

# Higgs Characterisation

in the effective field theory approach  
via the FeynRules and MadGraph5\_aMC@NLO frameworks

Kentarou Mawatari

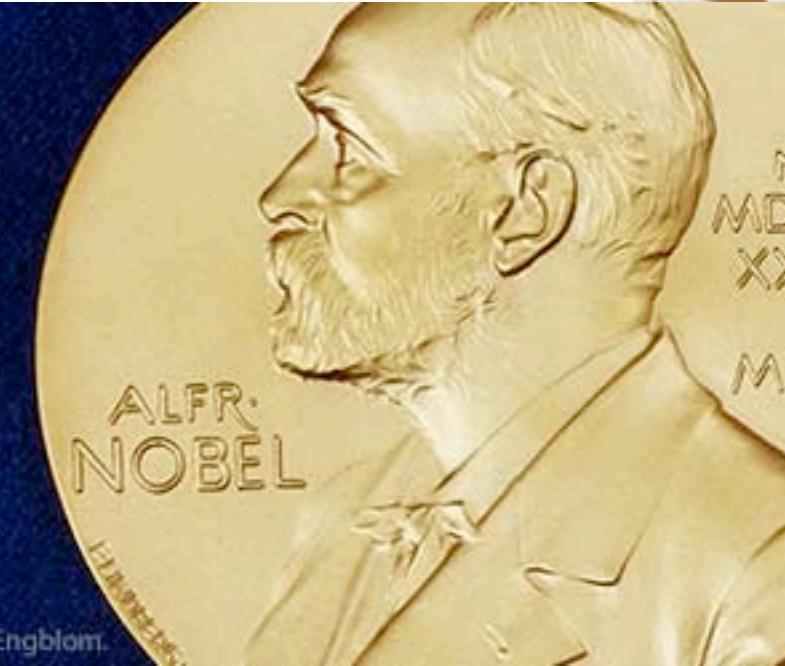
(Vrije Universities Brussel and International Solvay Institutes)

- ▶ Sec. II in YR3 of the LHC Higgs Cross Section Working Group [arXiv:1307.1347]
- ▶ Artoisenet, de Aquino, Demartin, Frederix, Frixione, Maltoni, Mandal, Mathews, KM, Ravindran, Seth, Torrielli, Zaro  
“A framework for Higgs characterisation” JHEP11(2013)043 [arXiv:1306.6464]
- ▶ Maltoni, KM, Zaro  
“Higgs characterisation via VBF/VH” EPJC74(2014)2710 [arXiv:1311.1829]
- ▶ Alloul, Fuks, Sanz  
“Phenomenology of the Higgs Effective Lagrangian via FR” JHEP04(2014)110 [arXiv:1310.5150]

# 2013 NOBEL PRIZE IN PHYSICS

# François Englert Peter W. Higgs

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## BROKEN SYMMETRY AND THE MASS OF GAUGE VECTOR MESONS\*

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(Received 26 June 1964)

## BROKEN SYMMETRIES, MASSLESS PARTICLES AND GAUGE FIELDS

P. W. HIGGS

Tait Institute of Mathematical Physics, University of Edinburgh, Scotland

Received 27 July 1964

"for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider" [<http://www.nobelprize.org>]



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- Introduction
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- Summary

# Spin/parity determination

$X \rightarrow VV^* \rightarrow 4l$

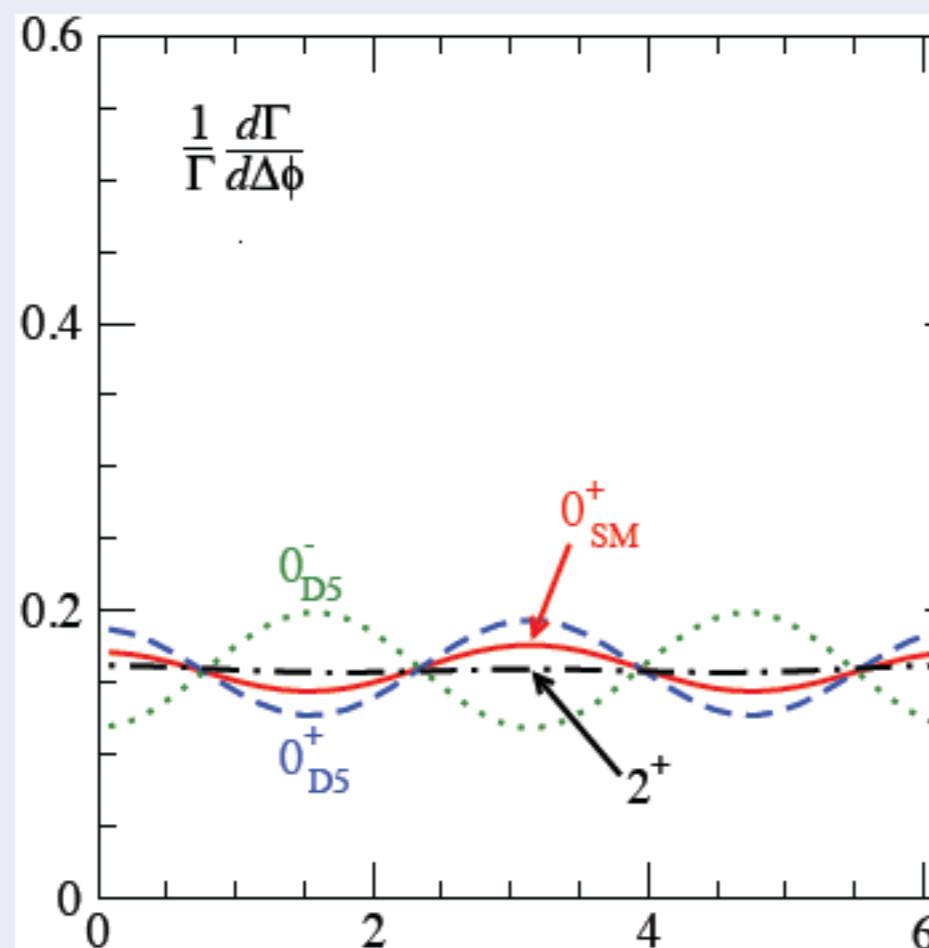
[Dell'Aquila, Nelson, PRD(1986)]

[Choi, Miller, Mühlleitner, Zerwas, PLB(2003)]

[Gao et al, PRD(2010)] ...

[Bolognesi et al, PRD(2012)]

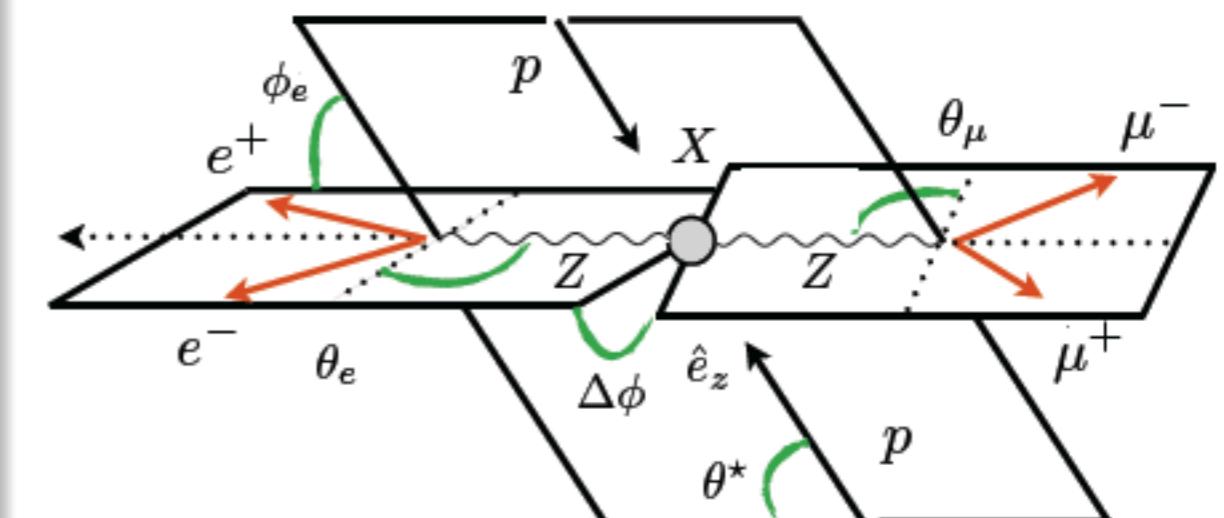
$X \rightarrow ZZ^* \rightarrow 4\ell$



$$\mathcal{L}_{0^+_{SM}} = g_{0^+_{SM}} V_\mu V^\mu X_0$$

$$\mathcal{L}_{0^+_{D5}} = g_{0^+_{D5}} V_{\mu\nu} V^{\mu\nu} X_0$$

$$\mathcal{L}_{0^-_{D5}} = g_{0^-_{D5}} V_{\mu\nu} \tilde{V}^{\mu\nu} X_0$$



$d\sigma/d\Delta\phi \sim \text{const.}$  for  $0^+_{SM}$ ,  $d\sigma/d\Delta\phi \sim 1 \pm A \cos 2\Delta\phi$  for  $0^\pm_{D5}$ .

# Spin/parity determination

## $X \rightarrow VV^* \rightarrow 4l$ vs. VBF

[Choi, Miller, Mühlleitner, Zerwas, PLB(2003)]

[Gao et al, PRD(2010)] ...

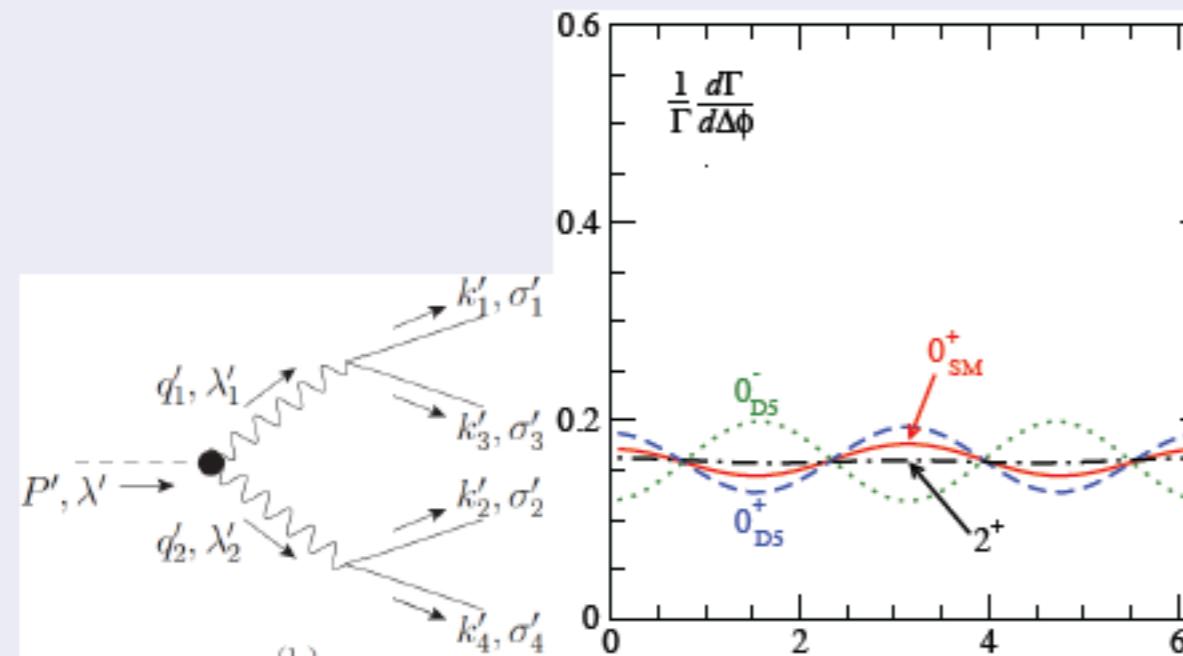
[Bolognesi et al, PRD(2012)]

[Plehn, Rainwater, Zeppenfeld, PRL(2002)]

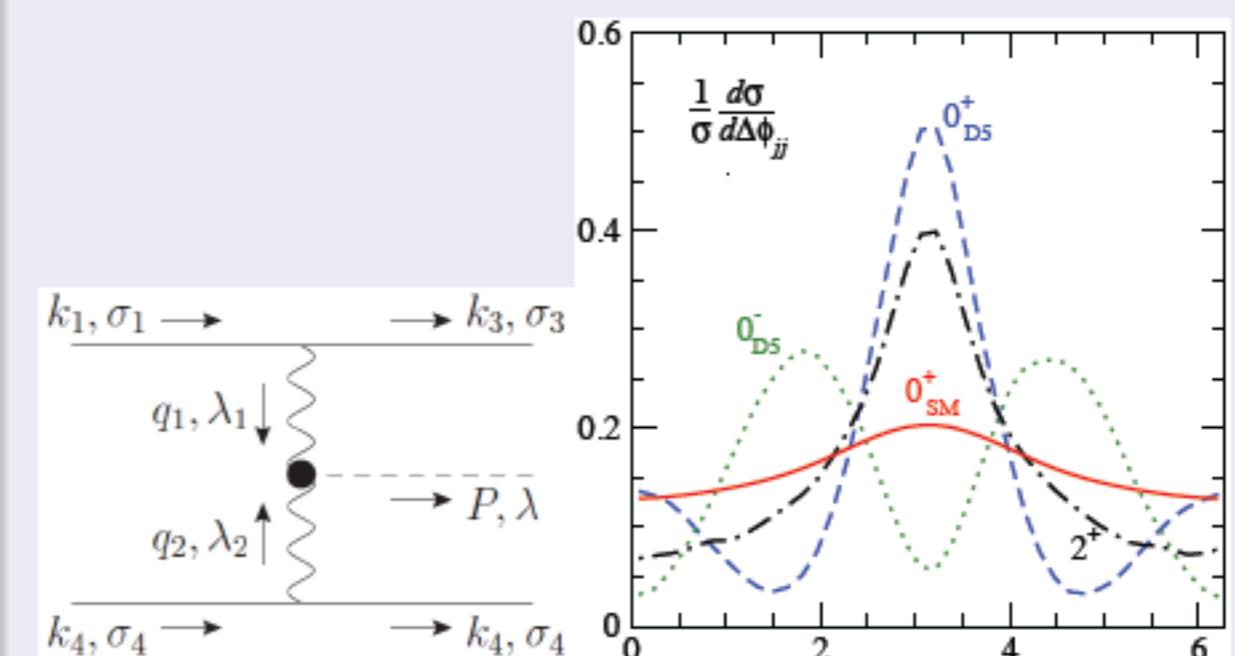
[Hagiwara, Li, KM, JHEP(2009)] ...

[Englert, Goncalves-Netto, KM, Plehn, JHEP(2013)]

$X \rightarrow ZZ^* \rightarrow 4\ell$



Vector boson fusion (VBF)



$$d\sigma/d\Delta\phi \sim \text{const. for } 0_{SM}^+, \quad d\sigma/d\Delta\phi \sim 1 \pm A \cos 2\Delta\phi \text{ for } 0_{D5}^\pm.$$

Nontrivial azimuthal angle correlations of the decay planes ( $X \rightarrow ZZ$ ) and the jets (VBF) can be explained as the quantum interference among different helicity states of the intermediate vector-bosons.

# The SM cannot be the ultimate theory !

- How can we find the BSM physics?
  - ✓ Find new particles/phenomena.
    - ➡ Top-down approach: SUSY, ED, 2HDM, ...
  - ✓ Find small deviations from the SM expectation.
    - ➡ Bottom-up approach: Effective field theory

# Is this the Standard Model scalar boson?

- Higgs boson precision measurement
- determination of the Higgs boson Lagrangian
  - the structure of the operators, linked to the spin/parity of a Higgs boson
    - ▶ distributions
    - the coupling strength
      - ▶ rate
- How do we approach to get them?

# Tools for Higgs Physics

## Cross Section

### ggF

- [HIGLU](#) (NNLO QCD+NLO EW)
- [iHixs](#) (NNLO QCD+NLO EW)
- [FeHiPro](#) (NNLO QCD+NLO EW)
- [HNNLO, HRes](#) (NNLO+NNLL QCD)
- [SusHi](#) (NNLO QCD)
- [RGHiggs](#) (NNLO+NNNLL QCD)
- [ggHiggs](#) (approx. NNNLO QCD)

### VBF

- [VV2H](#) (NLO QCD)
- [VBFNLO](#) (NLO QCD)
- [HAWK](#) (NLO QCD+EW)
- [VBF@NNLO](#) (NNLO QCD)

### WH/ZH

- [V2HV](#) (NLO QCD)
- [HAWK](#) (NLO QCD+EW)
- [VH@NNLO](#) (NNLO)

### ttH

- [HQQ](#) (LO QCD)

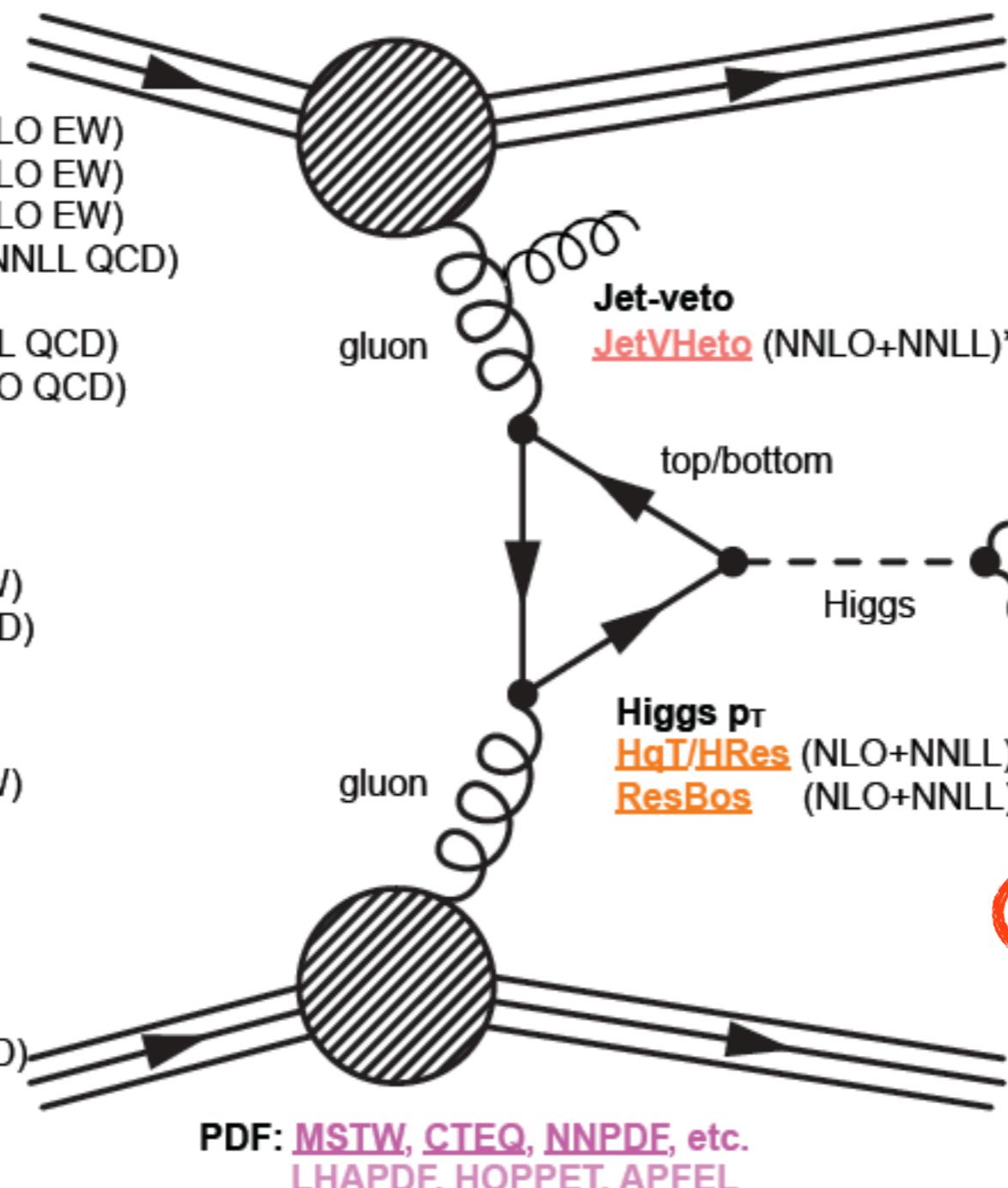
### bbH

- [bbh@NNLO](#) (NNLO QCD)

### HH

- [HPAIR](#) (NLO QCD)

+ private codes.



## NLO MC

- [POWHEG](#) [MiNLO](#)
- [MadGraph5](#) [aMC@NLO](#)
- [SHERPA](#) [MEPS@NLO](#)

## LO MC

- [gg2VV](#)

## NLO ME

- [MCFM](#), [MG5\\_aMC@NLO](#)

## W/Z

- Higgs Decay**
- [HDECAY](#) (NLO++)
- [Prophecy4f](#) (NLO)

## W/Z

## Higgs Properties

- [MELA/JHU](#), [MEKD](#)
- [MG5\\_aMC@NLO](#) (HC)

in this talk

## MSSM/2HDM

- [FeynHiggs](#), [CPSuperH](#)
- [SusHi+2HDMC](#)
- [HIGLU+HDECAY](#)

\* NLO+NNLL in differential

# Effective field theory approach

- Given the fact that only a 125 GeV SM-like boson and nothing else so far, the effective field theory approach is one of the best way to explore BSM effects.
- ▶ All new particles and phenomena are assumed to appear at some scale  $\Lambda$ .
- ▶ Not predictive at scales larger than  $\Lambda \rightarrow$  loss of unitarity
- ▶ Below  $\Lambda$ , all new physics effects are parametrized by higher dimensional gauge invariant operators made of SM fields.  $\rightarrow$  many parameters
- ▶ No assumption on the form of new physics  $\rightarrow$  model independent
- ▶ Renormalisable order by order in the scale  $\Lambda \rightarrow$  systematically improvable

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \frac{1}{\Lambda^2} \mathcal{L}_6 + \dots \quad \mathcal{L}_6 = \sum_i C_i Q_i$$

Buchmuller&Wyler 1986 ...  
Grzadkowski et al. 2010

# Higgs effective Lagrangian before vs. after EW symmetry breaking

- D6 (the gauge basis): **HEL** [Alloul, Fuks, Sanz, arXiv:1310.5150]
  - ▶ Only using Standard Model gauge-eigenstates
  - ▶ Several operators may be associated with a single coupling (in the mass basis)
  - ▶ One operator associated with several couplings (in the mass basis)
  - ▶ <https://feynrules.irmp.ucl.ac.be/wiki/HEL>
- D5 (the mass basis): **HC** [Artoisenet et al., arXiv:1306.6464]
  - ▶ Couplings of the physical Higgs boson to the Standard Model (physical) states
  - ▶ One operator associated with a single coupling (and Lorentz structure)
  - ▶ No assumption on the Higgs boson spin
  - ▶ <https://feynrules.irmp.ucl.ac.be/wiki/HiggsCharacterisation>

# D6 Higgs Effective Lagrangian

[ from Contino, Ghezzi, Grojean, Muhlleitner, Spira (JHEP '13) ]  
 [ Alloul, Fuks, Sanz (1310.5150) ]

$$\begin{aligned} \mathcal{L}_{\text{SILH}} = & \frac{\bar{c}_H}{2v^2} \partial^\mu [\Phi^\dagger \Phi] \partial_\mu [\Phi^\dagger \Phi] + \frac{\bar{c}_T}{2v^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] - \frac{\bar{c}_6 \lambda}{v^2} [H^\dagger H]^3 \\ & - \left[ \frac{\bar{c}_u}{v^2} y_u \Phi^\dagger \Phi \cdot \bar{Q}_L u_R + \frac{\bar{c}_d}{v^2} y_d \Phi^\dagger \Phi \bar{Q}_L d_R + \frac{\bar{c}_l}{v^2} y_\ell \Phi^\dagger \Phi \bar{L}_L e_R + \text{h.c.} \right] \\ & + \frac{ig \bar{c}_W}{m_W^2} [\Phi^\dagger T_{2k} \overleftrightarrow{D}^\mu \Phi] D^\nu W_{\mu\nu}^k + \frac{ig' \bar{c}_B}{2m_W^2} [\Phi^\dagger \overleftrightarrow{D}^\mu \Phi] \partial^\nu B_{\mu\nu} \\ & + \frac{2ig \bar{c}_{HW}}{m_W^2} [D^\mu \Phi^\dagger T_{2k} D^\nu \Phi] W_{\mu\nu}^k + \frac{ig' \bar{c}_{HB}}{m_W^2} [D^\mu \Phi^\dagger D^\nu \Phi] B_{\mu\nu} \\ & + \frac{\bar{g}'^2 c_\gamma}{m_W^2} \Phi^\dagger \Phi B_{\mu\nu} B^{\mu\nu} + \frac{\bar{g}_s^2 c_g}{m_W^2} \Phi^\dagger \Phi G_{\mu\nu}^a G_a^{\mu\nu}, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{CP} = & \frac{ig \bar{c}_{HW}}{m_W^2} D^\mu \Phi^\dagger T_{2k} D^\nu \Phi \widetilde{W}_{\mu\nu}^k + \frac{ig' \bar{c}_{HB}}{m_W^2} D^\mu \Phi^\dagger D^\nu \Phi \widetilde{B}_{\mu\nu} + \frac{g'^2 \bar{c}_\gamma}{m_W^2} \Phi^\dagger \Phi B_{\mu\nu} \widetilde{B}^{\mu\nu} \\ & + \frac{g_s^2 \bar{c}_g}{m_W^2} \Phi^\dagger \Phi G_{\mu\nu}^a \widetilde{G}_a^{\mu\nu} + \frac{g^3 \bar{c}_{3W}}{m_W^2} \epsilon_{ijk} W_{\mu\nu}^i W_{\rho\sigma}^{jk} \widetilde{W}^{\rho\mu k} + \frac{g_s^3 \bar{c}_{3G}}{m_W^2} f_{abc} G_{\mu\nu}^a G_{\rho\sigma}^b \widetilde{G}_{\mu\nu}^{\rho\sigma} \\ \mathcal{L}_G = & \frac{g^3 \bar{c}_{3W}}{m_W^2} \epsilon_{ijk} W_{\mu\nu}^i W_{\rho\sigma}^{jk} W^{\rho\mu k} + \frac{g_s^3 \bar{c}_{3G}}{m_W^2} f_{abc} G_{\mu\nu}^a G_{\rho\sigma}^b G^{\rho\mu c} + \frac{\bar{c}_{2W}}{m_W^2} D^\mu W_{\mu\nu}^k D_\rho W_k^{\rho\nu} \\ & + \frac{\bar{c}_{2B}}{m_W^2} \partial^\mu B_{\mu\nu} \partial_\rho B^{\rho\nu} + \frac{\bar{c}_{2G}}{m_W^2} D^\mu G_{\mu\nu}^a D_\rho G_a^{\rho\nu}, \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{F_1} = & \frac{i\bar{c}_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu Q_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\bar{c}'_{HQ}}{v^2} [\bar{Q}_L \gamma^\mu T_{2k} Q_L] [\Phi^\dagger T_2^k \overleftrightarrow{D}_\mu \Phi] \\ & + \frac{i\bar{c}_{Hu}}{v^2} [\bar{u}_R \gamma^\mu u_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{i\bar{c}_{Hd}}{v^2} [\bar{d}_R \gamma^\mu d_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] \\ & - \left[ \frac{i\bar{c}_{Hud}}{v^2} [\bar{u}_R \gamma^\mu d_R] [\Phi \cdot \overleftrightarrow{D}_\mu \Phi] + \text{h.c.} \right] \\ & + \frac{i\bar{c}_{HL}}{v^2} [\bar{L}_L \gamma^\mu L_L] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi] + \frac{4i\bar{c}'_{HL}}{v^2} [\bar{L}_L \gamma^\mu T_{2k} L_L] [\Phi^\dagger T_2^k \overleftrightarrow{D}_\mu \Phi] \\ & + \frac{i\bar{c}_{He}}{v^2} [\bar{e}_R \gamma^\mu e_R] [\Phi^\dagger \overleftrightarrow{D}_\mu \Phi], \end{aligned}$$

$$\begin{aligned} \mathcal{L}_{F_2} = & \left[ - \frac{2g' \bar{c}_{uB}}{m_W^2} y_u \Phi^\dagger \cdot \bar{Q}_L \gamma^{\mu\nu} u_R B_{\mu\nu} - \frac{4g \bar{c}_{uW}}{m_W^2} y_u \Phi^\dagger \cdot (\bar{Q}_L T_{2k}) \gamma^{\mu\nu} u_R W_{\mu\nu}^k \right. \\ & - \frac{4g_s \bar{c}_{uG}}{m_W^2} y_u \Phi^\dagger \cdot \bar{Q}_L \gamma^{\mu\nu} T_a u_R G_{\mu\nu}^a + \frac{2g' \bar{c}_{dB}}{m_W^2} y_d \Phi \bar{Q}_L \gamma^{\mu\nu} d_R B_{\mu\nu} \\ & + \frac{4g \bar{c}_{dW}}{m_W^2} y_d \Phi (\bar{Q}_L T_{2k}) \gamma^{\mu\nu} d_R W_{\mu\nu}^k + \frac{4g_s \bar{c}_{dG}}{m_W^2} y_d \Phi \bar{Q}_L \gamma^{\mu\nu} T_a d_R G_{\mu\nu}^a \\ & \left. + \frac{2g' \bar{c}_{eB}}{m_W^2} y_\ell \Phi \bar{L}_L \gamma^{\mu\nu} e_R B_{\mu\nu} + \frac{4g \bar{c}_{eW}}{m_W^2} y_\ell \Phi (\bar{L}_L T_{2k}) \gamma^{\mu\nu} e_R W_{\mu\nu}^k + \text{h.c.} \right] \end{aligned}$$

◆ The model file is publicly available. (<https://feynrules.irmp.ucl.ac.be/wiki/HEL>)

# Mapping between the D6 and D5 operators

HC [arXiv: 1306.6464]

$$\begin{aligned}\mathcal{L}_0^f &= - \sum_{f=t,b,\tau} \bar{\psi}_f (c_\alpha \kappa_{Hff} g_{Hff} + i s_\alpha \kappa_{Aff} g_{Aff} \gamma_5) \psi_f X_0 \\ \mathcal{L}_0^V &= \left\{ c_\alpha \kappa_{SM} \left[ \frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \right. \\ &\quad - \frac{1}{4} \left[ c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma} A_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma} A_{\mu\nu} \tilde{A}^{\mu\nu} \right] \\ &\quad - \frac{1}{2} \left[ c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma} Z_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{AZ\gamma} g_{AZ\gamma} Z_{\mu\nu} \tilde{A}^{\mu\nu} \right] \\ &\quad - \frac{1}{4} \left[ c_\alpha \kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + s_\alpha \kappa_{Agg} g_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu} \right] \\ &\quad \left. - \frac{1}{4} \frac{1}{\Lambda} \left[ c_\alpha \kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + s_\alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu} \right] \right. \\ &\quad - \frac{1}{2} \frac{1}{\Lambda} \left[ c_\alpha \kappa_{HWW} W_\mu^+ W^{-\mu\nu} + s_\alpha \kappa_{AWW} W_\mu^+ \tilde{W}^{-\mu\nu} \right] \\ &\quad - \frac{1}{\Lambda} c_\alpha [\kappa_{H\theta\gamma} Z_\nu \partial_\mu \gamma^\nu + (\kappa_{H\theta W} W)]\end{aligned}$$

HEL [arXiv: 1310.5150]

Eq. (2.25)	Ref. [46]	Section 2.1
$g_{hgg}$	$c_\alpha \kappa_{Hgg} g_{Hgg}$	$g_H - \frac{4\bar{c}_g g_s^2 v}{m_W^2}$
$\tilde{g}_{hgg}$	$s_\alpha \kappa_{Agg} g_{Agg}$	$- \frac{4\bar{c}_g g_s^2 v}{m_W^2}$
$g_{h\gamma\gamma}$	$c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma}$	$a_H - \frac{8\bar{c}_\gamma s_W^2}{m_W}$
$\tilde{g}_{h\gamma\gamma}$	$s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma}$	$- \frac{8\bar{c}_\gamma s_W^2}{m_W}$
$g_{hzz}^{(1)}$	$\frac{1}{\Lambda} c_\alpha \kappa_{HZZ}$	$\frac{2g}{c_W^2 m_W} \left[ \bar{c}_{HB} s_W^2 - 4\bar{c}_\gamma s_W^4 + c_W^2 \bar{c}_{HW} \right]$
$\tilde{g}_{hzz}$	$\frac{1}{\Lambda} s_\alpha \kappa_{AZZ}$	$\frac{2g}{c_W^2 m_W} \left[ \bar{c}_{HB} s_W^2 - 4c_\gamma s_W^4 + c_W^2 \bar{c}_{HW} \right]$
$g_{hzz}^{(2)}$	$\frac{1}{\Lambda} c_\alpha \kappa_{H\theta Z}$	$\frac{g}{c_W^2 m_W} \left[ (\bar{c}_{HW} + \bar{c}_W) c_W^2 + (\bar{c}_B + \bar{c}_{HB}) s_W^2 \right]$
$g_{hzz}^{(3)}$	$c_\alpha \kappa_{SM} g_{HZZ}$	$\frac{gm_W}{c_W^2} \left[ 1 - \frac{1}{2} \bar{c}_H - 2\bar{c}_T + 8\bar{c}_\gamma \frac{s_W^4}{c_W^2} \right]$
$g_{haz}^{(1)}$	$c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma}$	$\frac{gs_W}{c_W m_W} \left[ \bar{c}_{HW} - \bar{c}_{HB} + 8\bar{c}_\gamma s_W^2 \right]$
$\tilde{g}_{haz}$	$\frac{gs_W}{c_W m_W}$	$\left[ \bar{c}_{HW} - \bar{c}_{HB} + 8\bar{c}_\gamma s_W^2 \right]$
		$\left[ \bar{c}_{HW} - \bar{c}_{HB} - \bar{c}_B + \bar{c}_W \right]$

Two approaches equivalent as they can be mapped into one another.

$$\begin{aligned}V_{\mu\nu} &= \partial_\mu V_\nu - \partial_\nu V_\mu \quad (V = A, Z, W^\pm), \quad V_{\mu\nu} = \frac{1}{2} \epsilon_{\mu\nu\rho\sigma} V^{\rho\sigma} \\ G_{\mu\nu}^a &= \partial_\mu G_\nu^a - \partial_\nu G_\mu^a + g_s f^{abc} G_\mu^b G_\nu^c,\end{aligned}$$

$g_{hww}^{(2)}$	$\frac{1}{\Lambda} c_\alpha \kappa_{H\theta W}$	$\frac{g}{m_W} \bar{c}_W$
		$\left[ \bar{c}_W + \bar{c}_{HW} \right]$

# Higgs Characterisation model

- We implemented an effective Lagrangian featuring bosons
$$X(J^P=0^+, 0^-, 1^+, 1^-, 2^+)$$
in FeynRules.
- any-process, any-decay, any-observable
- Equally useful for theorists (it can be systematically improved, changed - easily) and experimentalists (MC's event generation directly in analyses).
- Adaptable to the present/future analyses and accuracy targets.

The parametrization is based on the recent work [Englert, Goncalves-Netto, KM, Plehn (2013)].

# Effective Lagrangian -- spin0

$$\mathcal{L}_0^f = - \sum_{f=t,b,\tau} \bar{\psi}_f (c_\alpha \kappa_{Hff} g_{Hff} + i s_\alpha \kappa_{Aff} g_{Aff} \gamma_5) \psi_f X_0$$

$$\begin{aligned} \mathcal{L}_0^V = & \left\{ c_\alpha \kappa_{\text{SM}} \left[ \frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \right. \\ & - \frac{1}{4} [c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma} A_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma} A_{\mu\nu} \tilde{A}^{\mu\nu}] \\ & - \frac{1}{2} [c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma} Z_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{AZ\gamma} g_{AZ\gamma} Z_{\mu\nu} \tilde{A}^{\mu\nu}] \\ & - \frac{1}{4} [c_\alpha \kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + s_\alpha \kappa_{Agg} g_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}] \\ & - \frac{1}{4} \frac{1}{\Lambda} [c_\alpha \kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + s_\alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu}] \\ & - \frac{1}{2} \frac{1}{\Lambda} [c_\alpha \kappa_{HWW} W_\mu^+ W^{-\mu\nu} + s_\alpha \kappa_{AWW} W_\mu^+ \tilde{W}^{-\mu\nu}] \\ & - \frac{1}{\Lambda} c_\alpha [\kappa_{H\theta\gamma} Z_\nu \partial_\mu A^{\mu\nu} + \kappa_{H\theta Z} Z_\nu \partial_\mu Z^{\mu\nu} \\ & \quad + (\kappa_{H\theta W} W_\nu^+ \partial_\mu W^{-\mu\nu} + h.c.)] \Big\} X_0 \end{aligned}$$

parameter	description
$\Lambda$ [GeV]	cutoff scale
$c_\alpha$ ( $\equiv \cos \alpha$ )	mixing between $0^+$ and $0^-$
$\kappa_i$	dimensionless coupling parameter

```
#####
## INFORMATION FOR FRBLOCK
#####
Block frblock
1 1.000000e+03 # Lambda
2 1.000000e+00 # ca
3 1.000000e+00 # kSM
4 1.000000e+00 # kHtt
5 1.000000e+00 # kAtt
6 1.000000e+00 # kHbb
7 1.000000e+00 # kAbb
8 1.000000e+00 # kHll
9 1.000000e+00 # kAll
10 1.000000e+00 # kHaa
11 1.000000e+00 # kAaa
12 1.000000e+00 # kHz
13 1.000000e+00 # kAza
14 1.000000e+00 # kHgg
15 1.000000e+00 # kAgg
16 0.000000e+00 # kHzz
17 0.000000e+00 # kAzz
18 0.000000e+00 # kHww
19 0.000000e+00 # kAww
20 0.000000e+00 # kHda
21 0.000000e+00 # kHdz
22 0.000000e+00 # kHdwR
23 0.000000e+00 # kHdwI
```

# Effective Lagrangian -- spin0

$$\mathcal{L}_0^f = - \sum_{f=t,b,\tau} \bar{\psi}_f (c_\alpha \kappa_{Hff} g_{Hff} + i s_\alpha \kappa_{Aff} g_{Aff} \gamma_5) \psi_f X_0$$

$$\begin{aligned} \mathcal{L}_0^V = & \left\{ c_\alpha \kappa_{\text{SM}} \left[ \frac{1}{2} g_{HZZ} Z_\mu Z^\mu + g_{HWW} W_\mu^+ W^{-\mu} \right] \right. \\ & - \frac{1}{4} [c_\alpha \kappa_{H\gamma\gamma} g_{H\gamma\gamma} A_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{A\gamma\gamma} g_{A\gamma\gamma} A_{\mu\nu} \tilde{A}^{\mu\nu}] \\ & - \frac{1}{2} [c_\alpha \kappa_{HZ\gamma} g_{HZ\gamma} Z_{\mu\nu} A^{\mu\nu} + s_\alpha \kappa_{AZ\gamma} g_{AZ\gamma} Z_{\mu\nu} \tilde{A}^{\mu\nu}] \\ & - \frac{1}{4} [c_\alpha \kappa_{Hgg} g_{Hgg} G_{\mu\nu}^a G^{a,\mu\nu} + s_\alpha \kappa_{Agg} g_{Agg} G_{\mu\nu}^a \tilde{G}^{a,\mu\nu}] \\ & \left. - \frac{1}{4} \frac{1}{\Lambda} [c_\alpha \kappa_{HZZ} Z_{\mu\nu} Z^{\mu\nu} + s_\alpha \kappa_{AZZ} Z_{\mu\nu} \tilde{Z}^{\mu\nu}] \right. \\ & - \frac{1}{2} \frac{1}{\Lambda} [c_\alpha \kappa_{HWW} W_\mu^+ W^{-\mu\nu} + s_\alpha \kappa_{AWW} W_\mu^+ \tilde{W}^{-\mu\nu}] \\ & - \frac{1}{\Lambda} c_\alpha [\kappa_{H\theta\gamma} Z_\nu \partial_\mu A^{\mu\nu} + \kappa_{H\theta Z} Z_\nu \partial_\mu Z^{\mu\nu} \\ & \quad + (\kappa_{H\theta W} W_\nu^+ \partial_\mu W^{-\mu\nu} + h.c.)] \Big\} X_0 \end{aligned}$$

Dimensionful **couplings g** are set as internal parameters so as to reproduce a **SM Higgs** for  $\kappa=1$ .

$g_{Xyy'} \times v$	$ff$	$ZZ/WW$	$\gamma\gamma$	$Z\gamma$	$gg$
$H$	$m_f$	$2m_{Z/W}^2$	$47\alpha_{\text{EM}}/18\pi$	$C(94\cos^2\theta_W - 13)/9\pi$	$-\alpha_s/3\pi$
$A$	$m_f$	0	$4\alpha_{\text{EM}}/3\pi$	$2C(8\cos^2\theta_W - 5)/3\pi$	$\alpha_s/2\pi$

HiggsCharacterisation – FeynRules

http://feynrules.irmp.ucl.ac.be/wiki/HiggsCharacterisation

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HiggsCharacterisation – FeynRules

## The Higgs Characterisation model

### Authors

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### Description of the model

This is a model file for the characterisation of the boson recently discovered at the LHC. Our effective lagrangian consists of the SM (except for the Higgs itself), expressed through the physical degrees of freedom present below the EWSB scale, plus a new bosonic state  $X(J^P)$  with spin/parity assignments  $J^P = 0^+, 0^-, 1^+, 1^-,$  or  $2^+$ . The new state can couple to SM particles via interactions of the lowest possible dimensions. In addition, the state  $0^+$  is allowed to mix with the  $0^-$  one, and can interact with SM particles with higher-dimensional operators beyond those of the SM. See more details in

- ↗ [1306.6464](#) : P. Artoisenet, P. de Aquino, F. Demartin, R. Frederix, S. Frixione, F. Maltoni, M. K. Mandal, P. Mathews, K. Mawatari, V. Ravindran, S. Seth, P. Torrielli, M. Zaro, "A framework for Higgs characterisation" (*JHEP11(2013)043*).
- ↗ [1307.5607](#) : P. de Aquino, K. Mawatari, "Characterising a Higgs-like resonance at the LHC" (*Proceedings for HPNP2013*).
- ↗ [1311.1829](#) : F. Maltoni, K. Mawatari, M. Zaro, "Higgs characterisation via vector-boson fusion and associated production: NLO and parton-shower effects" (*EPJC74(2014)2710*).

### Model files for LO

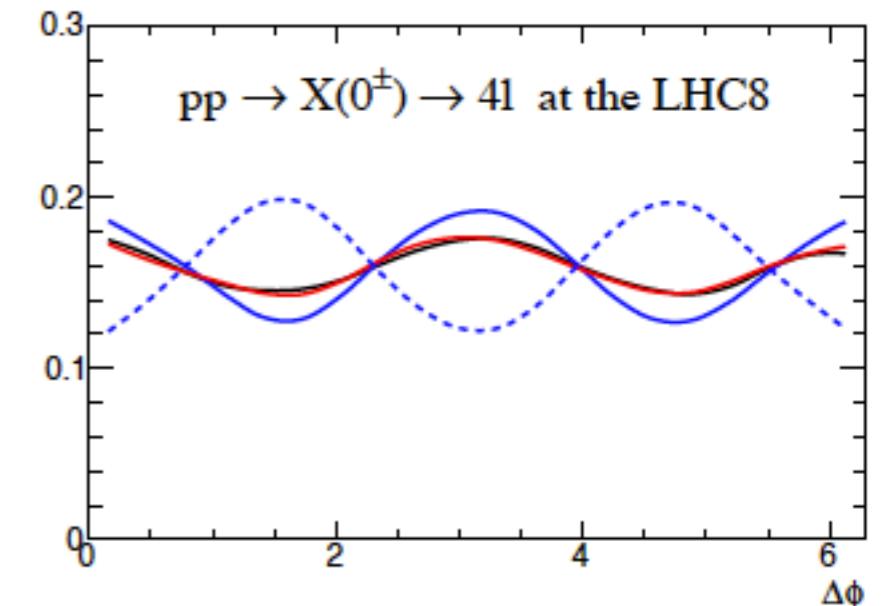
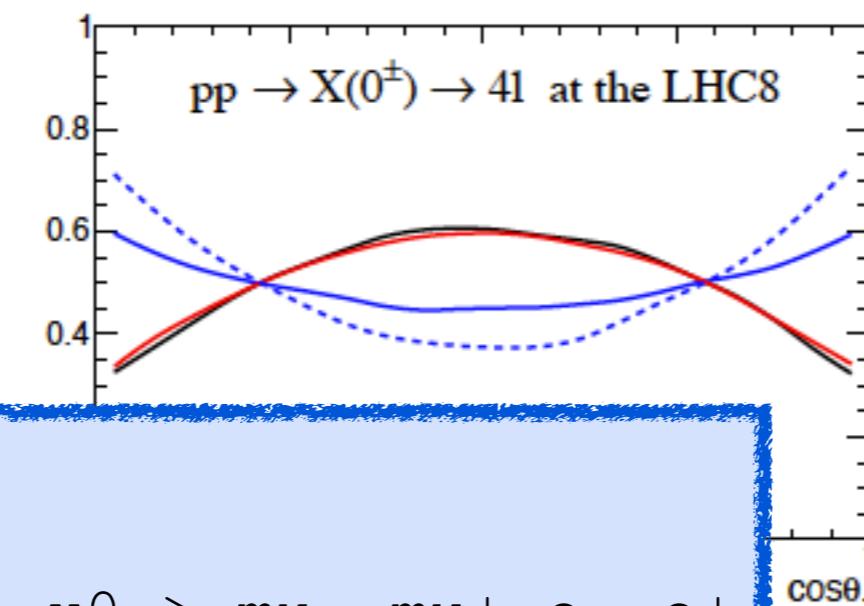
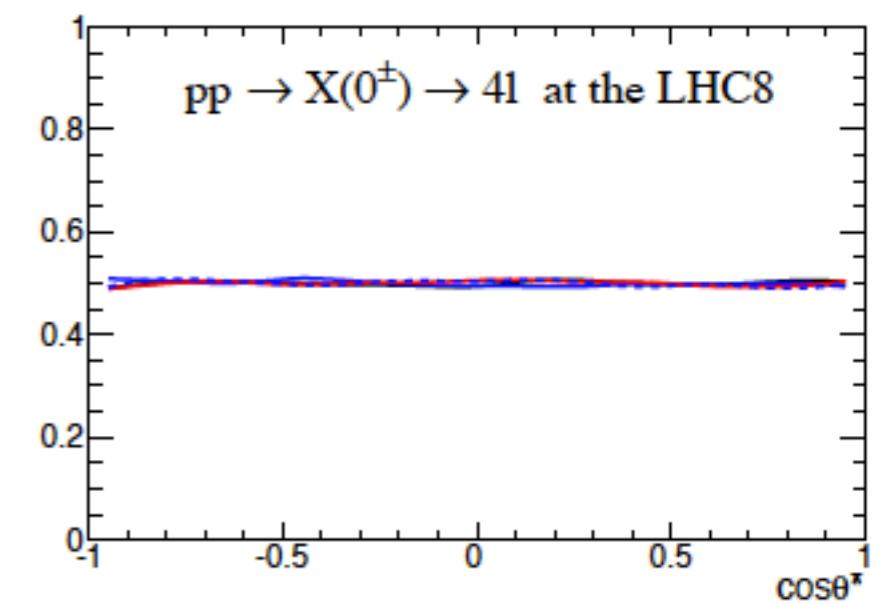
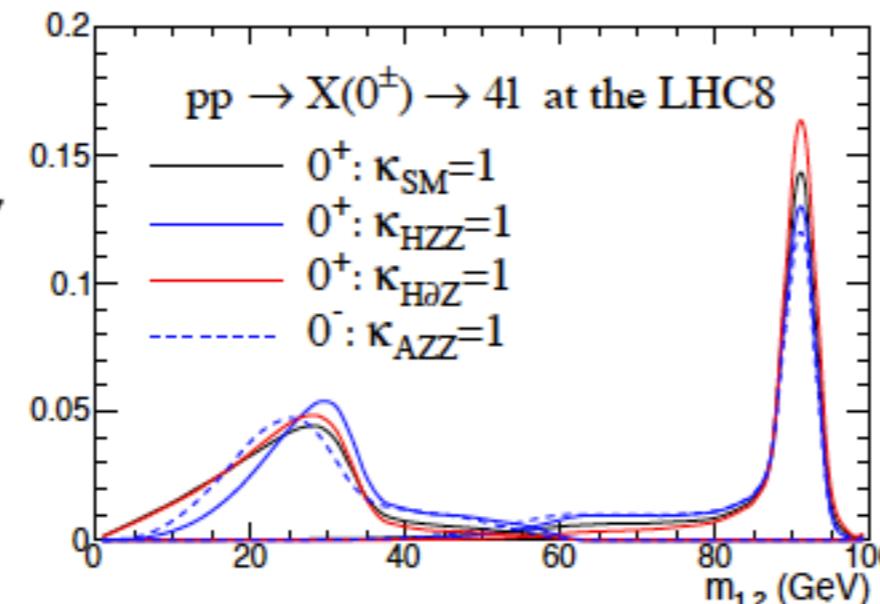
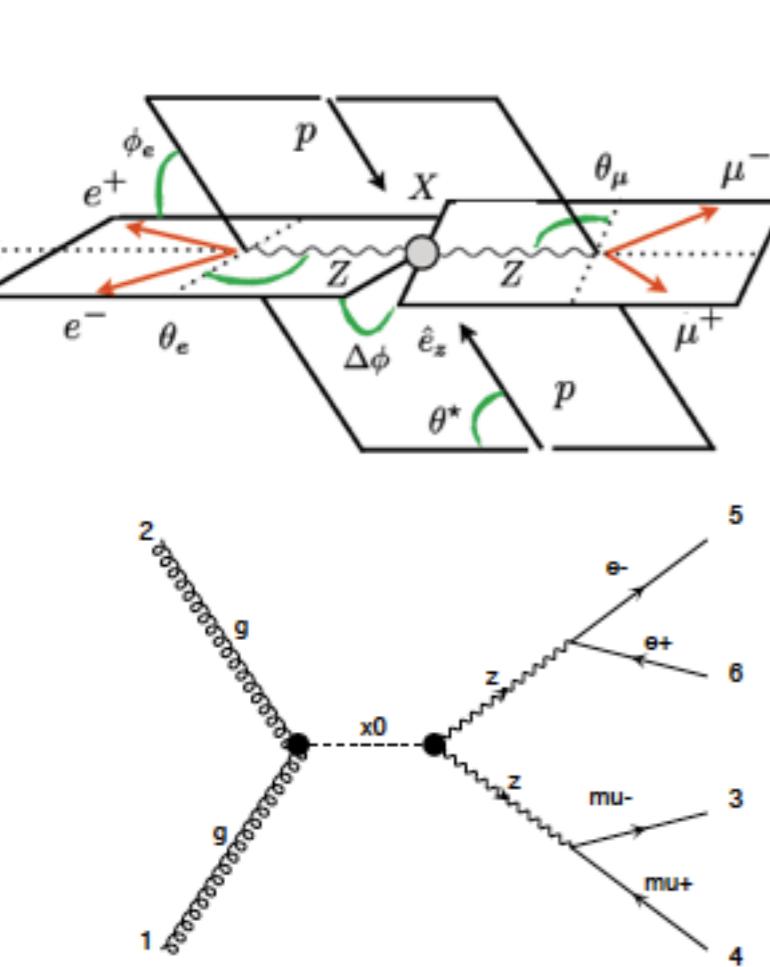
- [HC.fr](#) : the main model file.
- [SM\\_HC.fr](#) : This model requires the modified Standard Model implementation of [FeynRules](#).
- [Massless.rst](#), [Cabibbo.rst](#) : SM restriction files.
- [HC.nb](#) : this is an example Mathematica notebook that loads the model, calculates the Feynman rules and extract the model files within the UFO format.
- [HC\\_UFO.zip](#) : The model files in UFO format (for MadGraph5).

### Model files for NLO (only for the $J=0$ case)

A few remarks before use; see the README file in the model.

- [HC\\_NLO\\_X0\\_UFO.zip](#) : The model files in UFO format (for MadGraph5).

# Mass and angular distributions -- spin0



```
./bin/mg5_aMC
>import model HC
>generate p p > x0, x0 > mu- mu+ e- e+
>launch
```

# Effective Lagrangian -- spin 1

- The most general interactions at the lowest canonical dimension:

$$\mathcal{L}_1^f = \sum_{f=q,\ell} \bar{\psi}_f \gamma_\mu (\kappa_{f_a} a_f - \kappa_{f_b} b_f \gamma_5) \psi_f X_1^\mu$$

$$\begin{aligned} \mathcal{L}_1^W = & i\kappa_{W_1} g_{WWZ} (W_{\mu\nu}^+ W^{-\mu} - W_{\mu\nu}^- W^{+\mu}) X_1^\nu + i\kappa_{W_2} g_{WWZ} W_\mu^+ W_\nu^- X_1^{\mu\nu} \\ & - \kappa_{W_3} W_\mu^+ W_\nu^- (\partial^\mu X_1^\nu + \partial^\nu X_1^\mu) \\ & + i\kappa_{W_4} W_\mu^+ W_\nu^- \tilde{X}_1^{\mu\nu} - \kappa_{W_5} \epsilon_{\mu\nu\rho\sigma} [W^{+\mu} (\partial^\rho W^{-\nu}) - (\partial^\rho W^{+\mu}) W^{-\nu}] X_1^\sigma \end{aligned}$$

$$\mathcal{L}_1^Z = -\kappa_{Z_1} Z_{\mu\nu} Z^\mu X_1^\nu - \kappa_{Z_3} X_1^\mu (\partial^\nu Z_\mu) Z_\nu - \kappa_{Z_5} \epsilon_{\mu\nu\rho\sigma} X_1^\mu Z^\nu (\partial^\rho Z^\sigma)$$

- Parity conservation implies that
  - ▶ for  $X_1^-$      $\kappa_{f_b} = \kappa_{V_4} = \kappa_{V_5} = 0$
  - ▶ for  $X_1^+$      $\kappa_{f_a} = \kappa_{V_1} = \kappa_{V_2} = \kappa_{V_3} = 0$

# Effective Lagrangian -- spin2

- via the energy-momentum tensor of the SM fields, starting from D5:

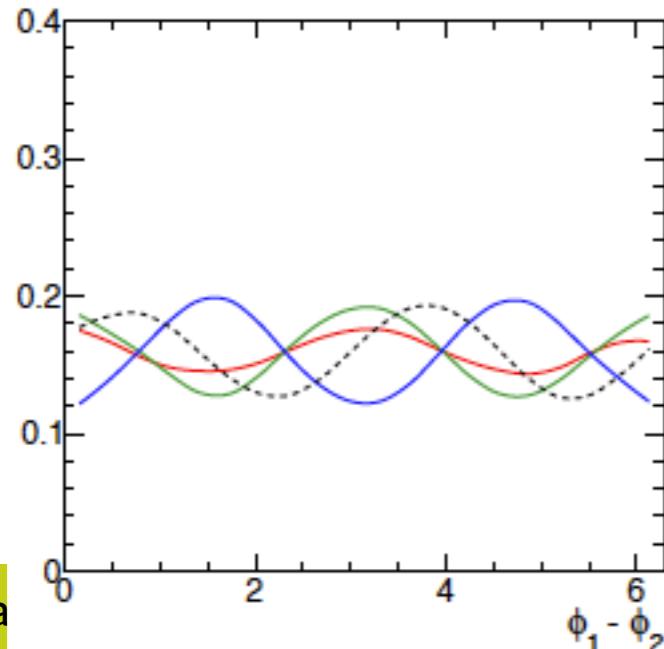
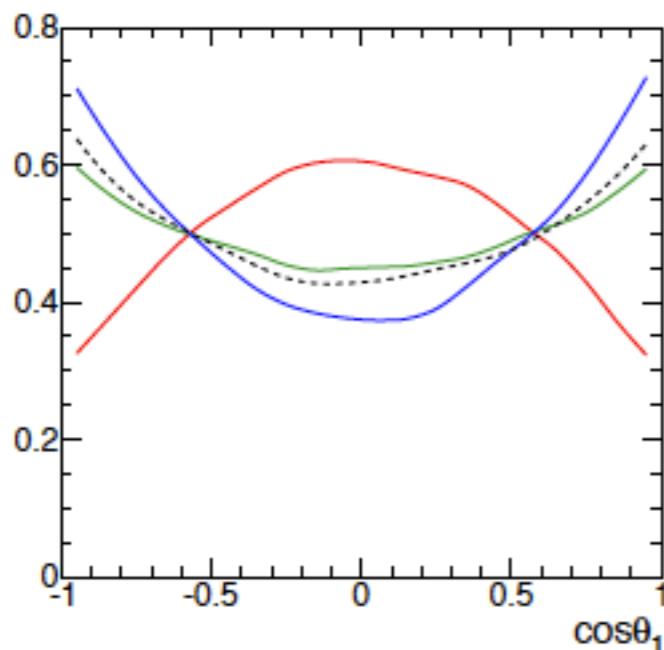
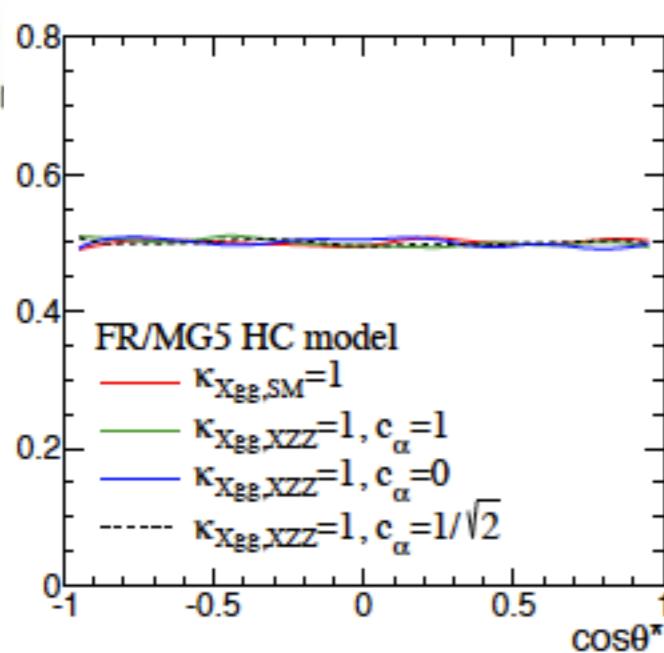
$$\mathcal{L}_2^f = -\frac{1}{\Lambda} \sum_{f=q,\ell} \kappa_f T_{\mu\nu}^f X_2^{\mu\nu}$$
$$\mathcal{L}_2^V = -\frac{1}{\Lambda} \sum_{V=Z,W,\gamma,g} \kappa_V T_{\mu\nu}^V X_2^{\mu\nu}$$

► The E-M tensor for QED:

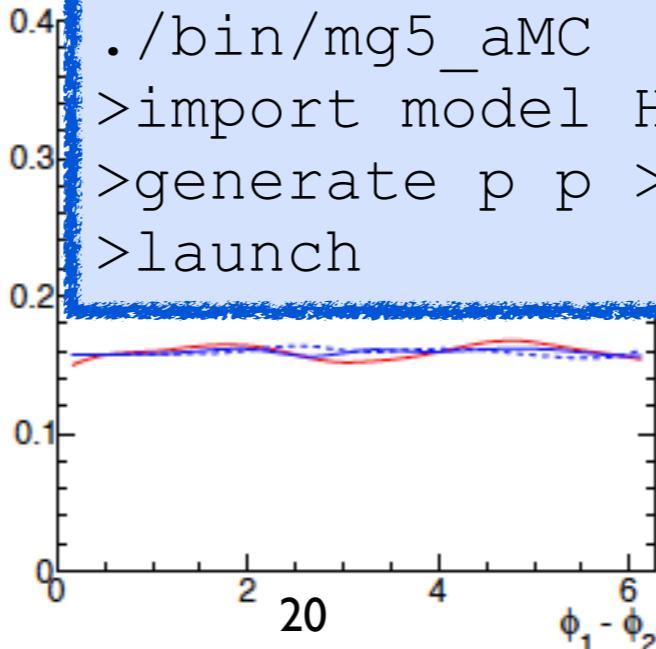
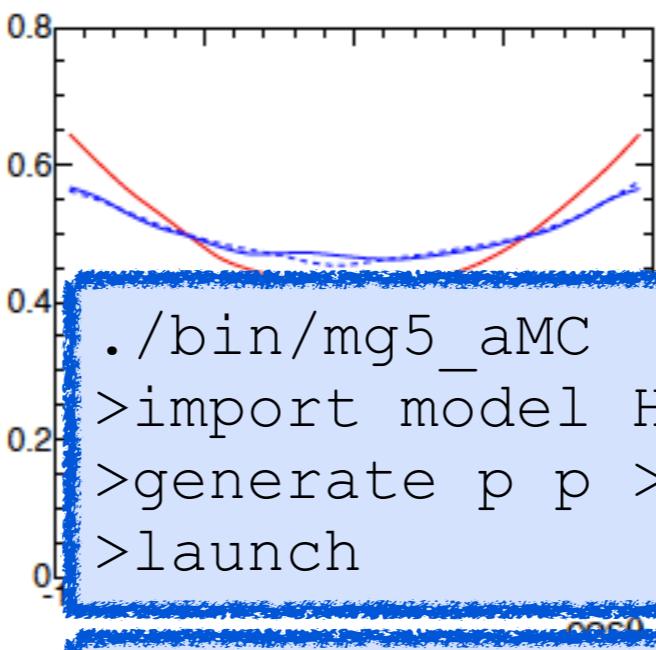
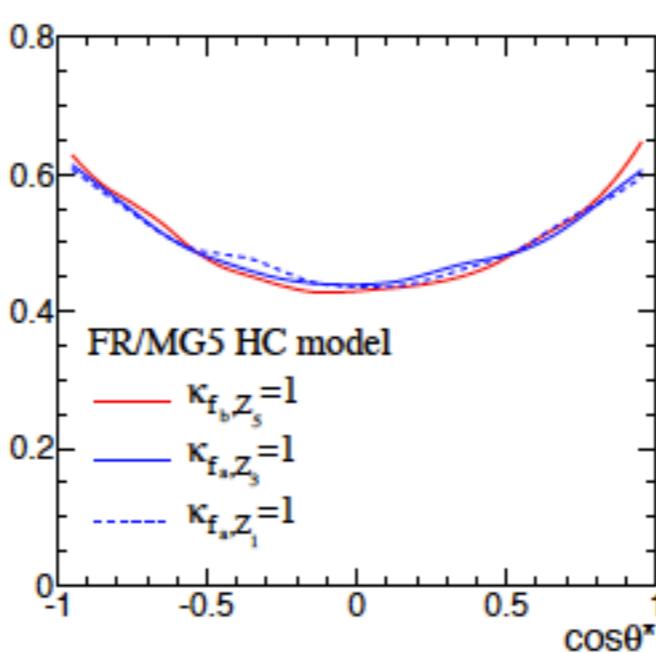
$$T_{\mu\nu}^f = -g_{\mu\nu} \left[ \bar{\psi}_f (i\gamma^\rho D_\rho - m_f) \psi_f - \frac{1}{2} \partial^\rho (\bar{\psi}_f i\gamma_\rho \psi_f) \right] + \left[ \frac{1}{2} \bar{\psi}_f i\gamma_\mu D_\nu \psi_f - \frac{1}{4} \partial_\mu (\bar{\psi}_f i\gamma_\nu \psi_f) + (\mu \leftrightarrow \nu) \right],$$
$$T_{\mu\nu}^\gamma = -g_{\mu\nu} \left[ -\frac{1}{4} A^{\rho\sigma} A_{\rho\sigma} + \partial^\rho \partial^\sigma A_\sigma A_\rho + \frac{1}{2} (\partial^\rho A_\rho)^2 \right] - A_\mu^\rho A_{\nu\rho} + \partial_\mu \partial^\rho A_\rho A_\nu + \partial_\nu \partial^\rho A_\rho A_\mu,$$



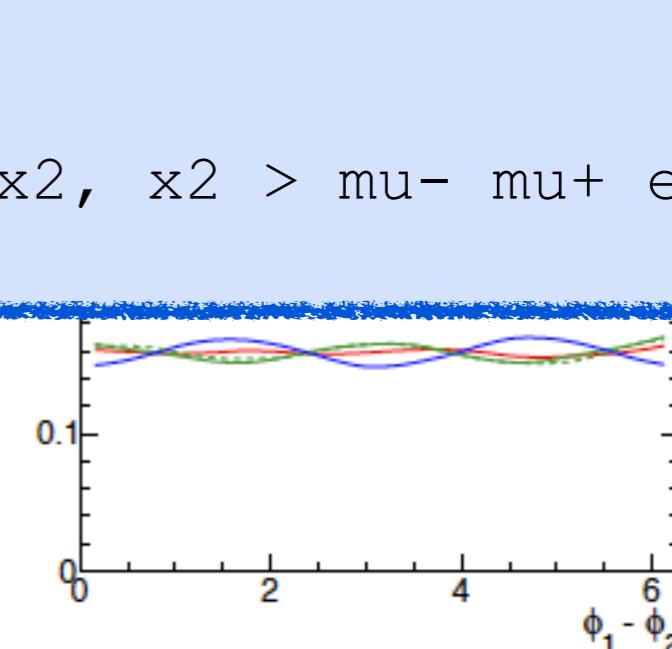
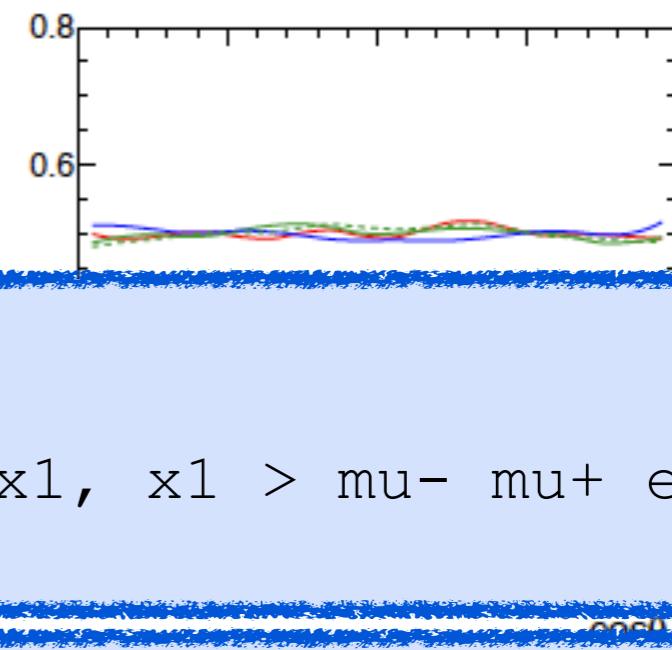
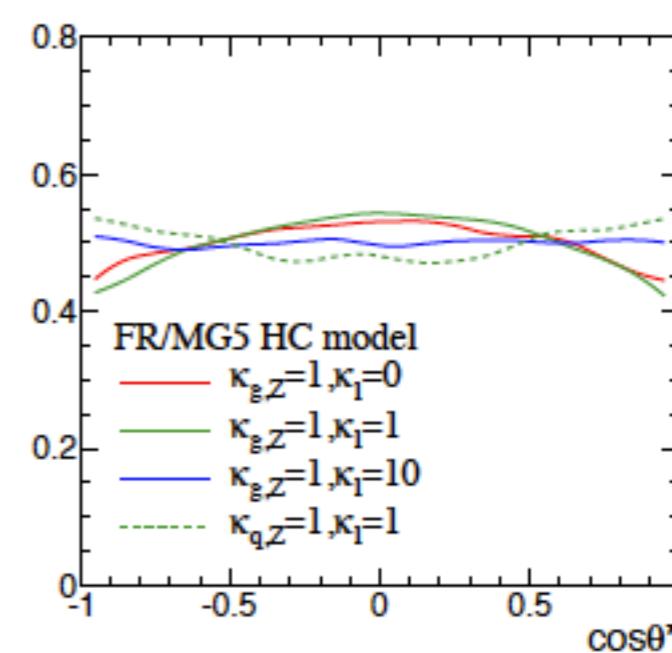
spin-0



spin-1



spin-2

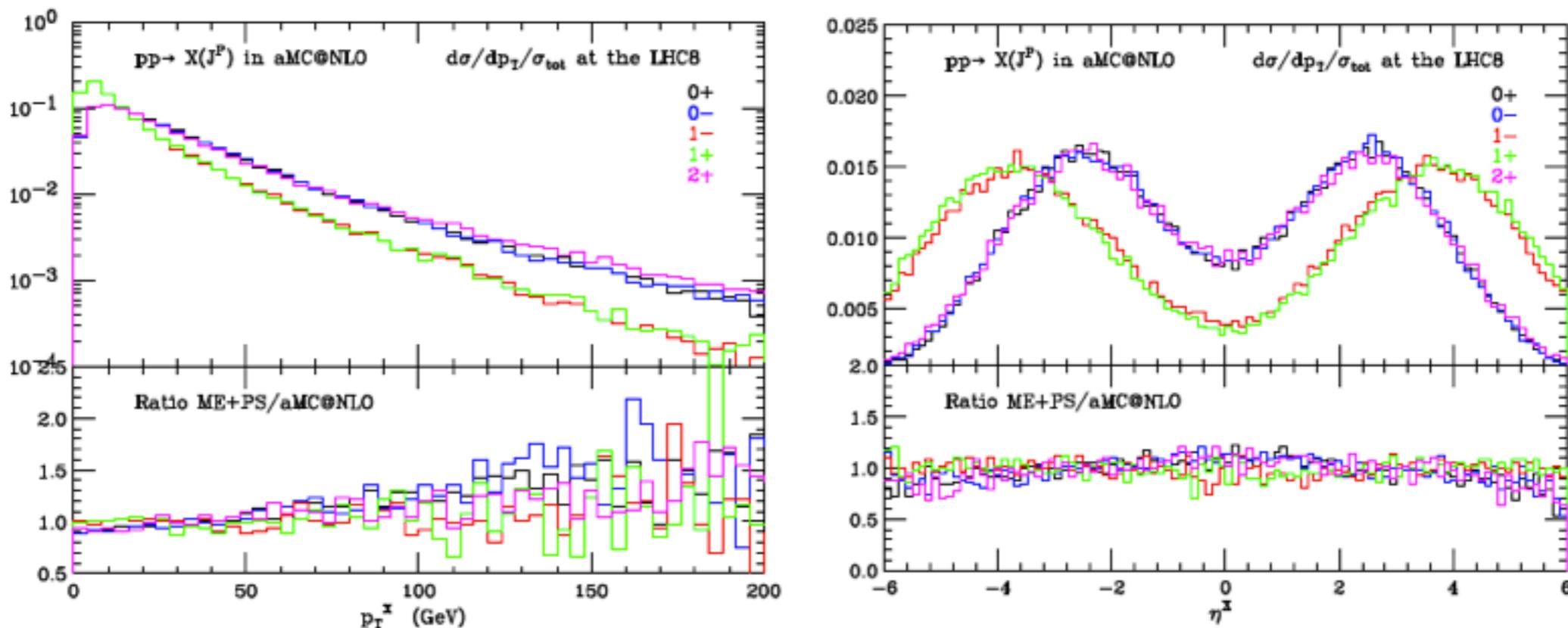


```
./bin/mg5_aMC  
>import model HC  
>generate p p > x1, x1 > mu- mu+ e- e+  
>launch
```

```
./bin/mg5_aMC  
>import model HC  
>generate p p > x2, x2 > mu- mu+ e- e+  
>launch
```

# Higher order effects in QCD

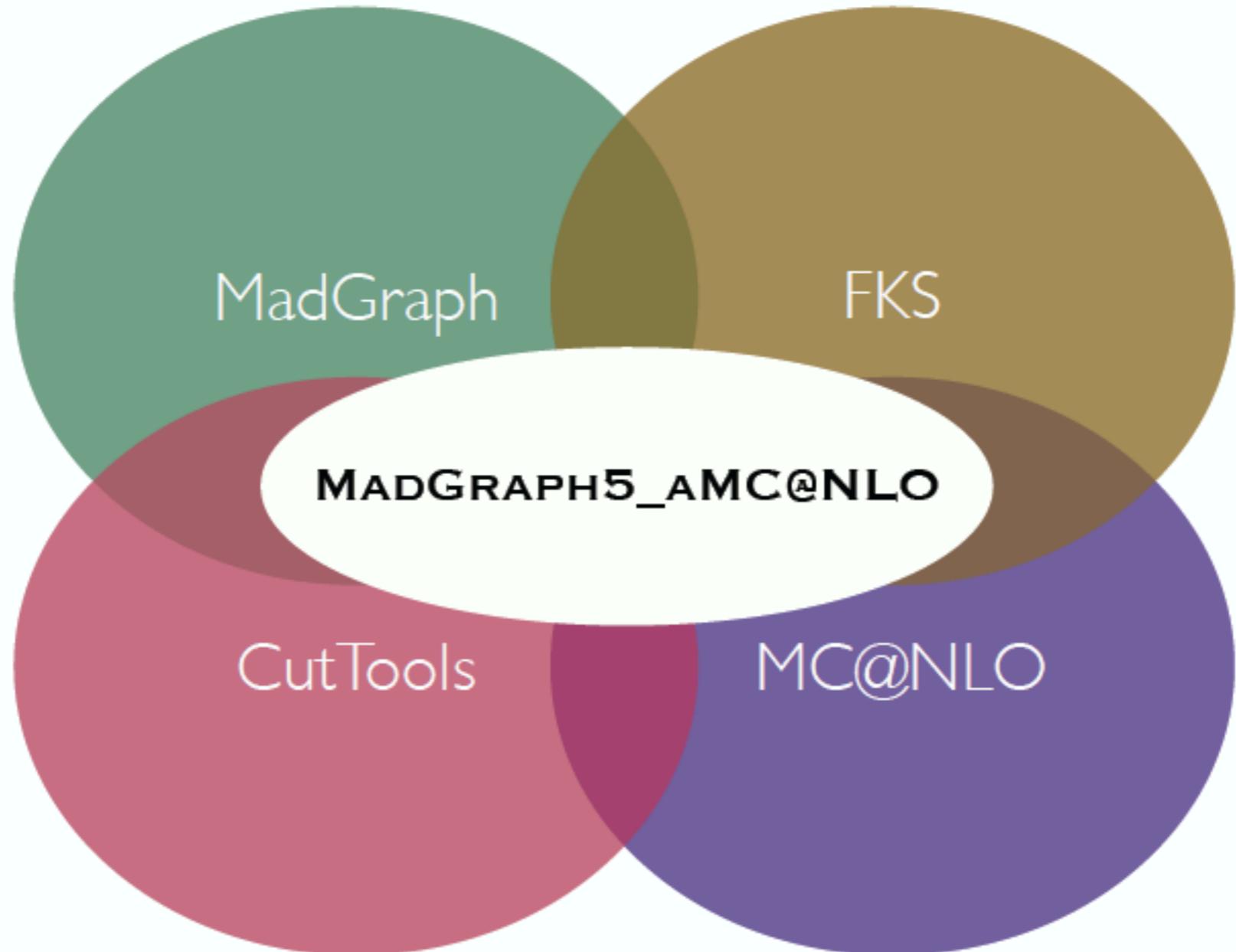
- The LO predictions can be systematically improved by including the effects due to the emission of QCD partons.
  - ▶ LO Matrix-Element/Parton-Shower merging [ME+PS]
  - ▶ full-NLO matrix element with parton-shower [aMC@NLO+Pythia/Herwig]



Good agreement between the ME+PS and aMC@NLO predictions for most observables.

# MADGRAPH5\_AMC@NLO

J. Alwall, R. Frederix, S. Frixione, F. Maltoni, O. Mattelaer, H. S. Shao, T. Stelzer, P. Torrielli, V. Hirschi, MZ  
*arXiv:1405.0301*



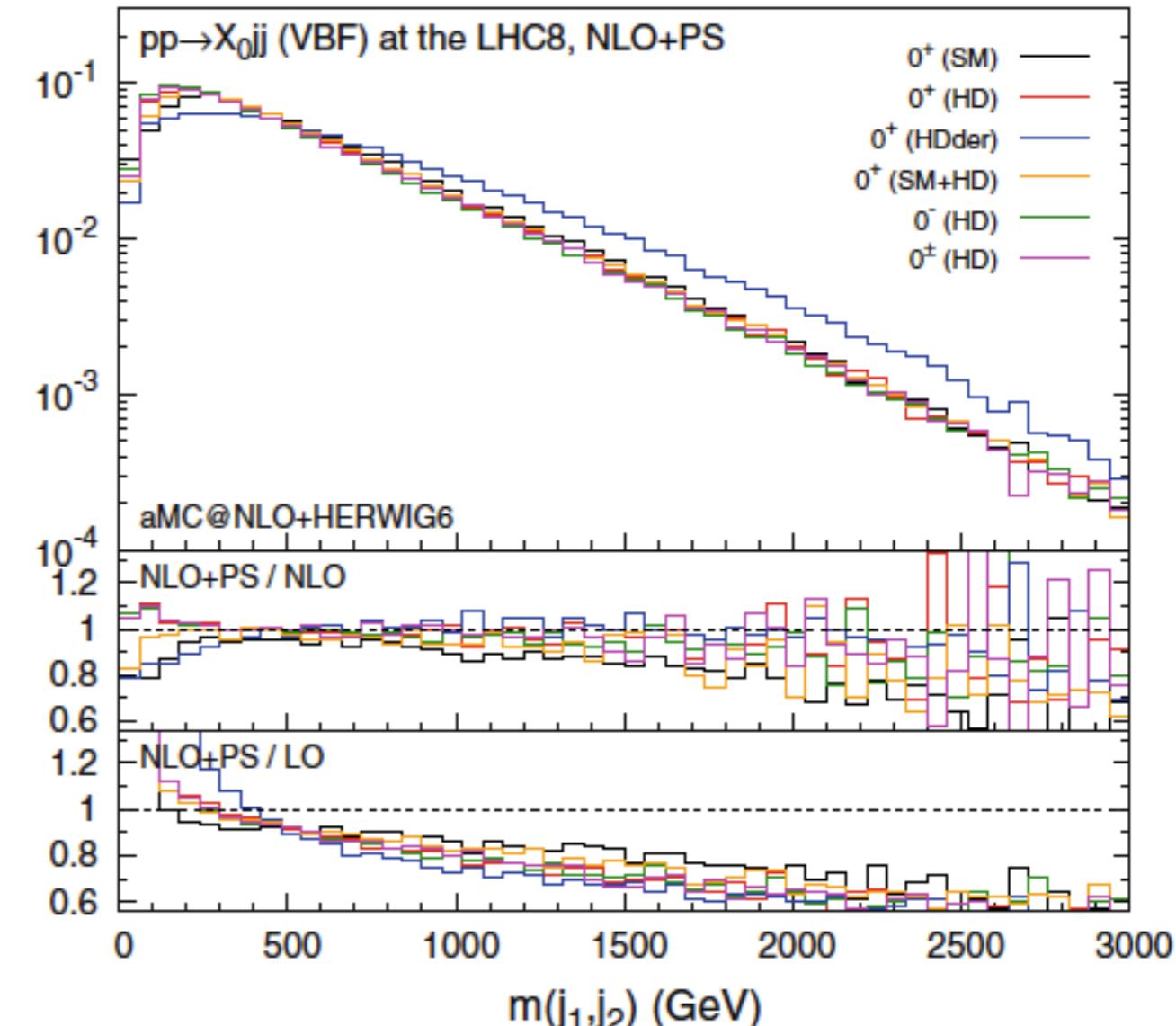


# Mjj distributions

```
./bin/mg5_aMC
>import model HC_NLO
>generate p p > x0 j j QCD=0 [QCD]
>launch
```

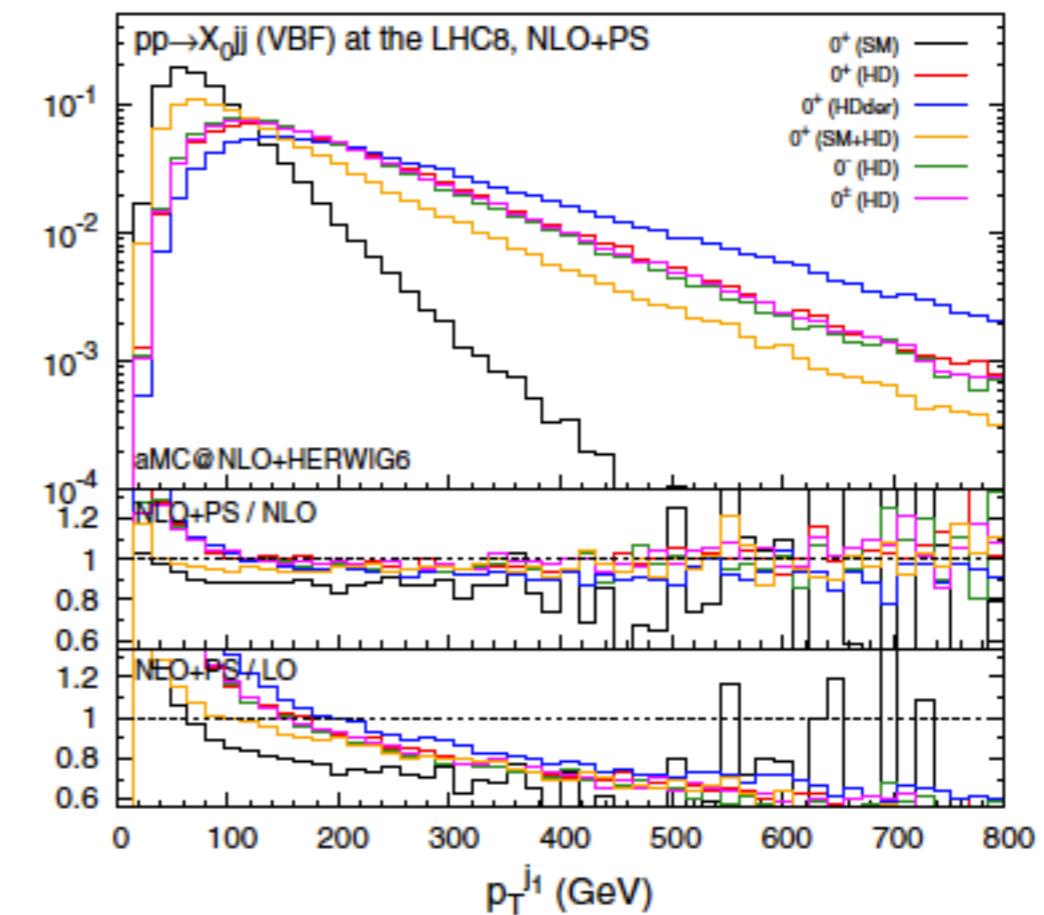
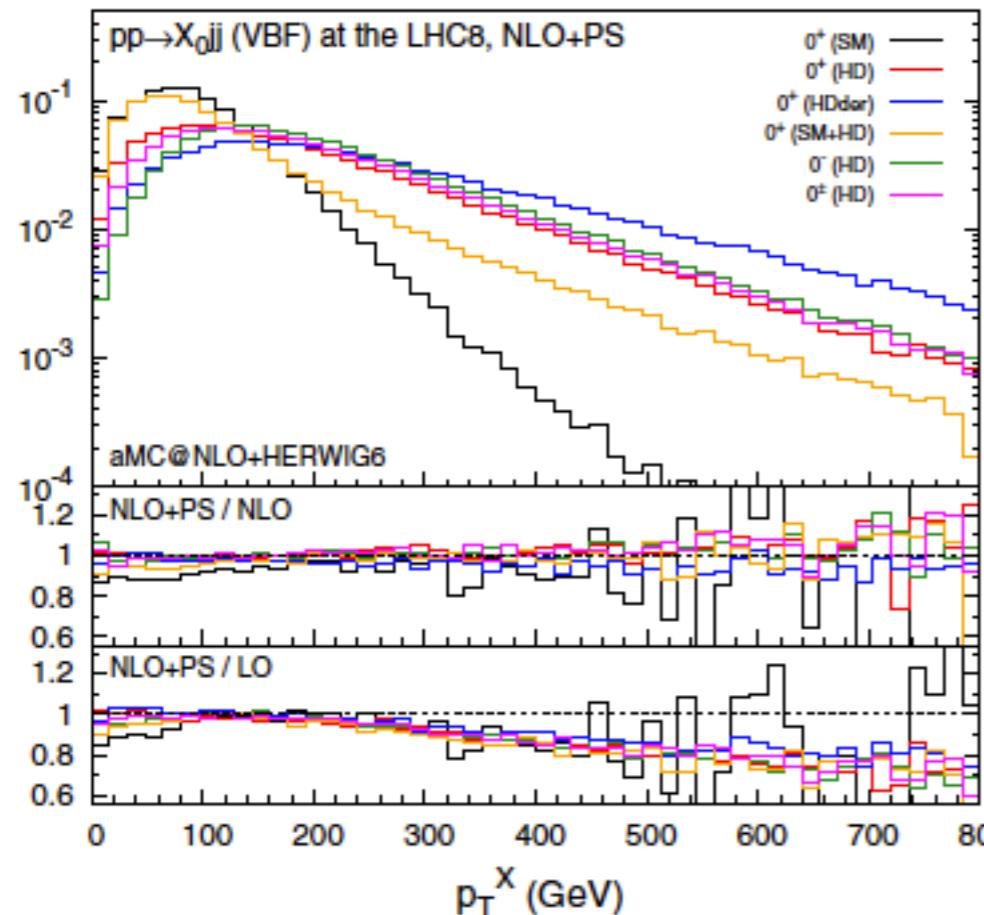
Scenario	HC parameter choice
0 <sup>+</sup> (SM)	$\kappa_{\text{SM}} = 1$ ( $c_\alpha = 1$ )
0 <sup>+</sup> (HD)	$\kappa_{HZZ, HWW} = 1$ ( $c_\alpha = 1$ )
0 <sup>+</sup> (HDder)	$\kappa_{H\partial Z, H\partial W} = 1$ ( $c_\alpha = 1$ )
0 <sup>+</sup> (SM+HD)	$\kappa_{SM, HZZ, HWW} = 1$ ( $c_\alpha = 1, \Lambda = v$ )
0 <sup>-</sup> (HD)	$\kappa_{AZZ, AWW} = 1$ ( $c_\alpha = 0$ )
0 <sup>±</sup> (HD)	$\kappa_{HZZ, AZZ, HWW, AWW} = 1$ ( $c_\alpha = 1/\sqrt{2}$ )

Scenario	$\sigma_{\text{LO}}$ (fb)	$\sigma_{\text{NLO}}$ (fb)	$K$
0 <sup>+</sup> (SM)	1509(1) $^{+4.7\%}_{-4.4\%}$	1633(2) $^{+2.0\%}_{-1.5\%}$	1.08
0 <sup>+</sup> (HD)	69.66(6) $^{+7.5\%}_{-6.6\%}$	67.08(13) $^{+2.2\%}_{-2.3\%}$	0.96
0 <sup>+</sup> (HDder)	721.9(6) $^{+11.0\%}_{-9.0\%}$	684.9(1.5) $^{+2.3\%}_{-2.8\%}$	0.95
0 <sup>+</sup> (SM+HD)	3065(2) $^{+5.6\%}_{-5.1\%}$	3144(5) $^{+1.6\%}_{-1.1\%}$	1.03
0 <sup>-</sup> (HD)	57.10(4) $^{+7.7\%}_{-6.7\%}$	55.24(11) $^{+2.1\%}_{-2.5\%}$	0.97
0 <sup>±</sup> (HD)	63.46(5) $^{+7.6\%}_{-6.7\%}$	61.07(13) $^{+2.3\%}_{-2.0\%}$	0.96



- The mjj distributions are all very similar (except the scenario with the derivative operator).
- The QCD corrections tend to make the tagging jets softer.

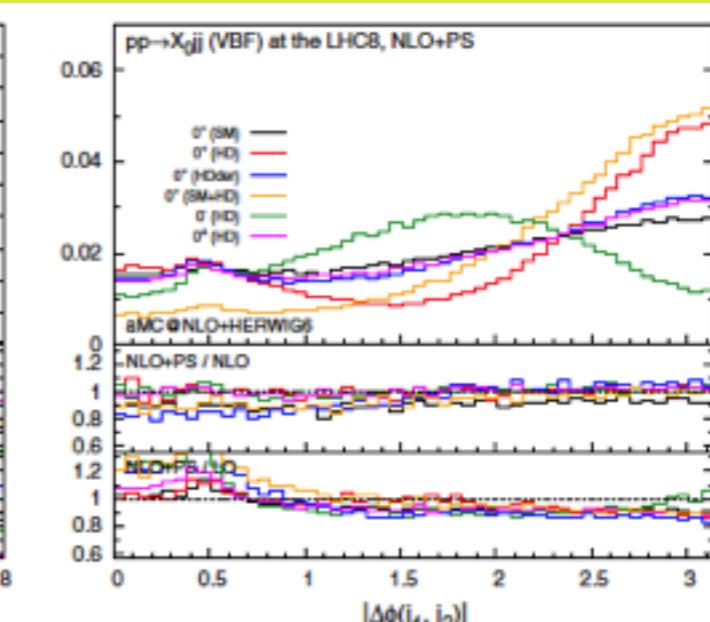
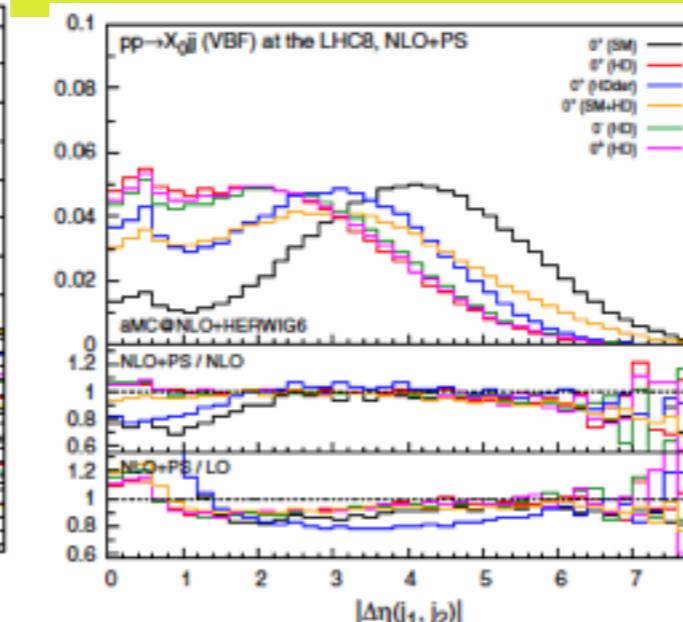
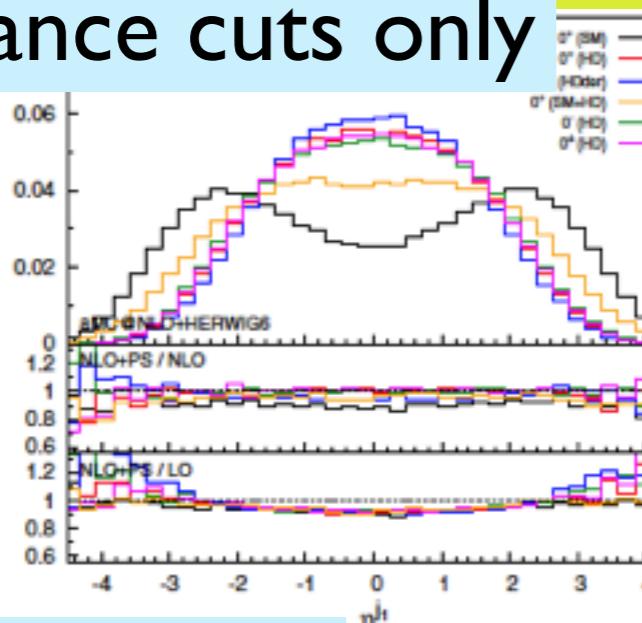
# pT distributions



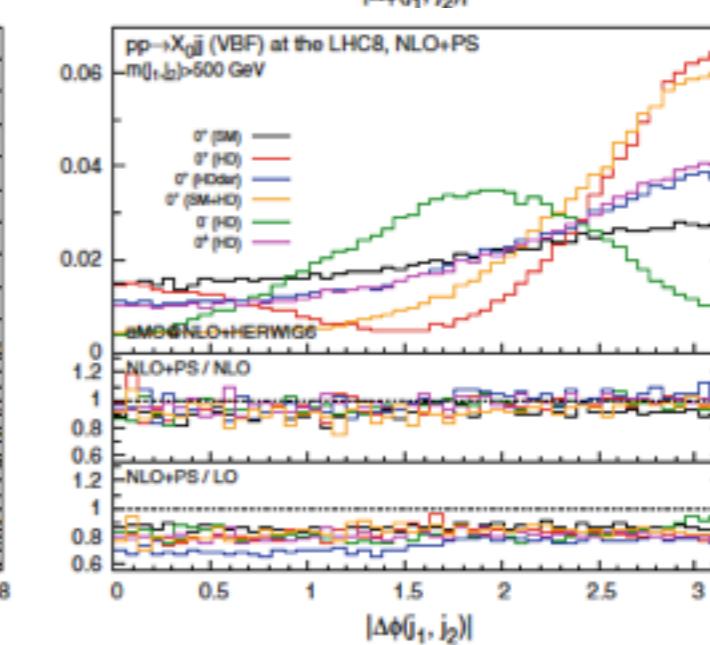
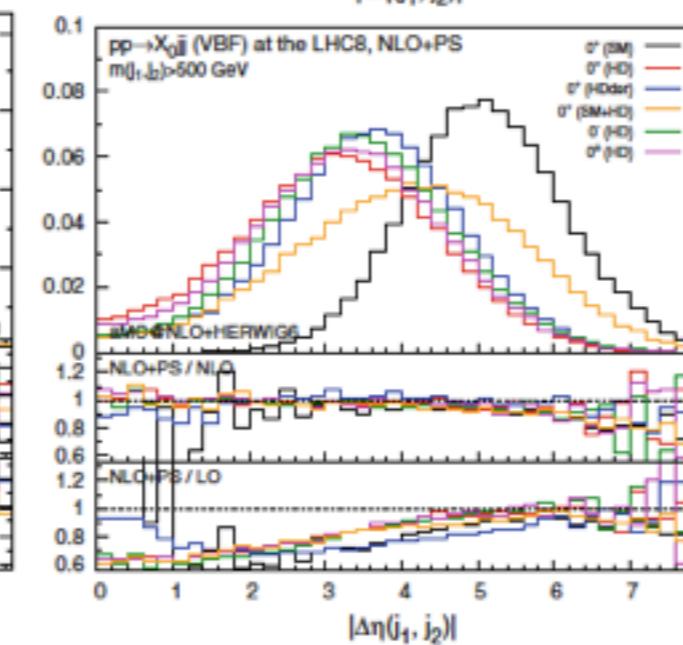
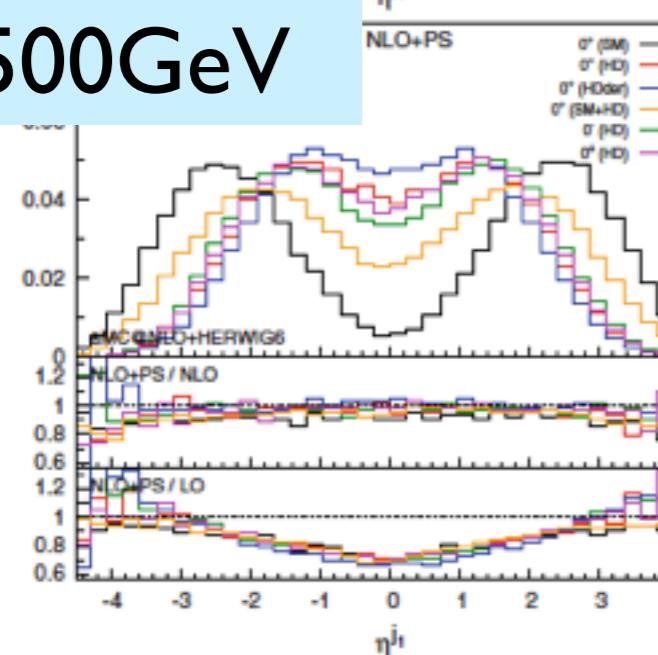
- The unitarity violating behavior of the HD interactions, especially HDder, clearly manifests itself.

# mjj cuts

acceptance cuts only



$m_{jj} > 500\text{GeV}$

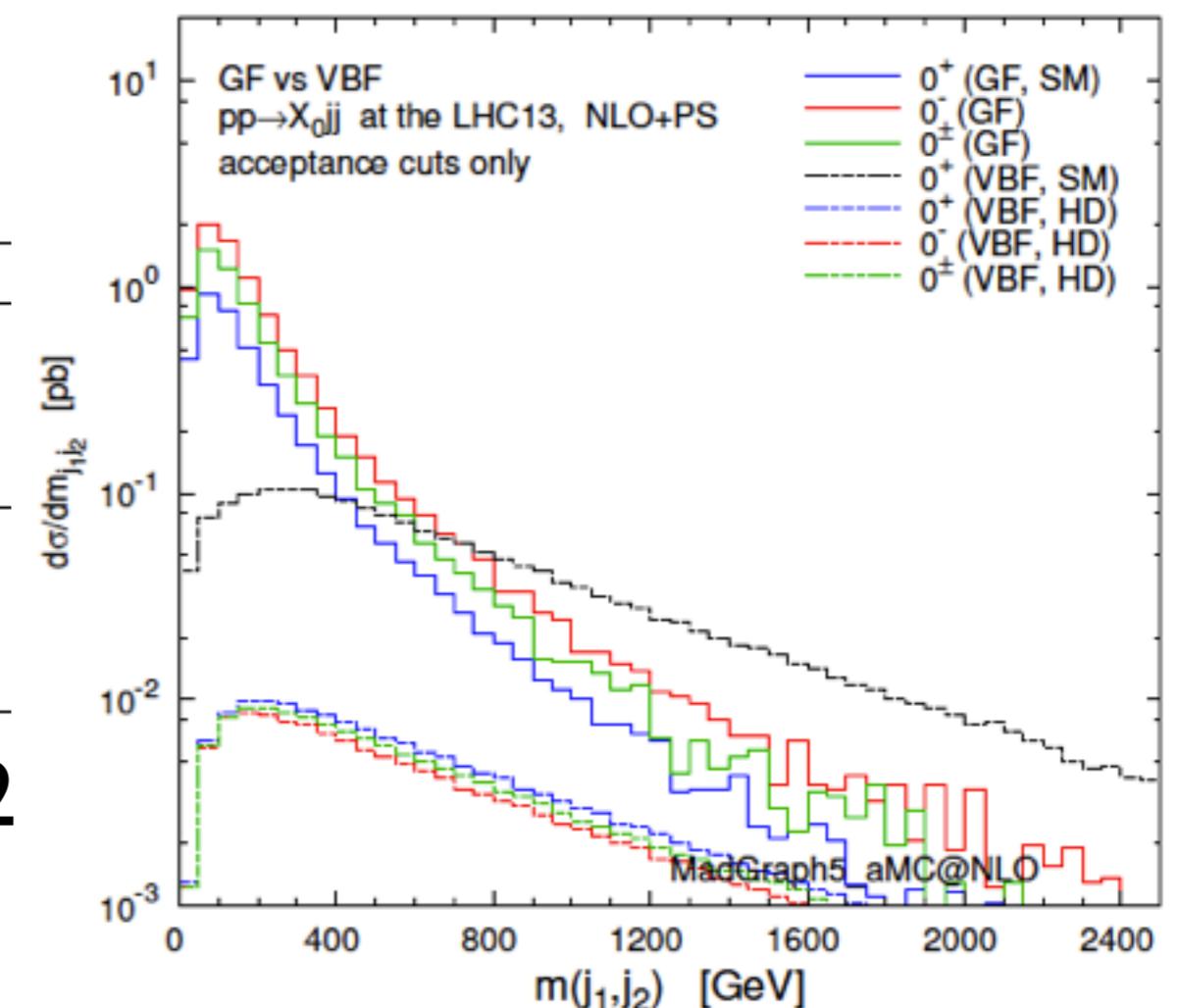


- The mjj cut effectively suppresses the central jet activity, especially for SM.
- The difference among the scenarios becomes more pronounced.
- NLO corrections cannot be described by an overall K factor, and also depends on the applied cuts.

# (weak) VBF vs. GF (plus 2jets)

scenario		$\sigma_{\text{LO}}$ (pb)	$\sigma_{\text{NLO}}$ (pb)
LHC 8 TeV	$0^+$	1.351(1) $^{+67.1}_{-36.8}$ $^{+4.3\%}_{-4.3\%}$	1.702(6) $^{+19.7}_{-20.8}$ $^{+1.7\%}_{-1.7\%}$
	$0^-$	2.951(3) $^{+67.2}_{-36.8}$ $^{+4.4\%}_{-4.4\%}$	3.660(15) $^{+19.1}_{-20.6}$ $^{+1.7\%}_{-1.7\%}$
	$0^\pm$	2.142(2) $^{+67.1}_{-36.8}$ $^{+4.4\%}_{-4.4\%}$	2.687(10) $^{+19.6}_{-20.8}$ $^{+1.7\%}_{-1.7\%}$
LHC 13 TeV	$0^+$	4.265(4) $^{+61.5}_{-34.9}$ $^{+3.3\%}_{-3.3\%}$	5.092(23) $^{+15.4}_{-17.9}$ $^{+1.2\%}_{-1.2\%}$
	$0^-$	9.304(9) $^{+61.6}_{-34.9}$ $^{+3.4\%}_{-3.4\%}$	11.29(4) $^{+16.0}_{-18.2}$ $^{+1.2\%}_{-1.2\%}$
	$0^\pm$	6.775(6) $^{+61.5}_{-34.9}$ $^{+3.3\%}_{-3.3\%}$	8.055(35) $^{+15.8}_{-18.2}$ $^{+1.2\%}_{-1.2\%}$

$K \sim 1.2$



# (weak) VBF vs. GF (plus 2jets)

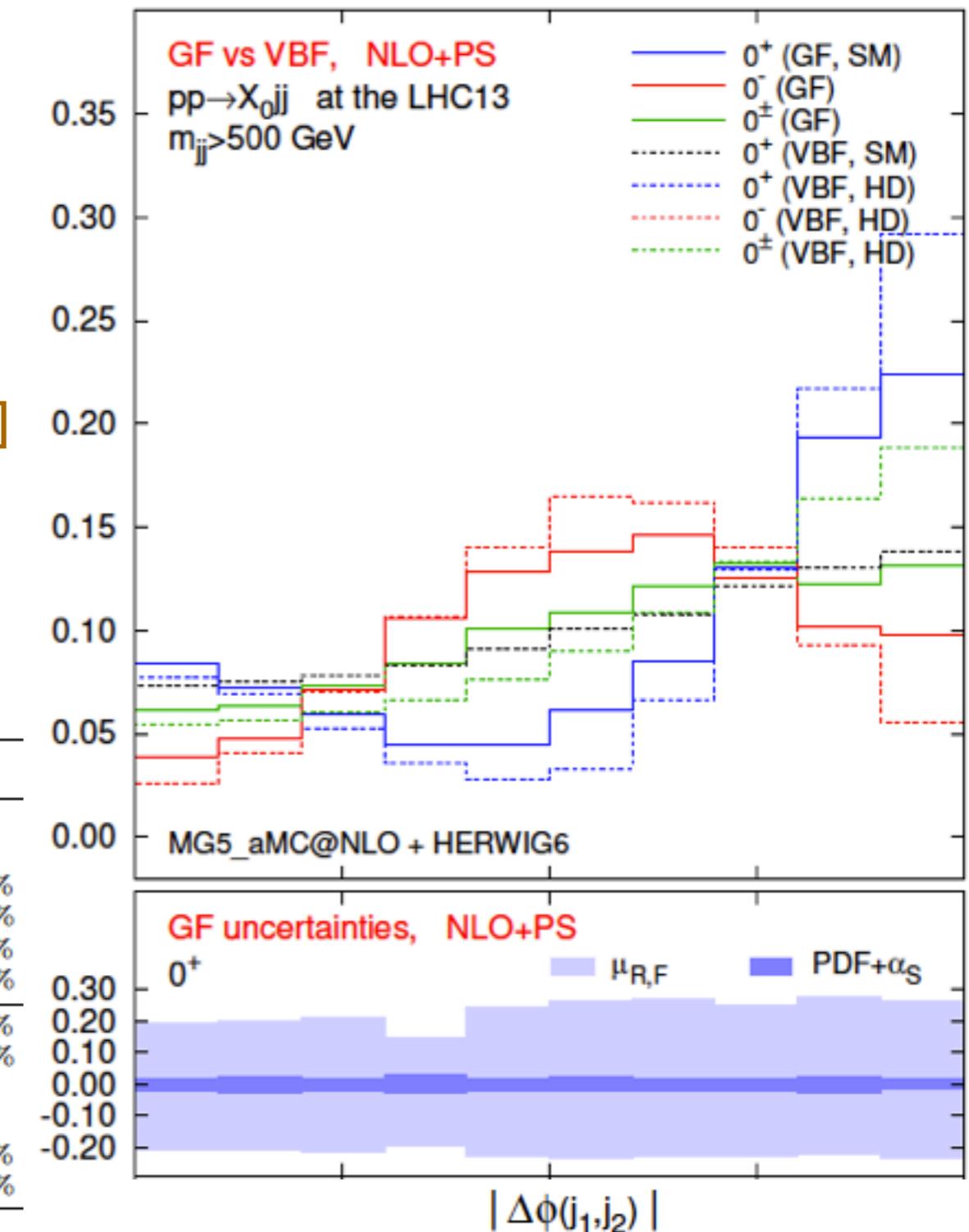
Maltoni, KM, Zaro [arXiv:1311.1829]

```
./bin/mg5_aMC
> import model HC_NLO
> generate p p > x0 j j QCD=0 [QCD]
> launch
```

Demartin, Maltoni, KM, Page, Pittau, Zaro [in progress]

```
./bin/mg5_aMC
> import model HC_NLO-heft
> generate p p > x0 j j / t [QCD]
> launch
```

	scenario	$\sigma_{\text{LO}}$ (pb)	$\sigma_{\text{NLO}}$ (pb)
LHC 8 TeV	$0^+$	1.351(1) $^{+67.1}_{-36.8}$ $^{+4.3\%}_{-4.3\%}$	1.702(6) $^{+19.7}_{-20.8}$ $^{+1.7\%}_{-1.7\%}$
	$0^-$	2.951(3) $^{+67.2}_{-36.8}$ $^{+4.4\%}_{-4.4\%}$	3.660(15) $^{+19.1}_{-20.6}$ $^{+1.7\%}_{-1.7\%}$
	$0^\pm$	2.142(2) $^{+67.1}_{-36.8}$ $^{+4.4\%}_{-4.4\%}$	2.687(10) $^{+19.6}_{-20.8}$ $^{+1.7\%}_{-1.7\%}$
LHC 13 TeV	$0^+$	4.265(4) $^{+61.5}_{-34.9}$ $^{+3.3\%}_{-3.3\%}$	5.092(23) $^{+15.4}_{-17.9}$ $^{+1.2\%}_{-1.2\%}$
	$0^-$	9.304(9) $^{+61.6}_{-34.9}$ $^{+3.4\%}_{-3.4\%}$	11.29(4) $^{+16.0}_{-18.2}$ $^{+1.2\%}_{-1.2\%}$
	$0^\pm$	6.775(6) $^{+61.5}_{-34.9}$ $^{+3.3\%}_{-3.3\%}$	8.055(35) $^{+15.8}_{-18.2}$ $^{+1.2\%}_{-1.2\%}$

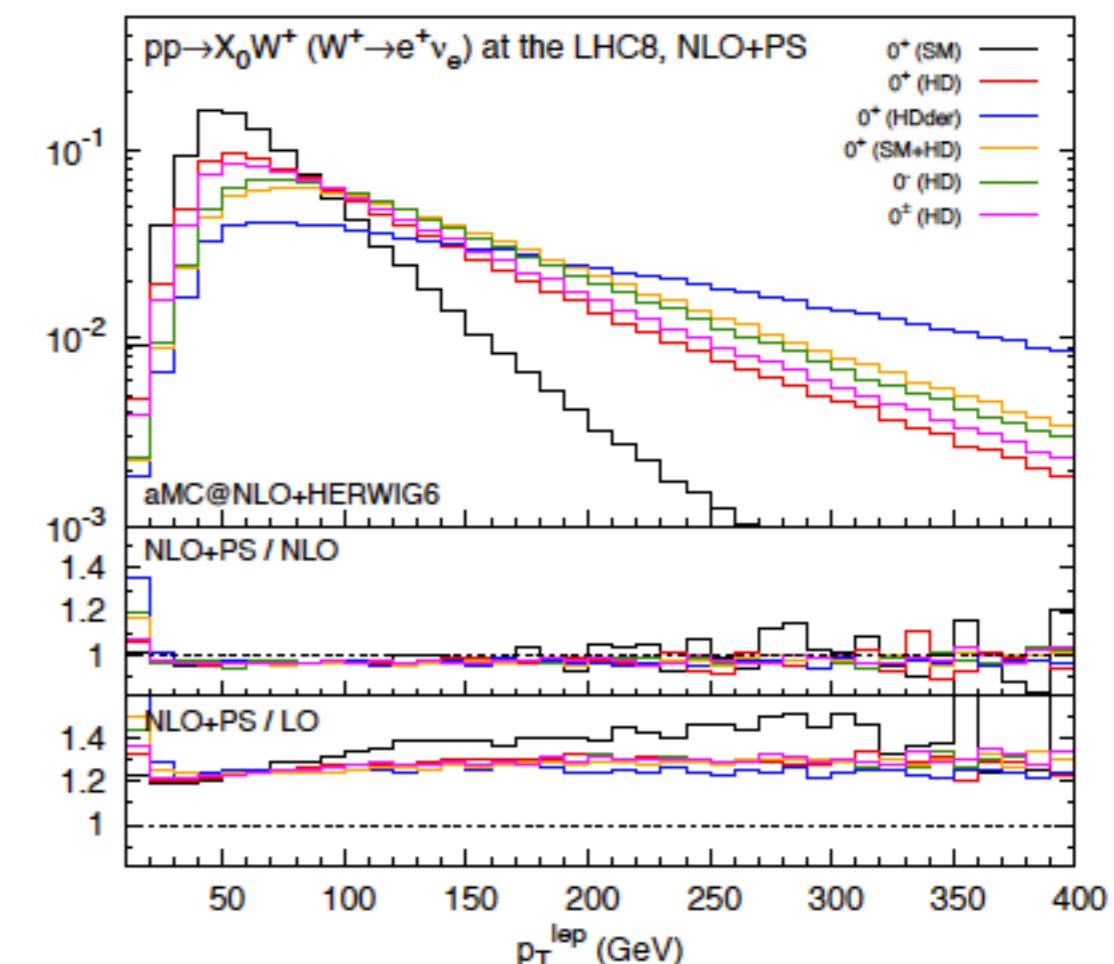
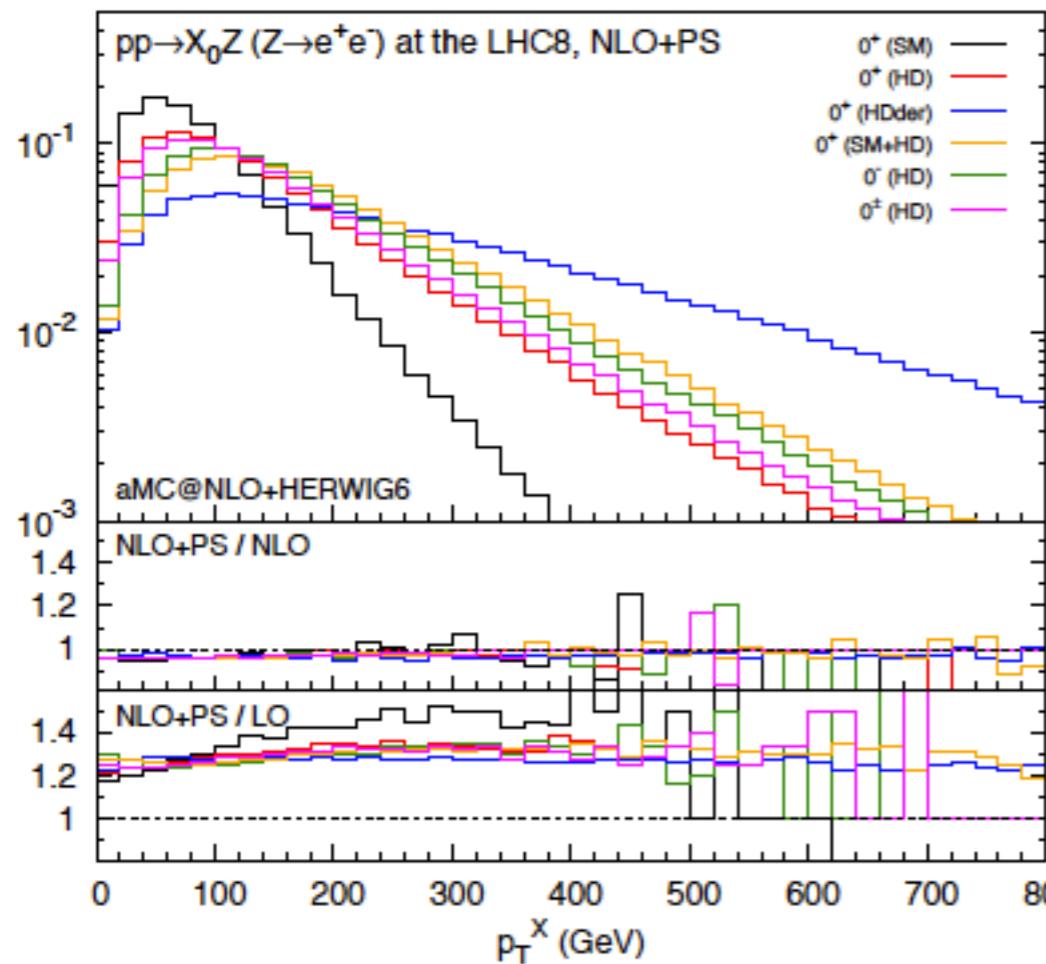


# VH

Maltoni, KM, Zaro [arXiv:1311.1829]

```
./bin/mg5_aMC
> import model HC_NLO
> generate p p > x0 e+ e- [QCD]
> launch
```

scenario	$\sigma_{\text{LO}}$ (fb)	$\sigma_{\text{NLO}}$ (fb)	$K$
$0^+$ (SM)	39.58(3)	$+0.1\%$ $-0.6\%$	51.22(5) $+2.2\%$ $-1.8\%$
$0^+$ (HD)	13.51(1)	$+1.5\%$ $-1.7\%$	17.51(1) $+1.9\%$ $-1.3\%$
$0^+$ (HDder)	324.2(2)	$+4.7\%$ $-4.3\%$	416.1(4) $+2.3\%$ $-2.1\%$
$0^+$ (SM+HD)	118.8(1)	$+3.0\%$ $-2.9\%$	154.2(1) $+1.8\%$ $-1.6\%$
$0^-$ (HD)	8.386(7)	$+2.6\%$ $-2.6\%$	10.89(1) $+1.8\%$ $-1.5\%$
$0^\pm$ (HD)	10.96(1)	$+1.9\%$ $-2.1\%$	14.22(1) $+1.8\%$ $-1.3\%$

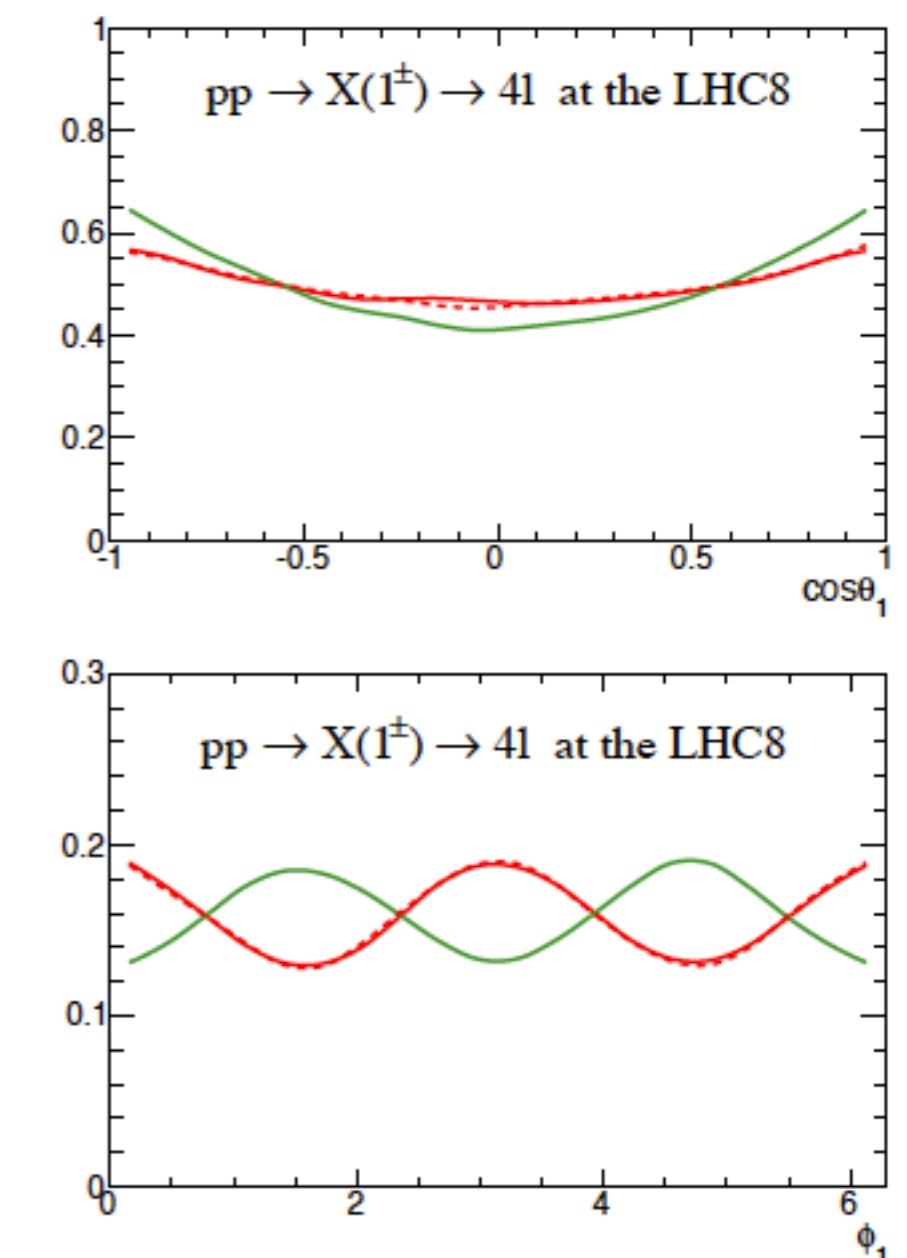
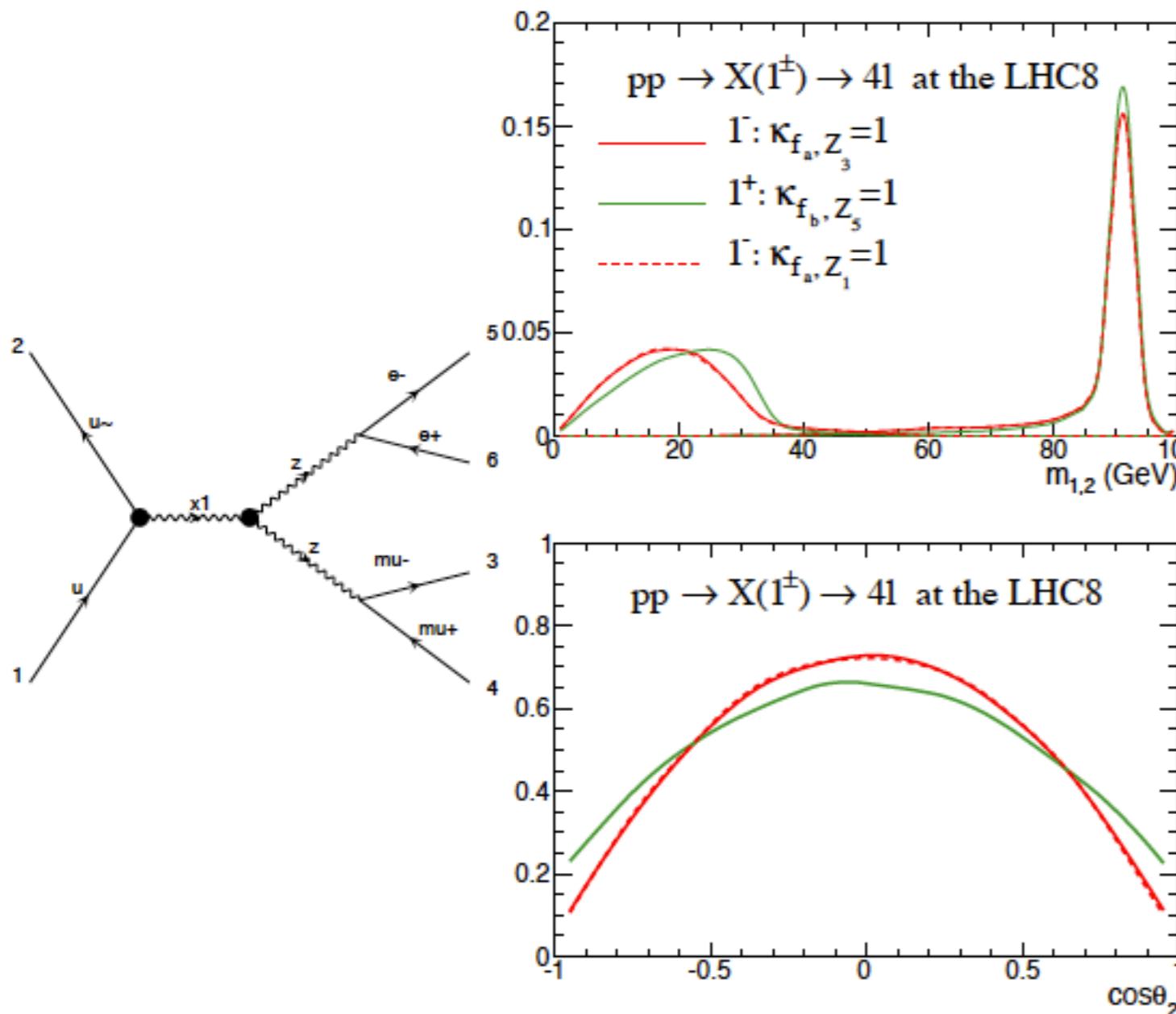


# Summary

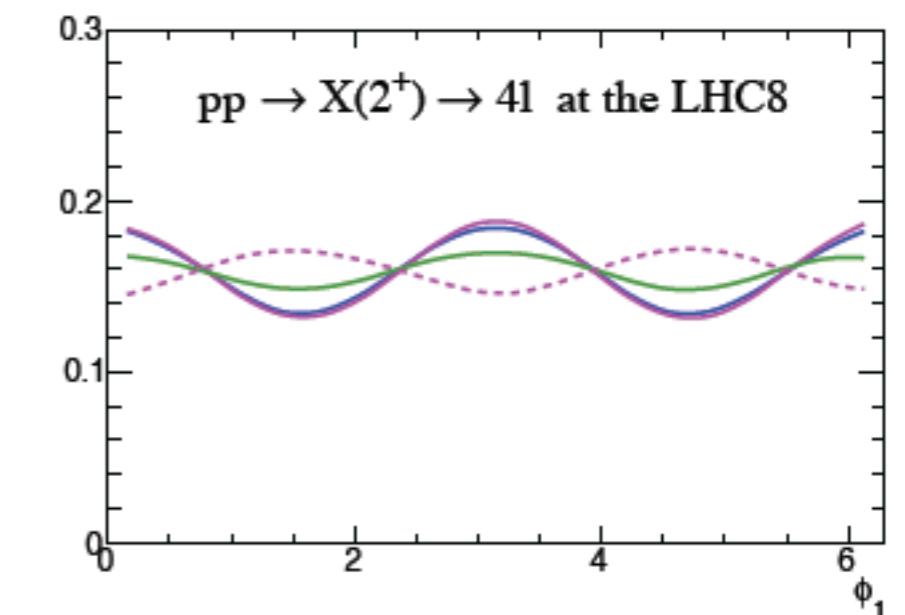
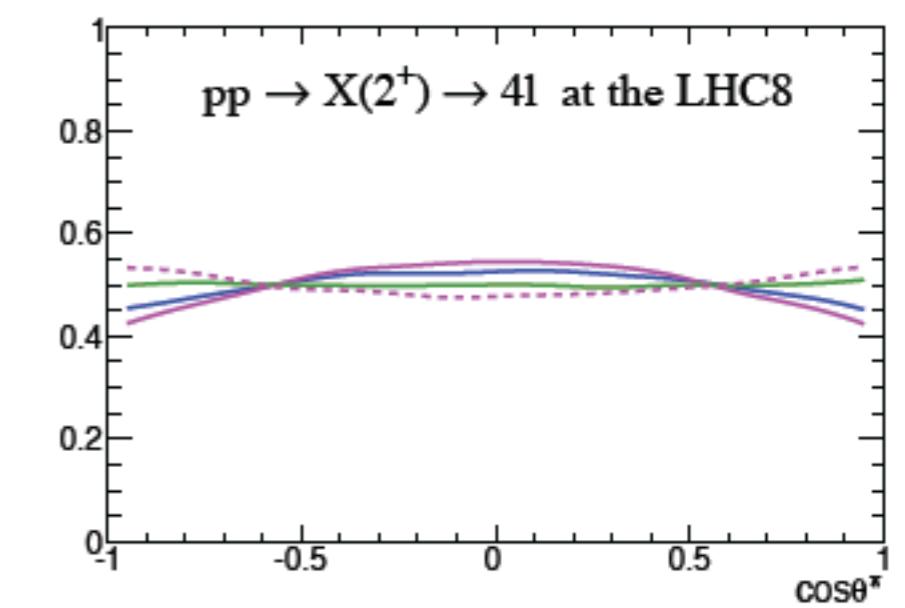
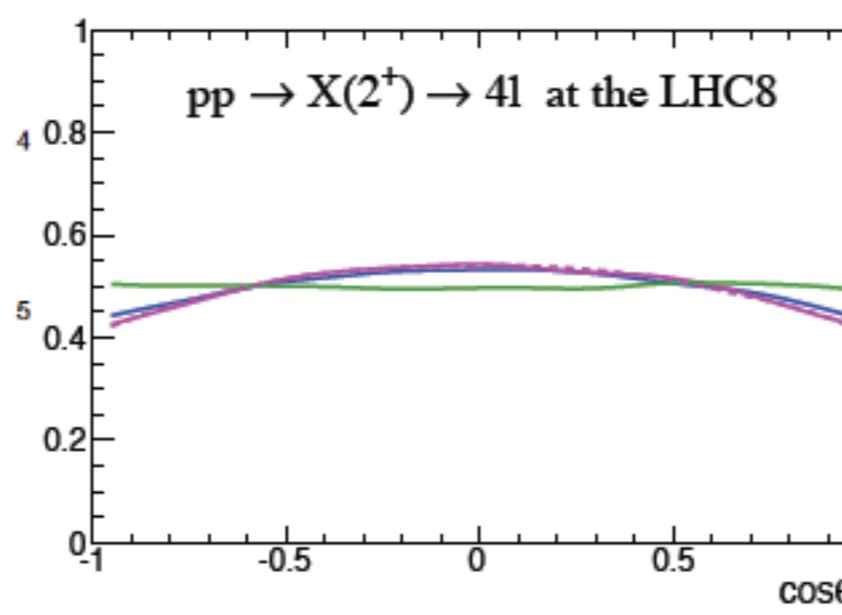
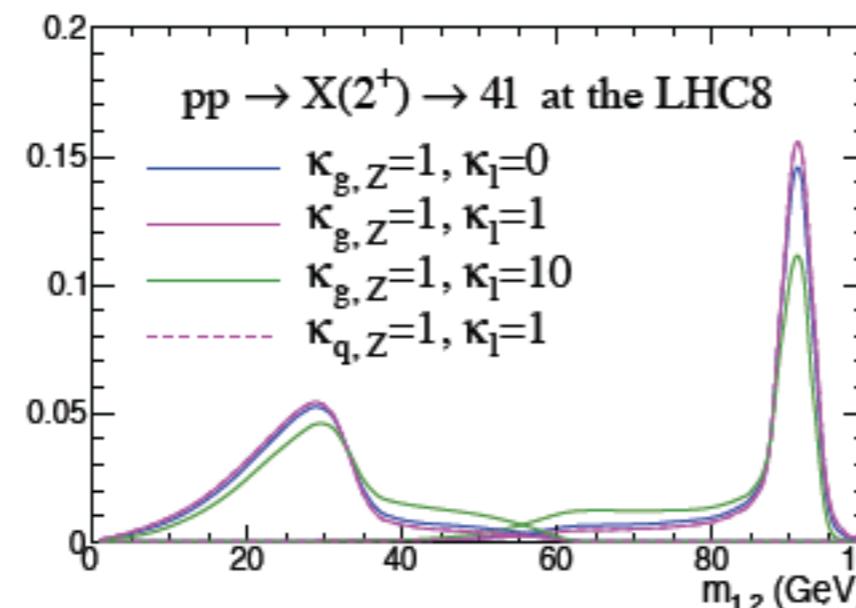
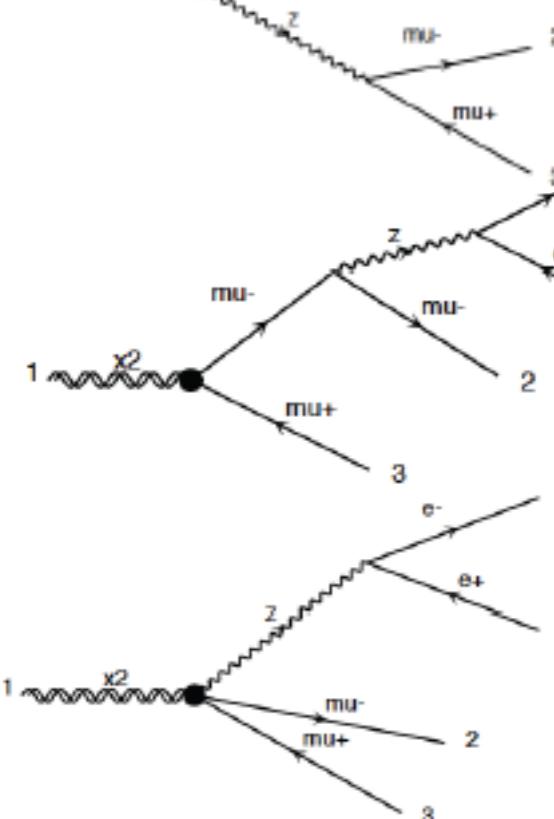
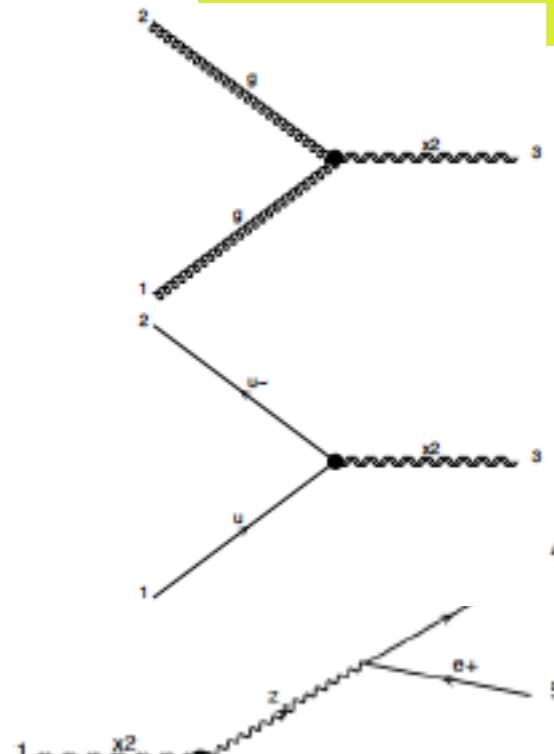
- After the discovery of a Higgs-like resonance at the LHC, the main focus of the analyses now is **the determination of the Higgs Lagrangian.**
- This includes
  - **the structure of the operators**, linked to the spin/parity of a ‘Higgs’ boson.
  - **the coupling strength.**
- MC tools to study the property of the SM-like boson are publicly available, e.g. HC, HEL (based on the EFT).
- Event generation at NLO is possible for (several) spin 0,1,2 hypothesis and can be used to validate merged samples.

# Back-up

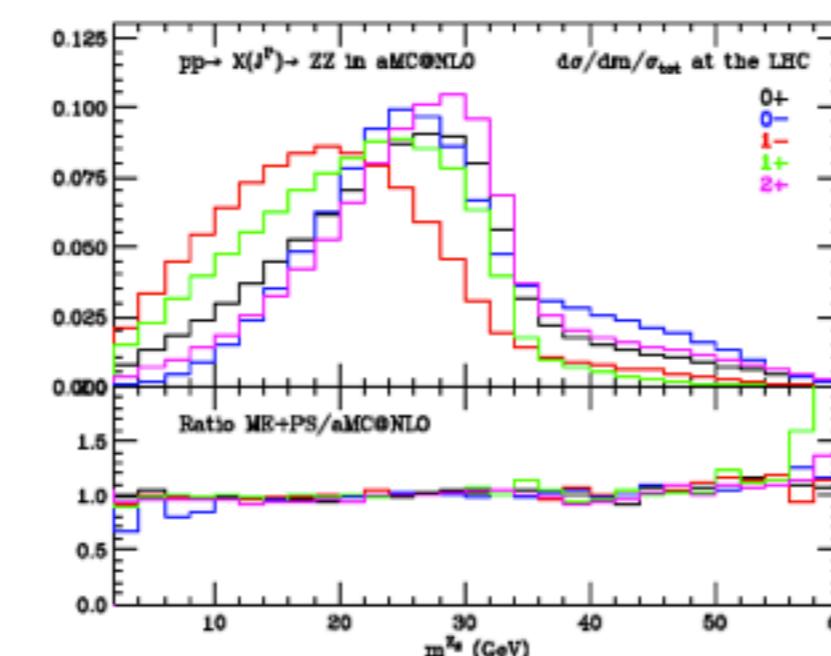
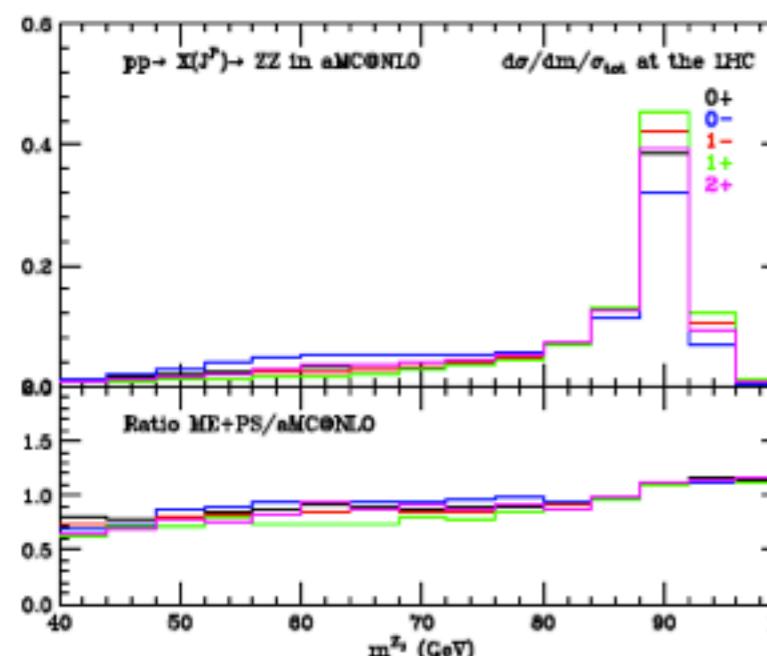
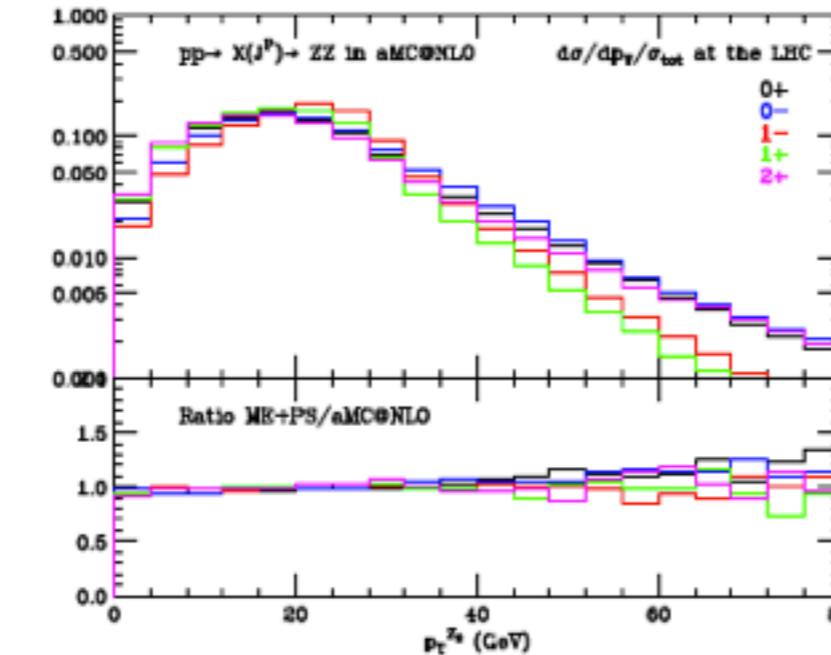
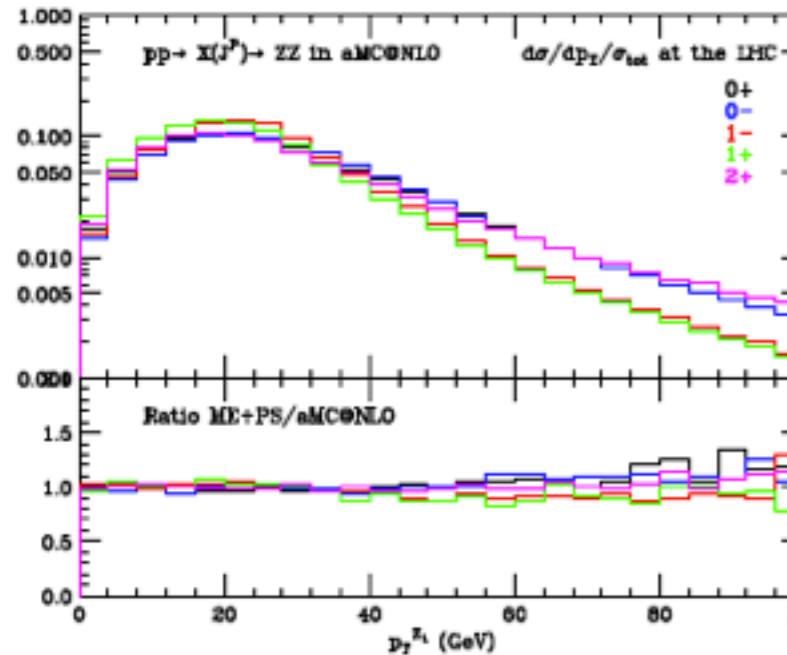
# Mass and angular distributions -- spin I



# Mass and angular distributions -- spin2



# aMC@NLO vs. ME+PS

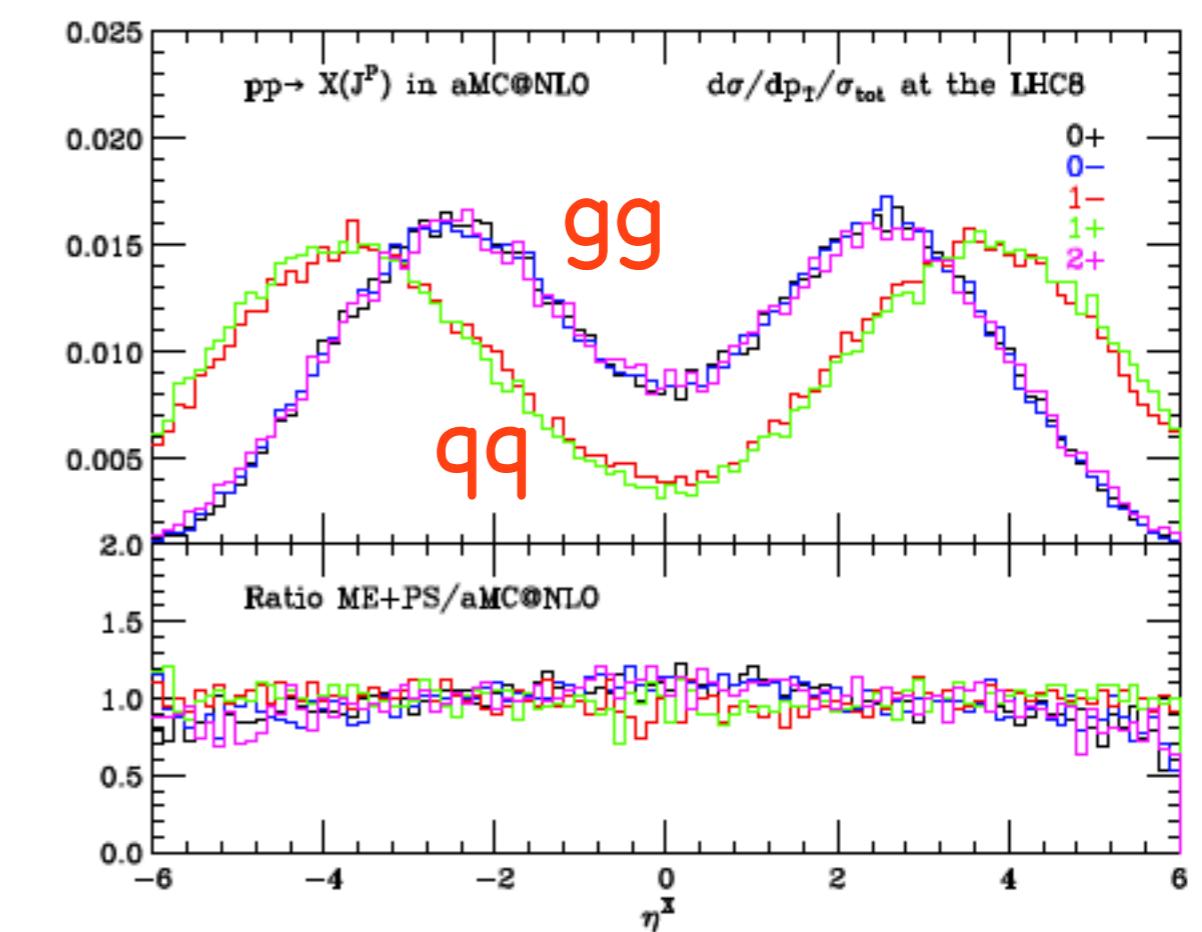
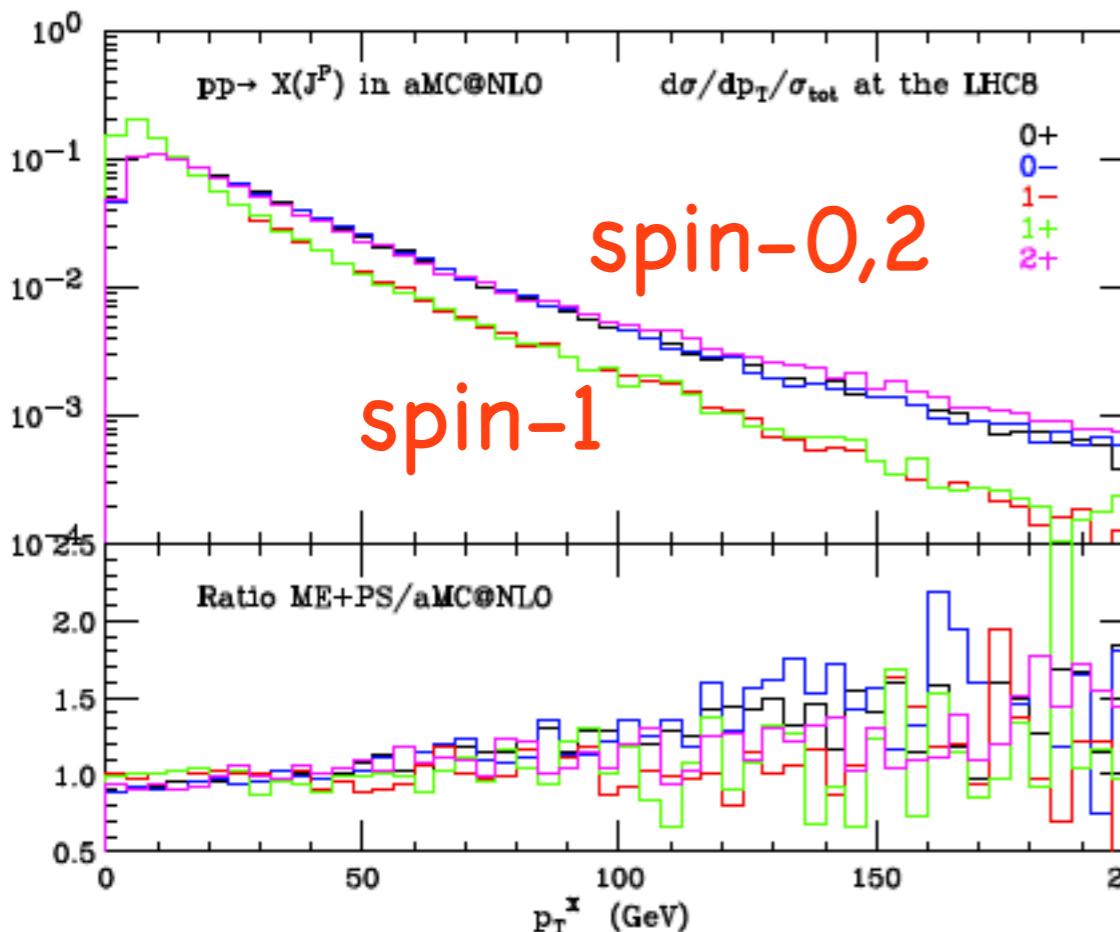


Good agreement between the ME+PS and aMC@NLO predictions for most observables.

For spin 0, the production and decay factorize, for spin 1 and 2 this does not happen and the full  $2 \rightarrow 4,5$  matrix elements need to be used.

# Higher order effects in QCD (I)

## inclusive production in $\text{pp} \rightarrow X(J^P)$



The matched sample is harder than aMC@NLO at large pT due to the extra 2 ME patrons in the matched sample.

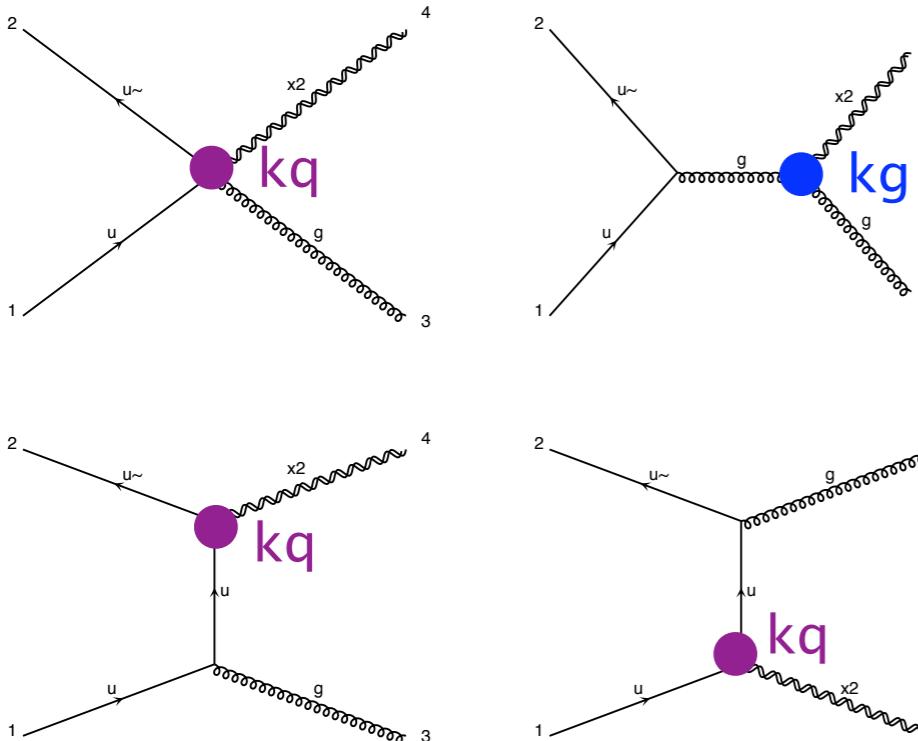
excellent agreement between  
ME+PS and aMC@NLO

The different shapes are due to the different initial state.

# Higher order effects in QCD (II)

## unitarity-violating behavior of models with a spin-2 state

$$\mathcal{L} = -\frac{1}{\Lambda} \kappa_q T_{\mu\nu}^q X_2^{\mu\nu} - \frac{1}{\Lambda} \kappa_g T_{\mu\nu}^g X_2^{\mu\nu}$$



$$|\mathcal{M}|^2 \propto s/\Lambda^2 \quad \text{for } \kappa_q = \kappa_g$$

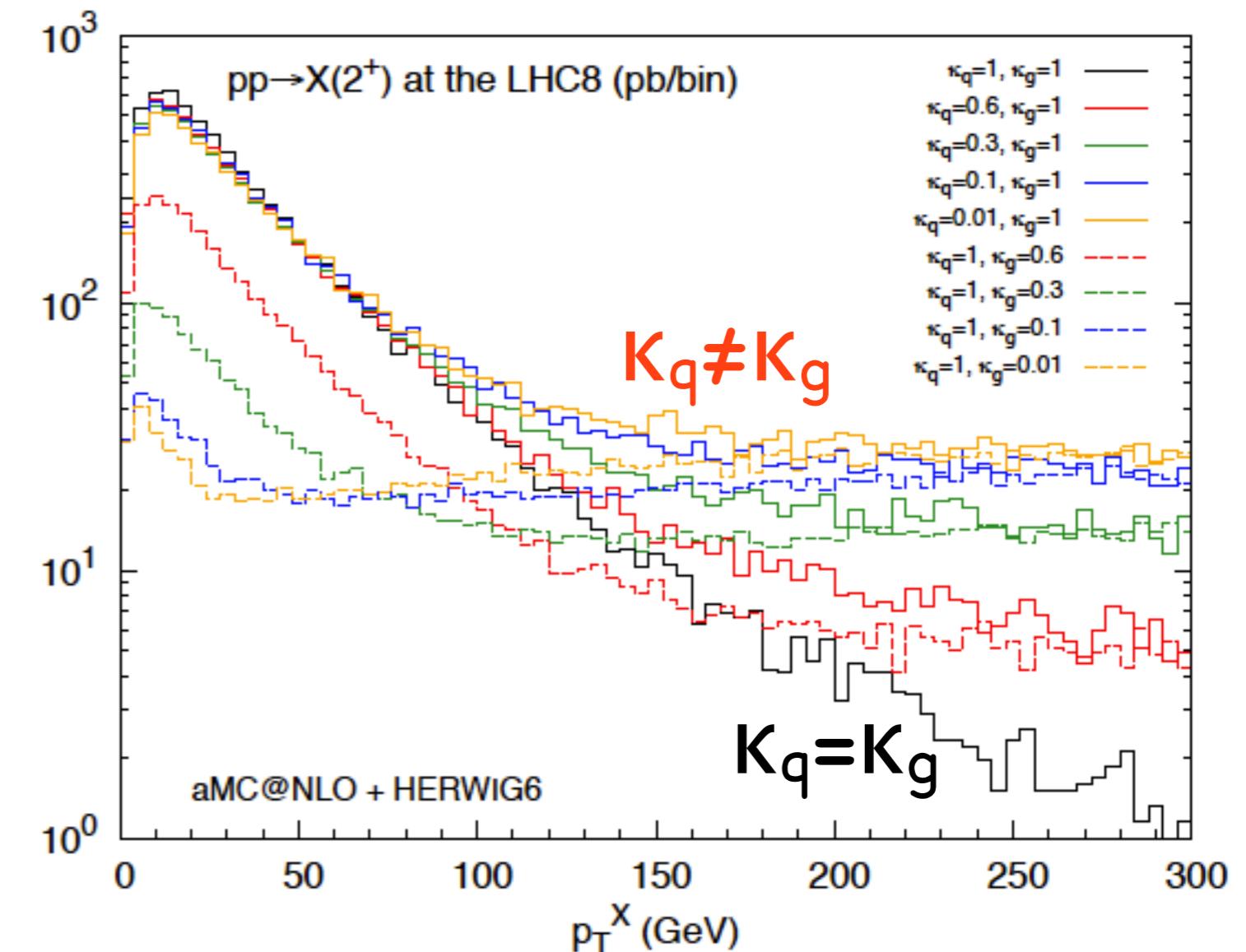
$$|\mathcal{M}|^2 \propto s^3/m^4 \Lambda^2 \quad \text{for } \kappa_q \neq \kappa_g$$

$$|\mathcal{M}|^2 = \frac{N}{\Lambda^2 s t u m^4} \left\{ 3\kappa_g^2 m^4 [2m^4 - 2m^2(t+u) + t^2 + u^2] [m^4 - m^2(t+u) + 4tu] \right. \\ + (\kappa_q - \kappa_g) 6 \kappa_g m^4 s [m^6 + m^2 s(s+2u) - 2su(s+u)] \\ + (\kappa_q - \kappa_g)^2 s [6m^{10} - 6m^8(t+u) + 3m^6(t^2 + u^2) - 12m^4 tu(t+u) \\ \left. + 2m^2 tu(t^2 + 12tu + u^2) - 2tu(t^3 + t^2u + tu^2 + u^3)] \right\}, \quad (4.2)$$

# Higher order effects in QCD (II)

## unitarity-violating behavior of models with a spin-2 state

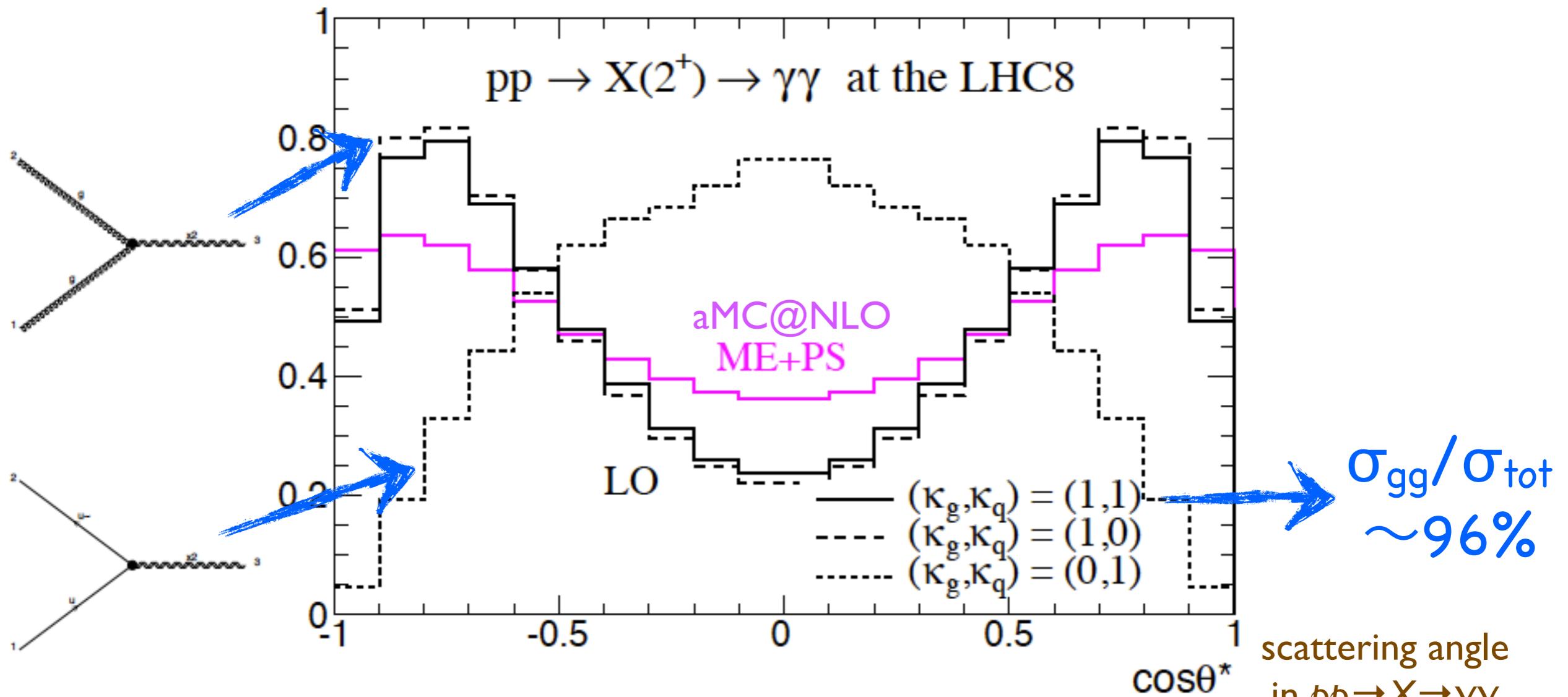
$$\mathcal{L} = -\frac{1}{\Lambda} \kappa_q T_{\mu\nu}^q X_2^{\mu\nu} - \frac{1}{\Lambda} \kappa_g T_{\mu\nu}^g X_2^{\mu\nu}$$



A model with non-universal couplings  
dramatically changes the  $p_T(X)$  spectrum.

# Higher order effects in QCD (III)

## on spin observables for a spin-2 state



$$\frac{d\sigma(gg)}{d\cos\theta^*} \propto |d_{22}^2(\theta^*)|^2 + |d_{2-2}^2(\theta^*)|^2 = \frac{1}{8}(1 + 6\cos^2\theta^* + \cos^4\theta^*),$$

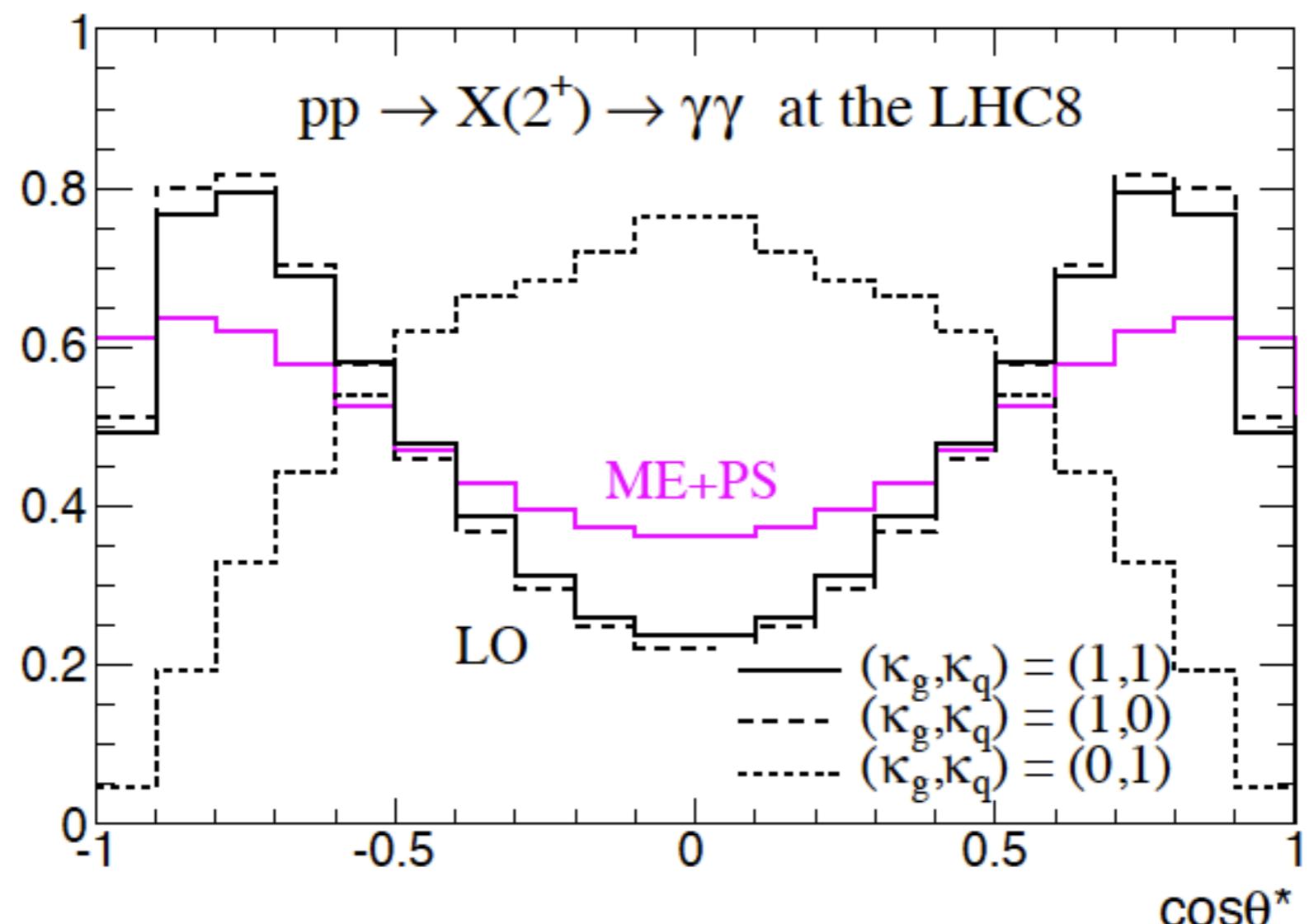
$$\frac{d\sigma(q\bar{q})}{d\cos\theta^*} \propto |d_{12}^2(\theta^*)|^2 + |d_{1-2}^2(\theta^*)|^2 = \frac{1}{2}(1 - \cos^4\theta^*).$$

# How can we get the spin/parity information?

1.  $X \rightarrow \gamma\gamma$
2.  $X \rightarrow VV^* \rightarrow 4l$
3.  $pp \rightarrow jjX$
4.  $pp \rightarrow vX$
5.  $X \rightarrow \tau\tau$

# Spin/parity determination

## I. $X \rightarrow \gamma\gamma$



$$\frac{d\sigma(gg)}{d\cos\theta^*} \propto |d_{22}^2(\theta^*)|^2 + |d_{2-2}^2(\theta^*)|^2 = \frac{1}{8}(1 + 6\cos^2\theta^* + \cos^4\theta^*) ,$$

$$\frac{d\sigma(q\bar{q})}{d\cos\theta^*} \propto |d_{12}^2(\theta^*)|^2 + |d_{1-2}^2(\theta^*)|^2 = \frac{1}{2}(1 - \cos^4\theta^*) .$$

# Spin/parity determination

## 2. $X \rightarrow VV^* \rightarrow 4l$

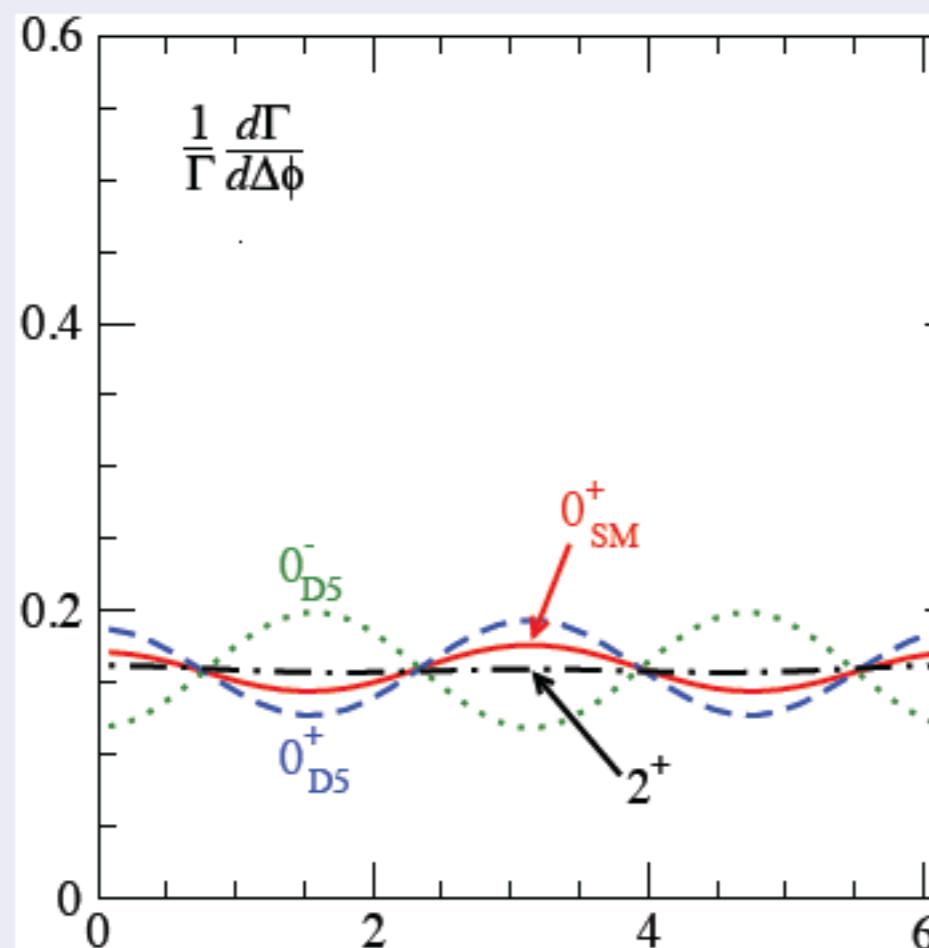
[Dell'Aquila, Nelson, PRD(1986)]

[Choi, Miller, Mühlleitner, Zerwas, PLB(2003)]

[Gao et al, PRD(2010)] ...

[Bolognesi et al, PRD(2012)]

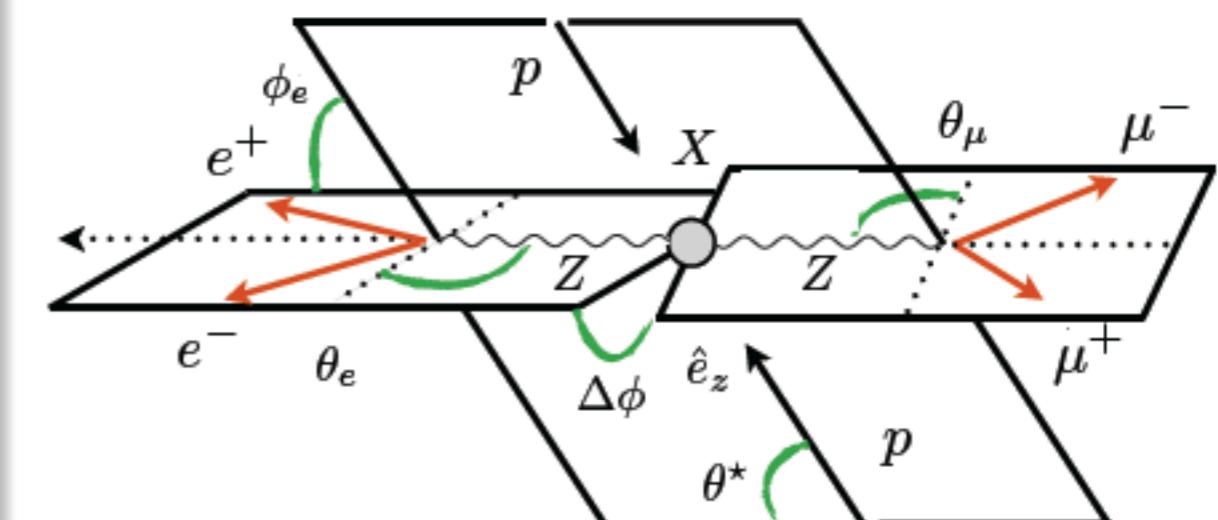
$X \rightarrow ZZ^* \rightarrow 4\ell$



$$\mathcal{L}_{0^+_{SM}} = g_{0^+_{SM}} V_\mu V^\mu X_0$$

$$\mathcal{L}_{0^+_{D5}} = g_{0^+_{D5}} V_{\mu\nu} V^{\mu\nu} X_0$$

$$\mathcal{L}_{0^-_{D5}} = g_{0^-_{D5}} V_{\mu\nu} \tilde{V}^{\mu\nu} X_0$$



$d\sigma/d\Delta\phi \sim \text{const.}$  for  $0^+_{SM}$ ,  $d\sigma/d\Delta\phi \sim 1 \pm A \cos 2\Delta\phi$  for  $0^\pm_{D5}$ .

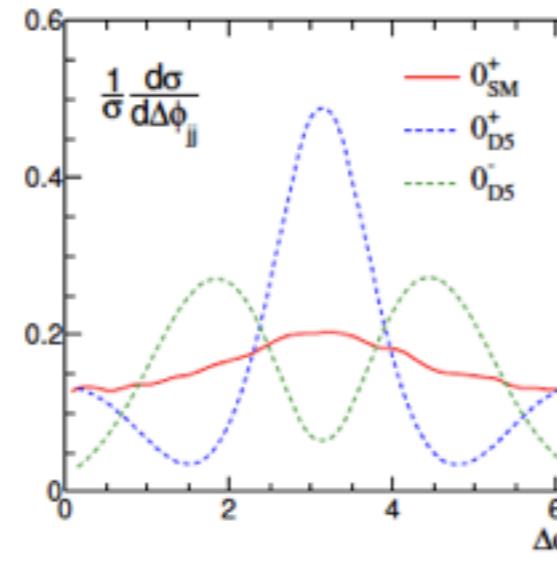
# Spin/parity determination

## 3. $p\bar{p} \rightarrow jjX$

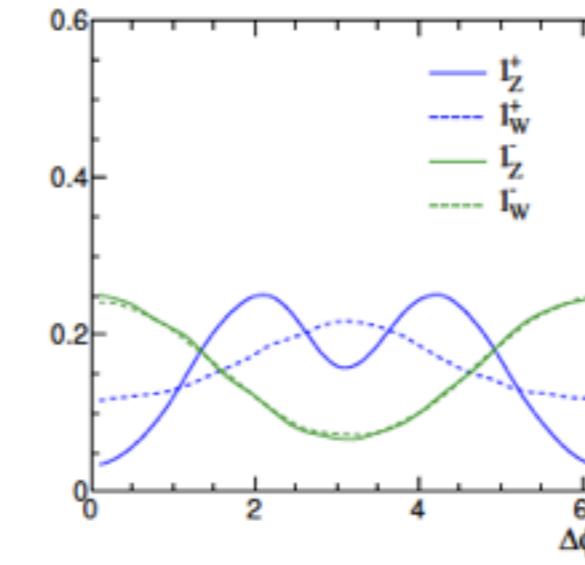
di-jet correlations

Englert, Goncalves-Netto, KM, Plehn (2013)

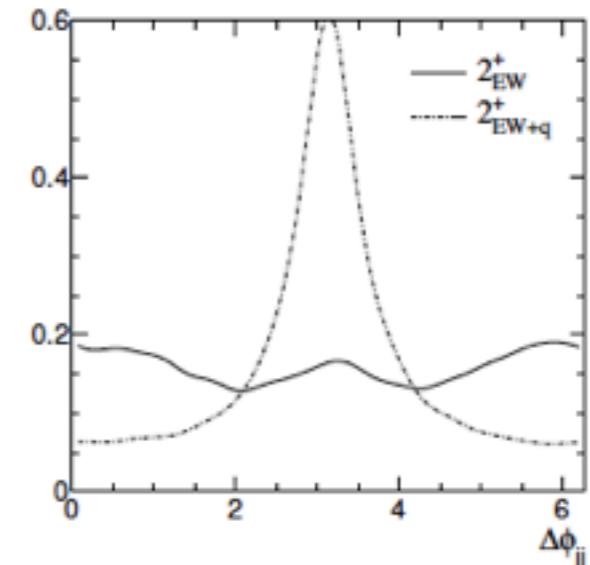
spin-0



spin-1



spin-2



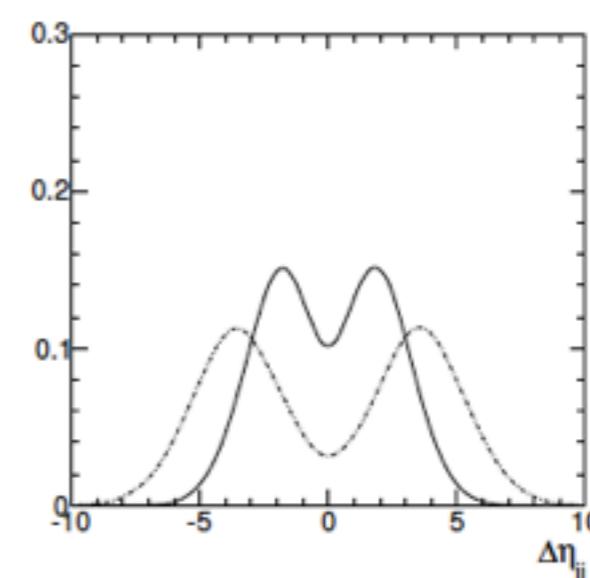
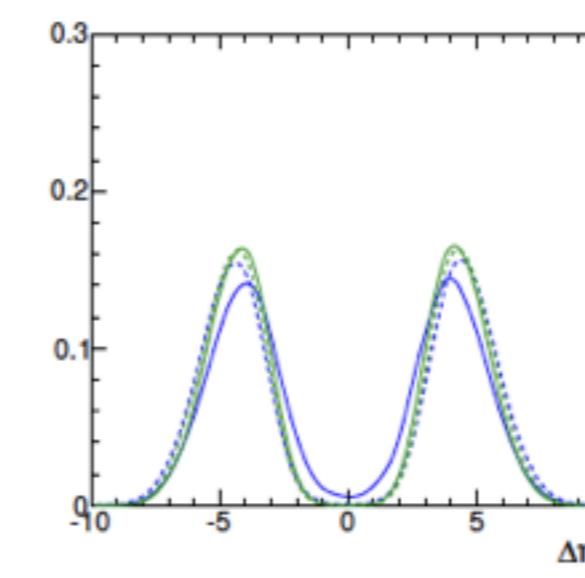
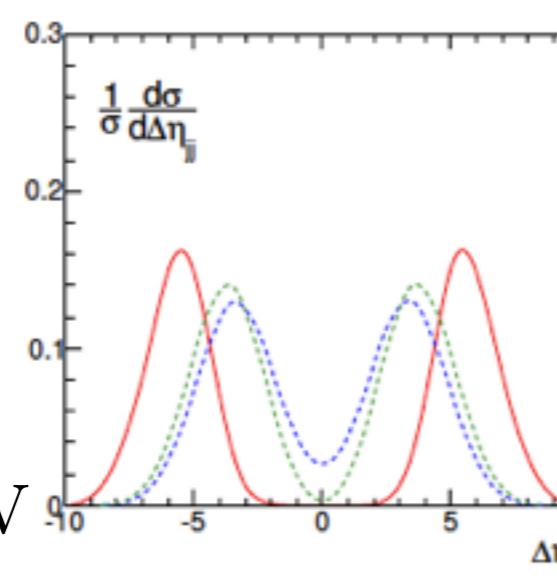
$\sqrt{s} = 14$  TeV

$p_{Tj} > 20$  GeV

$\Delta R_{jj} > 0.6$

$|\eta_j| < 5$

$m_{jj} > 600$  GeV

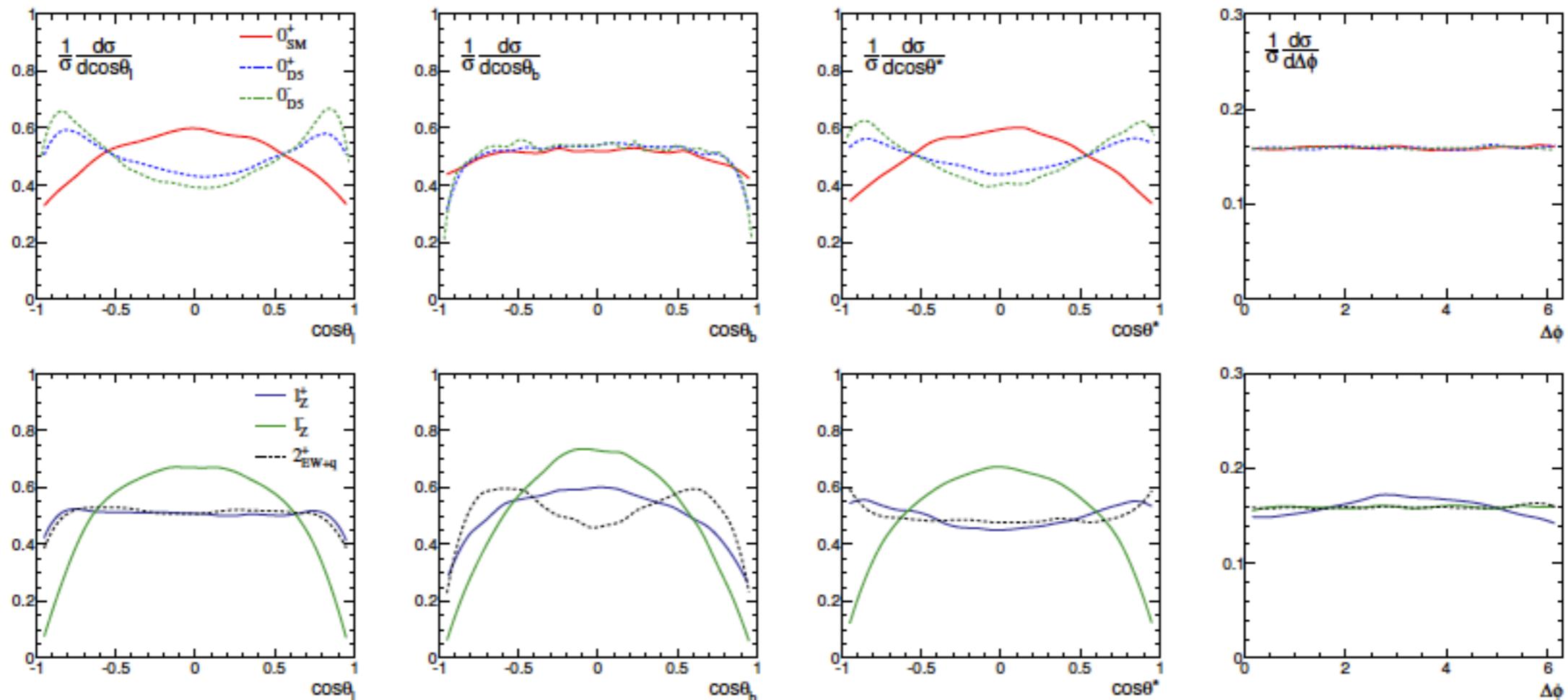


$\Delta\eta$  as well as  $\Delta\Phi$  are the powerful observables.

# Spin/parity determination

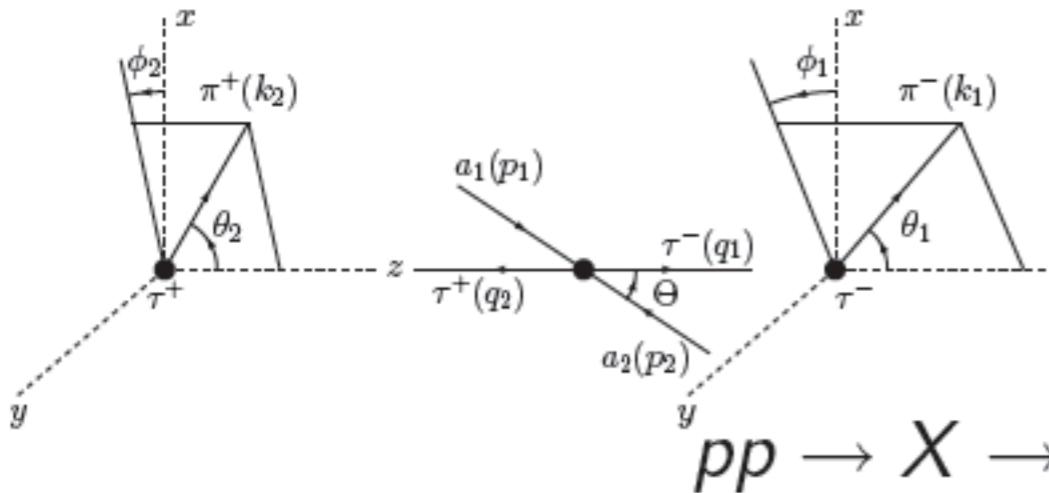
## 4. $p\bar{p} \rightarrow ZX$

Englert, Goncalves-Netto, KM, Plehn (2013)



# Spin/parity determination

## 5. $X \rightarrow \tau\tau$



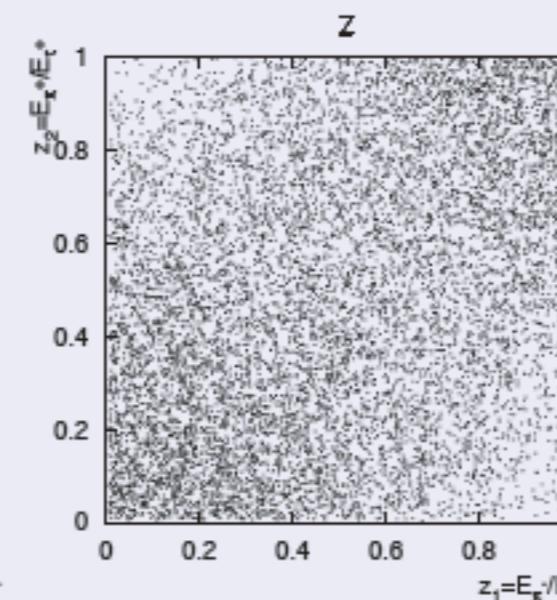
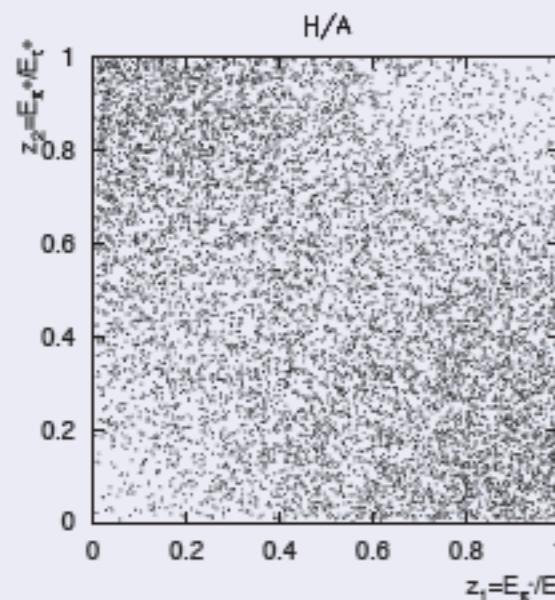
[Bullock, Hagiwara, Martin, NPB(1993)]

[Krämer, Kühn, Stong, Zerwas, ZPC(1994)]

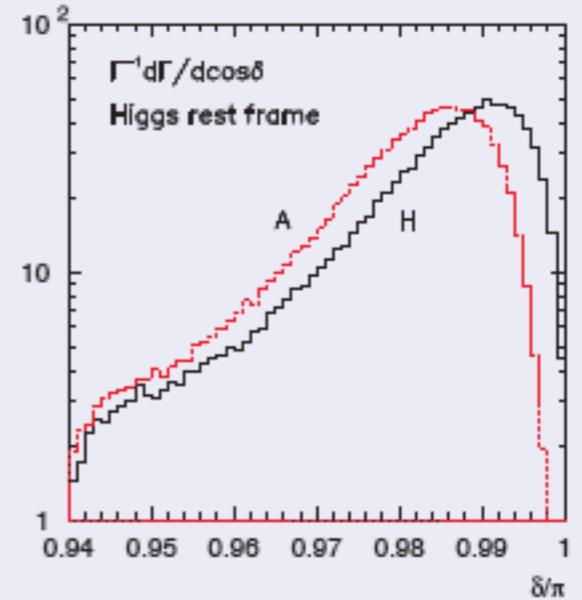
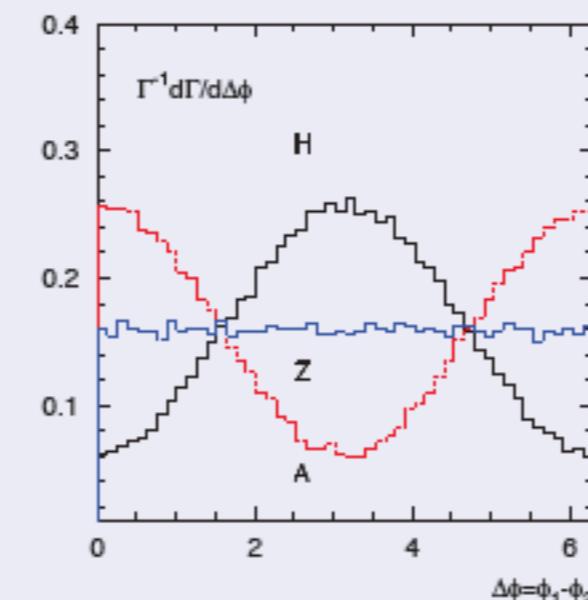
[Pierzchala, Richter-Was, Was, Worek, APPB(2001,2002,...)]

[Hagiwara, Li, KM, Nakamura, 1212.6247]

### Longitudinal spin (helicity) effect



### Transverse spin effect



$$d^2\Gamma/dz_1 dz_2 \sim 1 \mp z_1 z_2 \text{ for spin-0/1, } d\Gamma/d\Delta\phi \sim 1 \mp A \cos \Delta\phi \text{ for } 0^\pm$$

$\tau$  could be a spin/parity analyzer!