

$K^{(*)}\pi$ revisited

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Three messages from this talk:

- i) typical predictions of factorization in the infinite mass limit for $K\pi$ amplitudes are off by $\sim -30-40\%$ (Luca was almost right!)
- ii) $K\pi$ decays are not puzzling once subleading terms are included. Measured CP asymmetries are compatible with the Standard Model
- iii) $K^*\pi$ decays are a perfect playground for QCD challenges

MC, Franco, Martinelli, Pierini, Silvestrini,
in preparation

new physics in $K\pi$ CP asymmetries?

$$\mathcal{A}_{K^\pm \pi^\mp} \equiv \frac{N(\bar{B}^0 \rightarrow K^- \pi^+) - N(B^0 \rightarrow K^+ \pi^-)}{N(\bar{B}^0 \rightarrow K^- \pi^+) + N(B^0 \rightarrow K^+ \pi^-)} = -0.094 \pm 0.018 \pm 0.008$$

Belle collaboration
Nature 452,2008

$$\mathcal{A}_{K^\pm \pi^0} = +0.07 \pm 0.03 \pm 0.01$$

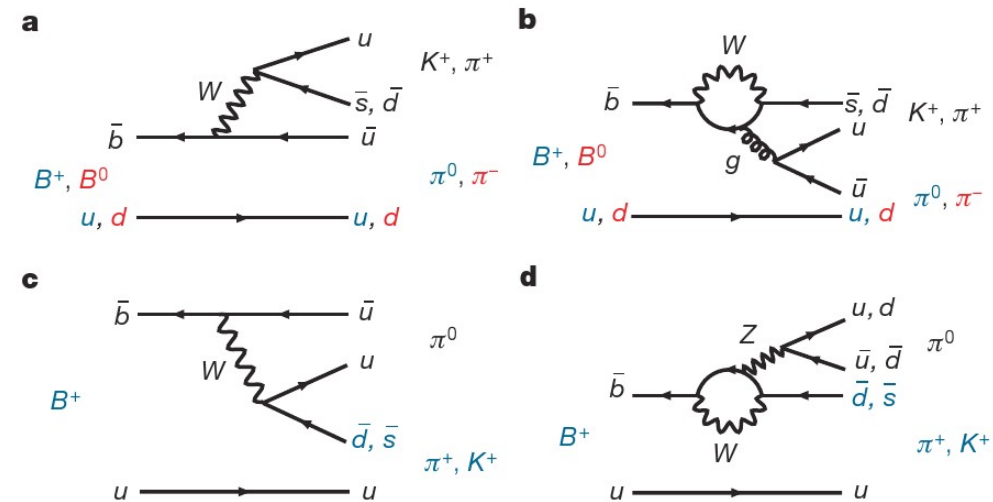
$$\Delta\mathcal{A} \equiv \mathcal{A}_{K^\pm \pi^0} - \mathcal{A}_{K^\pm \pi^\mp} = +0.164 \pm 0.037$$

4.4 σ away from 0

Is this new physics?

It could be but SM predictions depend on hadronic models

Silvestrini
arXiv:0705.1624
%



	QCDF [50]	PQCD [54, 55]	SCET [58]	GP [92]
$A_{CP}(\pi^0 K^-)$	$7.1^{+1.7+2.0+0.8+9.0}_{-1.8-2.0-0.6-9.7}$	-1^{+3}_{-5}	$-11 \pm 9 \pm 11 \pm 2$	3.4 ± 2.4
$A_{CP}(\pi^+ K^-)$	$4.5^{+1.1+2.2+0.5+8.7}_{-1.1-2.5-0.6-9.5}$	-9^{+6}_{-8}	$-6 \pm 5 \pm 6 \pm 2$	-8.9 ± 1.6

Amplitude Parametrization

general parametrization
 *one simplification only:
 isospin breaking in the
 hadronic ME neglected
 can be reintroduced
 if need be

$$\begin{aligned}
 A(B^+ \rightarrow K^0 \pi^+) &= -V_{ts} V_{tb}^* P + V_{us} V_{ub}^* A, \\
 A(B^+ \rightarrow K^+ \pi^0) &= \frac{1}{\sqrt{2}} \left(V_{ts} V_{tb}^* (P + \Delta P_1 + \Delta P_2) \cdot \right. \\
 &\quad \left. V_{us} V_{ub}^* (E_1 + E_2 + A) \right), \\
 A(B^0 \rightarrow K^+ \pi^-) &= V_{ts} V_{tb}^* (P + \Delta P_1) - V_{us} V_{ub}^* E_1 \\
 A(B^0 \rightarrow K^0 \pi^0) &= -\frac{1}{\sqrt{2}} \left(V_{ts} V_{tb}^* (P - \Delta P_2) + V_{us} V_{ub}^* E_2 \right)
 \end{aligned}$$

$$\begin{aligned}
 E_1 &= E_1^F + F \left(r(E_1) e^{i\delta(E_1)} - r(P_1^{\text{GIM}}) e^{i\delta(P_1^{\text{GIM}})} \right) \\
 &= E_1^F + F R(E_1) e^{i\Delta(E_1)}, \\
 E_2 &= E_2^F + F \left(r(E_2) e^{i\delta(E_2)} + r(P_1^{\text{GIM}}) e^{i\delta(P_1^{\text{GIM}})} \right) \\
 &= E_2^F + F R(E_2) e^{i\Delta(E_2)}, \\
 A &= A^F + F \left(r(A) e^{i\delta(A)} - r(P_1^{\text{GIM}}) e^{i\delta(P_1^{\text{GIM}})} \right), \\
 &= A^F + F R(A) e^{i\Delta(A)}, \\
 P &= P^F + F r(P) e^{i\delta(P)}, \\
 \Delta P_1 &= \Delta P_1^F + F \alpha_{\text{em}} r(\Delta P_1) e^{i\delta(\Delta P_1)}, \\
 \Delta P_2 &= \Delta P_2^F + F \alpha_{\text{em}} r(\Delta P_2) e^{i\delta(\Delta P_2)},
 \end{aligned}$$

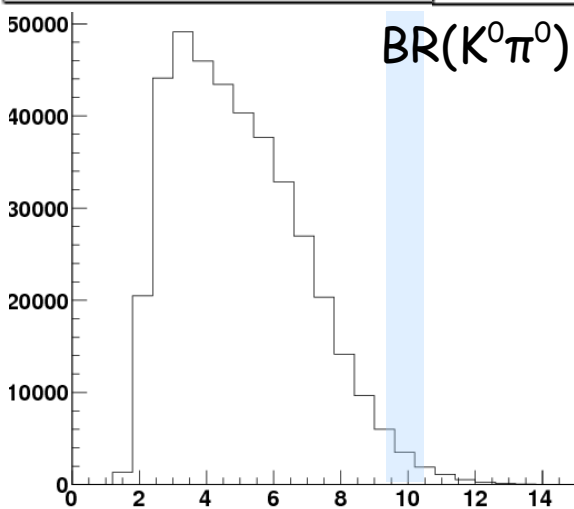
deviations from factorization:
 $R(X) \exp[i\Delta(X)]$ in units of $F = \bar{A}_{K\pi}$

related to Buras, Silvestrini,
[hep-ph/9806278](https://arxiv.org/abs/hep-ph/9806278)

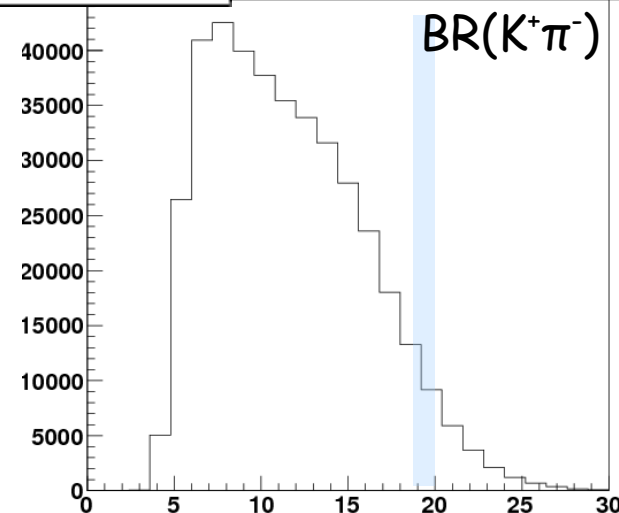
$$\begin{aligned}
 E_1 &= E_1(s, q, q; B, K, \pi) - P_1^{\text{GIM}}(s, q; B, K, \pi) \\
 E_2 &= E_2(q, q, s; B, \pi, K) + P_1^{\text{GIM}}(s, q; B, K, \pi) \\
 A &= A_1(s, q, q; B, K, \pi) - P_1^{\text{GIM}}(s, q; B, K, \pi) \\
 P &= P_1(s, d; B, K, \pi), \\
 \Delta P_1 &= P_1(s, u; B, K, \pi) - P_1(s, d; B, K, \pi), \\
 \Delta P_2 &= P_2(s, u; B, \pi, K) - P_2(s, d; B, \pi, K).
 \end{aligned}$$

Step #0: try throwing away all these ugly parameters

Decay Mode	$\text{BR}^{\text{exp}} \times 10^6$		$\mathcal{A}_{\text{CP}}^{\text{exp}} = -C$	
$K^+ \pi^-$	19.4 ± 0.6	11.6 ± 4.4	-0.097 ± 0.012	0.065 ± 0.016
$K^+ \pi^0$	12.9 ± 0.6	8.4 ± 3.2	0.050 ± 0.025	0.086 ± 0.014
$K^0 \pi^+$	23.1 ± 1.0	14.8 ± 5.7	0.009 ± 0.025	0.008 ± 0.001
$K^0 \pi^0$	9.9 ± 0.6	5.0 ± 2.0	-0.14 ± 0.11	-0.03 ± 0.01
		S	0.38 ± 0.19	0.73 ± 0.01



QCD factorization
 predictions for $m_b \rightarrow \infty$
 (including χ -enhanced terms)



Two non-contradictory statements:

- typical factorized $K\pi$ amplitudes are off by $\sim -30-40\%$
- factorized amplitudes can reproduce the $K\pi$ data

Old method, new perspective

* old idea: use data to determine the subleading terms, but

11 real unknowns

9 measurements

too many parameters! One can:

- reduce the parameter set

(like in the good old charming-penguin days)

- vary all the parameters in theoretically

sensible ranges (we take $r \in [0, 0.5]$, $\delta \in [-\pi, \pi]$)

Final goal: find "upper bounds" to the theoretical errors compatible with data and the $1/m_b$ expansion

Quick facts on charming penguins

- first appearance
Colangelo, Nardulli, Paver, Riazuddin
Z. Phys. C45 (1990) 575
- christening
MC, Franco, Martinelli, Silvestrini
hep-ph/9703353
- revisited (I)
MC, Franco, Martinelli, Pierini,
Silvestrini, hep-ph/0104126
- revisited (II)
Bauer, Pirjol, Rothstein, Stewart
hep-ph/0401188

Results of the fit to the $K\pi$ data

“global fit”:

results obtained fitting the whole data set

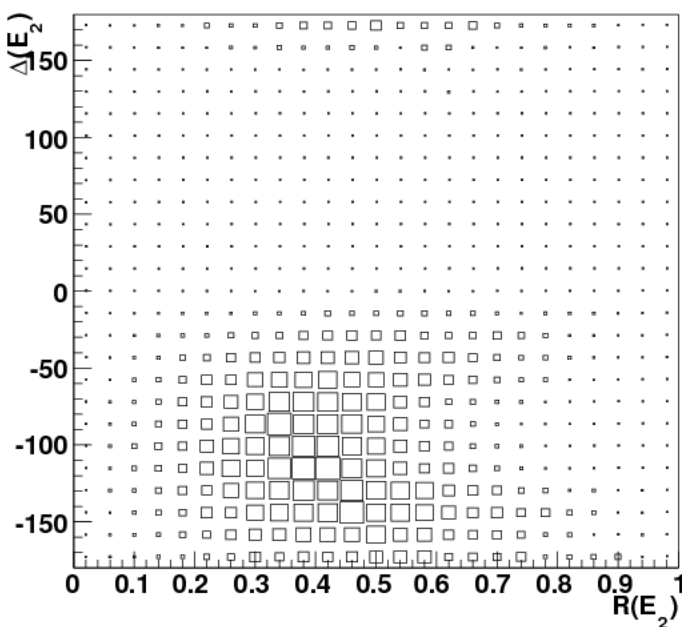
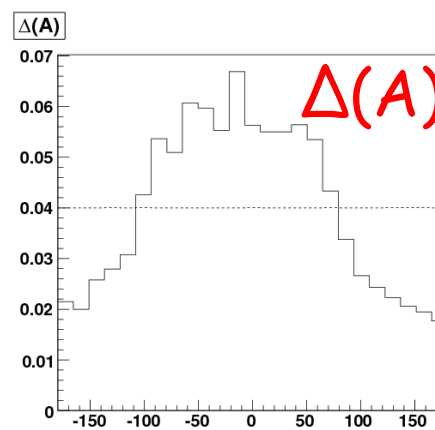
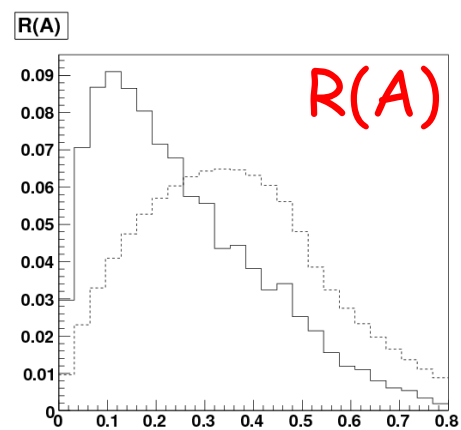
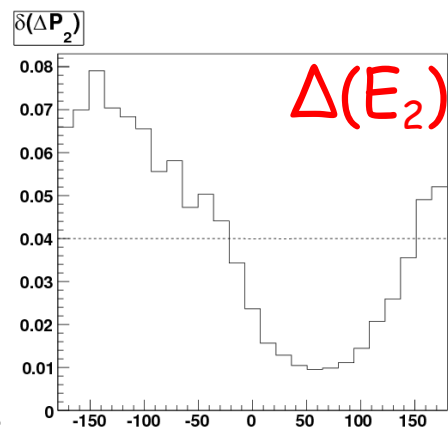
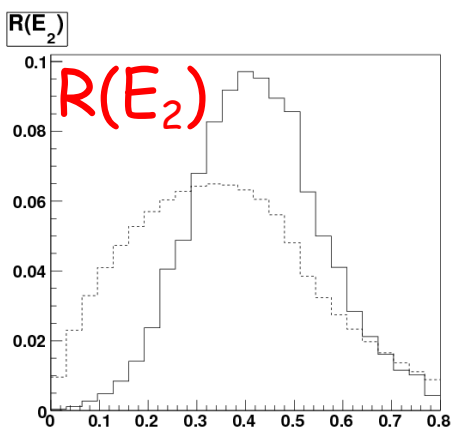
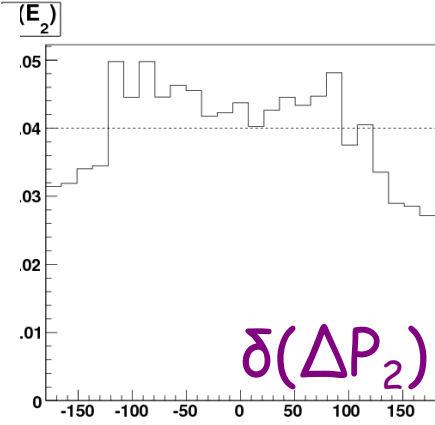
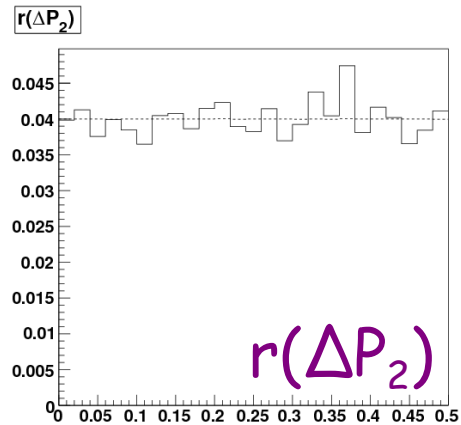
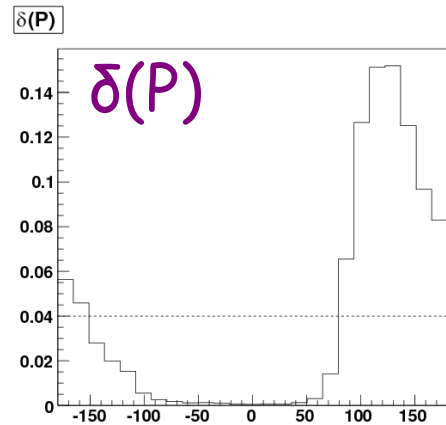
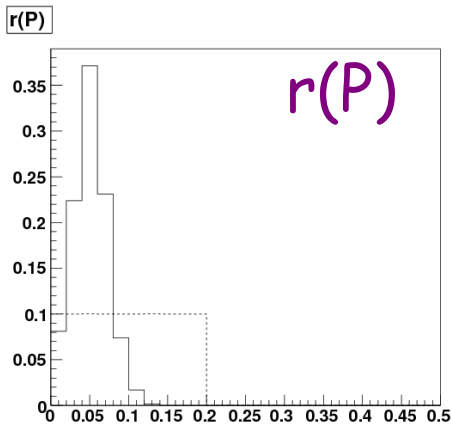
“fit predictions”:

results obtained fitting the whole data set but the “prediction”

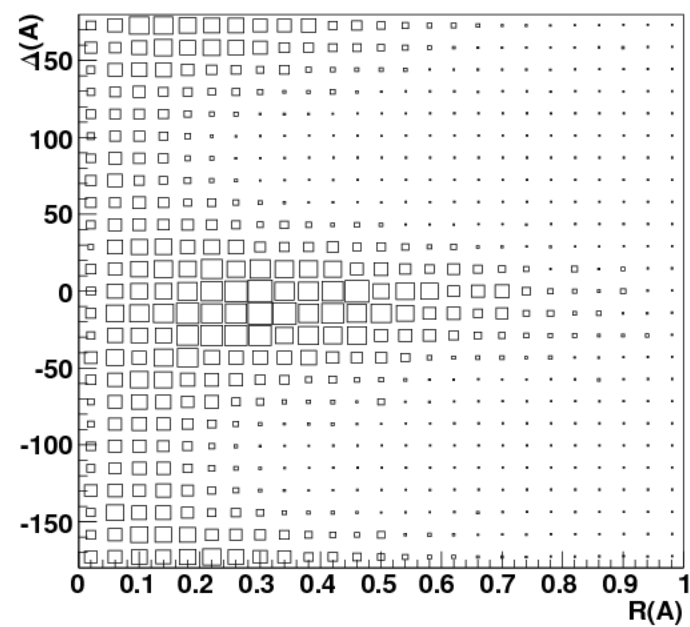
	global fit	fit prediction
$\text{BR}(K^+ \pi^-) \times 10^6$	19.6 ± 0.5	20.1 ± 1.0
$\text{BR}(K^+ \pi^0) \times 10^6$	12.7 ± 0.5	12.4 ± 0.7
$\text{BR}(K^0 \pi^+) \times 10^6$	23.7 ± 0.8	24.6 ± 1.2
$\text{BR}(K^0 \pi^0) \times 10^6$	9.2 ± 0.4	8.6 ± 0.6
$\mathcal{A}_{\text{CP}}(K^+ \pi^-)$	-0.095 ± 0.012	-0.01 ± 0.08
$\mathcal{A}_{\text{CP}}(K^+ \pi^0)$	0.043 ± 0.024	-0.02 ± 0.07
$\mathcal{A}_{\text{CP}}(K^0 \pi^+)$	0.010 ± 0.023	0.02 ± 0.06
$C(K_S \pi^0)$	0.12 ± 0.04	0.12 ± 0.04
$S(K_S \pi^0)$	0.702 ± 0.067	0.74 ± 0.06

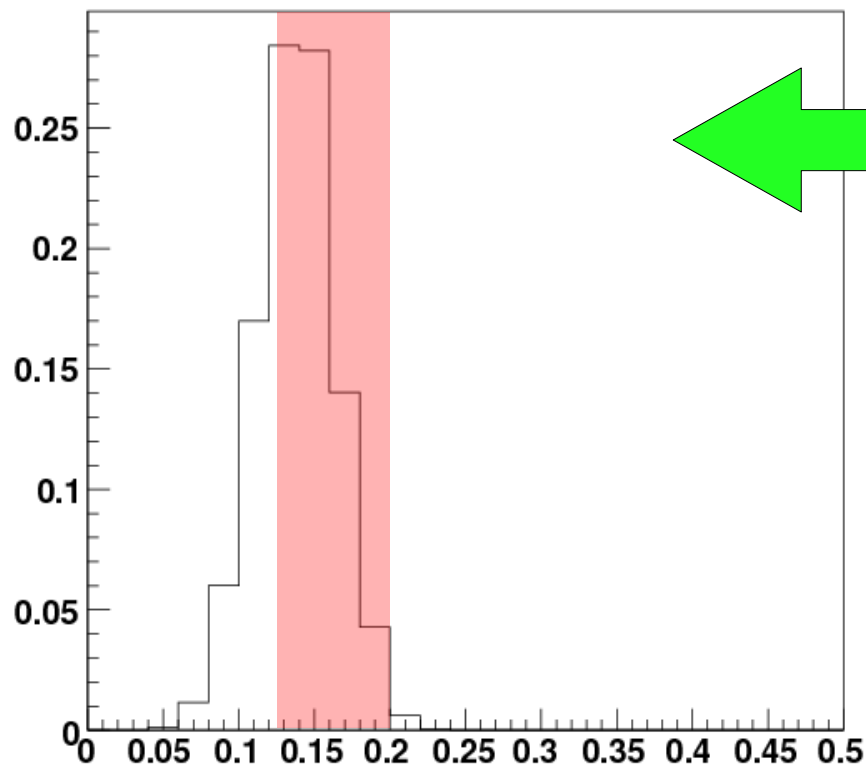
Decay Mode	$\text{BR}^{\text{exp}} \times 10^6$	$\mathcal{A}_{\text{CP}}^{\text{exp}} = -C$	S
$K^+ \pi^-$	19.4 ± 0.6	-0.097 ± 0.012	–
$K^+ \pi^0$	12.9 ± 0.6	0.050 ± 0.025	–
$K^0 \pi^+$	23.1 ± 1.0	0.009 ± 0.025	–
$K^0 \pi^0$	9.9 ± 0.6	-0.14 ± 0.11	0.38 ± 0.19

- BR's OK and fairly insensitive to the “ Λ/m_b noise”
- \mathcal{A}_{CP} can be reproduced thanks to the “ Λ/m_b noise”
- $S(K_S \pi^0)$ cannot be “satisfactorily” reproduced



* charming penguins $r(P)$
are well determined
* Λ/m_b corrections $R(A)$
 $\sim 0.1-0.2$ and $R(E_2)$
 $\sim 0.3-0.4$ are selected
* data do not fix values
for the other terms



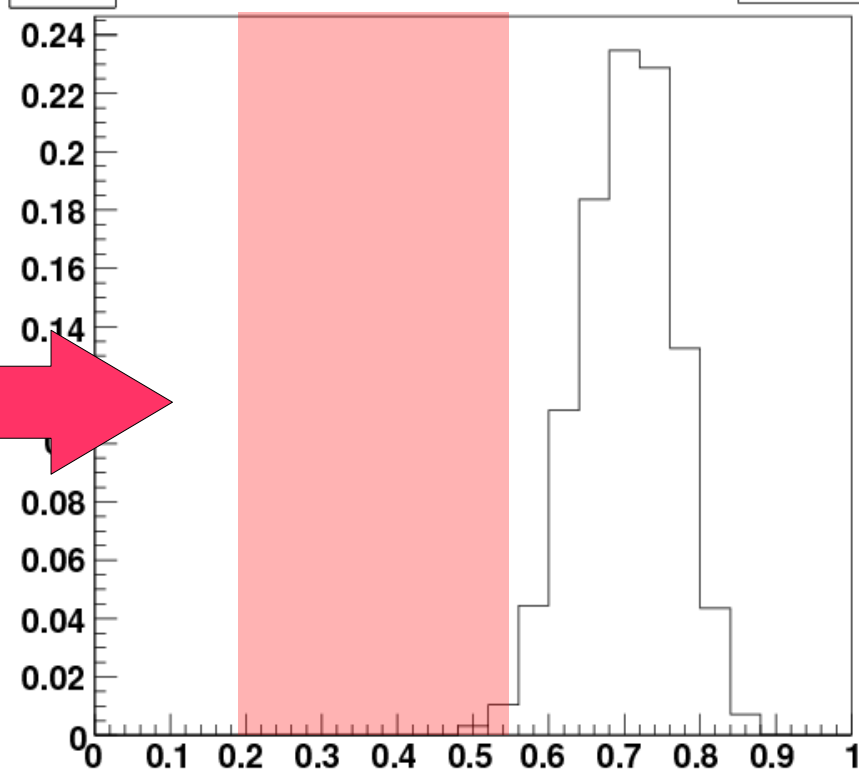
ΔA_{CP} Mean 0.1378
RMS 0.02633

preliminary

No clash with
the Belle measurement
of the CP asymmetries

However...

$$S(K_S\pi^0)_{exp} = 0.38 \pm 0.19$$

 $S_{K_S\pi^0}$ Mean 0.7023
RMS 0.06381

The $K^*\pi$ playground for QCD challenges

- * 11 real hadronic parameters as in the $K\pi$ case
- * 11 observables \Rightarrow fully determined in the SM

1. $K^+ \pi^- \pi^0$ Dalitz plot: (3)

$$|A(K^{*+} \pi^-)|, |A(K^{*0} \pi^0)|, \\ \arg A(K^{*+} \pi^-) - \arg A(K^{*0} \pi^0)$$

2. $K^- \pi^+ \pi^0$ Dalitz plot: (3)

$$|A(K^{*-} \pi^+)|, |A(\bar{K}^{*0} \pi^0)|, \\ \arg A(K^{*-} \pi^+) - \arg A(\bar{K}^{*0} \pi^0)$$

3. $K_S \pi^- \pi^+$ Dalitz plot: (1)

$$|A(K^{*+} \pi^-)|, |A(K^{*-} \pi^+)|, \\ \arg A(K^{*+} \pi^-) - \arg A(K^{*-} \pi^+)$$

4. $K_S \pi^0 \pi^0$ Dalitz plot: (0)

$$|A(K^{*0} \pi^0)|, |A(\bar{K}^{*0} \pi^0)|, \\ \arg A(K^{*0} \pi^0) - \arg A(\bar{K}^{*0} \pi^0)$$

5. $K_S \pi^+ \pi^0$ Dalitz plot: (3)

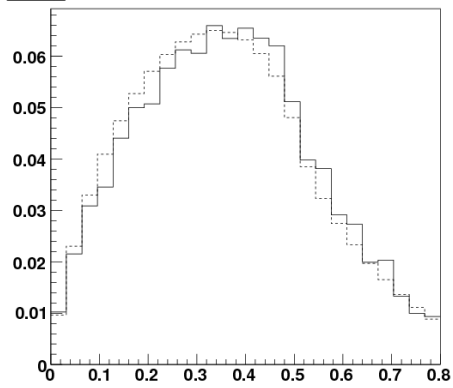
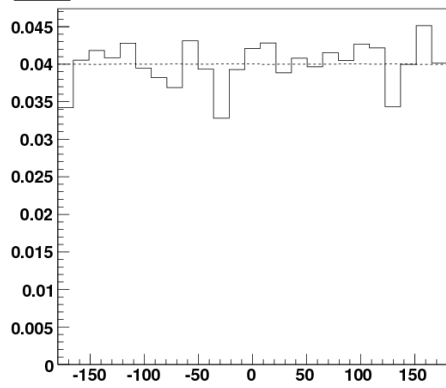
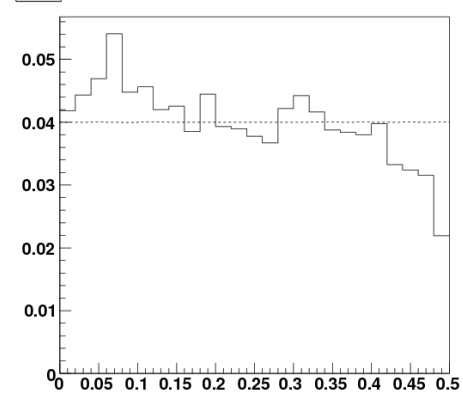
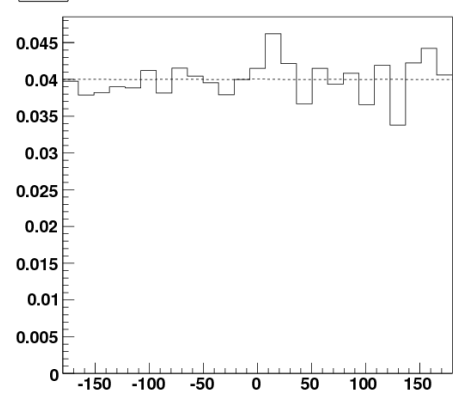
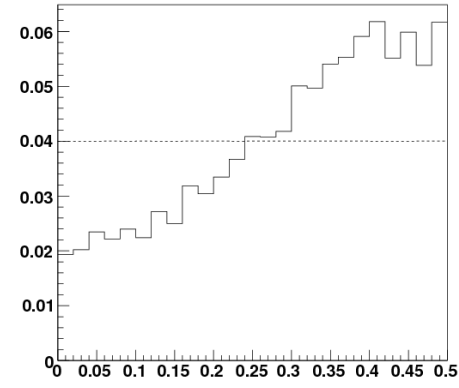
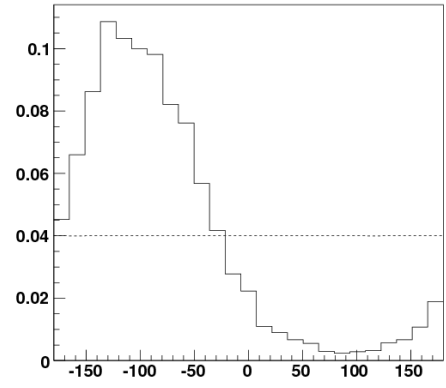
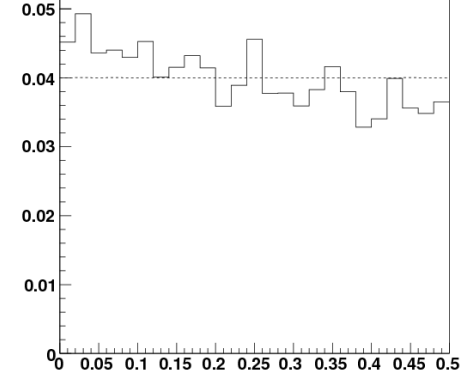
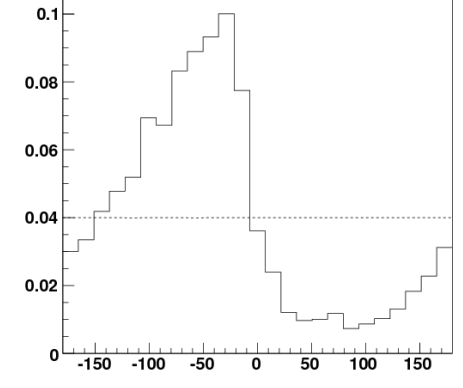
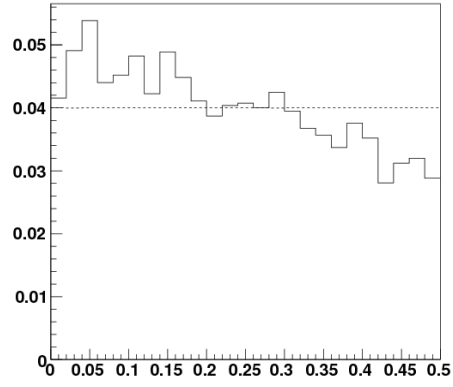
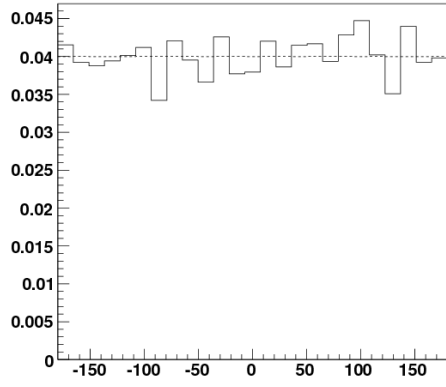
$$|A(K^{*+} \pi^0)|, |A(K^{*0} \pi^+)|, \\ \arg A(K^{*+} \pi^0) - \arg A(K^{*0} \pi^+)$$

6. $K_S \pi^- \pi^0$ Dalitz plot: (3)

$$|A(K^{*-} \pi^0)|, |A(K^{*0} \pi^-)|, \\ \arg A(K^{*-} \pi^0) - \arg A(K^{*0} \pi^-)$$

amplitudes satisfy 2 isospin quadrangular relations (-2)

Spare Slides

$R(E_1)$  $\Delta(E_1)$  $r(A)$  $\delta(A)$  $r(E_2)$  $\delta(E_2)$  $r(P_1^{GIM})$  $\delta(P_1^{GIM})$  $r(\Delta P_1)$  $\delta(\Delta P_1)$ 

$$E_1^F = A_{\pi K} \left(-\alpha_1 - \alpha_4^u + \alpha_4^c - \alpha_{4,EW}^u + \alpha_{4,EW}^c \right)$$

$$E_2^F = A_{K\pi} \left(-\alpha_2 - \frac{3}{2}(\alpha_{3,EW}^u - \alpha_{3,EW}^c) \right) \\ + A_{\pi K} \left(\alpha_4^u - \alpha_4^c - \frac{1}{2}(\alpha_{4,EW}^u - \alpha_{4,EW}^c) \right),$$

$$A^F = A_{\pi K} \left(-\alpha_4^u + \alpha_4^c + \frac{1}{2}(\alpha_{4,EW}^u - \alpha_{4,EW}^c) \right),$$

$$P^F = A_{\pi K} \left(-\alpha_4^c + \frac{1}{2}\alpha_{4,EW}^c \right),$$

$$\Delta P_1^F = -A_{\pi K} \frac{3}{2}\alpha_{4,EW}^c,$$

$$\Delta P_2^F = -A_{K\pi} \frac{3}{2}\alpha_{3,EW}^c,$$

$$+ m_s = (98 \pm 6 \pm 12) \text{ MeV}$$

$$A_{\pi K} = G_F / \sqrt{2} m_B^2 f_k F_\pi(0)$$

$$A_{K\pi} = G_F / \sqrt{2} m_B^2 f_\pi F_k(0)$$

f_π	0.1307 GeV	f_K	0.1598 GeV
$F^{B \rightarrow \pi}$	0.27 ± 0.08	$F^{B \rightarrow K} / F^{B \rightarrow \pi}$	1.20 ± 0.10
τ_{B^0}	$1.546 \cdot 10^{-12}$ ps	τ_{B^+}	$1.674 \cdot 10^{-12}$ ps
m_B	$5.2794 \text{ GeV}/c^2$	f_B	$0.189 \pm 0.027 \text{ GeV}$
m_π	$0.14 \text{ GeV}/c^2$	m_K	$0.493677 \text{ GeV}/c^2$

Conclusions

Flavour physics is a unique tool for searching and studying NP complementary to the LHC

There is a first evidence for NP in $b \leftrightarrow s$ transitions. Confirmation in Summer

From $\Delta F=2$ transitions, a pattern of flavour violation in NP emerges:

$2 \leftrightarrow 3: O(1)$, $1 \leftrightarrow 3: < O(0.1)$, $1 \leftrightarrow 2$ strong suppr.

The next 15 years of flavour physics are well motivated and clearly planned:
exciting times ahead