

# THEORY ISSUES IN MEASURING $\gamma$

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Introduction

$\gamma$  from tree  $B_{(s)}$  decays

$\gamma$  from penguin  $B_{(s)}$  decays

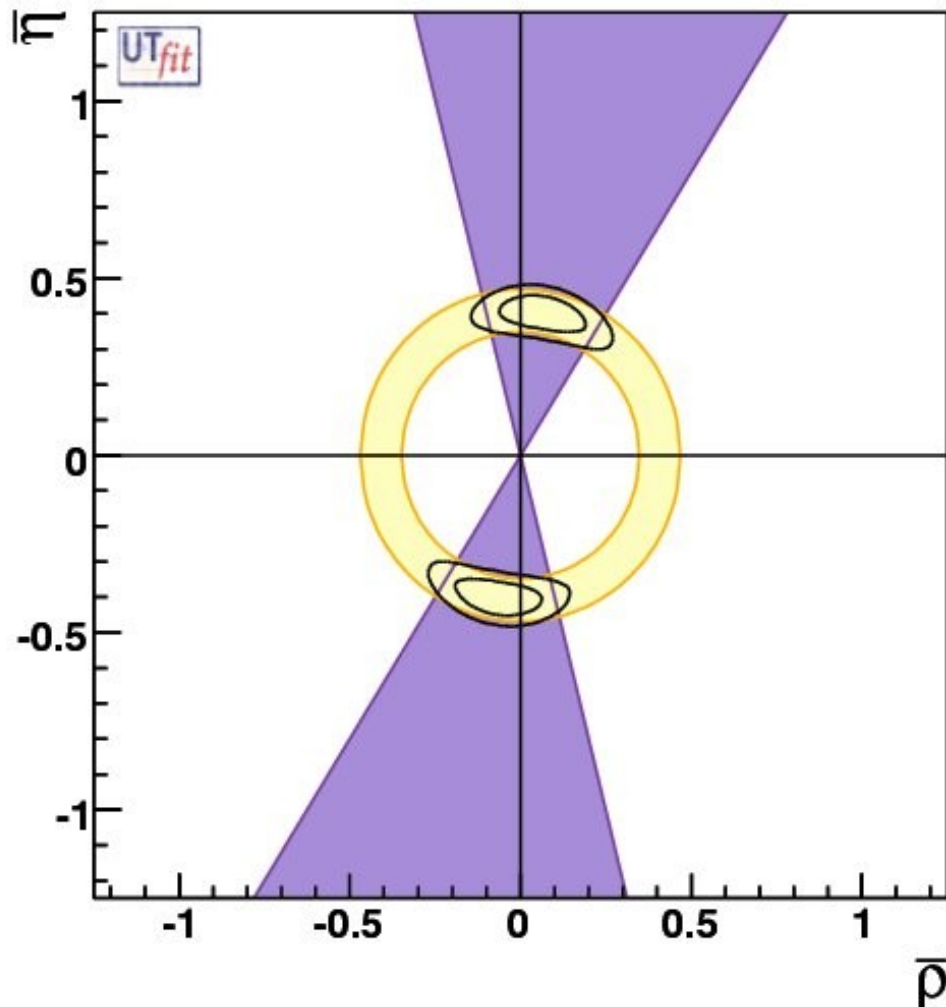
Extracting the tree from penguin  $B_s$  decays

Conclusions

# INTRODUCTION - I

- The extraction of  $\gamma$  from tree-level decays is of fundamental importance for NP searches:
  - Fix  $\rho$  and  $\eta$  from tree-level  $V_{ub}$  and  $\gamma$ : insensitive to loop-mediated NP
  - Determine NP contributions to loop-mediated FCNC and CPV (meson mixing, penguin decays,...)
- Precision measurement of  $\gamma$  needed to fully exploit the constraining power of K and B mixing

# INTRODUCTION - II



- Impressive achievement of B-factories: tree-level determination of the UT
- Allows to constrain new CPV in K and  $B_d$  mixing at the 10% level

# INTRODUCTION - III

- The extraction of  $\gamma$  from penguin decays is sensitive to NP
- Some knowledge of hadronic parameters is necessary
- Several strategies based on hadronic models and/or flavour symmetries
- Quantifying hadronic uncertainties is the main issue

# $\gamma$ FROM TREE DECAYS

- Basic idea: use interference between  $b \rightarrow c\bar{u}q$  and  $b \rightarrow \bar{u}c q$  decays, with  $q=d,s$ .
- No penguins present: no loop-mediated NP
- Methods can be classified according to the charge of the decaying B meson and to the final state  $f$  in which the D meson is reconstructed

# $\gamma$ FROM TREE DECAYS - $B^\pm$

- Use interference between  $B^\pm \rightarrow D_f K^\pm$  and  $B^\pm \rightarrow \bar{D}_f K^\pm$ , with  $f$  a common final state for  $D$  and  $\bar{D}$
- Define  $A(DK^+)/A(\bar{D}K^+) = r_B e^{i(\delta_B + \gamma)}$ , with  $r_B = 0.37 |\bar{C} + A| / |T + C|$ .
- Ratio of hadronic matrix elements difficult to estimate, expect  $r_B \sim 0.1$ .

# $\gamma$ FROM TREE DECAYS - $B^\pm$

- GLW: use  $f$   $CP^\pm$  eigenstate, measure

$$R_{CP^\pm} = \frac{\Gamma(B^+ \rightarrow D_{CP^\pm}^0 K^+) + \Gamma(B^- \rightarrow D_{CP^\pm}^0 K^-)}{\Gamma(B^+ \rightarrow D^0 K^+) + \Gamma(B^- \rightarrow \bar{D}^0 K^-)} = 1 + r_B^2 \pm 2r_B \cos \gamma \cos \delta_B$$

$$A_{CP^\pm} = \frac{\Gamma(B^+ \rightarrow D_{CP^\pm}^0 K^+) - \Gamma(B^- \rightarrow D_{CP^\pm}^0 K^-)}{\Gamma(B^+ \rightarrow D_{CP^\pm}^0 K^+) + \Gamma(B^- \rightarrow D_{CP^\pm}^0 K^-)} = \frac{\pm 2r_B \sin \gamma \sin \delta_B}{R_{CP^\pm}}$$

extract  $r_B$  and  $\gamma$  with a four-fold ambiguity

Gronau & Wiler 91; Gronau & London 91

- ADS:  $f$  non  $CP$  eigenstate. Define

Atwood, Dunietz & Soni 97, 01

$A(D \rightarrow f)/A(D \rightarrow \bar{f}) = r_D e^{i\delta_D}$  and measure

$$R_{ADS} = \frac{\Gamma(B^+ \rightarrow \bar{f} K^+) + \Gamma(B^- \rightarrow f K^-)}{\Gamma(B^+ \rightarrow f K^+) + \Gamma(B^- \rightarrow \bar{f} K^-)} = r_D^2 + r_B^2 + 2r_B r_D \cos \gamma \cos(\delta_B + \delta_D)$$

$$A_{ADS} = \frac{\Gamma(B^- \rightarrow f K^-) - \Gamma(B^+ \rightarrow \bar{f} K^+)}{\Gamma(B^- \rightarrow f K^-) + \Gamma(B^+ \rightarrow \bar{f} K^+)} = r_B r_D [\cos(\delta + \gamma) + \cos(\delta - \gamma)] / R_{ADS}$$

# $\gamma$ FROM TREE DECAYS - $B^\pm$

- $\gamma$ ,  $r_B$  and  $\delta_B$  are common to all final states  $f$
- $r_{Df}$  can be measured from  $D$  decays, so for each  $f$  add one unknown ( $\delta_{Df}$ ) and two observables ( $R_f$  and  $A_f$ )
- The same analysis can be repeated for  $D^*$  or  $K^*$
- Add as many final states as possible!



# $\gamma$ FROM TREE DECAYS - $B^\pm$

- GGSZ: use Dalitz analysis of multi-body CP eigenstate final states (equivalent of measuring a continuum of final states)
- model-dependent approach: fit  $A(D \rightarrow f)$  from the Dalitz plot
- model-independent approach: bin the Dalitz plot and determine the D decay strong phase from the B decay data. Need much more statistics and/or  $D\bar{D}$  data.

Giri, Grossman, Soffer & Zupan 03

# $\gamma$ FROM TREE DECAYS - $B_d$

- Can use DK final states also for  $B_d$  decays
- Define  $A(DK^0)/A(\bar{D}K^0)=r_{B0}e^{i(\delta_{B0}+\gamma)}$ , with  $r_{B0}=0.37 |\bar{C}|/|C|$  expected around  $r_{B0}\sim 0.4$
- If final state is self-tagging (e.g.  $K^{*0} \rightarrow K^+\pi^-$ )  
no  $B_d$  mixing
- If final state admits mixing (e.g.  $B^0 \rightarrow D^-K^0\pi^+$ )  
need a time-dependent analysis, measure

$$2\beta + \gamma$$

Kayser & London 00; Atwood & Soni 03; Fleischer 03; Gronau et al. 04

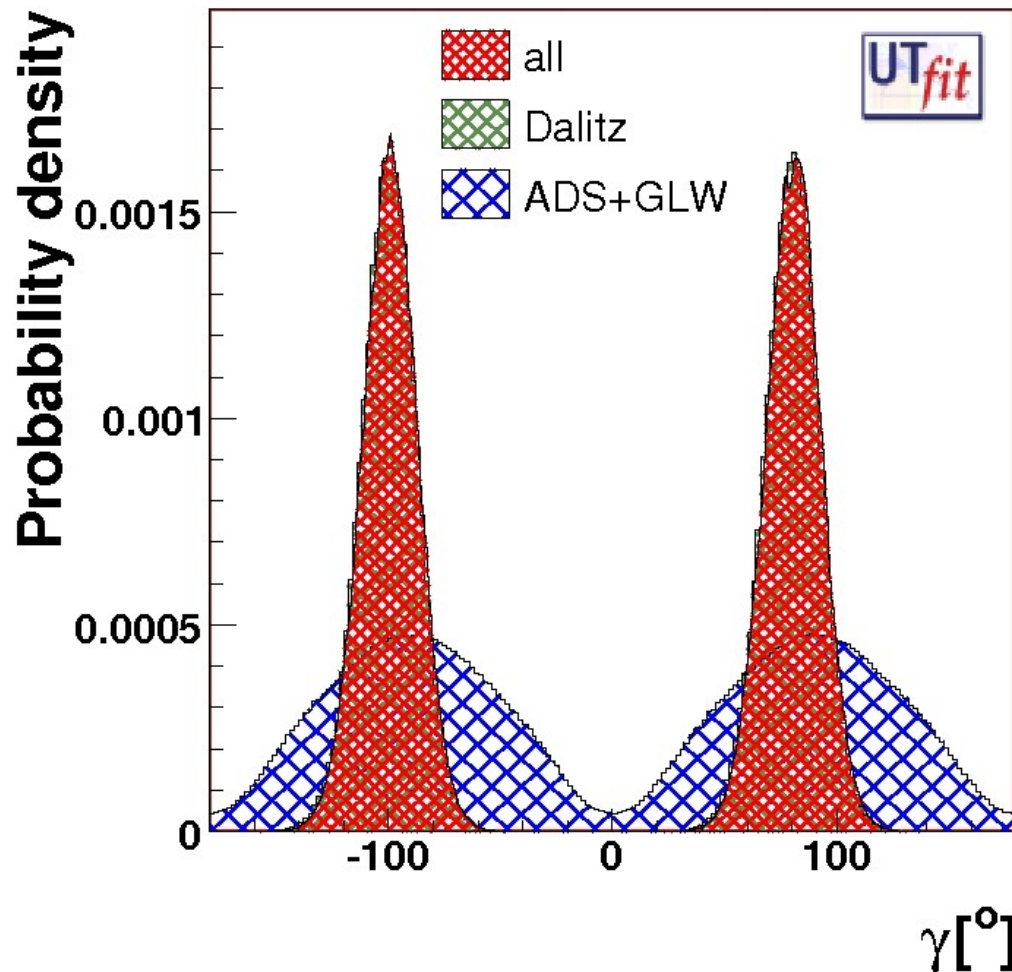
# $\gamma$ FROM TREE DECAYS - $B_s$

- A time-dependent analysis of  $B_s \rightarrow D_s^\pm K^\pm$  gives access to  $\phi_s + \gamma$

Dunietz & Sachs 88; Aleksan, Dunietz & Kaiser 92; Dunietz 98

- Combining this with the measurement of the  $B_s$  mixing phase  $\phi_s$  one can extract  $\gamma$
- There is however an eightfold ambiguity (fourfold if  $\Delta\Gamma_s$  effects are observable)
- The ambiguity can be resolved using  $SU(3)$  related  $B_d \rightarrow D^\pm \pi^\pm$  decays

# $\gamma$ FROM TREE DECAYS TODAY



- Combining all results gives

$$\gamma = (80 \pm 13)^\circ$$

$$r_B = 0.10 \pm 0.02$$

$$r_B^* = 0.09 \pm 0.04$$

- Theory error is (and will remain in the future) negligible

# $\gamma$ FROM PENGUIN DECAYS

- Idea: extract  $\gamma$  using the interference of tree ( $V_{ub}$ ) and penguin ( $V_{tb}$ ) amplitudes
- **Sensitive to NP contributions to penguins**
- **Need some control of hadronic matrix elements**
- The main issue is the determination of hadronic uncertainties: theoretical errors are dominant in this case

# $\gamma$ FROM $B \rightarrow K\pi$

- Quark level transition is  $b \rightarrow s q \bar{q}$  ( $q=u,d$ )
- $V_{ub} V_{us}^*$  contribution is doubly Cabibbo suppressed w.r.t.  $V_{tb} V_{ts}$ : interference effects expected at the level of 10% or smaller
- Six independent hadronic matrix elements  $\Rightarrow$  11 real hadronic parameters +  $\gamma$
- 9 observables, 2 quadrangular isospin relations  $\Rightarrow$  need some theory input

# $\gamma$ FROM $B \rightarrow K\pi$ - II

- Most popular approaches: neglect of some contributions (annihilations, electroweak penguins, ...) and/or  $B \rightarrow \pi\pi$  plus flavour symmetries  
Fleischer 96; Fleischer & Mannel 98; Fleischer 99; Atwood & Soni 98; Gronau & Rosner 98, 02; Neubert & Rosner 98; Neubert 99; Buras & Fleischer 99, 00; Kim, Oh & Yoon 07; ...
- Theoretical error can only be guessed: no way to check it on data
- Cannot conclude from vanishing CP asymmetries that one amplitude dominates: can be just due to vanishing strong phases

# $\gamma$ FROM $B_s \rightarrow KK$ and $B \rightarrow \pi\pi$

- Using  $SU(3)$  can combine  $B_s \rightarrow KK$  and  $B \rightarrow \pi\pi$  to obtain a determination of  $\gamma$
- The same considerations on hadronic uncertainties apply here: to estimate the theory error need to quantify the amount of  $SU(3)$  (U-spin) breaking

Fleischer 99; Fleischer & Matias 02  
London, Matias & Virto 05; Baek,  
London, Matias & Virto 06; ...



# $\gamma$ FROM $B_s \rightarrow K^* \pi$

- Idea: isolate tree decay amplitude using a Dalitz plot analysis and extract its phase
- Quark level transition is  $b \rightarrow d \bar{q} q$  ( $q=u,d$ )
  - tree contribution ( $V_{ub} V_{ud}^*$ ) is not Cabibbo suppressed w.r.t.  $V_{tb} V_{td}^*$
  - electroweak penguins are negligible
- Analogous to the extraction of  $\alpha$  from  $B_d \rightarrow \rho \pi$  decays

Ciuchini, Pierini & L.S. 06;  
Gronau, Pirjol, Soni & Zupan 06, 07

# $\gamma$ FROM $B_s \rightarrow K^* \pi$ - II

- Extract  $\gamma$  using the following amplitude isospin relations:

$$A_s = A(B_s \rightarrow K^{*-} \pi^+) + \sqrt{2} A(B_s \rightarrow \bar{K}^{*0} \pi^0) = -V_{ub}^* V_{ud} (E_1 + E_2)$$

$$\bar{A}_s = A(\bar{B}_s \rightarrow K^{*+} \pi^-) + \sqrt{2} A(\bar{B}_s \rightarrow K^{*0} \pi^0) = -V_{ub} V_{ud}^* (E_1 + E_2)$$

$$R_s = \frac{\bar{A}_s}{A_s} = \frac{V_{ub} V_{ud}^*}{V_{ub}^* V_{ud}} = e^{-2i\gamma}$$

- $A(B_s \rightarrow K^{*-} \pi^+)$  and  $A(B_s \rightarrow \bar{K}^{*0} \pi^0)$  can be extracted from  $B_s \rightarrow K^- \pi^+ \pi^0$  Dalitz plot, and the conjugate amplitudes from  $\bar{B}_s \rightarrow K^+ \pi^- \pi^0$

# $\gamma$ FROM $B_s \rightarrow K^* \pi$ - III

- To obtain the relative phase of the  $B_s \rightarrow K^- \pi^+ \pi^0$  and  $\bar{B}_s \rightarrow K^+ \pi^- \pi^0$  Dalitz plots, use the  $B_s \rightarrow K_S \pi^+ \pi^-$  Dalitz plot, exploiting interference of  $K^{*+} \pi^-$  and  $K^{*-} \pi^+$  with  $\rho^0 K_S$  and other  $\pi^+ \pi^-$  resonances.
- At hadron colliders, the sensitivity is given by the  $\text{Re } \lambda \Delta\Gamma_s / \Gamma_s$  term in the time-integrated rate ( $\lambda = q/p \bar{A}/A$ ). Of course, a time-dependent analysis would also help.

# CONCLUSIONS

- The extraction of  $\gamma$  from tree-level amplitudes plays a crucial role in looking for NP in the flavour sector
- Theory uncertainties are negligible in tree-level decays
- Tree-level amplitudes can be successfully isolated in  $B_s \rightarrow K^* \pi$  decays - same procedure as in  $B_d \rightarrow \rho \pi$

# CONCLUSIONS - II

- The extraction of  $\gamma$  from the interference of tree and penguin amplitudes looks much more problematic
- Some theoretical input needed to reduce the number of unknown hadronic parameters
- Theoretical uncertainties very difficult to estimate, cannot be verified using data
- Use penguin channels to look for NP in penguin amplitudes, taking  $\gamma$  as input!