## pp→ttyy 23 years after  $\overline{\phantom{a}}$

#### Zoltán Trócsányi



University of Debrecen and MTA-DE Particle Physics Research Group in collaboration with A. Kardos

> Zoltan 70, Zürich May 23, 2014





## **Outline**

- Motivation
- Method
- Predictions
- Conclusions



Zoltan's excitement about a clear signal of observing a Higgs-boson of intermediate mass  $(70 \times m_H c^2 / GeV \times 140)$ 

irreducible background ttγγ final-state was unknown −<br>+

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- message: works with excellent resolution of photons
- … actually achieved by ATLAS and CMS

Physics Letters B 271 (1991) 247-255 PHYSICS LETTERS B North-Holland

#### **Clear signal of intermediate mass Higgs boson production at LHC and SSC**

Z. Kunszt, Z. Trócsányi<sup>1</sup> *Theoretical Physics, ETH, CH-8093 Zurich, Switzerland* 

and

W.J. Stirling *Departments of Physics and Mathematical Sciences, University of Durham, Durham DH I 3LE, UK* 

Received 11 July 1991

We compute the event rates for two potentially important background processes  $pp \rightarrow t\bar{t}\gamma\gamma X$  and  $pp \rightarrow b\bar{b}\gamma\gamma X$  to the recently proposed signature for Higgs production of one isolated lepton and two isolated photons in the intermediate-mass range. We find that the background can be suppressed assuming good ( $\approx$  (2–3)%) mass resolution in the invariant mass of the two photons, and assuming that the isolation criteria for the lepton and photons can be efficiently implemented experimentally. We reanalyse the signal to background ratios using realistic experimental cuts and find that by measuring the inclusive production of one isolated lepton and two isolated photons at the LHC or the SSC we can obtain a clear experimental signal for the production of the Higgs boson in the mass region  $70 < M_H < 140$  GeV.

## Higgs boson has been discovered



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 $m_H$  [GeV]=125.5±0.2<sub>stat</sub>±0.6<sub>syst</sub> (ATLAS 2013) 125.7±0.3stat±0.3syst (CMS 2013)

- All measured properties are consistent with SM expectations within experimental uncertainties
	- branching ratios as predicted
	- spin zero
	- parity +
	- $\circ$  couples to masses of W and Z (with  $c_v=1$  within experimental uncertainty)

## t-quark: potential tool for discovery

The t-quark is heavy, Yukawa coupling ∼1  $m_f$  [GeV]=173.34 $\pm$ 0.64 (LHC+TeVatron, 2014)  $( \Rightarrow y_1 = 0.997 \pm 0.003)$ 

⇒ plays important role in Higgs physics (more tantalizing:  $m_1 m_2 = (125.7 \pm 0.3)^2$  GeV<sup>2</sup>)  $\circ$  y<sub>t</sub> cannot be measured in H  $\rightarrow$  tT decay (m<sub>t</sub> > m<sub>H</sub>)

## How to measure yt?

- $\Theta$  H  $\rightarrow$  yy is sensitive to y<sub>t</sub> through t-quark loop,
	- but rates are small and W loop also contributes



- $g \circ g g \to H$  is sensitive to y<sub>t</sub> through t-quark loop
	- if only SM model particles contribute (so far xsec is consistent with SM) 000000



- $\theta$  gg  $\rightarrow$  H is sensitive to BSM physics
	- if  $y_t$  is measured separately

## tTH hadroproduction

 $\bullet$  y<sub>t</sub> can be measured in pp  $\rightarrow$  tTH through many decay channels (all very difficult): 000000

- hadrons with single lepton:  $t \to b \ell \nu, \, \overline{t} \to \bar{b} j j, \, H \to b \bar{b}$
- hadrons with dilepton:  $t \to b \ell \nu, \bar{t} \to \bar{b} \ell \nu, H \to b \bar{b}$
- hadrons with hadronic tau:
- $\bullet$  diphoton with lepton:
- $\bullet$  diphoton with hadrons:
- same sign dilepton:
- 3 leptons with di, trilepton:  $t \to b\ell\nu, \bar{t} \to \bar{b}jj, H \to \ell[\nu]\ell[\nu]$
- 4 lepton with di, trilepton:  $t \to b\ell\nu, \bar{t} \to \bar{b}\ell[\nu], H \to \ell[\nu]\ell[\nu]$



 $t \to b \; jj, \; \overline{t} \to b \; jj, \; H \to \gamma \gamma$  $t \to b \; jj, \; \bar{t} \to \bar{b}jj, \; H \to \ell \nu \ell[\nu]$ 

 $t \to b\ell\nu, \bar{t} \to \bar{b}jj, H \to \gamma\gamma$ 

 $t \to b\ell\nu, \bar{t} \to \bar{b}jj, H \to \tau_h^+ \tau_h^-$ 

## tTH hadroproduction

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## The importance of being top

These require precise predictions of distributions at hadron level for  $pp \rightarrow$  +T+hard X, X = H, W, Z,  $\gamma$ , j, bB, 2 j...

…with decays: the t-quark is not detected because it decays before hadronization  $b$ -ie  $|V_{tb}|^2 \gg |V_{ts}|^2$ ,  $|V_{td}|^2$ 





#### ...to distributions, full of pitfalls & difficulties



#### There is a long way from loops and legs...

SMC idea: use probabilistic picture of parton splitting in the collinear approximation, iterate splitting to high orders

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Standard MC first emission:

$$
d\sigma_{\rm SMC} = B(\Phi_n) d\Phi_n \left[ \Delta_{\rm SMC}(t_0) + \Delta_{\rm SMC}(t) \frac{\alpha_{\rm s}(t)}{2\pi} \frac{1}{t} P(z) \Theta(t - t_0) d\Phi_{\rm rad}^{\rm SMC} \right]
$$

$$
= \lim_{k_{\perp} \to 0} R(\Phi_{n+1}) / B(\Phi_n)
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$$
  

$$
\int B(\Phi_n) d\Phi_n = \sigma_{\rm LO}
$$
  

$$
\circ
$$
 **POWHEG MC first emission:**

$$
d\sigma = \bar{B}(\Phi_n) d\Phi_n \left[ \Delta(\Phi_n, p_{\perp}^{\min}) + \Delta(\Phi_n, k_{\perp}) \frac{R(\Phi_{n+1})}{B(\Phi_n)} \Theta(k_{\perp} - p_{\perp}^{\min}) d\Phi_{\text{rad}} \right]
$$

$$
\bar{B}(\Phi_n) = B(\Phi_n) + V(\Phi_n) + \int \left[ R(\Phi_{n+1}) - A(\Phi_{n+1}) \right] d\Phi_{\text{rad}}
$$

[Frixione, Nason, Oleari arXiv: 0709.2092]

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$$
  
\n[Frixione, Nason, Oleari  
\n
$$
a\text{rXiv: O7O9.2092}
$$

#### POWHEG-BOX framework





15 Les Houches file of Born and Born+1st radiation events (LHE) ready for processing with SMC followed by almost arbitrary experimental analysis

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16 event generation is, faster than an NLO computation (once the code is ready!) ...but we deliver the events on request

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- •Three solutions:
	- 1. use inclusive photons 1 attractive only 1. use inclusive photons 1 attractive or<br>2. use smooth isolation 5 theoretically
	-
	- 17 3. include photon fragmentation cumbersome theoretically and photon fragmentation is not, known well

#### Second paper with Zoltan

Nuclear Physics B394 (1993) 139–168 North-Holland

NUCLEAR<br>PHYSICS B

#### **QCD** corrections to photon production in association with hadrons in  $e^+e^-$  annihilation  $*$

Zoltan Kunszt and Zoltán Trócsányi<sup>1</sup>

*Theoretical Physics, ETH, Zurich, Switzerland*

Received 17 July 1992 Accepted for publication 19 October 1992

Next-to-leading order QCD corrections for inclusive photon production in  $e^+e^-$  annihilation are derived. We emphasize that in a well-defined perturbative analysis — with or without isolation — it is always necessary to subtract the photon—quark collinear singularity. The subtraction term is absorbed into the non-perturbative fragmentation functions of the photon. The  $Q^2$ -dependence of the photon fragmentation functions is determined by inhomogeneous The  $Q$  -dependence of the photon inaginemation functions is determined by impointigeneous discussed. We also analyse the validity of the approaches where the non-perturbative contribudiscussed. We also analyse the validity of the approaches where the non-perturbative contributions are neglected. Using a general purpose next-to-leading order Monte Carlo program, we calculate various physical quantities that were measured in LEP experiments recently.

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• Solution: isolate photon from partons with cone radius decreasing as parton energy inside the cone decreases

 $E_{\perp\,\text{,had}} = \sum E_{\perp\,,i} \Theta\left(\delta - R(p_{\gamma},p_{i})\right) < E_{\perp\,,\gamma}$  $i \in$ tracks  $\sqrt{1-\cos\delta}$  $1 - \cos \delta_0$ ◆

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- Experimenters prefer cone with fixed radius, with reduced hadronic activity inside the cone  $E_{\perp\,\,\mathrm{,had}} = \sum_{}~E_{\perp\,\,,i} \Theta\left(R_\gamma - R(p_\gamma,p_i)\right) < E_{\perp\,\,\mathrm{,had}}^{\max}$

 $i \in$ tracks

## Experimental cone vs. smooth cone



# Preliminary

tTγγ hadroproduction at NLO

#### PowHel can produce distributions at NLO



Cross sections agree with predictions of MadGraph5 if  $m_b \rightarrow O$ : (1052±10) pb with given cuts

#### PowHel can produce distributions at NLO



NLO K-factor is in the range 1.2-1.5 for fixed default scale, improves with dynamic scale

## NLO+PS matching makes possible (almost) inclusive event sample with photons

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- •Event generation requires generation cuts, chosen much smaller than physical cuts  $\Rightarrow$
- measurable cross sections are independent of the generation cuts
- •Applicable to processes without final-state light patrons in the Born cross section

## Formal accuracy of the POWHEG MC

$$
\langle O \rangle = \int \mathrm{d}\Phi_\mathrm{B} \widetilde{B} \left[ \Delta(p_{\perp\,\mathrm{min}}) O(\Phi_\mathrm{B}) + \int \mathrm{d}\Phi_\mathrm{rad} \Delta(p_{\perp}) \frac{R}{B} O(\Phi_\mathrm{R}) \right] =
$$

 $\bullet\bullet\bullet$ 

$$
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$$
  
...  

$$
= \left\{ \int d\Phi_{\rm B} \left[ B + V \right] O(\Phi_{\rm B}) + \int d\Phi_{\rm R} RO(\Phi_{\rm R}) \right\} (1 + O(\alpha_{\rm S}))
$$
  
Useful for checking

#### σLHE agree with σNLO w smooth isolation



#### NLO and LHE predictions agree

#### σLHE agree with σNLO w smooth isolation



#### NLO and LHE predictions agree

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#### Exception: Sudakov damping

## Message: we can trust the LHE's, so can make



## Four possible forms of predictions

**LHE:** distributions from events at BORN+1st radiation

**Decay:** on-shell decays of heavy particles (t-quarks), shower and hadronization effects turned off

**PS:** parton showering (PYTHIA or HERWIG) included (t-quarks kept stable)

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**Full SMC:** decays, parton showering and hadronization are included by using PYTHIA or HERWIG

Number and type of particles are very different => to study the effect of SMC we employ selection cuts to keep the cross section fixed

## Predictions after full SMC with physical cuts are independent of generation cuts



## Predictions after full SMC with physical cuts are independent of generation cuts



## Message: we can trust the (almost) inclusive LHE's, so can use experimental isolation

# Conclusions |

## Conclusions

- First computation of pp *→* ttγγ at NLO + SMC accuracy −<br>+
- NLO cross sections agree with predictions of public codes
- Experimentally preferred cone isolation is possible on (almost) inclusive event sample
- Effects of parton shower, SMC can be quantified
- LHE event files for  $pp \rightarrow \, \texttt{tt}$ ,  $\texttt{ttH}$ ,  $\texttt{ttW}$ ,  $\texttt{ttZ}$ ,  $\texttt{ttjet}$ ,  $\texttt{ttbb}$ , t $\overline{\textsf{ty}}$ , t $\overline{\textsf{ty}}$ , WWbb processes available, to put into SMC and perform experimental analyses on events with hadrons − − − − − − − − − −

# $An example: pp \rightarrow \overline{f} \rightarrow W^*W^*b\overline{b}$



#### details in arXiv:1405.5859

#### Processes available in PowHel

✓tT  $\sqrt{1T+Z}$  $\sqrt{1T + W}$  $\sqrt{1T}$  +  $H/A$  $\sqrt{1}T + j$ ✓WWbB  $\sqrt{1T + bB}$  $\sqrt{1}T + y$  $\sqrt{1T + v v}$ [Kardos et al, arXiv: 1111.0610,1111.1444, 1208.2665, 1108.0387, 1101.2672, PoS LL2012 057, 1405.5659 1303.6291 ready to submit to be published soon]

#### Kedves Zoli

### Köszönöm az útnak indítást és Isten éltessen sokáig erőben, egészségben!



# tTγ hadroproduction at NLO

#### $NLO$  vs. PS at 8TeV,  $\mu = H_T/2$



#### $NLO$  vs. PS at 8TeV,  $\mu = H_T/2$



## Scale dependence after full SMC at 8TeV



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