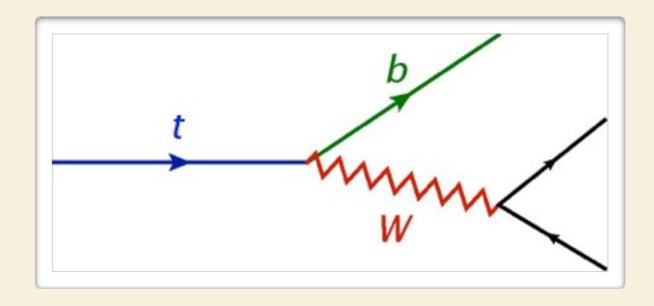
The importance of being top

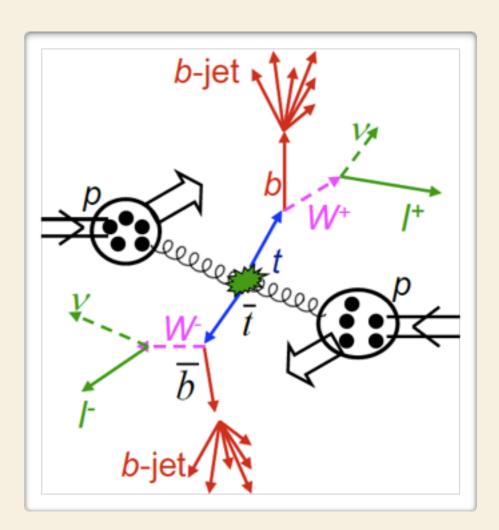
These require precise predictions of distributions at hadron level for $pp \rightarrow tT+hard X$, $X = H,W,Z,\gamma,j,bB,2j...$

...with decays: the t-quark is not detected because it

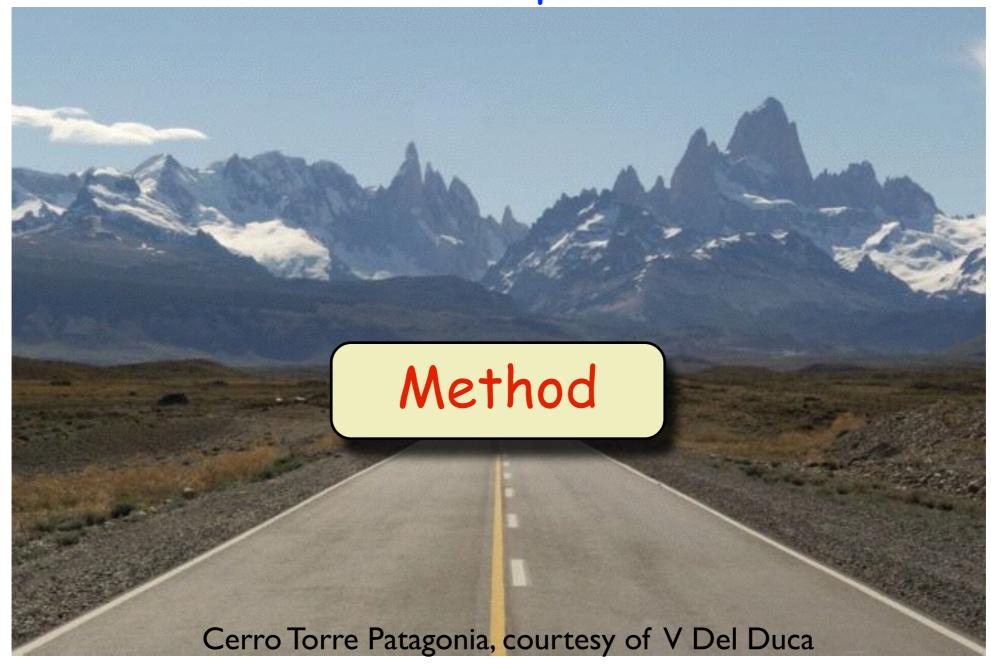
decays before hadronization

$$|V_{tb}|^2 \gg |V_{ts}|^2, |V_{td}|^2$$





... to distributions, full of pitfalls & difficulties



There is a long way from loops and legs...

SMC idea: use probabilistic picture of parton splitting in the collinear approximation, iterate splitting to high orders

SMC idea: use probabilistic picture of parton splitting in the collinear approximation, iterate splitting to high orders

Standard MC first emission:

$$d\sigma_{\text{SMC}} = B(\Phi_n)d\Phi_n \left[\Delta_{\text{SMC}}(t_0) + \Delta_{\text{SMC}}(t) \frac{\alpha_s(t)}{2\pi} \frac{1}{t} P(z) \Theta(t - t_0) d\Phi_{\text{rad}}^{\text{SMC}} \right]$$

$$= \lim_{k_{\perp} \to 0} R(\Phi_{n+1}) / B(\Phi_n)$$

SMC idea: use probabilistic picture of parton splitting in the collinear approximation, iterate splitting to high orders

Standard MC first emission:

$$d\sigma_{\text{SMC}} = B(\Phi_n) d\Phi_n \left[\Delta_{\text{SMC}}(t_0) + \Delta_{\text{SMC}}(t) \frac{\alpha_s(t)}{2\pi} \frac{1}{t} P(z) \Theta(t - t_0) d\Phi_{\text{rad}}^{\text{SMC}} \right]$$

$$= \lim_{k_{\perp} \to 0} R(\Phi_{n+1}) / B(\Phi_n)$$

POWHEG MC first emission:

$$d\sigma = \bar{B}(\Phi_n)d\Phi_n \left[\Delta(\Phi_n, p_{\perp}^{\min}) + \Delta(\Phi_n, k_{\perp}) \frac{R(\Phi_{n+1})}{B(\Phi_n)} \Theta(k_{\perp} - p_{\perp}^{\min}) d\Phi_{\text{rad}} \right]$$
$$\bar{B}(\Phi_n) = B(\Phi_n) + V(\Phi_n) + \int \left[R(\Phi_{n+1}) - A(\Phi_{n+1}) \right] d\Phi_{\text{rad}}$$

[Frixione, Nason, Oleari arXiv: 0709.2092]

SMC idea: use probabilistic picture of parton splitting in the collinear approximation, iterate splitting to high orders

Standard MC first emission:

$$d\sigma_{\text{SMC}} = B(\Phi_n) d\Phi_n \left[\Delta_{\text{SMC}}(t_0) + \Delta_{\text{SMC}}(t) \underbrace{\frac{\alpha_s(t)}{2\pi} \frac{1}{t} P(z) \Theta(t - t_0) d\Phi_{\text{rad}}^{\text{SMC}}}_{\text{rad}} \right]$$

$$= \lim_{k_{\perp} \to 0} R(\Phi_{n+1}) / B(\Phi_n)$$

POWHEG MC first emission:

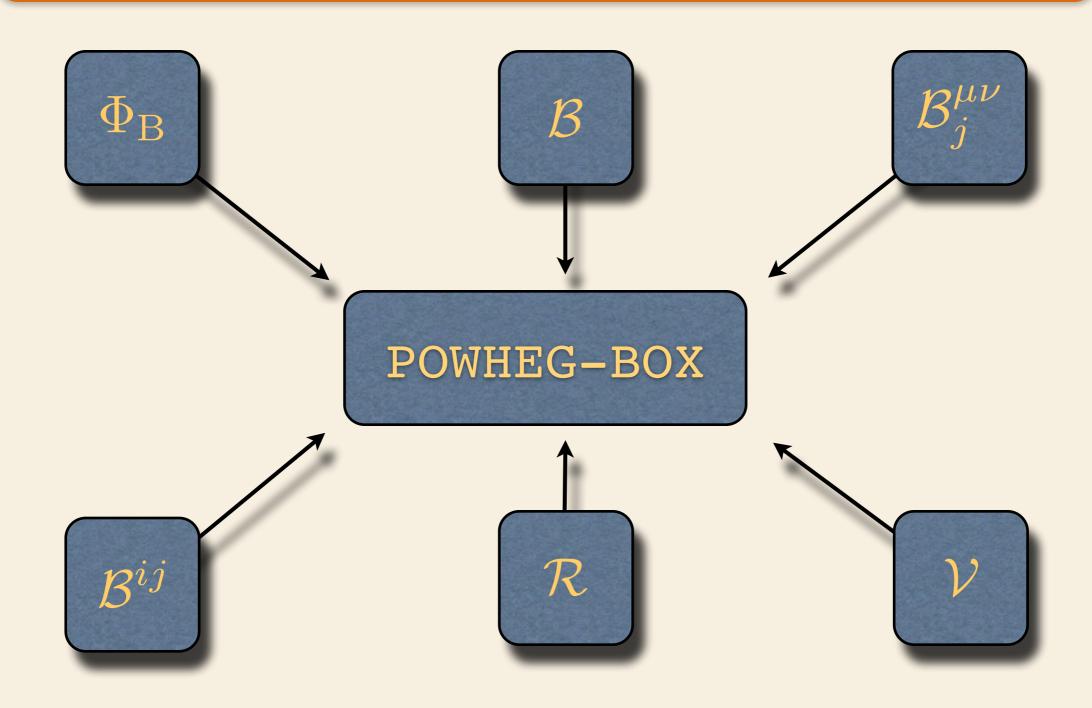
$$\mathrm{d}\sigma = \bar{B}(\Phi_n)\mathrm{d}\Phi_n \left[\Delta(\Phi_n, p_\perp^{\mathrm{min}}) + \Delta(\Phi_n, k_\perp) \frac{R(\Phi_{n+1})}{B(\Phi_n)} \, \Theta(k_\perp - p_\perp^{\mathrm{min}}) \, \mathrm{d}\Phi_{\mathrm{rad}}\right]$$

$$\bar{B}(\Phi) = B(\Phi_n) + V(\Phi_n) + \int \left[R(\Phi_{n+1}) - A(\Phi_{n+1})\right] \mathrm{d}\Phi_{\mathrm{rad}}$$

$$\bar{B}(\Phi_n) \, \mathrm{d}\Phi_n = \sigma_{\mathrm{NLO}}$$
[Frixione, Nason, Oleani

[Frixione, Nason, Oleari arXiv: 0709.2092]

POWHEG-BOX framework



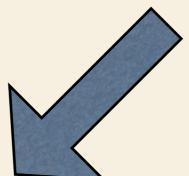
PowHel framework

HELAC-NLO

[Bevilaqua et al, arXiv: 1110.1499]



POWHEG-BOX



[Alioli, Nason, Oleari, Re, arXiv: 1002.2581]

PowHel

RESULT of PowHel:

Les Houches file of Born and Born+1st radiation events (LHE) ready for processing with SMC followed by almost arbitrary experimental analysis

- · Hadrons in final state
- •Closer to experiments, realistic analysis becomes feasible

- · Hadrons in final state
- •Closer to experiments, realistic analysis becomes feasible
- Decayed tops
- Parton shower can have significant effect
 (e.g. in Sudakov regions)

- · Hadrons in final state
- •Closer to experiments, realistic analysis becomes feasible
- Decayed tops
- Parton shower can have significant effect (e.g. in Sudakov regions)
- For the user:

event generation is, faster than an NLO computation (once the code is ready!)

- · Hadrons in final state
- •Closer to experiments, realistic analysis becomes feasible
- Decayed tops
- Parton shower can have significant effect (e.g. in Sudakov regions)
- For the user:

event generation is, faster than an NLO computation

(once the code is ready!)

...but we deliver the events on request

- Photons have to be isolated
- Easy at LO in perturbation theory: complete isolation in a cone around the y

- Photons have to be isolated
- Easy at LO in perturbation theory: complete isolation in a cone around the y
- Problematic at NLO:
 a completely isolated photon is not IR safe

- ·Photons have to be isolated
- Easy at LO in perturbation theory: complete isolation in a cone around the y
- Problematic at NLO:
- a completely isolated photon is not IR safe
- · Three solutions:
 - 1. use inclusive photons attractive only
 - 2. use smooth isolation I theoretically
- attractive only theoretically
 - 3. include photon fragmentation cumbersome theoretically and photon fragmentation is not, known well

Second paper with Zoltan

Nuclear Physics B394 (1993) 139–168 North-Holland NUCLEAR PHYSICS B

QCD corrections to photon production in association with hadrons in e⁺e⁻ annihilation *

Zoltan Kunszt and Zoltán Trócsányi 1

Theoretical Physics, ETH, Zurich, Switzerland

Received 17 July 1992 Accepted for publication 19 October 1992

Next-to-leading order QCD corrections for inclusive photon production in e^+e^- annihilation are derived. We emphasize that in a well-defined perturbative analysis – with or without isolation – it is always necessary to subtract the photon–quark collinear singularity. The subtraction term is absorbed into the non-perturbative fragmentation functions of the photon. The Q^2 -dependence of the photon fragmentation functions is determined by inhomogeneous evolution equations. The modification of the evolution equations due to photon isolation is discussed. We also analyse the validity of the approaches where the non-perturbative contributions are neglected. Using a general purpose next-to-leading order Monte Carlo program, we calculate various physical quantities that were measured in LEP experiments recently.

Second paper with Zoltan

Nuclear Physics B394 (1993) 139–168 North-Holland NUCLEAR PHYSICS B

QCD corrections to photon production in association with hadrons in e⁺e⁻ annihilation *

Zoltan Kunszt and Zoltán Trócsányi 1

Theoretical Physics, ETH, Zurich, Switzerland

Received 17 July 1992 Accepted for publication 19 October 1992

Next-to-leading order QCD corrections for inclusive photon production in e^+e^- annihilation are derived. We emphasize that in a well-defined perturbative analysis – with or without isolation – it is always necessary to subtract the photon–quark collinear singularity. The subtraction term is absorbed into the non-perturbative fragmentation functions of the photon. The Q^2 -dependence of the photon fragmentation functions is determined by inhomogeneous evolution equations. The modification of the evolution equations due to photon isolation is discussed. We also analyse the validity of the approaches where the non-perturbative contributions are neglected. Using a general purpose next-to-leading order Monte Carlo program, we calculate various physical quantities that were measured in LEP experiments recently.

 Problem with cone isolation at NLO: isolation cone cuts into phase space of soft gluons, hence not infrared safe

- Problem with cone isolation at NLO: isolation cone cuts into phase space of soft gluons, hence not infrared safe
- Solution: isolate photon from partons with cone radius decreasing as parton energy inside the cone decreases

$$E_{\perp,\text{had}} = \sum_{i \in \text{tracks}} E_{\perp,i} \Theta \left(\delta - R(p_{\gamma}, p_{i}) \right) < E_{\perp,\gamma} \left(\frac{1 - \cos \delta}{1 - \cos \delta_{0}} \right)$$

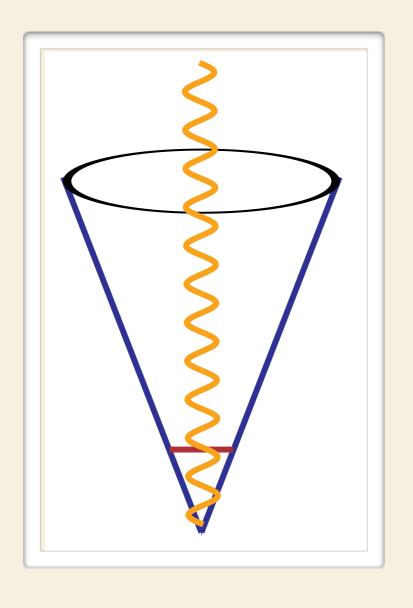
- Problem with cone isolation at NLO: isolation cone cuts into phase space of soft gluons, hence not infrared safe
- Solution: isolate photon from partons with cone radius decreasing as parton energy inside the cone decreases

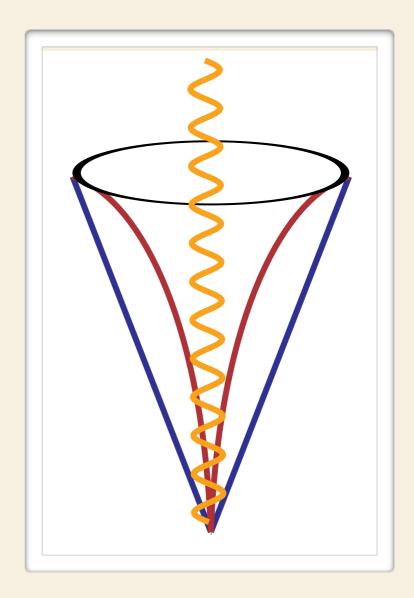
$$E_{\perp,\text{had}} = \sum_{i \in \text{tracks}} E_{\perp,i} \Theta\left(\delta - R(p_{\gamma}, p_{i})\right) < E_{\perp,\gamma} \left(\frac{1 - \cos \delta}{1 - \cos \delta_{0}}\right)$$

• Experimenters prefer cone with fixed radius, with reduced hadronic activity inside the cone

$$E_{\perp,\text{had}} = \sum_{i \in \text{tracks}} E_{\perp,i} \Theta \left(R_{\gamma} - R(p_{\gamma}, p_i) \right) < E_{\perp,\text{had}}^{\text{max}}$$

Experimental cone vs. smooth cone

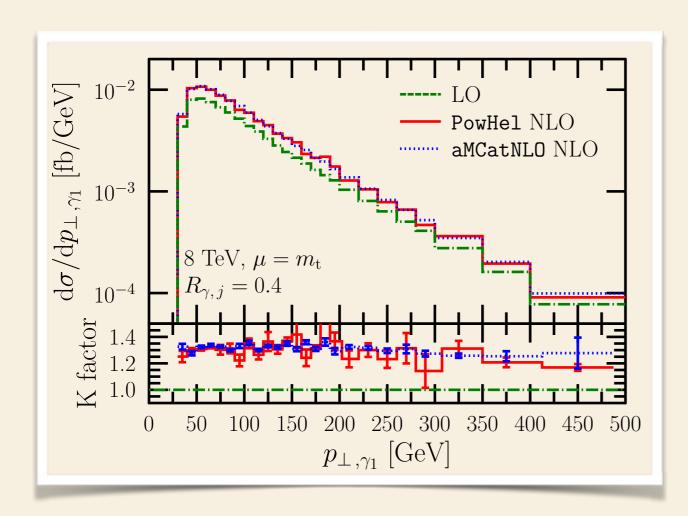


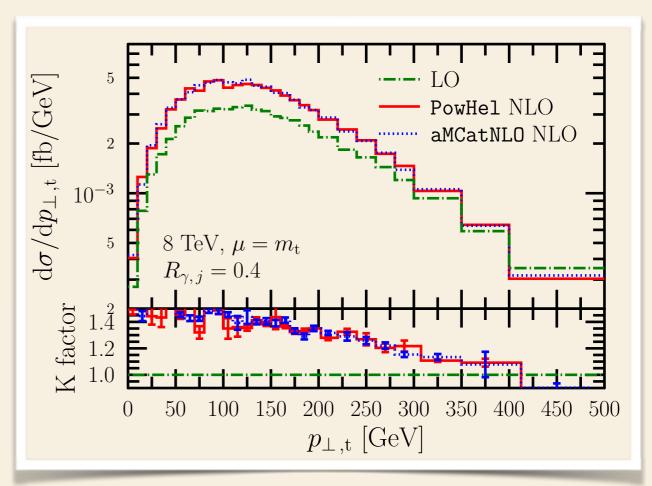


tTyy hadroproduction at NLO

Preliminary

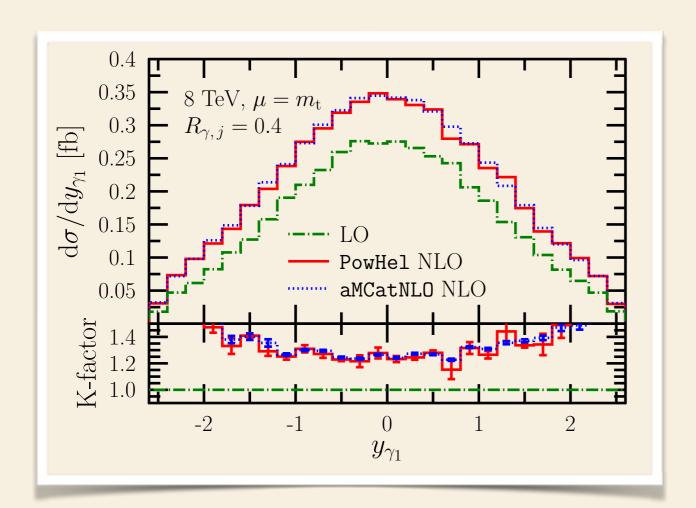
PowHel can produce distributions at NLO

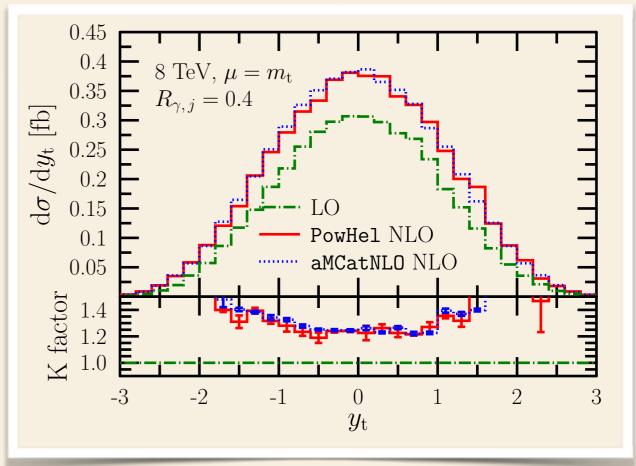




Cross sections agree with predictions of MadGraph5 if $m_b \rightarrow 0$: (1052±10) pb with given cuts

PowHel can produce distributions at NLO





NLO K-factor is in the range 1.2-1.5 for fixed default scale, improves with dynamic scale

NLO+PS matching makes possible (almost) inclusive event sample with photons

NLO+PS matching makes possible (almost) inclusive event sample with photons

- •Event generation requires generation cuts, chosen much smaller than physical cuts \Rightarrow
- measurable cross sections are independent of the generation cuts
- Applicable to processes without final-state light patrons in the Born cross section

Formal accuracy of the POWHEG MC

$$\langle O \rangle = \int d\Phi_{\rm B} \widetilde{B} \left[\Delta(p_{\perp,\rm min}) O(\Phi_{\rm B}) + \int d\Phi_{\rm rad} \Delta(p_{\perp}) \frac{R}{B} O(\Phi_{\rm R}) \right] =$$

• • •

Formal accuracy of the POWHEG MC

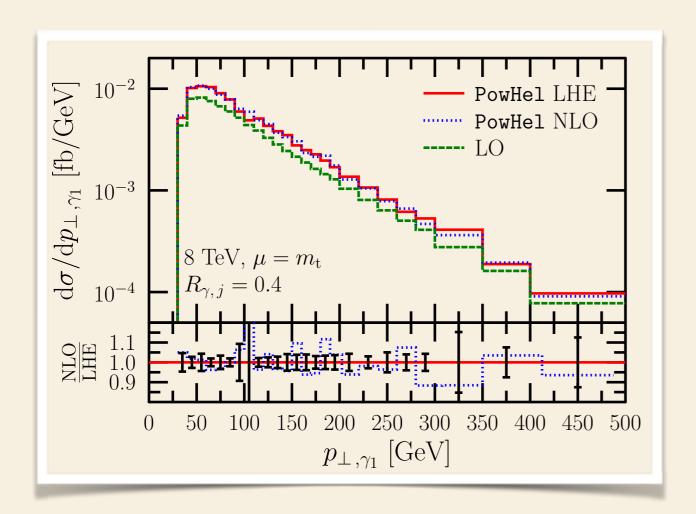
$$\langle O \rangle = \int d\Phi_{\rm B} \widetilde{B} \left[\Delta(p_{\perp, \rm min}) O(\Phi_{\rm B}) + \int d\Phi_{\rm rad} \Delta(p_{\perp}) \frac{R}{B} O(\Phi_{\rm R}) \right] =$$

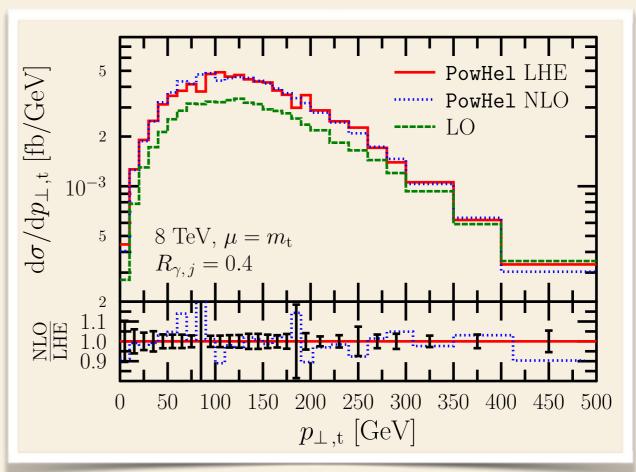
• • •

$$= \left\{ \int d\Phi_{\rm B} \left[B + V \right] O(\Phi_{\rm B}) + \int d\Phi_{\rm R} R O(\Phi_{\rm R}) \right\} (1 + \mathcal{O}(\alpha_{\rm S}))$$

Useful for checking

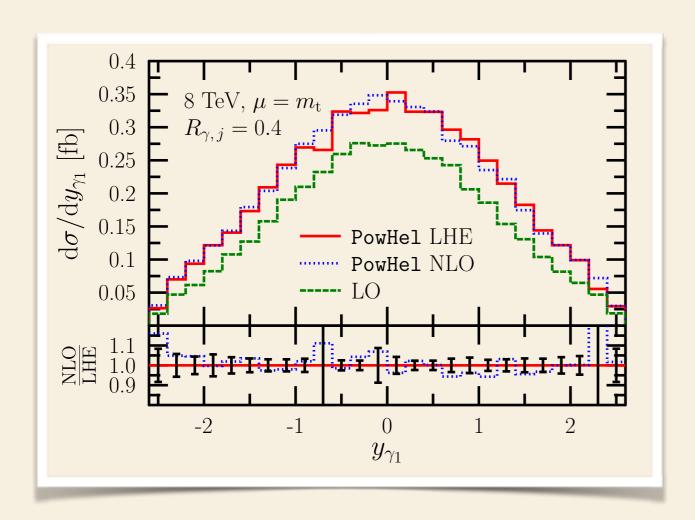
OLHE agree with ONLO w smooth isolation

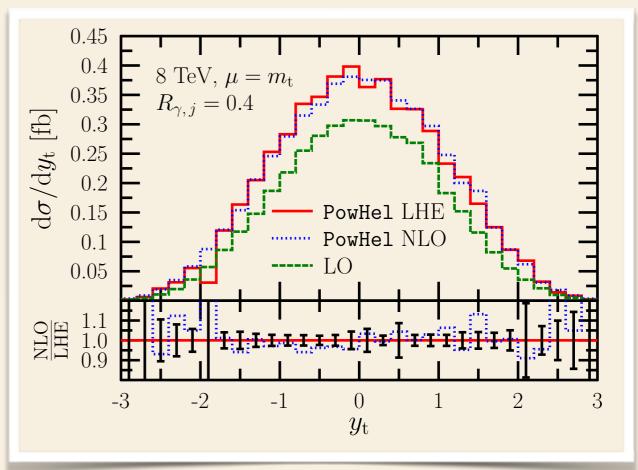




NLO and LHE predictions agree

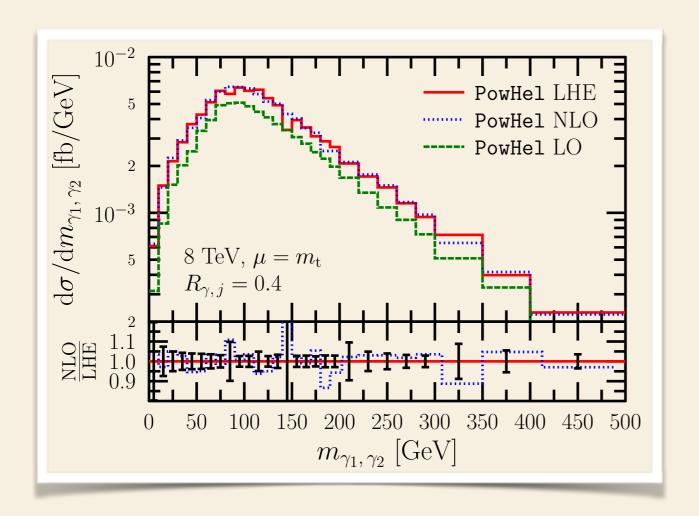
olhe agree with onlo w smooth isolation

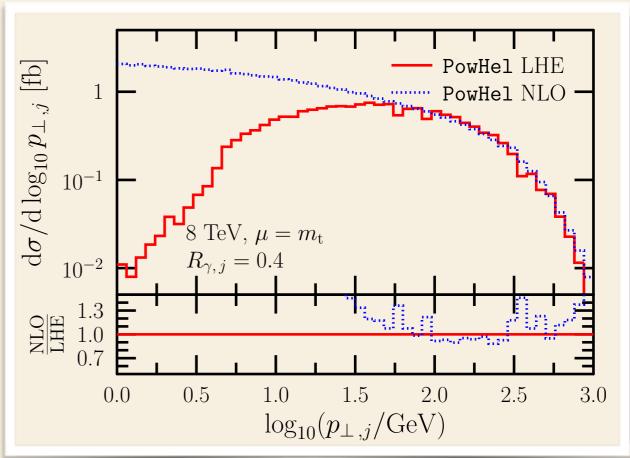




NLO and LHE predictions agree

olhe agree with onlo w smooth isolation





Exception: Sudakov damping

Message: we can trust the LHE's, so can make

Predictions

Four possible forms of predictions

LHE: distributions from events at BORN+1st radiation

Decay: on-shell decays of heavy particles (t-quarks), shower and hadronization effects turned off

PS: parton showering (PYTHIA or HERWIG) included (t-quarks kept stable)

Four possible forms of predictions

LHE: distributions from events at BORN+1st radiation

Decay: on-shell decays of heavy particles (t-quarks), shower and hadronization effects turned off

PS: parton showering (PYTHIA or HERWIG) included (t-quarks kept stable)

Full SMC: decays, parton showering and hadronization are included by using PYTHIA or HERWIG

Four possible forms of predictions

LHE: distributions from events at BORN+1st radiation

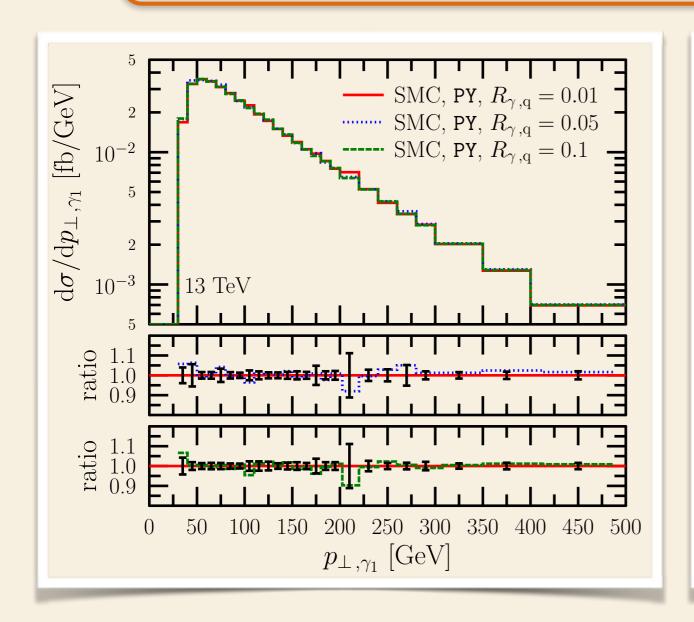
Decay: on-shell decays of heavy particles (t-quarks), shower and hadronization effects turned off

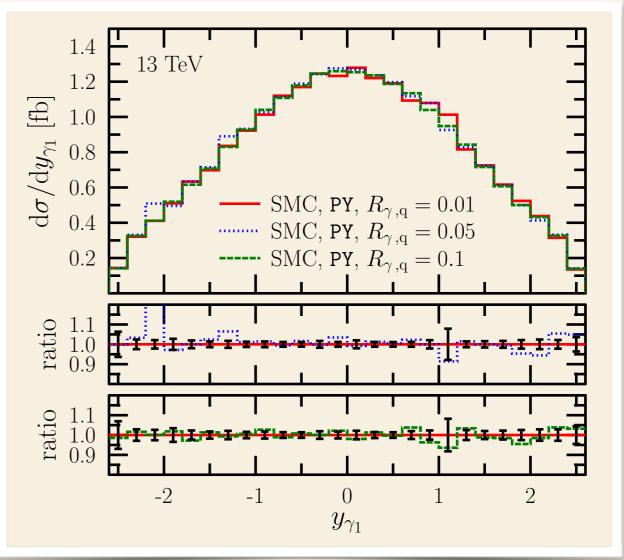
PS: parton showering (PYTHIA or HERWIG) included (t-quarks kept stable)

Full SMC: decays, parton showering and hadronization are included by using PYTHIA or HERWIG

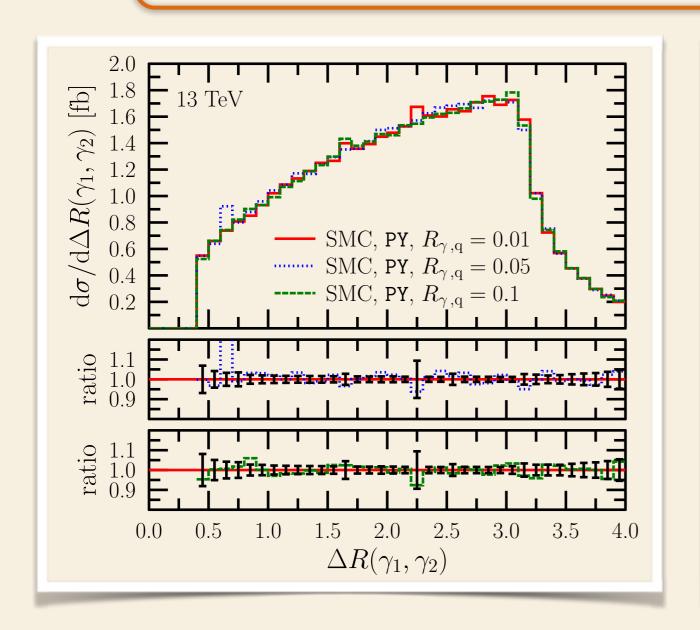
Number and type of particles are very different => to study the effect of SMC we employ selection cuts to keep the cross section fixed

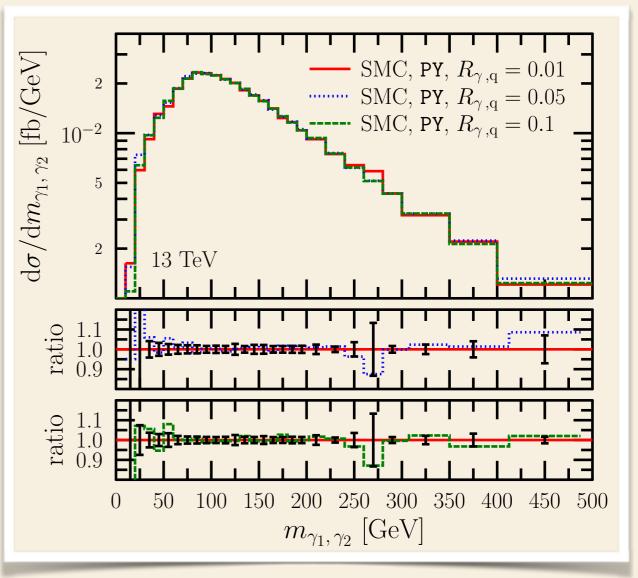
Predictions after full SMC with physical cuts are independent of generation cuts





Predictions after full SMC with physical cuts are independent of generation cuts





Message:

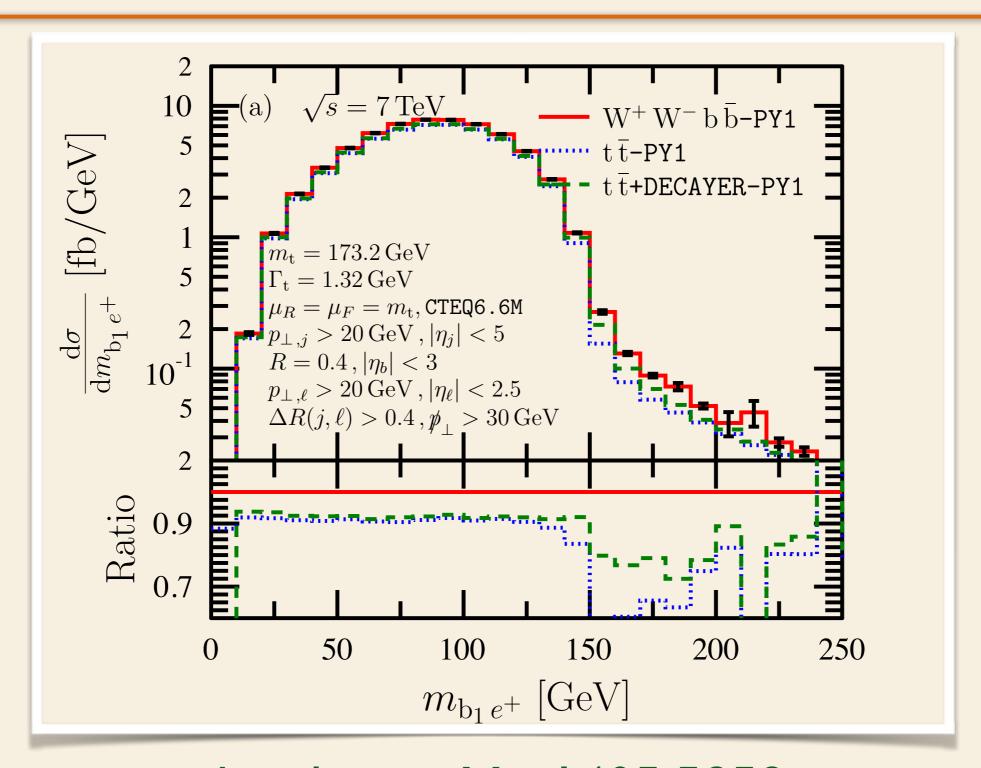
we can trust the (almost) inclusive LHE's, so can use experimental isolation

Conclusions

Conclusions

- First computation of pp \rightarrow ttyy at NLO + SMC accuracy
- NLO cross sections agree with predictions of public codes
- Experimentally preferred cone isolation is possible on (almost) inclusive event sample
- Effects of parton shower, SMC can be quantified
- LHE event files for pp → tt, ttH, ttW, ttZ, ttjet, ttbb, tty, ttyy, WWbb processes available, to put into SMC and perform experimental analyses on events with hadrons

An example: $pp \rightarrow t\bar{t} \rightarrow W^{\dagger}W^{\dagger}b\bar{b}$



details in arXiv:1405.5859

Processes available in PowHel

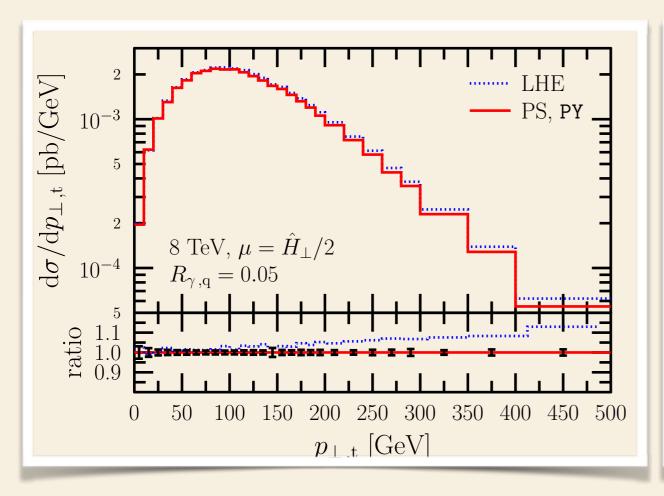
```
[Kardos et al, arXiv:
√†T
              1111.0610,1111.1444,
\sqrt{T+Z}
√†T + W
              1208.2665,
              1108.0387,
√+T + H/A
\sqrt{T+j}
              1101.2672,
              PoS LL2012 057, 1405.5659
√ WWbB
              1303.6291
√tT+bB
              ready to submit
\sqrt{T} + \gamma
\sqrt{T+VV}
              to be published soon]
```

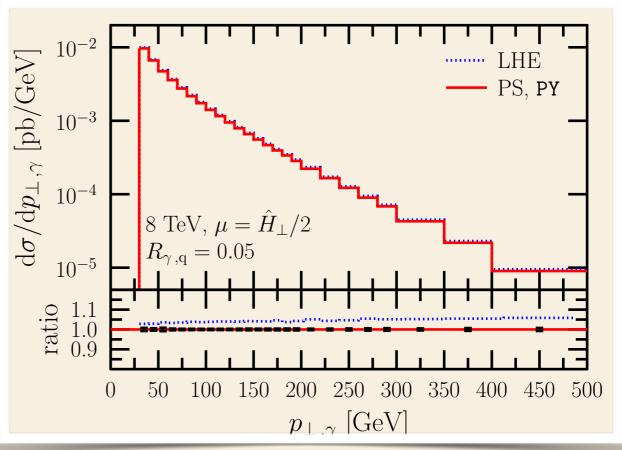
Kedves Zoli

Köszönöm az útnak indítást és Isten éltessen sokáig erőben, egészségben! Appendix

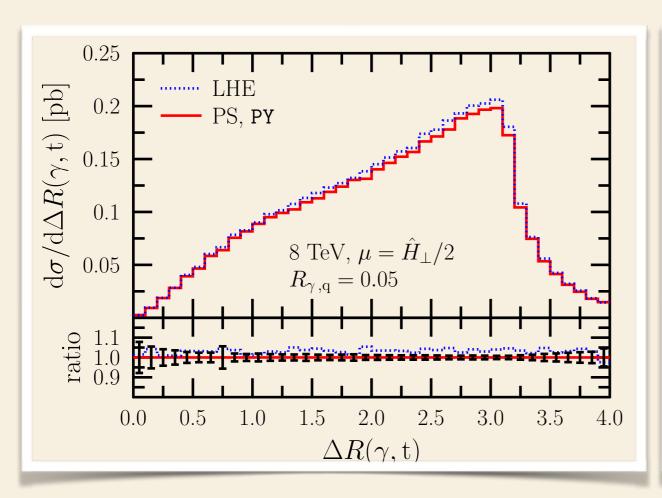
tTy hadroproduction at NLO

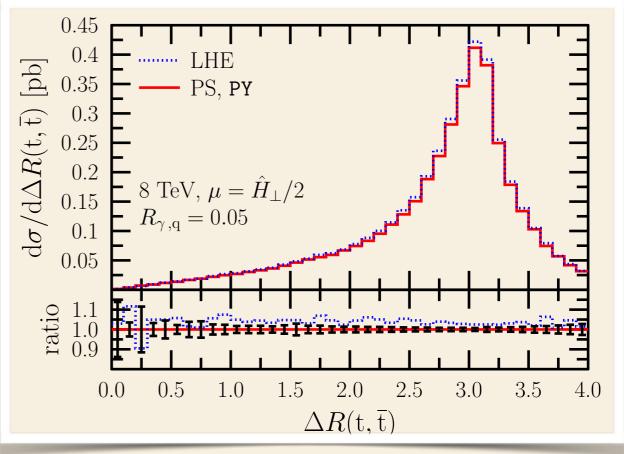
NLO vs. PS at 8TeV, $\mu = H_T/2$



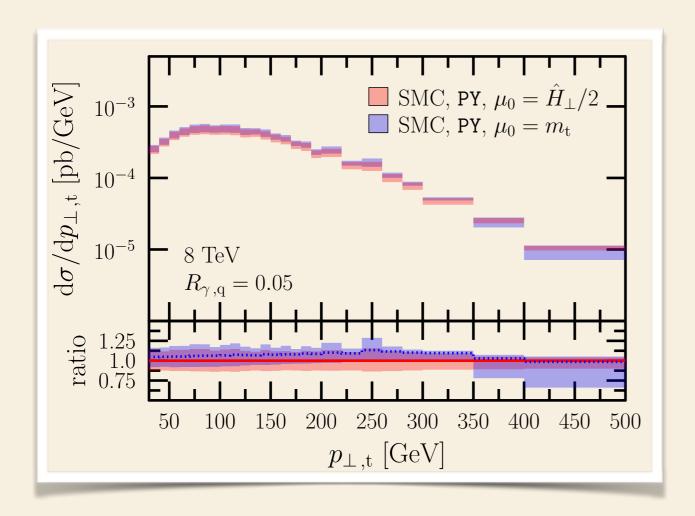


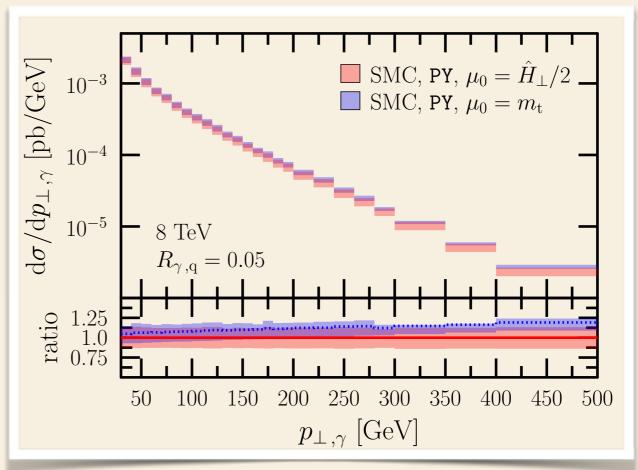
NLO vs. PS at 8TeV, $\mu = H_T/2$





Scale dependence after full SMC at 8TeV





Scale dependence after full SMC at 8TeV

