

① SU(8) FAMILY UNIFICATION WITH
BOSON-FERMION BALANCE

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MOTIVATION:

SM HAS BOSONS AND FERMIONS → SUPERSYMMETRY?

BUT: • NO DEFINITIVE MECHANISM FOR SUSY BREAKING

• SO FAR, SUSY NOT FOUND

TRY WEAKER PRINCIPLE THAN EXACT SUSY:

BOSON-FERMION BALANCE

THAT IS, EQUAL NUMBERS OF MASSLESS BOSON
AND FERMION HELICITY STATES

② STARTING POINT: $SO(8)$ $N=8$ SUPERGRAVITY

PARTICLE	HELICITIES
1 GRAVITON	2
8 GRAVITINOS	16
28 VECTORS	56
56 MAJORANA FERMIONS	112
70 SCALARS	70

$16+112=128$

$2+56+70=128$

28 \equiv ADJOINT OF $SO(8)$ TOO SMALL TO CONTAIN SM

REDISTRIBUTE 70 SCALAR HELICITIES TO VECTORS

$\frac{70}{2} = 35$

$35 + 28 = 63$

CAN BE ADJOINT OF $SU(8)$ BIG ENOUGH!

	SPIN \rightarrow 2	WEYL $\frac{3}{2}$	1	WEYL $\frac{1}{2}$	COMPLEX 0
" $SO(8)$ GRAVITON MULTIPLY" $^{\circ}$	1	8_L	63	56_L	$\leftarrow SU(8)$ REPS
" $SU(8)$ MATTER MULTIPLY" $^{\circ}$				$\bar{28}_L$ 28_L	56

③ SU(8) ANOMALY CANCELLATION

CHIRAL ANOMALY OF SPIN $3/2$

= 3 TIMES CHIRAL ANOMALY OF SPIN $1/2$

$$3 \times \text{ANOMALY}(8_L) = 3 \times 1 = 3$$

$$2 \times \text{ANOMALY}(\overline{28}_L) = 2 \times (-4) = -8$$

$$\text{ANOMALY}(56_L) = 5$$

$$\text{TOTAL ANOMALY} = 3 - 8 + 5 = 0$$

[THIS COUNTING FIRST NOTED BY MARCUS (1985)]

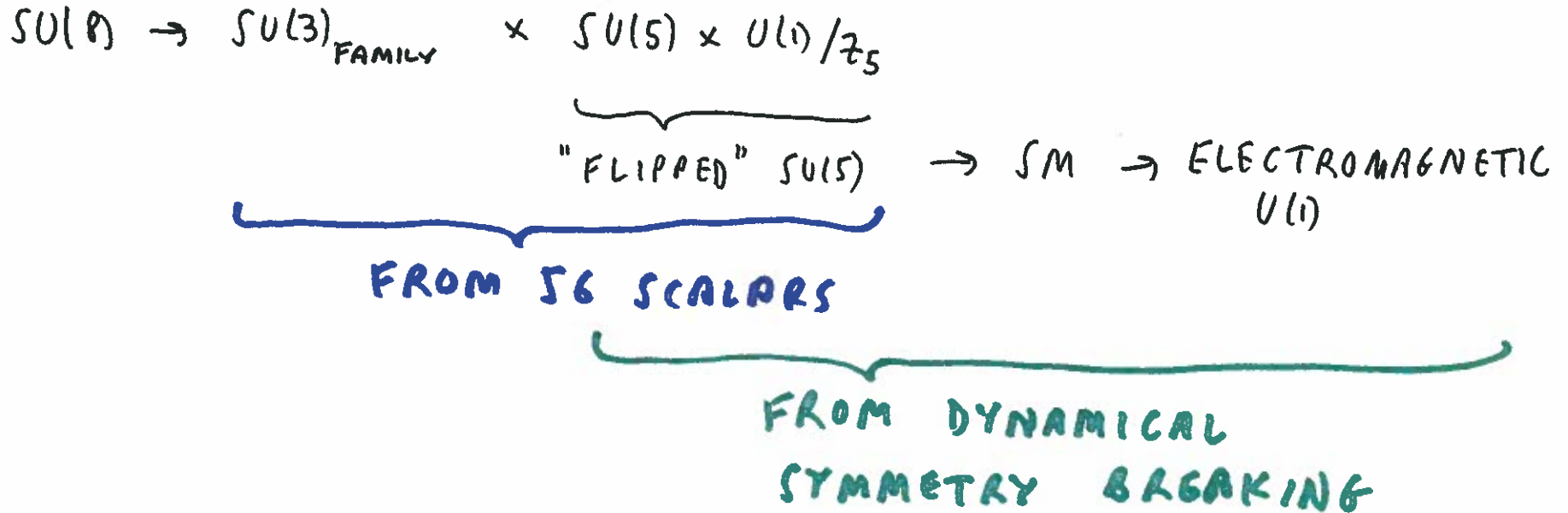
[SPIN $3/2$ ANOMALY DUFF (1982)
NIELSEN + RÖMER (1985)]

SO THE MODEL IS VIABLE

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WILL NOW SHOW:

- MODEL CONTAINS FIELDS OF SM
- MODEL HAS A SYMMETRY BREAKING PATH TO SM



UNDER $SU(8) \rightarrow SU(3) \times SU(5) \times U(1)$

$$56 \rightarrow (1, 1) (-15) + (1, 10) (9) + (\bar{3}, 5) (-7) + (3, 10) (1)$$

\uparrow
 $U(1)$ GENERATOR $\neq 0$

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MODULAR GROUND STATE

$|\Omega\rangle =$ GROUND STATE BEFORE DYNAMICAL SYMMETRY BREAKING

NEED $\langle \Omega | (1,1) (-15) | \Omega \rangle \neq 0$

$\rightarrow |\Omega\rangle$ MUST BE A SUPERPOSITION OF $U(1)$ EIGENSTATES
DIFFERING BY 15 UNITS

MODULAR GROUND STATE MODULO DIVISOR OF 15

ASSUME MOD 5 IS NATURE'S CHOICE

$$|\Omega\rangle = \sum_{h=-\infty}^{\infty} e^{i h \omega} |5h\rangle = |0\rangle + e^{i\omega} |5\rangle + e^{2i\omega} |10\rangle + e^{3i\omega} |15\rangle + e^{-i\omega} |-5\rangle + \dots$$

$\{ \} =$ MODULO 5 EQUIVALENT TO $()$

$$56 \rightarrow (1,1) (-15) \{0\} + (1,10) (9) \{-1\} + (3,5) (-7) \{-2\} + (3,10) (1) \{1\}$$

WILL SHOW MODULO 5 REPRESENTATIVES THAT CORRESPONDS
TO PATH TO SM: ASSUME UNIQUE MOD 5 REPRESENTATIVE IS
PICKED BY DYNAMICAL SYMMETRY BREAKING

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FIELD CONTENT OF MODEL

<u>FIELD</u>	<u>SPIN</u>	<u>SU(8) REP</u>	<u>HELICITIES</u>	<u>BRANCHING \rightarrow SU(3) x SU(5) x U(1)</u>
$h_{\mu\nu}$	2	1	2	1
ψ_{μ}^{α}	WEYL $\frac{3}{2}$	8_L	16	$(3,1)(-5) \{0\} + (1,5)(3) \{-2\}$
A_{μ}^A	1	63	126	$(1,1)(0) \{0\} + (8,1)(0) \{0\} + (1,24)(0) \{0\}$ $+ (3,3)(-1) \{2\} + (\bar{3},5)(1) \{-2\}$
$\chi^{[\alpha\beta\gamma]}$	WEYL $\frac{1}{2}$	56_L	112	$(1,1)(-15) \{0\} + (1, \bar{10})(9) \{-1\} + (\bar{3},5)(-7) \{3\} + \underline{(3,10)(1) \{1\}}$
$\lambda_1[\alpha\beta]$	WEYL $\frac{1}{2}$	$\bar{28}_L$	56	$(3,1)(10) \{5\} + (1, \bar{10})(-6) \{-1\} + (\bar{3}, \bar{5})(2) \{-3\}$
$\lambda_2[\alpha\beta]$	WEYL $\frac{1}{2}$	$\bar{28}_L$	56	$\underline{(3,1)(10) \{5\}} + (1, \bar{10})(-6) \{-1\} + \underline{(\bar{3},5)(2) \{-3\}}$
$\phi^{[\alpha\beta\gamma]}$	COMPLEX 0	56	112	$(1,1)(-15) \{0\} + \underline{(1, \bar{10})(9) \{-1\}} + \underline{(3,10)(1) \{1\}}$ $+ (\bar{3},5)(-7) \{-2\}$

↔ STATES WITH FLIPPED SU(5) REPRESENTATIONS

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SYMMETRY BREAKING $SU(9) \rightarrow SU(3) \times SU(5) \times U(1)$

BREAK USING SCALAR $\phi^{(AB)}$

A_μ^A

$$(1,1)(0)\{0\} + (8,1)(0)\{0\} + (1,24)(0)\{0\}$$

$\begin{matrix} \nearrow \\ U(1) \end{matrix}$
 $\begin{matrix} \nearrow \\ SU(3) \end{matrix}$
 $\begin{matrix} \nearrow \\ SU(5) \end{matrix}$

$$+ (3, \bar{5})(-8)\{2\} + (\bar{3}, 5)(8)\{-2\}$$

ABSORB 30 REAL COMPONENTS OF $\phi^{(AB)}$
TO BECOME MASSIVE

$\phi^{(AB)}$

$$(1,1)(-15)\{0\} + \underline{(1, \bar{10})(9)\{-1\}} + \underline{(3, 10)(1)\{1\}}$$

$\begin{matrix} \nearrow \\ \text{SETS VACUUM} \\ \text{EXPECTATION} \end{matrix}$
 $\underbrace{\hspace{10em}}$
RESIDUAL SCALARS

$$+ (\bar{3}, 5)(-7)\{-2\}$$

30 REAL COMPONENTS THAT ARE ABSORBED
BY VECTOR $(\bar{3}, 5)(8)\{-2\}$ AND CONJUGATE

NOTE: $8 \equiv -7 \pmod{15}$

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"STANDARD" SU(5)

GEORGI + GLASHOW
PRL 32, 438 (1974)

$SU(5) \supset SM$

$Q = T_L^3 + Y \leftarrow \text{(LANGACKER DEFINITION)}$

$Y = \frac{1}{6}$ SLANSKY $U(1)_{SM}$ GENERATOR
↑
PHYS. REP. 79, 1 (1981)

SU(5) REPS:

10
$\bar{5}$

"FLIPPED" SU(5)

BARR
PLB 112, 219 (1982) } FROM
SO(10)

DERENDINGER, KIM, NANOPOULOS
PLB 139, 176 (1984)

ANTONIADIS, ELLIS, HAGELIN, NANOPOULOS
PLB 194, 231 (1987)

$SU(5) \times U(1)_X \supset SM$

$Q = T_L^3 + Y$

$Y = \frac{1}{5} (X - Z)$

$Z = \frac{1}{6}$ SLANSKY $U(1)_{SM}$ GENERATOR

SU(5) REPS:

10
$\bar{5}$
1

X VALUES:

1
-3
5

DESCENT FROM SO(10) 16 RE P.

⇒ ALL $SU(5) \times U(1)_X$ ANOMALIES CANCEL

ALSO, ALL $SU(2)$ FAMILY $\times SU(5) \times U(1)_X$ " " "

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STATE ASSIGNMENTS IN STANDARD; FLIPPED SU(5)

<u>$(SU(2)_L, SU(3)_{\text{color}})_Y$</u>	<u>NAME</u>	<u>STD. SU(5)</u>	<u>FLIPPED SU(5)</u>	
$(1, \bar{3})_{-2/3}$	u_L^c	10	$\bar{5}$	FLIP 10 to $\bar{5}$
$(2, 1)_{-1/2}$	$\begin{pmatrix} e_L \\ \nu_L \end{pmatrix} = l_L$	$\bar{5}$	$\bar{5}$	
$(1, 1)_0$	N_L	ABSENT	10	
$(1, \bar{3})_{1/3}$	d_L^c	$\bar{5}$	10	FLIP $\bar{5}$ to 10
$(2, 3)_{1/6}$	$\begin{pmatrix} u_L \\ d_L \end{pmatrix} = q_L$	10	10	
$(1, 1)_1$	e_L^c	10	1	

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SYMMETRY BREAKING TO SM

$$SU(5) \rightarrow SU(2) \times SU(3) \times U(1)$$

$$\bar{5} = (2, 1) (-3) + (1, \bar{3}) (2)$$

$$10 = (1, 1) (6) + (1, \bar{3}) (-4) + (2, 3) (1)$$

$$24 = (1, 1) (0) + (3, 1) (0) + (1, 8) (0) + (2, 3) (-5) + (2, \bar{3}) (5)$$

STANDARD $SU(5)$: USES $(1, 1) (0)$ OF REAL SCALAR 24

FLIPPED $SU(5)$: USES $(1, 1) (6)$ OF COMPLEX SCALAR 10

↑ $U(1)$ GENERATOR $\neq 0$

MODULAR GROUND STATE MOD (6)

IN 24 VECTOR, $(2, 3) (-5) + (2, \bar{3}) (5)$ ABSORB 12 REAL

COMPONENTS OF COMPLEX 10 TO BECOME MASSIVE

THESE COME FROM $(2, 3) (1)$ AND CONJUGATE

$$\begin{array}{c} \downarrow \\ (2, 3) (-5) \quad | \equiv -5 \text{ MOD } (6) \end{array}$$

②

SYMMETRY BREAKING TO SM - CONTINUED

$$SU(8) \rightarrow SU(3)_{\text{FAMILY}} \times \underbrace{SU(5) \times U(1)}_{\substack{\uparrow \\ U(1)_X \\ \text{FLIPPED } SU(5)}}$$

↑
USE SCALAR ϕ_6

↳ RESIDUAL $(1, \bar{10}) \{-1\} + (3, 10) \{1\}$

CORRECT ASSIGNMENTS TO BREAK FLIPPED $SU(5) \rightarrow SM$

$$10 \supset (1, 1) (6) \quad Z = \frac{1}{6} (6) = 1$$

$$X = 1$$

$$\Rightarrow Y = \frac{1}{5} (X - Z) = 0$$

$$T_L^3 = 0 \text{ (SINGLET)} \Rightarrow Q = 0$$

$(3, 10)$ CAN ALSO BREAK FAMILY SYMMETRY $SU(3) \rightarrow SU(2) \times \frac{U(1)}{Z_2}$

AFTER $SU(5)$ BREAKING 10 DECOMPOSES INTO THREE COMPONENTS $10_{A,B,C}$

$(3, 10_{A,B,C})$ CAN COMPLETELY BREAK FAMILY $SU(3)$

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FERMION STATE SUMMARY

(SU(3) FAMILY, SU(5))

χ	<u>$(1, 1) \{0\}$</u> + $(1, \bar{10}) \{-1\}$ + $(\bar{3}, 5) \{3 \equiv -2\}$ + <u>$(3, 10) \{1\}$</u>
λ_1	$(3, 1) \{5\}$ + $(1, \bar{10}) \{-1\}$ + $(\bar{3}, \bar{5}) \{-3 \equiv 2\}$
λ_2	<u>$(3, 1) \{5\}$</u> + $(1, \bar{10}) \{-1\}$ + <u>$(\bar{3}, \bar{5}) \{-3\}$</u>

 3 SETS WITH FLIPPED SU(5) ASSIGNMENTS

 3 VECTOR-LIKE SETS BARR ^{avXw:} 1307.5770 ROTATE AWAY PROTON DELAY?

 DARK MATTER CANDIDATE?

POSSIBLE CONDENSATES:

- AFTER FAMILY BREAKING, THREE $(1, \bar{10}) \{-1\}$ WITH $(3, 10) \{1\}$
- $3 \times \bar{3} \supset 1$ PAIRINGS

$\chi (\bar{3}, 5) \{3\}$	$\lambda (3, 1) \{5\}$	$\supset (1, 5) \{3 \equiv -2\}$	}	MECHANISM TO GET SM HIGGS (MORE TO SAY ON THIS)
$\chi (3, 10) \{1\}$	$\lambda (\bar{3}, \bar{5}) \{-3\}$	$\supset (1, 5) \{-2\}$		
$\lambda_1 (3, 1) \{5\}$	$\lambda_2 (\bar{3}, \bar{5}) \{-3\}$	$\supset (1, \bar{5}) \{2\}$		

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ASYMPTOTIC FREEDOM AND GLOBAL SYMMETRIES

SU(8) IS ASYMPTOTICALLY FREE: $C(S) = \text{INDEX OF } SU(8) \text{ REP. OF SPIN } S$

$$\frac{1}{3} \left[11 C(1) - 26 C(\text{WEYL } \frac{3}{2}) - 2 C(\text{WEYL } \frac{1}{2}) - C(\text{COMPLEX } 0) \right]$$

$$= \frac{1}{3} \left[11 \times 16 - 26 \times 1 - 2 \times (15 + 2 \times 6) - 15 \right] = 27 > 0$$

GLOBAL CHIRAL SYMMETRIES

- OVERALL $U(1)$ REPHASING OF ALL FERMIONS - BROKEN BY USUAL INSTANTON + ANOMALY MECHANISM
- RELATIVE $U(1)$ REPHASING OF $56_L \chi$
- RELATIVE $U(2)$ REPHASING AND MIXING OF TWO $\overline{28}_L \lambda_{1,2}$

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DYNAMICAL SYMMETRY BREAKING: GENERAL FORMALISM

- FROM LEFT CHIRAL SPINORS, CAN FORM TWO TYPES OF SCALARS

$$\bar{\Psi}_{L1} \Psi_{L2} \equiv 0 \quad \left. \begin{aligned} \bar{\Psi}_{L1}^c \Psi_{L2} &= \Psi_{L1}^T i \gamma^0 \Psi_{L2} \\ &= \Psi_{L2}^T i \gamma^0 \Psi_{L1} \end{aligned} \right\} \neq 0 \quad \begin{array}{l} \text{ALLOWED} \\ \text{FORM OF} \\ \text{CONDENSATES} \end{array}$$

GROUP REP. CONTENT OF CONDENSATE IS DIRECT PRODUCT OF REP. CONTENT OF Ψ_1 AND Ψ_2

- VECTOR GLUON POTENTIAL FOR $A+B \rightarrow A+B$

$$V = \frac{g^2 K(A+B; A, B)}{2\pi}$$

$$K(A+B; A, B) = C_2(A+B) - C_2(A) - C_2(B)$$

$$C_2 = \text{CASIMIR} > 0$$

$$C_2(1) = 0$$

SPECIAL CASES:

$$K(1; A, \bar{A}) = C_2(1) - C_2(A) - C_2(\bar{A}) = -2C_2(A) < 0$$

$$K(A; \bar{A}, B) = C_2(A) - C_2(\bar{A}) - C_2(B) = -C_2(B) < 0$$

$$K(A; A, 1) = C_2(A) - C_2(A) - C_2(1) = 0$$

BOTH
ATTRACTIVE

(15) DYNAMICAL SYMMETRY BREAKING TO GET SM HIGGS

$SU(3): \quad 3 \times \bar{3} = \underline{1} + 8$

- = MOST ATTRACTIVE

$SU(5): \quad 10 \times \bar{5} = \underline{5} + 45$

$1 \times 5 = 5, \quad 1 \times \bar{5} = \bar{5}$

ALL FAMILY $3 \times \bar{3} > 1$ PAIRINGS:

$\chi_{(3,5)} \{3\}$	$\lambda_a (3,1) \{5\}$	$\supset (1,5) \{3 \equiv -2\}$	} SO GET CONDENSATE (1,5) {-2} AND ITS CONJUGATE (1,5) {2}
$\chi_{(3,10)} \{1\}$	$\lambda_a (\bar{3}, \bar{5}) \{-3\}$	$\supset (1,5) \{-2\}$	
$\lambda_1 (3,1) \{5\}$	$\lambda_2 (\bar{3}, \bar{5}) \{-3\}$	$\supset (1, \bar{5}) \{2\}$	

GLOBAL $U(2)$ CHIRAL SYMMETRY OF λ_a BROKEN

BY CONDENSATE \Rightarrow GET 4 GOLDSTONES WITH REPRESENTATION

(1,5) {-2} THIS IS WHAT IS NEEDED IN
FLIPPED $SU(5)$ TO BREAK SM ELECTROWEAK

ELIMINATION OF HIGGS TRIPLET (ANTONIADES ET AL)

$10 \{1\} \times 10 \{1\} \times 5 \{-2\} \supset (2,3)_{\frac{1}{6}} \times (2,3)_{\frac{1}{6}} \times (1,3)_{-\frac{1}{3}} \supset (1,1)_0$

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ABSENCE OF SCALAR-FERMION YUKAWA COUPLINGS

CHIRALITY \Rightarrow YUKAWA COUPLINGS OF SPIN $\frac{1}{2}$ FERMIONS
HAVE GENERAL FORM

$$\bar{\Psi}_{L1}^T i \gamma^0 \Psi_{L2} \Phi$$

$\Psi_{L1,2}$ ANY OF χ, λ

Φ EITHER ϕ OR ϕ^*

NO WAY TO CONTRACT $SU(8)$
INDICES TO FORM A SINGLET:

$\chi \chi$ 6 UPPER INDICES

$\lambda \lambda$ 4 LOWER INDICES

$\chi \lambda$ 3 UPPER, TWO LOWER INDICES

ϕ THREE UPPER INDICES

ϕ^* THREE LOWER INDICES

NO YUKAWA COUPLINGS INVOLVING SPIN $\frac{1}{2}$ FERMIONS
MUST BE GENERATED BY RADIATIVE CORRECTIONS

SO AFTER $SU(8)$ BREAKING, SPIN $\frac{1}{2}$ FERMIONS
REMAIN MASSLESS

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SUPERSYMMETRIES IN LIMIT OF ZERO GAUGE COUPLING g

TWO CONSERVED REP. 8 SUPERCURRENTS

$$J_a^{\mu\alpha} = \gamma^\nu (\partial_\nu \phi^{[\alpha\beta\gamma]}) \gamma^\mu \chi_a^{[\beta\gamma]} \quad a=1,2$$

ONE CONSERVED SINGLET SUPERCURRENT

$$J^\mu = \gamma^\nu (\partial_\nu \phi_{[\alpha\beta\gamma]}^*) \gamma^\mu \chi^{[\alpha\beta\gamma]}$$

$$\partial_\mu J_a^{\mu\alpha} = 0$$

$$\partial_\mu J^\mu = 0$$

NOT INVARIANCES OF SCALAR SECTOR SELF-COUPPLINGS

$$S_{\text{SELF}}(\phi) = \phi_{[\rho\kappa\tau]}^* \phi_{[\alpha\beta\gamma]}^* \left(g_1 \phi^{[\rho\kappa\tau]} \phi^{[\alpha\beta\gamma]} + g_2 \phi^{[\alpha\kappa\tau]} \phi^{[\rho\beta\gamma]} \right)$$

 $\Rightarrow g_{1,2}$ ORDER g^2 OR HIGHER

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DISCUSSION

THREE NOVEL INGREDIENTS

- BOSON-FERMION BALANCE WITHOUT FULL SUSY
- CANCELING ANOMALIES BETWEEN SPIN $\frac{1}{2}$ AND SPIN $\frac{3}{2}$
- BREAKING GAUGE SYMMETRY WITH SCALAR REPRESENTATION
WITH NONZERO $U(1)$ GENERATOR: MODULAR GROUND STATE

PROMISING FEATURES

- THREE FAMILIES
- ROUTE TO FLIPPED $SU(5)$
- $SU(8) \rightarrow SU(3)_{\text{FAMILY}} \times SU(5) \times U(1) \rightarrow SM$ USING SCALAR 56
 $SM \rightarrow ELECTROMAGNETIC U(1)$ USING HIGGS GENERATED
BY DYNAMICAL SYMMETRY
BREAKING FROM MOST
ATTRACTIVE CHANNELS
- VANISHING BARE YUKAWA COUPLINGS
ZERO GAUGE COUPLING SUMS

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EXPERIMENTAL SIGNATURES

- STERILE NEUTRINOS (GENERIC)
- SINGLET H1665 + SU(2) TRIPLET H1665 SAME QUANTUM NUMBERS

OPEN ISSUES

- CHOICE OF MOD(5) REPRESENTATIVE - DISCRETE ANALOG OF VACUUM ALIGNMENT CONDITION?

- DYNAMICAL ISSUES - RUNNING COUPLINGS
 EXTRA FERMION STATES BEYOND USUAL FAMILIES | PROTON DECAY
 MASSES AND MIXINGS OF SM
 CP VIOLATION
 FLAVOR CHANGING NEUTRAL CURRENTS
 MONOPOLES

FURTHER UNIFICATION

"SU(8) GRAVITY MULTIPLY" 128 BOSON AND FERMION HELICITIES

"SU(8) MATTER MULTIPLY" 112 " " "

MATCH NUMBERS OF HALF-INTEGER, INTEGER ROOTS OF E(8)

IS THERE A CONNECTION TO THE E(8) ROOT LATTICE?