

*Winter School on Trapped Charged Particles 2015,
Les Houches, France*

**Tutorial exercises related to the
guiding center approximation**

Problem 1 (2014.2.1)

Single particle orbit in the low magnetic field limit.

In the lecture, we found that particles drift across the magnetic field at the $\mathbf{E} \times \mathbf{B}$ drift velocity. In the inertial frame moving at this velocity, the electric field vanishes. This is valid for $E/B < c$. Imagine what happens in the opposite case, $E/B > c$. Note that this occurs for a finite E-field, when B is sufficiently small, so this is not an unphysical situation per se.

- a) If B is in the z direction and E is in the x direction, and you start a particle at zero velocity at $(x,y)=(0,0)$, will the particle ever come back to the $x=0$ plane?
- b) Find an inertial frame in which the particle experiences a simple linear acceleration. In order to do this, you need to consider the Lorentz transformation of the electric and magnetic fields from a stationary frame to a frame moving at velocity v . Note that when $E/B < c$, the frame moving at the $\mathbf{E} \times \mathbf{B}$ velocity is precisely the frame in which the electric field is zero. However, in the situation you now need to consider, the $\mathbf{E} \times \mathbf{B}$ velocity is larger than the speed of light so you cannot transform to the $\mathbf{E} \times \mathbf{B}$ frame anymore.

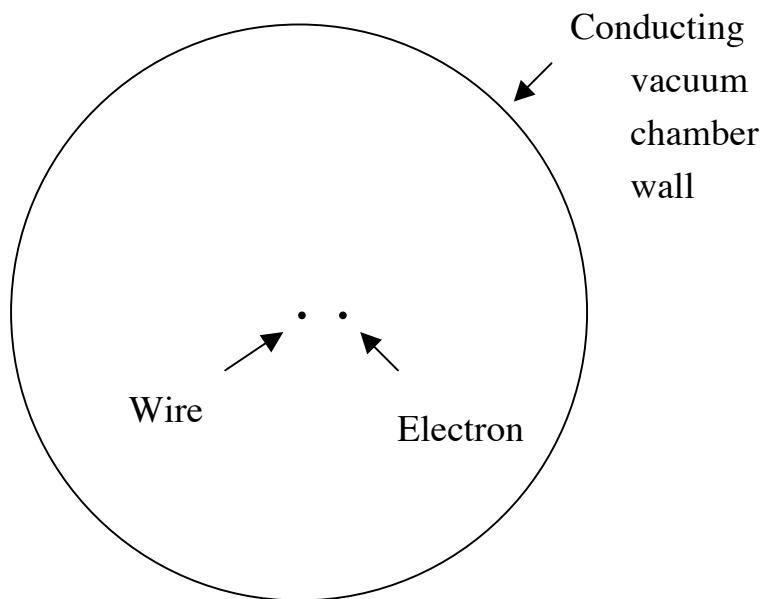
Background information: The electric and magnetic fields \mathbf{E}' and \mathbf{B}' in a frame moving at velocity v are related to the electric and magnetic fields \mathbf{E} and \mathbf{B} in a stationary frame by the Lorentz transformation (SI units):

$$\vec{E}' = \gamma(\vec{E} + \vec{v} \times \vec{B}) - \frac{\gamma^2}{\gamma+1} \frac{\vec{v}}{c} (\frac{\vec{v}}{c} \cdot \vec{E})$$

$$\vec{B}' = \gamma(\vec{B} - \frac{\vec{v}}{c^2} \times \vec{E}) - \frac{\gamma^2}{\gamma+1} \frac{\vec{v}}{c} (\frac{\vec{v}}{c} \cdot \vec{B})$$

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Hint: Note that there is some symmetry between E and B in the Lorentz transform. You may be able to guess what frame of reference eliminates the B -field in this case, by considering this symmetry.



Problem 2 (2014.2.2)

An electron lies, initially motionless, in the magnetic field of an infinite straight wire carrying a current \mathbf{I} . The wire is in the center of an infinite cylindrical conducting grounded vacuum chamber (see end-on view above). At $t=0$, the wire is suddenly charged to a positive potential ϕ without affecting \mathbf{I} . The electron gains energy from the electric field and begins to drift.

- (a) Draw a diagram showing the orbit of the electron and the relative directions of \mathbf{I} , \mathbf{B} , \mathbf{v}_E , \mathbf{v}_{VB} , \mathbf{v}_{R_c} .
- (b) Calculate the magnitudes of these drifts if the electron is at a radius of 1 cm and $\mathbf{I} = 500$ Amperes, $\phi = 460$ Volts, and the radius of the wire is 1 mm. Assume that ϕ is held at 0 Volts on the vacuum chamber walls 10 cm away.

Problem 3 (2014.3.1)

An electron (charge $-e$, mass m_e) is in a cylindrically symmetric mirror trap with an applied, uniform electric field. Use cylindrical coordinates (r, θ, z) . The particle has its guiding center located at $r=0$. On the $r=0$ line, the electric and magnetic fields are both time-independent and point in the z -direction and have the following magnitudes:

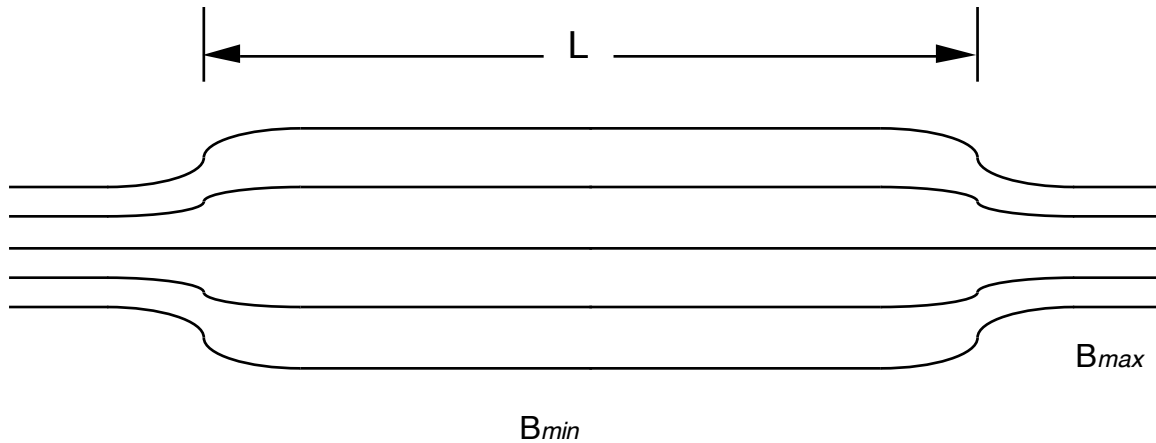
$$E_z = E_0, \quad B_z = B_0(1 + z^2/(z_0)^2).$$

Here E_0 , B_0 and z_0 are positive constants. The electron has $v_\perp = v_0$ and $v_\parallel = 0$ and is located at $z=0$. You may assume that v_0 is nonzero but small enough that the Larmor radius is small compared to all other spatial scales in the problem.

- Show that the particle is going to perform sinusoidal oscillations in the z direction, and express the oscillation frequency in terms of the parameters given in the problem.
- What is the oscillation amplitude for the electron's motion?
- Find the z -location at which an electron with perpendicular velocity v_0 and zero parallel velocity, and its guiding center at $r=0$, will remain at rest at that z -location – ie. find the equilibrium z -location for such an electron.

Problem 4 (2014.3.2)

Consider a charged particle trapped in a cylindrically symmetric magnetic mirror system with highly localized mirror regions where the magnetic field changes from B_{\min} to B_{\max} with $B_{\max} = 3 B_{\min}$ over a distance $\ll L$:



- (a) If a proton in the region where $B=B_{\min}$ has $v_{\perp} = 2v_{\parallel} = 2v_0$, compute the value of the magnetic moment, μ , and estimate the second adiabatic invariant, $J = \int v_{\parallel} dl$
- (b) If we slowly increase the magnetic field preserving the ratio of B_{\max}/B_{\min} , will the proton escape from the mirror system? If so, by what factor must the field be increased to allow the proton to escape from the mirror?

Problem 5 (2014.3.3)

When we derived the mirror force for the simple cylindrical system in class, we also found a force in the θ direction. Show that this force changes the particle's perpendicular velocity in just the right way as to conserve μ .