

Search for Direct Time-Reversal Violation

An Experimentalist-Theorist Collaboration

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Heavy Baryon LHCb Workshop

I– Introduction

II– T-odd Observables and Asymmetries

III– Kinematics and Dynamics

IV– Observable Measurements

I- Introduction

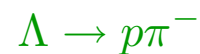
- **CPT** (theorem) \equiv Exact Symmetry, but CP violated \Rightarrow Indirect Violation of T
- Why Direct Search for Time-Reversal Violation (TRV) ??

TR symmetry can be tested **independently** of CP
 \Rightarrow CP conjugated channel not required

Example : **Electric Dipole Moment** of an elementary particle : electron, neutron ...

$$\vec{D}_e = d_e \vec{s} \quad d_e \neq 0 \implies \text{TRV}$$

- Historically, TR tested in Hyperon weak decays just after **Parity** Violation ([R. Gatto](#), 1958) like :



Also in Weak β -Decay of the Neutron ([Jackson et al, 1958](#))

\implies Essential to produce a **Polarized** particle (Resonance) in order to test TRV.

- Our Method : *Phenomenological and Intuitive approach.*

II- T-Odd Observables and Polarization

- Action of TR :

$$\vec{p} \rightarrow -\vec{p} \text{ (1)}, \quad \vec{s} \rightarrow -\vec{s} \text{ (2)}, \quad \text{and } |i\rangle \leftrightarrow |f\rangle \text{ (3)}$$

But :

- Last condition not (easy) realized in Nature \implies TR limited to transformations(1) and (2)

\Downarrow

Naive Time-Reversal operation.

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T-odd Observable \equiv Physical quantity changing sign by TR transformation

$$O \implies -O$$

- If TR is a **Good Symmetry** $\implies O = 0$, or $\langle O \rangle = 0$.

★ **But :**

Every T-odd Observable $\neq 0$ is not a sign of TRV, because of **Final State Interaction (FSI)**.

$$A + B \rightarrow C + D \implies TR \implies C' + D' \rightarrow A' + B' \quad X' = X(-\vec{p}, -\vec{s})$$

Normally, Strong Interactions between C and D \neq between $A^{(')}$ and $B^{(')}$

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FSI are **Unavoidable** Systematic Effects in Direct TRV searches

- Example : Transverse Polarization of Hyperons produced in p-p or p-nucleus collisions

$$\vec{t} = \frac{\vec{p}_i \times \vec{p}_h}{|\vec{p}_i \times \vec{p}_h|}, \quad P_t = \vec{t} \cdot \vec{s} \rightarrow \vec{t} \cdot (-\vec{s}) = -P_t$$

Usually, $P_t \neq 0$ and no sign of TRV.

Wolfenstein “Theorem” (1990-2000) :

If the FSI are negligible in a given process or if they can be subtracted from the scattering amplitude, any T-odd Observable is a *Sign of TRV*.

III- Kinematics and Dynamics

- General expression of a T-odd Observable :

$$C = \vec{v}_1 \cdot (\vec{v}_2 \times \vec{v}_3), \quad \vec{v}_i = \vec{s}_i \text{ or } \vec{p}_i$$

- Searching for T-odd Observables in Heavy Baryon Decays :

$$\mathcal{B}_b \implies R_1(\rightarrow a_1 + b_1) + R_2(\rightarrow a_2 + b_2), \quad R_1 = (1/2)^+ \quad R_2 = 0^-, 1^-$$

“Best Example” , $\Lambda_b \implies \Lambda(\rightarrow p\pi^-) \quad J/\psi(\rightarrow \mu^+\mu^-)$

- Kinematics based on the **Helicity Formalism** of Jacob-Wick-Jackson.

1-Frames

Transversity Frame constructed from the standard laboratory frame :

$$\mathbf{e}_x = \hat{\mathbf{p}}_{proton}, \quad \mathbf{e}_z = \frac{\mathbf{e}_x \times \mathbf{p}_b}{|\mathbf{e}_x \times \mathbf{p}_b|}, \quad \mathbf{e}_y = \mathbf{e}_z \times \mathbf{e}_x.$$

Then boosted to the \mathcal{B}_b rest-frame along \mathbf{p}_b , the quantization axis being taken parallel to \mathbf{e}_z . Then R_i *helicity frame* is defined as :

$$\mathbf{e}_L = \hat{\mathbf{p}}_R, \quad \mathbf{e}_T = \frac{\mathbf{e}_z \times \mathbf{e}_L}{|\mathbf{e}_z \times \mathbf{e}_L|}, \quad \mathbf{e}_N = \mathbf{e}_T \times \mathbf{e}_L,$$

\Rightarrow Analysis of the Angular Distributions of particles $a_i(b_i)$

2-Transformation under TR

\mathbf{e}_x and \mathbf{e}_y Change Sign under TR.

\mathbf{e}_z Unchanged.

In each R_i rest frame, \mathbf{e}_L and \mathbf{e}_T are **T-odd**, whereas \mathbf{e}_N is **T-even**.

3-Decay Amplitude

$$A_0(M_i) = \langle 1/2, M_i | \mathcal{S}^{(0)} | p, \theta, \phi; \lambda_1, \lambda_2 \rangle = \mathcal{A}_{(\lambda_1, \lambda_2)}(\Lambda_b \rightarrow R_1 R_2) D_{M_i M_f}^{1/2\star}(\phi, \theta, 0), \quad (1)$$

$$D_{M_i M_f}^j(\phi, \theta, 0) = d_{M_i M_f}^j(\theta) \exp(-i M_i \phi), \quad (2)$$

★ Weighting by the PDM of Initial Resonance Λ_b and Summation over all the states (intermediate and helicity final) :

$$\frac{d\Gamma}{d\Omega} = N^J \sum_{M_i, M'_i} \sum_{\lambda_1, \lambda_2} \rho_{M_i M'_i}^{\Lambda_b} |\mathcal{A}_{(\lambda_1, \lambda_2)}(\Lambda_b \rightarrow R_1 R_2)|^2 d_{M_i \lambda}^{1/2} d_{M'_i \lambda}^{1/2} \exp i(M'_i - M_i)\phi . \quad (3)$$

★ $A_{(\lambda_1, \lambda_2)} \neq A_{(-\lambda_1, -\lambda_2)}$ because of Parity violation.

★ Dynamics Model and 3 real parameters to fully determine $W(\theta, \phi) = \frac{d\Gamma}{d\Omega}$: \mathcal{P}^b , and complex ρ_{+-}

4-Asymmetry and Polarization(s)

• *Helicity Asymmetry Parameter*, α_{As} , computed according to :

$$\alpha_{As} = \frac{|\mathcal{B}_b(+)|^2 - |\mathcal{B}_b(-)|^2}{|\mathcal{B}_b(+)|^2 + |\mathcal{B}_b(-)|^2}, \quad (4)$$

with

$$|\mathcal{B}_b(\pm)|^2 = |A_{\pm 1/2, 0}^J|^2 + |A_{\mp 1/2, \mp 1}^J|^2. \quad (5)$$

• Polarization-vector of each Resonance R_i : Standard Relations (JWJ)

$$\vec{\mathcal{P}}_i = \frac{Tr(\rho_i^f \vec{S})}{W(\theta, \phi)}, \quad \rho_i^f \text{ being final PDM of } R_i$$

- Transformation of $\vec{\mathcal{P}}_i$ under TR :

$$\mathcal{P}_N = \vec{\mathcal{P}}_{(i)} \cdot \vec{e}_N \implies (-\vec{\mathcal{P}}_{(i)}) \cdot \vec{e}_N = -\mathcal{P}_N$$

\implies The Normal Component \mathcal{P}_N is **T-Odd**. So,
If $\mathcal{P}_N \neq 0 \implies$ *Sign of TRV*.

IV- Observable Measurements

- Weak Decaying Resonance $X \rightarrow (1/2)^+ + 0^-(1^-)$

$X = \Lambda_b$ or Λ , α^X its asymmetry parameter

$$\frac{dN}{d\Omega} \propto 1 + \alpha^X \mathbf{P}^X \cdot \hat{\mathbf{p}}, \quad (6)$$

\implies Polar and Azimuthal angle distributions :

$$f(\cos \theta) = \frac{1}{2}(1 + \alpha^X P_z^X \cos \theta) \quad (7)$$

and

$$g(\phi) = \frac{1}{2\pi} \left(1 + \frac{\pi}{4} \alpha^X (P_x^X \cos \phi + P_y^X \sin \phi) \right) \quad (8)$$

- Inferring P_x^X, P_y^X and P_z^X for the primary resonance, \mathcal{B}_b ,
- Secondary (Baryon) decay $R_1 \rightarrow a_1 + b_1$, P_x^X identified with the Normal Component $P_N^{R_1}$, a **T-odd observable**.
- Resonance R_2 Decay which conserves Parity, like $J/\psi \rightarrow \mu^- \mu^+$

$$\frac{dN}{d\cos\theta} \propto (1 - 3\rho_{00}) \cos^2\theta + (1 + \rho_{00}) \propto 1 + \gamma \cos^2\theta, \quad (9)$$

ρ_{00} is the probability for J/ψ to be longitudinally polarized, and

$$\gamma = \frac{1-3\rho_{00}}{1+\rho_{00}}.$$

★ Azimuthal Distribution of μ^- in J/ψ rest-frame is **Uniform**

How handling FSI (?)

- How to disentangle **FSI** from experimental results :

New Calculation Techniques in QCD, [QCDF](#)(Factorization).

- Old argument of [Bjorken](#) (1990) : *FSI are very weak* if Transition between

Heavy Quarks, like :

$$Q \implies Q' + X \dots \text{ like } \Lambda_b \rightarrow J/\psi + \Lambda$$

- Independently of *any argument or model*, we notice that FSI coming from Strong Interactions are **CP Invariant**

\implies Comparing \mathcal{B}_b and $\bar{\mathcal{B}}_b$ decays. (Valencia, London).

- Increasing the Number of TR tests as follows :

$$P_\alpha^{(1)} P_\beta^{(2)} \quad \alpha, \beta = L, T, N$$

- Time-Reversal Asymmetry parameter defined by :

$$\mathcal{A}_T = \frac{O - \bar{O}}{O + \bar{O}} \quad \text{with } \bar{O} = CP|O \rangle$$

Usually, we have **Interference** of two different amplitudes like *Spin-Orbit* term and another one.

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$$\mathcal{A}_T = 2A_1 A_2 \cos \delta_s \sin \phi_W$$

Even for small FSI, $\equiv \delta_s \rightarrow 0$, TR Asymmetry will be amplified.

Looking for New Physics (?)

- Defining new T-odd observables and Comparing them to their CP Conjugate ones.
- Eventual TRV effects independently of CPT theorem (?)

Thank you for your attention

References :

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E.DiSalvo, Z.J.Ajaltouni, *Baryon Polarization, Phases of Amplitudes and Time Reversal Odd Observables*
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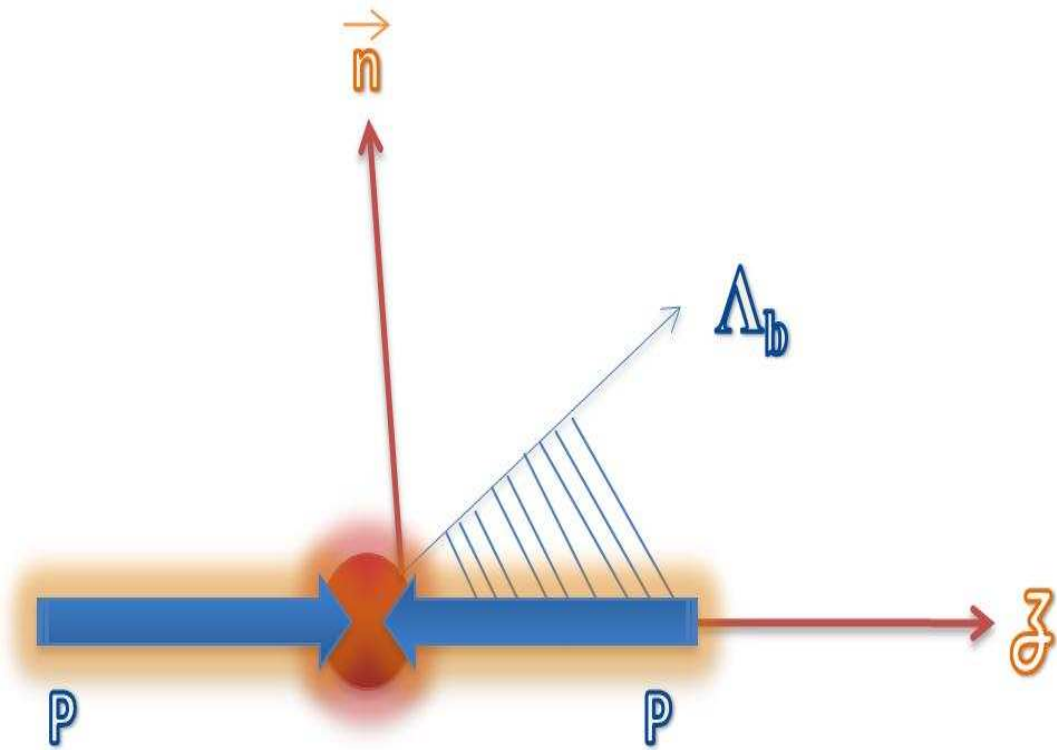


Figure 1: Λ_b in the Standard LHCb Frame

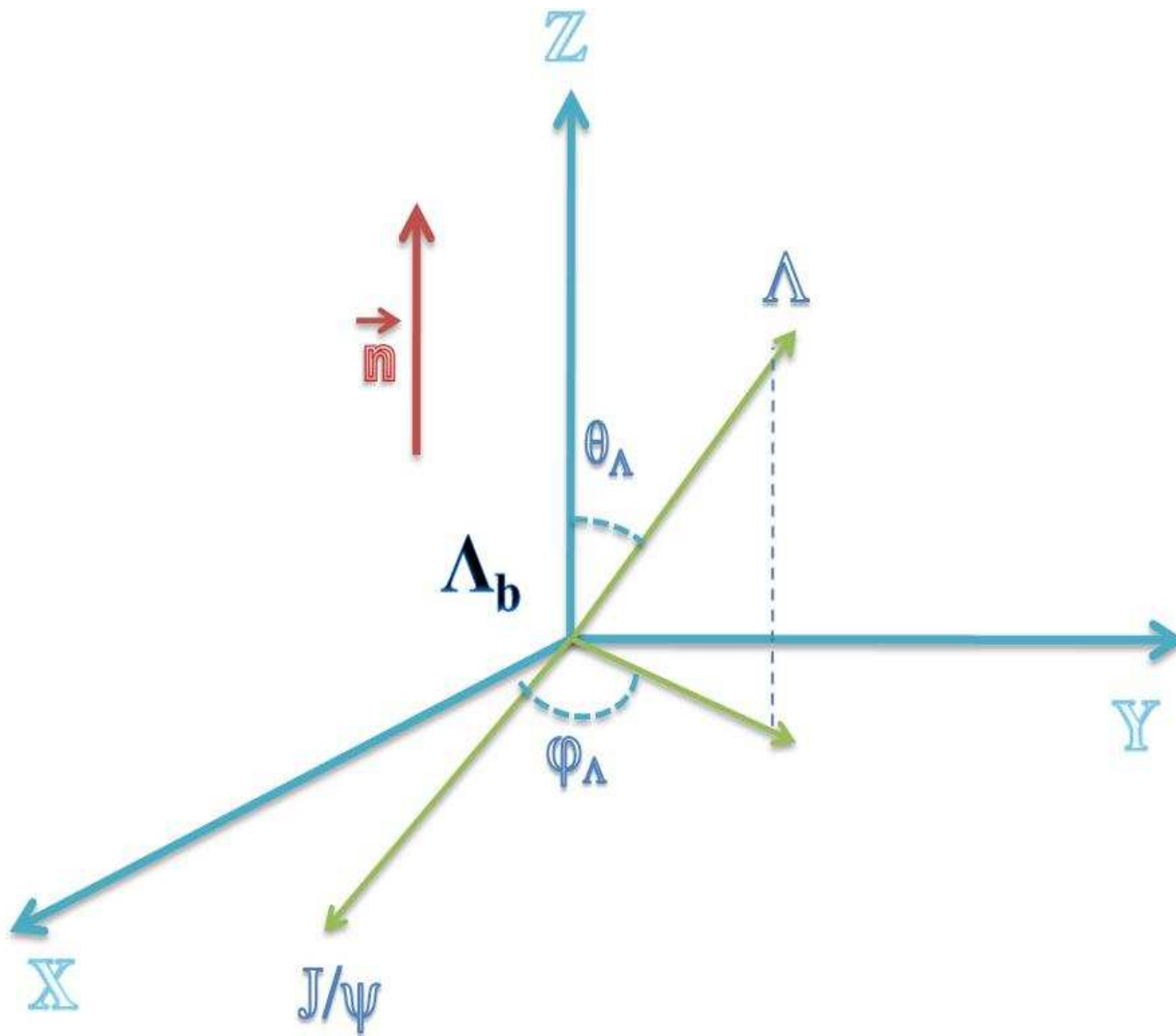
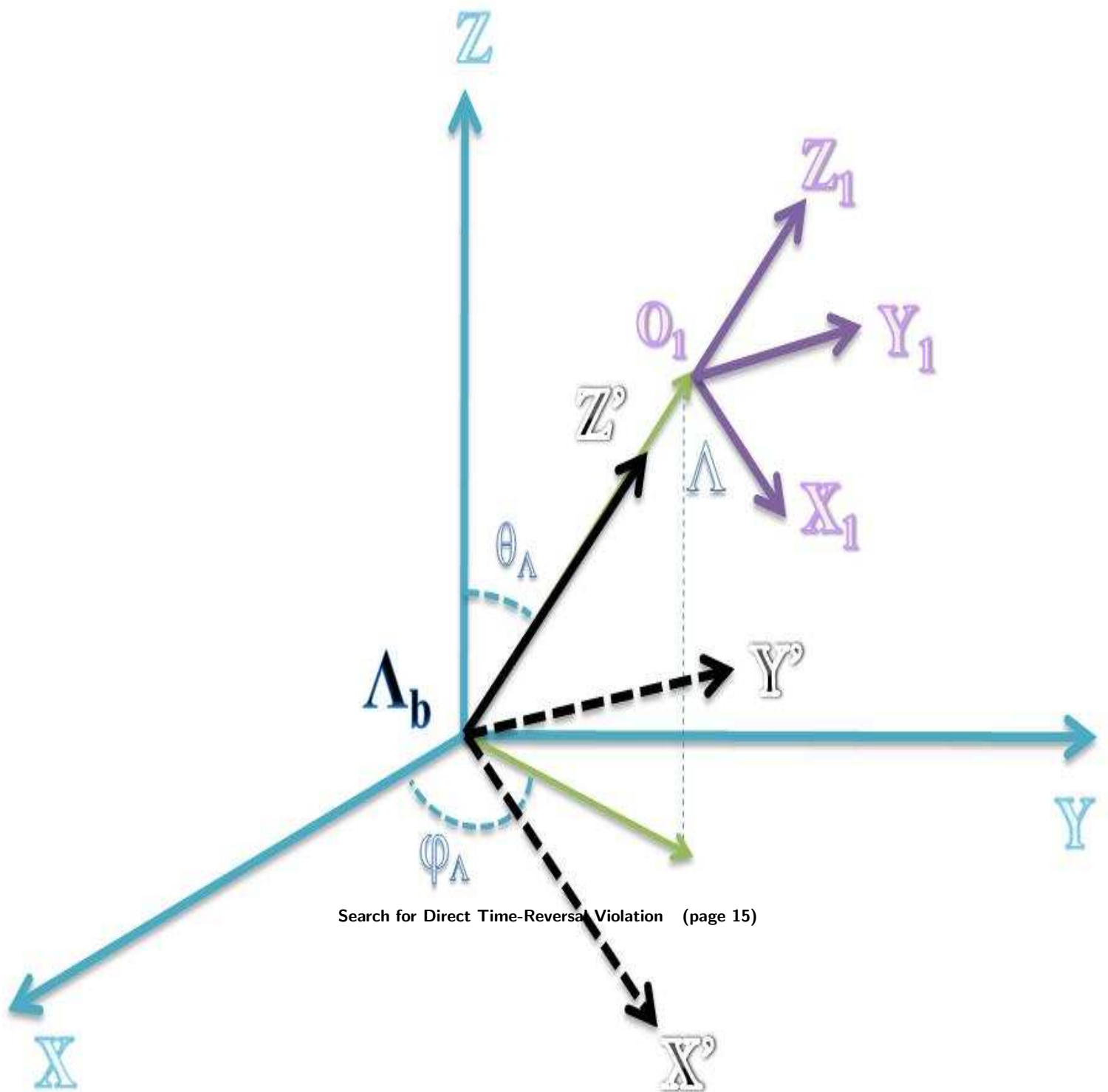
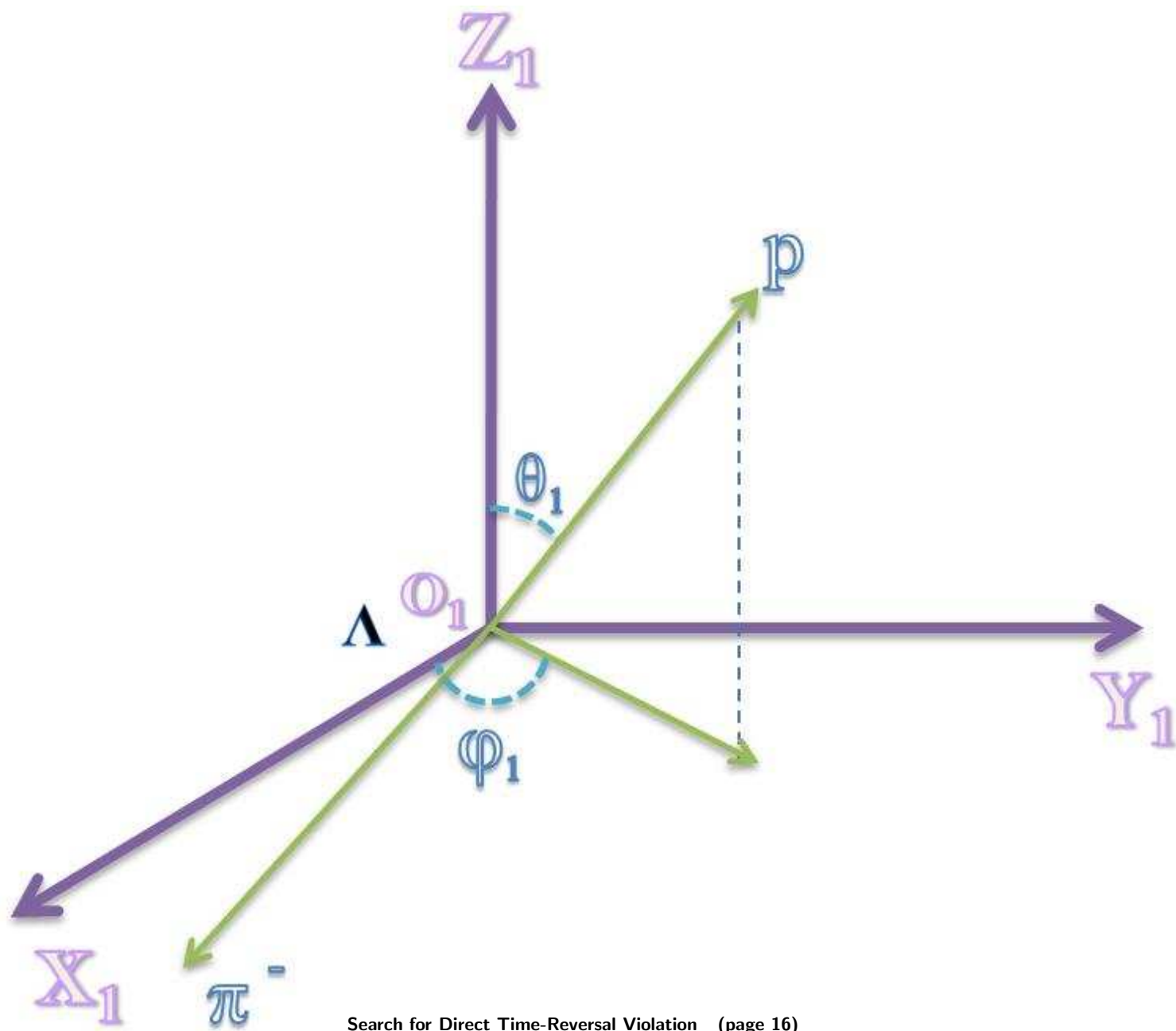


Figure 2: Λ_b Transversity Rest-Frame



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Figure 3: From Λ_b rest-frame to Λ^0 Helicity rest-frame



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Figure 4: Λ^0 Helicity Rest-Frame

Marwa Jahjah-Hussein,
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