Introduction to the SM (2)

Yuval Grossman

Cornell



The SM (2)

Contact details

- Email: yg73@cornell.edu
- Office while at CERN: 4-1.016



Happy 2nd birthday





The SM (2)

Yesterday...

- What is mechanics: $x(t) \Rightarrow L \Rightarrow$ symmetries
 - Q: the implication of $x_1 \rightarrow -3x_2$ and $x_2 \rightarrow -x_1/3$?
 - A: $q_1 = x_1 3x_2$, $q_2 = x_2 + 3x_1$ and we have $V(q_1)$
- \checkmark What is field theory: $\phi(x,t)$ are the coordinates, not x
- Harmonic oscillator
 - Leading term in the Taylor expansion
 - Normal modes and higher order couplings
 - We also expand around the minimum for fields

Today: QM, QFT, PT and Feynman diagrams

Y. Grossman

The SM (2)

Quantum mechanics



The SM (2)

What is QM?

- Many ways to formulate QM
- For example, we promote $x \to \hat{x}$
- We solve QM if we know the wave function $\psi(x,t)$
- Then we know things like $\langle x \rangle(t)$
- The wave function is mathematically a field

The quantum SHO

$$H = \frac{p^2}{2m} + \frac{m\omega^2 x^2}{2}$$
 $E_n = (n+1/2)\hbar\omega$

We also like to use

$$H = (a^{\dagger}a + 1/2)\hbar\omega \qquad a^{(\dagger)} \sim x \pm ip$$

• We call a^{\dagger} and a creation and annihilation operators

$$|a|n\rangle \propto |n-1\rangle \qquad a^{\dagger}|n\rangle \propto |n+1\rangle$$

So far this is abstract. What can we do with it?

Fields

- With many DOFs, $a \rightarrow a_i \rightarrow a(k)$
- And the states

$$|n\rangle \rightarrow |n_i\rangle \rightarrow |n(k)\rangle$$

And the energy

$$(n+1/2)\hbar\omega \to \sum (n_i+1/2)\hbar\omega_i \to \int [n(k)+1/2]\hbar\omega dk$$

- Just like in mechanics, we expand around the minimum of the fields, and to leading order we have SHOs
- In QFT fields are operators while x and t are not

SHO and photons

I have two questions:

- What is the energy that it takes to excite a SHO by one level?
- What is the energy of the photon?



The SM (2)

SHO and photons

I have two questions:

- What is the energy that it takes to excite a SHO by one level?
- What is the energy of the photon?

Same answer

$\hbar\omega$

Why the answer to both question is the same? Can we learn anything from it?

The SM (2)

What is a particle?

Excitations of SHOs are particles





The SM (2)

What about fermions?

- We see how photons are related to SHO
- We can construct a fermion SHO

$$[a, a^{\dagger}] = 1 \quad \rightarrow \quad \{b, b^{\dagger}\} = 1$$

- No classical analogue since $b^2 = 0$
- We can then think of fermionic fields. They can generate only one particle in a given state

What about mass and couplings?

A SHO give a "free" Lagrangian

$$\mathcal{L} = \frac{1}{2} \left(\partial_{\mu} \phi \right)^2$$

We can add "potential" terms (without derivatives)

$$\mathcal{L} = \frac{1}{2} \left(\partial_{\mu} \phi \right)^2 + m^2 \phi^2$$

- \checkmark Here *m* is the mass of the particle. Still free particle
- We can add terms. How do we choose what to add?

What about mass and couplings?

A SHO give a "free" Lagrangian

$$\mathcal{L} = \frac{1}{2} \left(\partial_{\mu} \phi \right)^2$$

We can add "potential" terms (without derivatives)

$$\mathcal{L} = \frac{1}{2} \left(\partial_{\mu} \phi \right)^2 + m^2 \phi^2$$

- \checkmark Here m is the mass of the particle. Still free particle
- We can add terms. How do we choose what to add?
 - Must be invariant under the symmetries
 - We keep some leading terms (usually, up to ϕ^4)

A short summary

- Particles are excitations of fields
- The fundeumental Lagrangian is giving in term of fields
- Our aim is to find it



Perturbation theory



The SM (2)

Perturbation theory

$H = H_0 + H_1 \qquad H_1 \ll H_0$

- In many cases pertubation theory is a mathematical tool
- There are cases, however, that PT is a better way to describe the physics
- Many times we prefer to work with EV of H (why?)
- Yet, at times it is better to work with EV of H_0 (why?)

PT for 2 SHOs

$$V(x,y) = \frac{kx^2}{2} + \frac{4ky^2}{2} + \alpha x^2 y$$

- Classically α moves energy between the two modes
- How it goes in QM?
- Recall the Fermi golden rule

$$P \propto |\mathcal{A}|^2 \times \text{P.S.}$$

If the initial state is $|0,1\rangle$ what transitions are alowed?

PT for 2 SHOs again

$$V(x,y) = \frac{kx^2}{2} + \frac{4ky^2}{2} + \alpha x^2 y$$

• since it is x^2y all we have is

$$\mathcal{A} \sim \langle 2, 0 | \alpha x^2 y | 0, 1 \rangle$$

- a_y annihilates the y "particle" and $(a_x^{\dagger})^2$ creates two x "particles"
- We have

$$P \propto \alpha^2$$

It is a "decay" of "particle" y into two x particles

Even More PT

$$V' = \alpha x^2 z + \beta x y z \qquad \omega_z = 10, \ \omega_y = 3, \ \omega_x = 1$$

• Calculate $y \to 3x$ using 2nd order PT

$$\mathcal{A} \sim \langle 3, 0, 0 | \mathcal{O} | 0, 1, 0 \rangle \qquad \mathcal{O} \sim \sum \frac{\langle 3, 0, 0 | V' | n \rangle \langle n | V' | 0, 1, 0 \rangle}{E_n - E_{0,1,0}}$$

What intermediate states contribute? only $|1,0,1\rangle$

$$P \propto \frac{\alpha^2 \beta^2}{64}$$

• What is the meaning of the $1/\Delta E$?

Y. Grossman

The SM (2)

Feynman diagrams



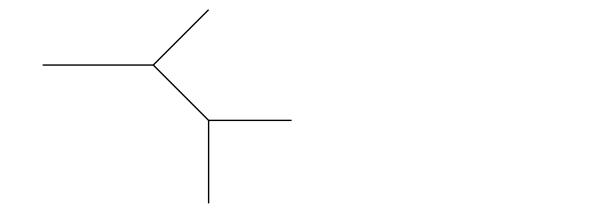
The SM (2)

Feynman diagrams

- A graphical way to do perturbation theory with fields
- Unlike SHOs before, a particle can have any energy as long as $E \ge m$
- Operators with 3 or more fields generate transitions between states. They give decay and scattering
- Decay rates and scattering cross sections are calculated using the Golden Rule
- Each coupling contributes to the amplitude linearly
- Each internal line correspond the the off-shellness of 2nd order perturbation theory and give $1/(\Delta E)^2$



 $\mathcal{L} = \lambda_1 X Y Z + \lambda_2 X^2 Z$ and we look for $Y \to 3X$



Homework: with $\mathcal{L} = \lambda X^2 Z$ estimate $\sigma(2X \to 2X)$

Y. Grossman

The SM (2)