# Introduction to the SM (2)

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## Happy 2nd birthday





### Yesterday...

- What is mechanics:  $x(t) \Rightarrow L \Rightarrow$  symmetries
	- Q: the implication of  $x_1 \rightarrow -3x_2$  and  $x_2 \rightarrow -x_1/3$ ?
	- A:  $q_1 = x_1 3x_2$ ,  $q_2 = x_2 + 3x_1$  and we have  $V(q_1)$
- What is field theory:  $\phi(x,t)$  are the coordinates, not  $x$
- Harmonic oscillator
	- **Leading term in the Taylor expansion**
	- Normal modes and higher order couplings
	- We also expand around the minimum for fields

Today: QM, QFT, PT and Feynman diagrams

## Quantum mechanics



## What is QM?

- Many ways to formulate QM $\bullet$
- For example, we promote  $x\rightarrow \hat{x}$
- We solve QM if we know the wave function  $\psi(x,t)$
- Then we know things like  $\langle x \rangle (t)$
- The wave function is mathematically <sup>a</sup> field

## The quantum SHO

$$
H = \frac{p^2}{2m} + \frac{m\omega^2 x^2}{2} \qquad E_n = (n+1/2)\hbar\omega
$$

• We also like to use

$$
H = (a^{\dagger}a + 1/2)\hbar\omega \qquad a^{(\dagger)} \sim x \pm ip
$$

We call  $a^\dagger$  and  $a$  creation and annihilation operators

$$
a|n\rangle \propto |n-1\rangle \qquad a^{\dagger}|n\rangle \propto |n+1\rangle
$$

So far this is abstract. What can we do with it?

#### Fields

- With many DOFs,  $a \rightarrow a_i \rightarrow a(k)$
- **And the states**

$$
|n\rangle \to |n_i\rangle \to |n(k)\rangle
$$

• And the energy

$$
(n+1/2)\hbar\omega \to \sum (n_i+1/2)\hbar\omega_i \to \int [n(k)+1/2]\hbar\omega dk
$$

- **O** Just like in mechanics, we expand around the minimum of the fields, and to leading order we have SHOs
- In QFT fields are operators while *<sup>x</sup>* and *<sup>t</sup>* are not

# SHO and photons

I have two questions:

- What is the energy that it takes to excite a SHO by one level?
- What is the energy of the photon?



# SHO and photons

I have two questions:

- What is the energy that it takes to excite a SHO by one level?
- What is the energy of the photon?

Same answer

#### $\hbar\,\omega$

Why the answer to both question is the same? Can welearn anything from it?

### What is a particle?

## Excitations of SHOs are particles





#### What about fermions?

- We see how photons are related to SHO
- We can construct <sup>a</sup> fermion SHO

$$
[a, a^{\dagger}] = 1 \rightarrow \{b, b^{\dagger}\} = 1
$$

- No classical analogue since*b*2 $^{2} = 0$
- We can then think of fermionic fields. They cangenerate only one particle in <sup>a</sup> given state

## What about mass and couplings?

A SHO give <sup>a</sup> "free" Lagrangian

$$
\mathcal{L}=\frac{1}{2}\left(\partial_{\mu}\phi\right)^2
$$

We can add "potential" terms (without derivatives)

$$
\mathcal{L}=\frac{1}{2}\left(\partial_{\mu}\phi\right)^{2}+m^{2}\phi^{2}
$$

- Here*m* is the mass of the particle. Still free particle
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- We can add terms. How do we choose what to add?
	- Must be invariant under the symmetries
	- We keep some leading terms (usually, up to  $\phi^4$  $^{4}$

### A short summary

- **•** Particles are excitations of fields
- The fundeumental Lagrangian is giving in term of fields $\bullet$
- Our aim is to find it



# Perturbation theory



### Perturbation theory

#### $H=H_0+H_1$  *H*<sub>1</sub>≪*H*<sub>0</sub>

- In many cases pertubation theory is <sup>a</sup> mathematical tool
- There are cases, however, that PT is a better way to describe the physics
- Many times we prefer to work with EV of *H* (why?)
- Yet, at times it is better to work with EV of  $H_0$  (why?)

#### PT for 2 SHOs

$$
V(x,y) = \frac{kx^2}{2} + \frac{4ky^2}{2} + \alpha x^2 y
$$

- Classically  $\alpha$  moves energy between the two modes
- How it goes in QM?  $\bullet$
- Recall the Fermi golden rule $\bullet$

$$
P \propto |\mathcal{A}|^2 \times \text{P.S.}
$$

If the initial state is  $|0,1\rangle$  what transitions are alowed?

# PT for 2 SHOs again

$$
V(x, y) = \frac{kx^2}{2} + \frac{4ky^2}{2} + \alpha x^2 y
$$

since it is  $x$ 2 $^{2}y$  all we have is

$$
\mathcal{A} \sim \langle 2, 0 | \alpha x^2 y | 0, 1 \rangle
$$

- $a_{y}$  annihilates the  $y$  "particle" and  $(a_{x}^{\dagger})^{2}$  creates two  $x$ "particles"
- We have

$$
P \propto \alpha^2
$$

It is a "decay" of "particle"  $y$  into two  $x$  particles  $\bullet$ 

#### Even More PT

$$
V' = \alpha x^2 z + \beta xyz \qquad \omega_z = 10, \ \omega_y = 3, \ \omega_x = 1
$$

Calculate  $y \rightarrow 3x$  using 2nd order PT

$$
\mathcal{A} \sim \langle 3, 0, 0 | \mathcal{O} | 0, 1, 0 \rangle \qquad \mathcal{O} \sim \sum \frac{\langle 3, 0, 0 | V' | n \rangle \langle n | V' | 0, 1, 0 \rangle}{E_n - E_{0,1,0}}
$$

What intermediate states contribute? only  $|1,0,1\rangle$ 

$$
P \propto \frac{\alpha^2 \beta^2}{64}
$$

What is the meaning of the <sup>1</sup>*/*∆*E*?

## Feynman diagrams



## Feynman diagrams

- A graphical way to do perturbation theory with fields
- Unlike SHOs before, <sup>a</sup> particle can have any energy aslong as  $E\geq m$
- Operators with 3 or more fields generate transitionsbetween states. They give decay and scattering
- Decay rates and scattering cross sections arecalculated using the Golden Rule
- Amplitude are calculated from  $\mathcal L$
- Each coupling contributes to the amplitude linearly
- Each internal line corrspnod the the off-shellness of 2nd order perturbation theory and give1*/*(∆*E*)2



$$
\mathcal{L} = \lambda \phi^3 \qquad \qquad \bigwedge \lambda
$$

 ${\cal L}$  $\mathcal{L} = \lambda_1 XYZ + \lambda_2 X^2 Z$  and we look for  $Y \to 3X$ 



Homework: with  $\mathcal{L} = \lambda X^2 Z$  estimate  $\sigma(2X \to 2X)$