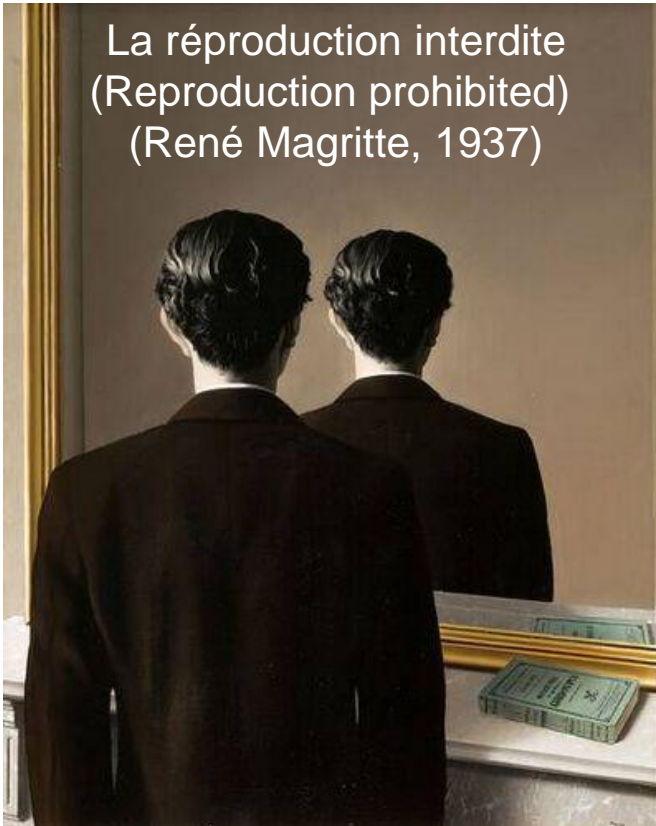




Reflections on the atomic nucleus, Liverpool

La réproduction interdite
(Reproduction prohibited)
(René Magritte, 1937)

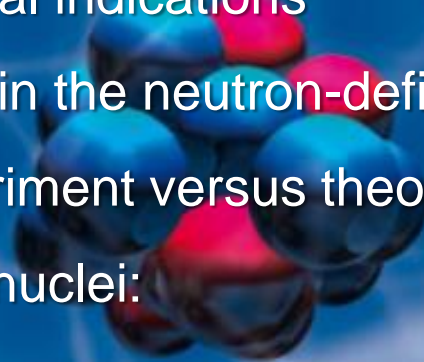


Coexisting Quadrupole Shapes in Heavy Exotic Nuclei

Piet Van Duppen
Department of Physics and Astronomy
KU Leuven, Belgium

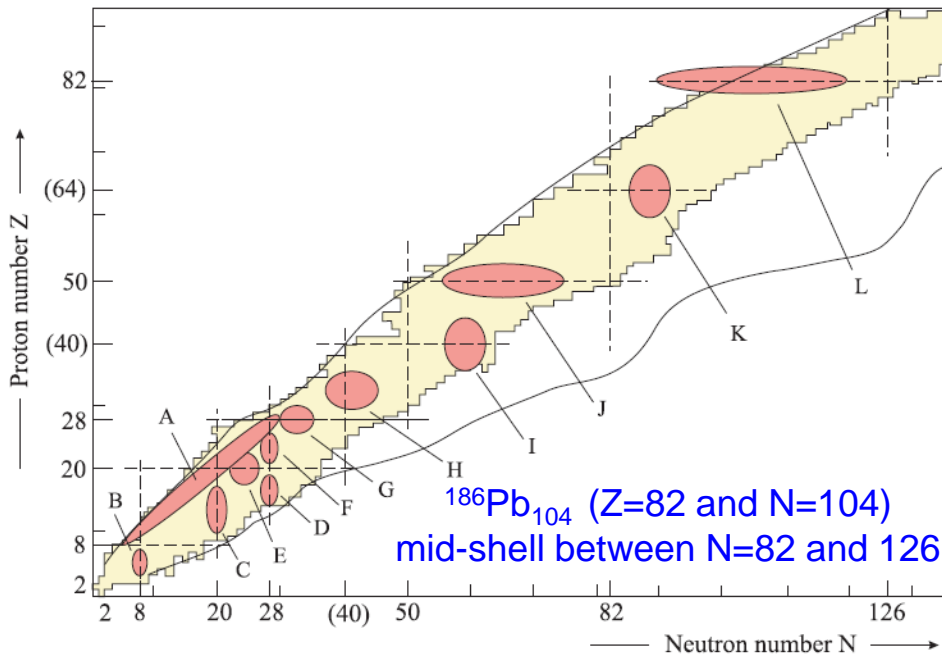
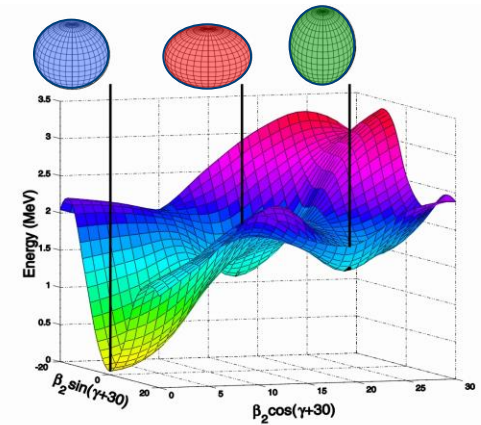
KU LEUVEN

Piet Van Duppen

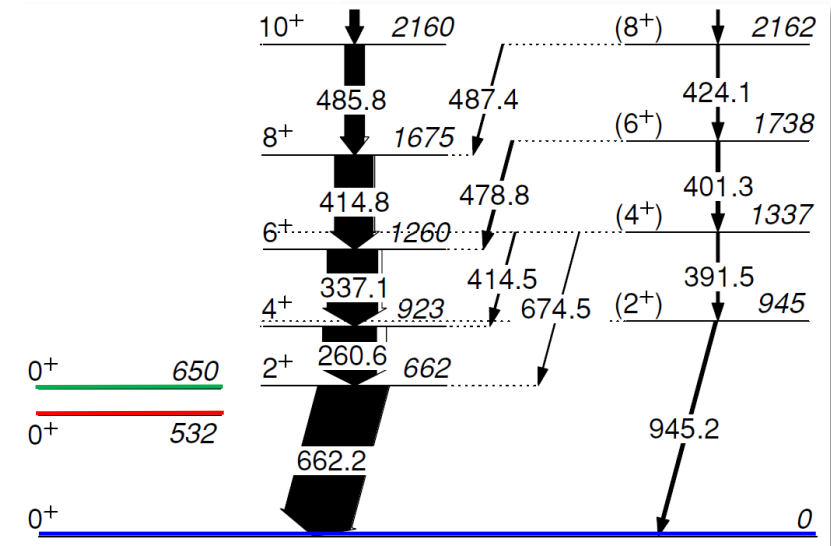
- 
- Shape coexistence in heavy nuclei: initial indications
 - Coulex and laser spectroscopy studies in the neutron-deficient nuclei around $Z=82$ (lead region): experiment versus theory
 - New approaches towards the heaviest nuclei:
 - In-gas laser ionization spectroscopy
 - Conclusion and outlook reflections

• Shape Coexistence in Atomic Nuclei

- states with different (*quadrupole*) shapes at similar binding energy
 - the result of the interplay between two opposing tendencies:
 - stabilizing effect of closed shells (and subshells):
 - residual proton-neutron interaction: correlation energy
- (Heyde and Wood, Review Modern Physics (2011))



Potential Energy Surface for ^{186}Pb

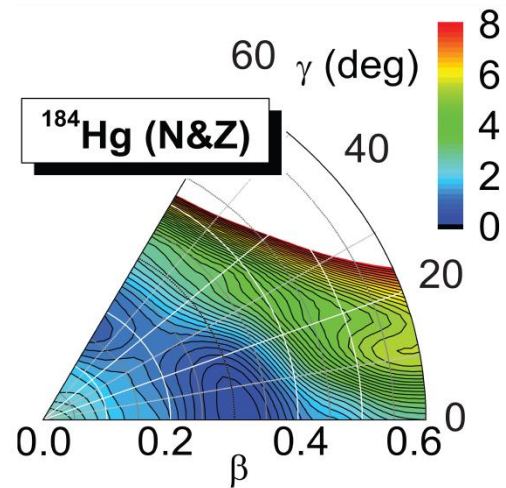
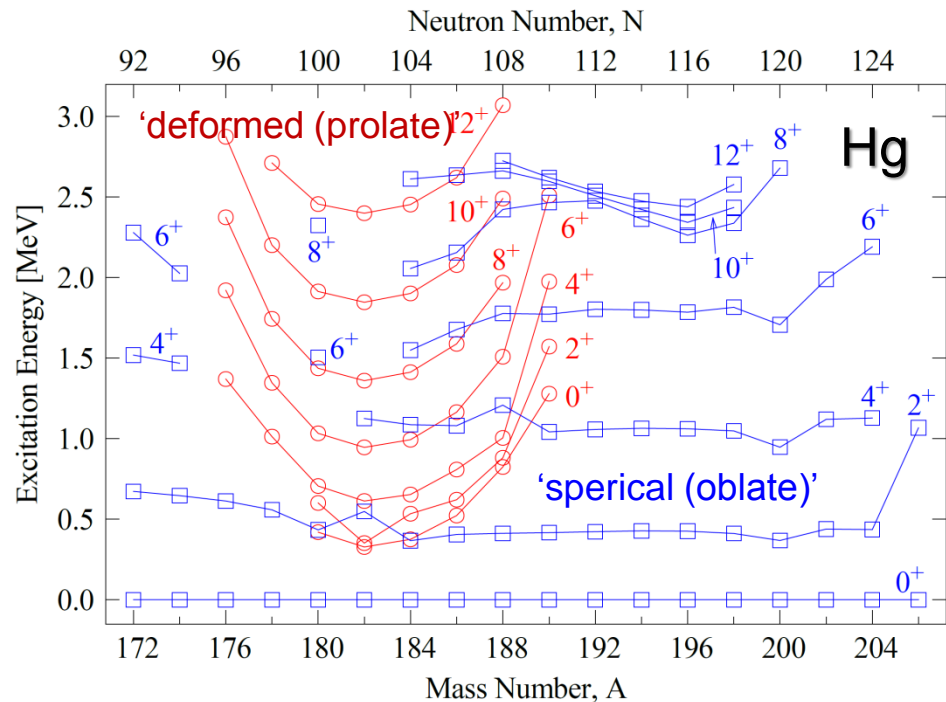
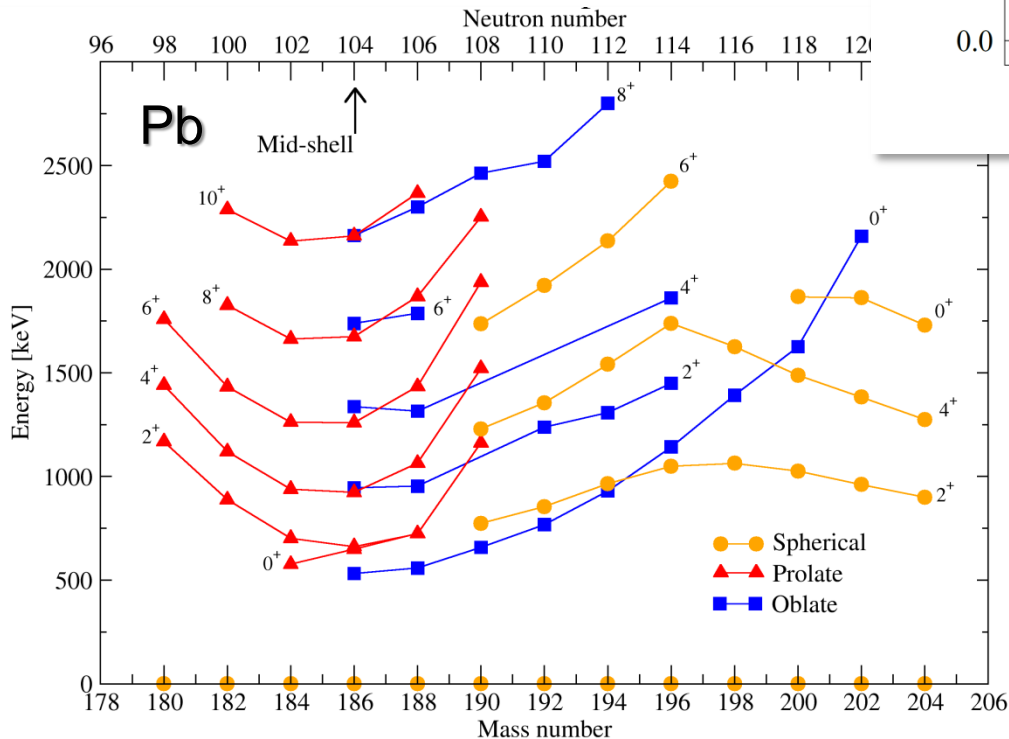


$^{186}\text{Pb}_{104}$



see also talks by R. Julin and J. Pakarinen

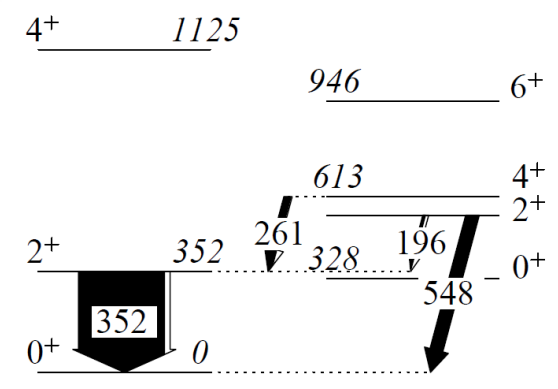
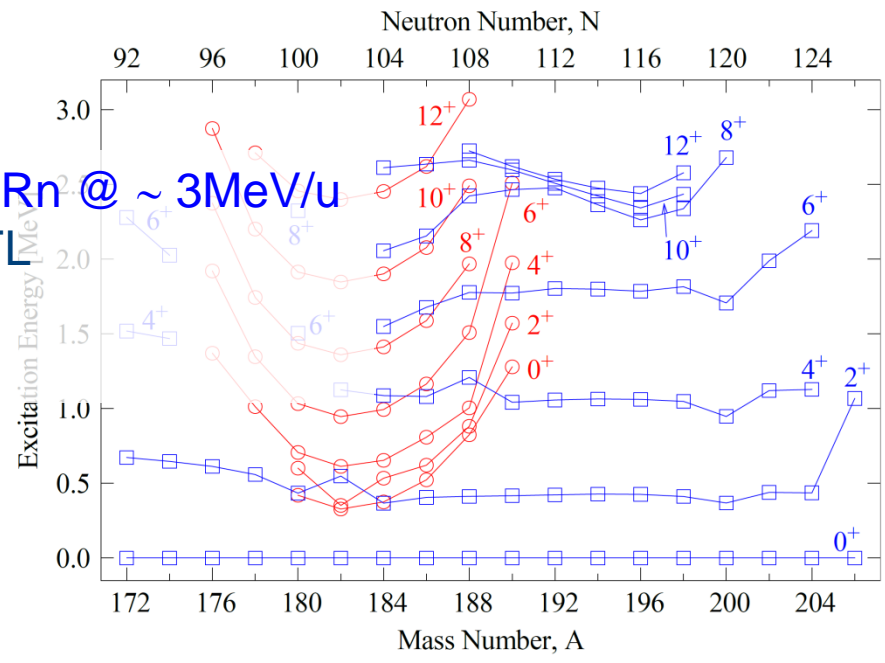
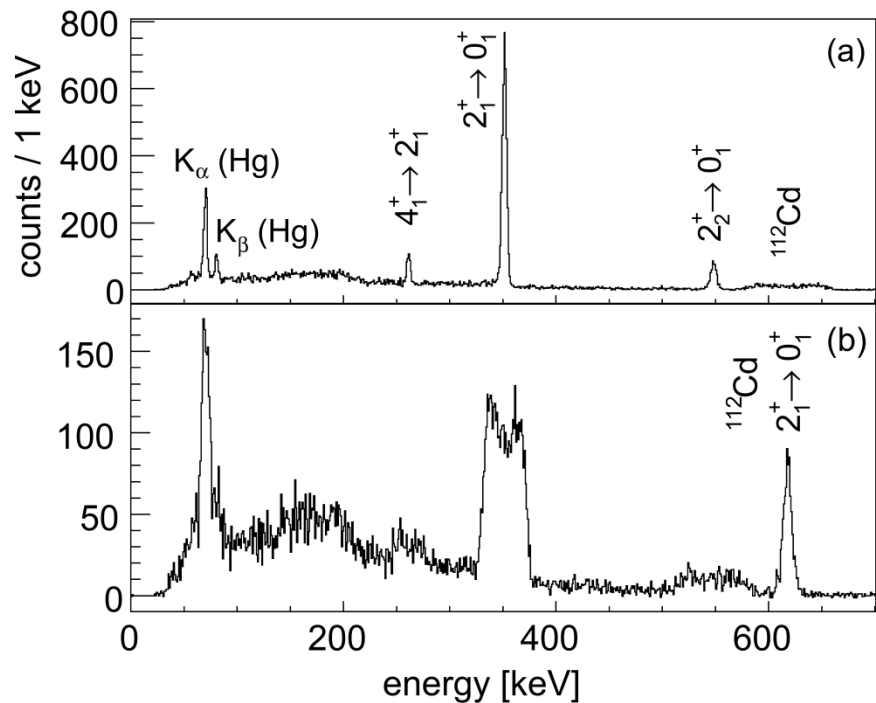
- Level energy systematics



Yao PRC87 (2013)
Egido PRL93 (2004)

KU LEUVEN

- Coulex of $^{182-188}\text{Hg}$, $^{188-198}\text{Pb}$, $^{196-202}\text{Po}$, $^{202-204}\text{Rn}$ @ $\sim 3\text{MeV/u}$
- Life time measurements: $^{184-188}\text{Hg}$: ANL, JYFL
- β -decay studies: $^{182,184}\text{Tl} - ^{182,184}\text{Hg}$
- γ -branchings, conversion coefficients



^{182}Hg ($T_{1/2}=10.8\text{ s}$; 3500 pps)

$\langle I_i E2 I_f \rangle$ (eb)	^{182}Hg	^{184}Hg	^{186}Hg	^{188}Hg
$\langle 0_1^+ E2 2_1^+ \rangle$	$1.29^{+0.04}_{-0.03}$	1.27 (3)	$1.25^{+0.10}_{-0.07}$	1.31 (10)

• Quadrupole Sum Rules

- relate experimental ME2's with quadrupole deformation parameters

(D. Cline, *Ann Rev Nucl Part Sci* 36 (1986) 683)

Overall deformation $\sim \beta \sim$

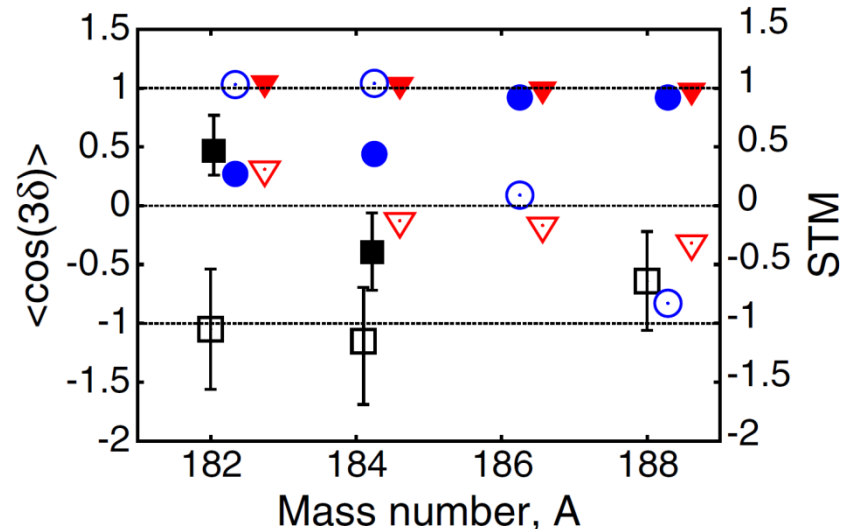
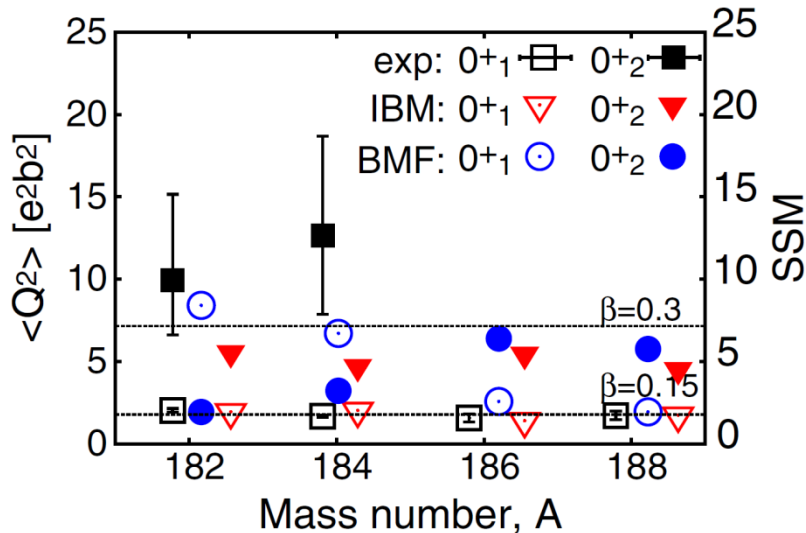
$$\frac{1}{\sqrt{5}} \langle Q^2 \rangle = \langle i | [E2 \times E2]_0 | i \rangle =$$

$$= \frac{1}{\sqrt{(2I_i + 1)}} \sum_t \langle i || E2 || t \rangle \langle t || E2 || i \rangle \begin{Bmatrix} 2 & 2 & 0 \\ I_i & I_i & I_t \end{Bmatrix}$$

Triaxiality $\sim \gamma \sim \langle \cos 3\delta \rangle :$

$$\sqrt{\frac{2}{35}} \langle Q^3 \cos 3\delta \rangle = \langle i | [[E2 \times E2]_2 \times E2]_0 | i \rangle =$$

$$= \frac{\mp 1}{(2I_i + 1)} \sum_{t,u} \langle i || E2 || u \rangle \langle u || E2 || t \rangle \langle t || E2 || i \rangle \begin{Bmatrix} 2 & 2 & 2 \\ I_i & I_t & I_u \end{Bmatrix}$$



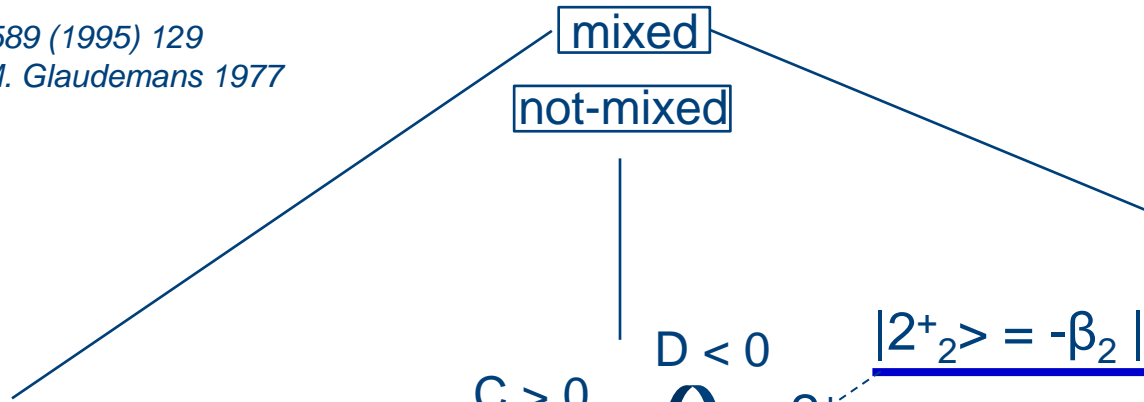
→ need more precise experimental matrix elements (e.g. HIE ISOLDE - 5 MeV/u)

→ need theoretical developments (including tri-axiality, odd-mass nuclei):

microscopic origin of shape coexistence

Two level mixing calculations

G.J. Lane et al., NPA 589 (1995) 129
 P. J. Brussaard, P.W.M. Glaudemans 1977

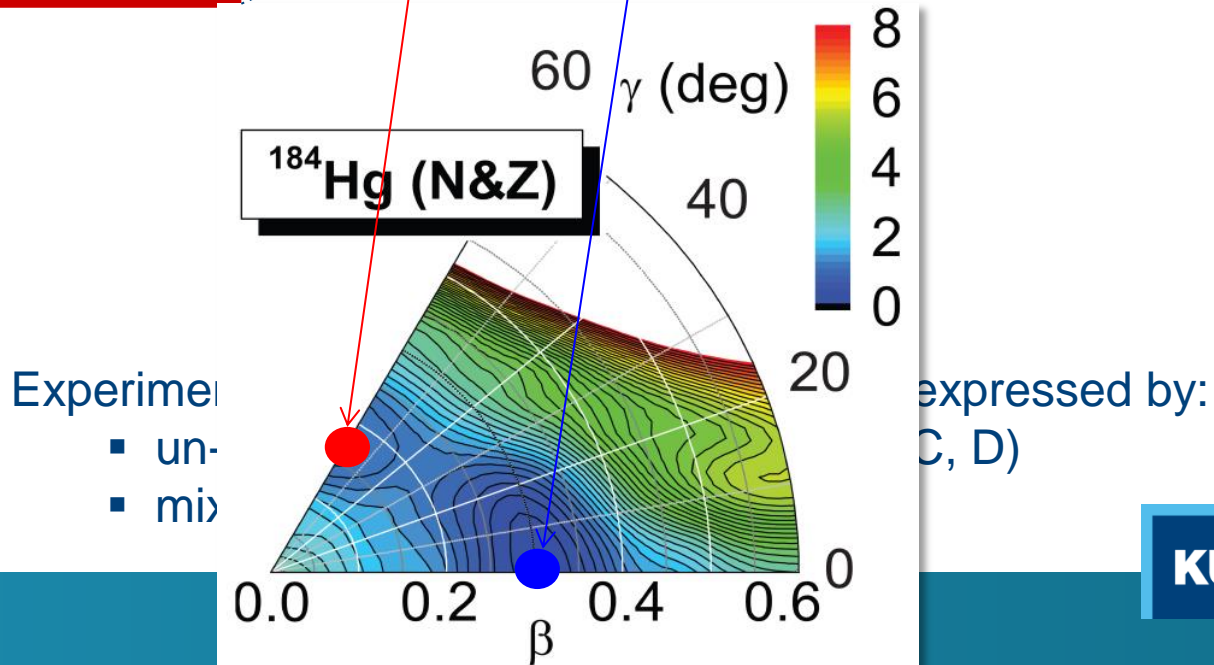


$$|2^+_{1>} = \alpha_2 |2^+_S> + \beta_2 |2^+_D>$$

$$|0^+_{1>} = \alpha_0 |0^+_S> + \beta_0 |0^+_D>$$

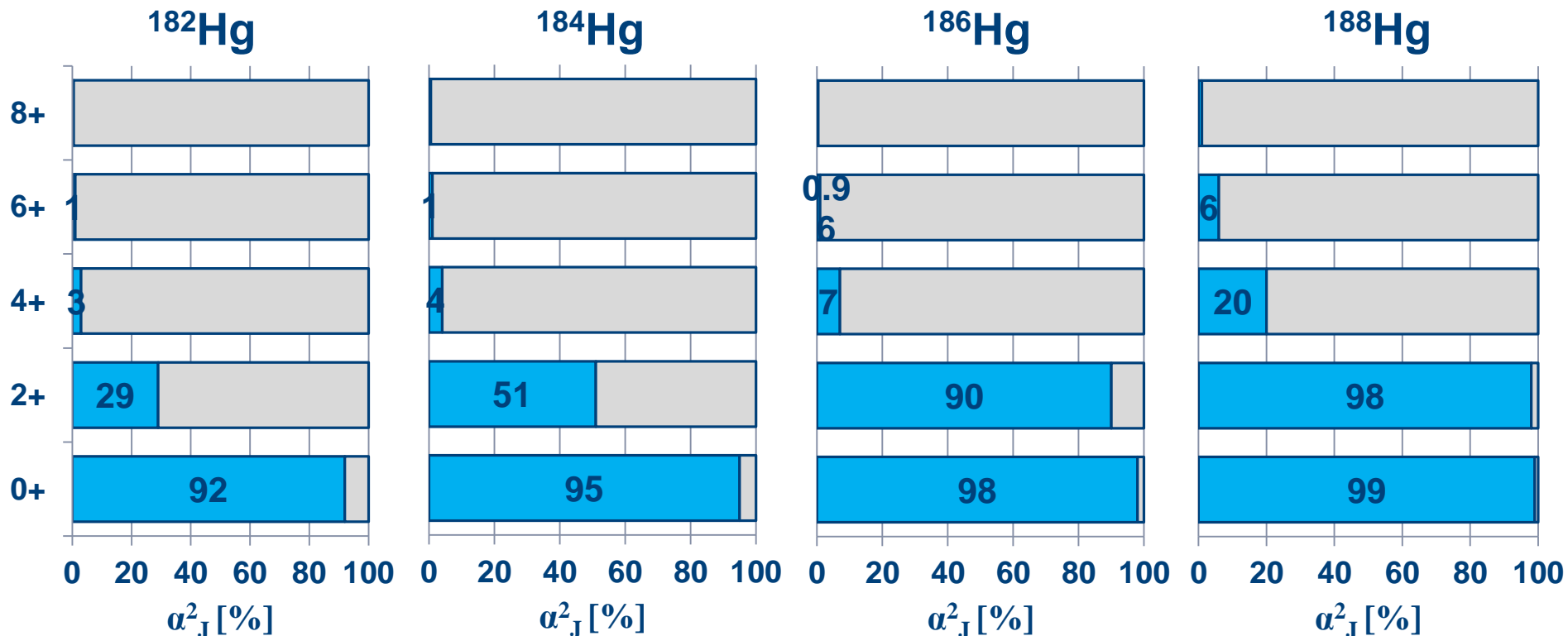
$$|2^+_{2>} = -\beta_2 |2^+_S> + \alpha_2 |2^+_D>$$

$$|0^+_{2>} = \beta_0 |0^+_S> - \alpha_0 |0^+_D>$$



- Two level mixing calculations

Deduced from fitting the higher-lying levels using the variable moment of inertia model



- 0+ states: only weakly mixed for all $^{182-188}\text{Hg}$
- 2+ states: mixing is changing from 29% up to 98%
- 4+ states: dominant deformed configuration for all $^{182-188}\text{Hg}$

- Two level mixing calculations

^{182}Hg : $\alpha_2^2 = 29\%$ $\alpha_4^2 = 3\%$
 ^{184}Hg : $\alpha_2^2 = 51\%$ $\alpha_4^2 = 4\%$
 ^{186}Hg : $\alpha_2^2 = 90\%$ $\alpha_4^2 = 7\%$
 ^{188}Hg : $\alpha_2^2 = 98\%$ $\alpha_4^2 = 20\%$

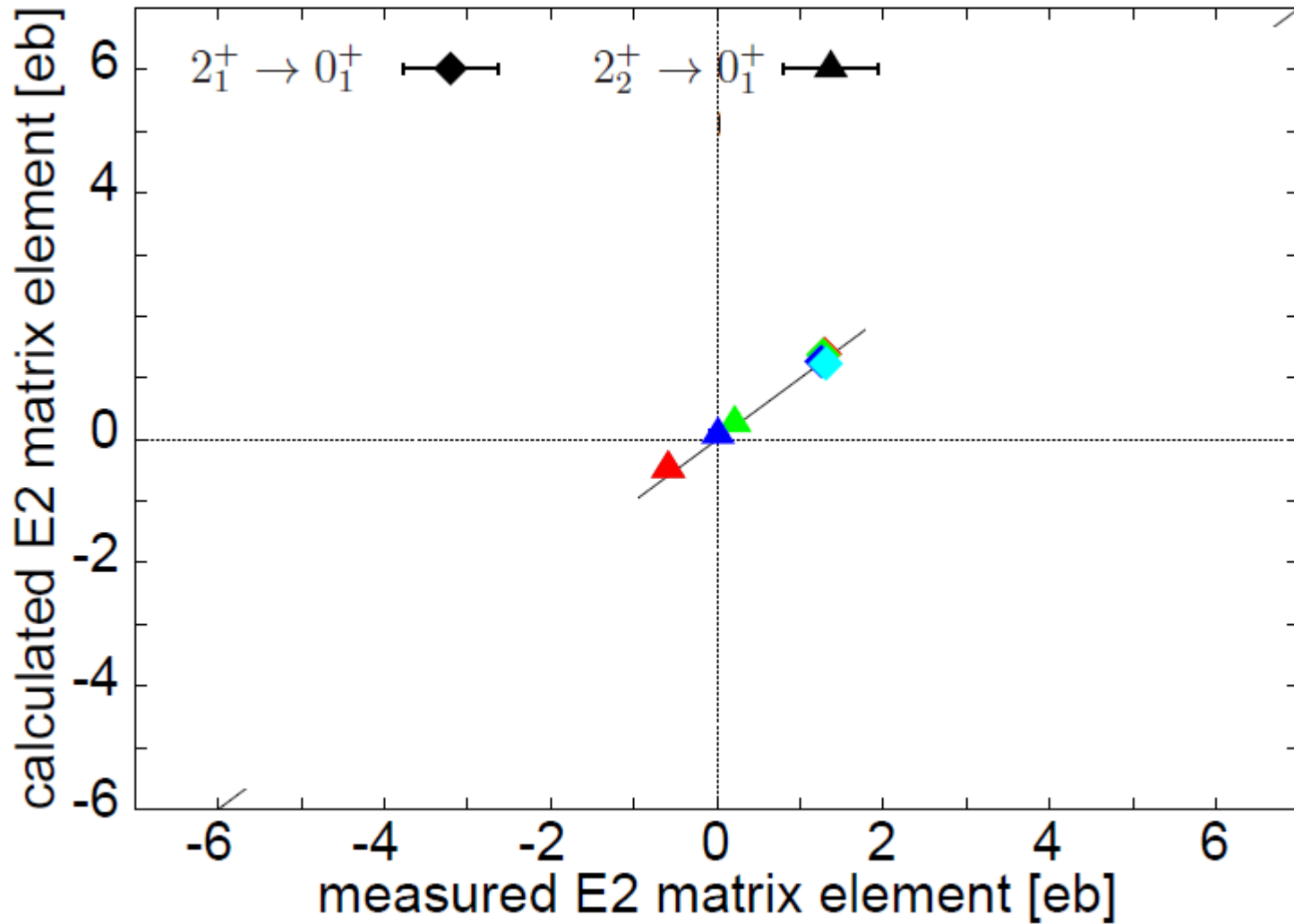
fitted un-mixed ME2's:

$$A = \langle 0^+_S || E2 || 2^+_S \rangle = 1.2 \text{ eb}$$

$$B = \langle 0^+_D || E2 || 2^+_D \rangle = 3.3 \text{ eb}$$

$$C = \langle 2^+_S || E2 || 2^+_S \rangle = 1.8 \text{ eb}$$

$$D = \langle 2^+_D || E2 || 2^+_D \rangle = -4.0 \text{ eb}$$



- Two level mixing calculations

^{182}Hg : $\alpha_2^2 = 29\%$ $\alpha_4^2 = 3\%$
 ^{184}Hg : $\alpha_2^2 = 51\%$ $\alpha_4^2 = 4\%$
 ^{186}Hg : $\alpha_2^2 = 90\%$ $\alpha_4^2 = 7\%$
 ^{188}Hg : $\alpha_2^2 = 98\%$ $\alpha_4^2 = 20\%$

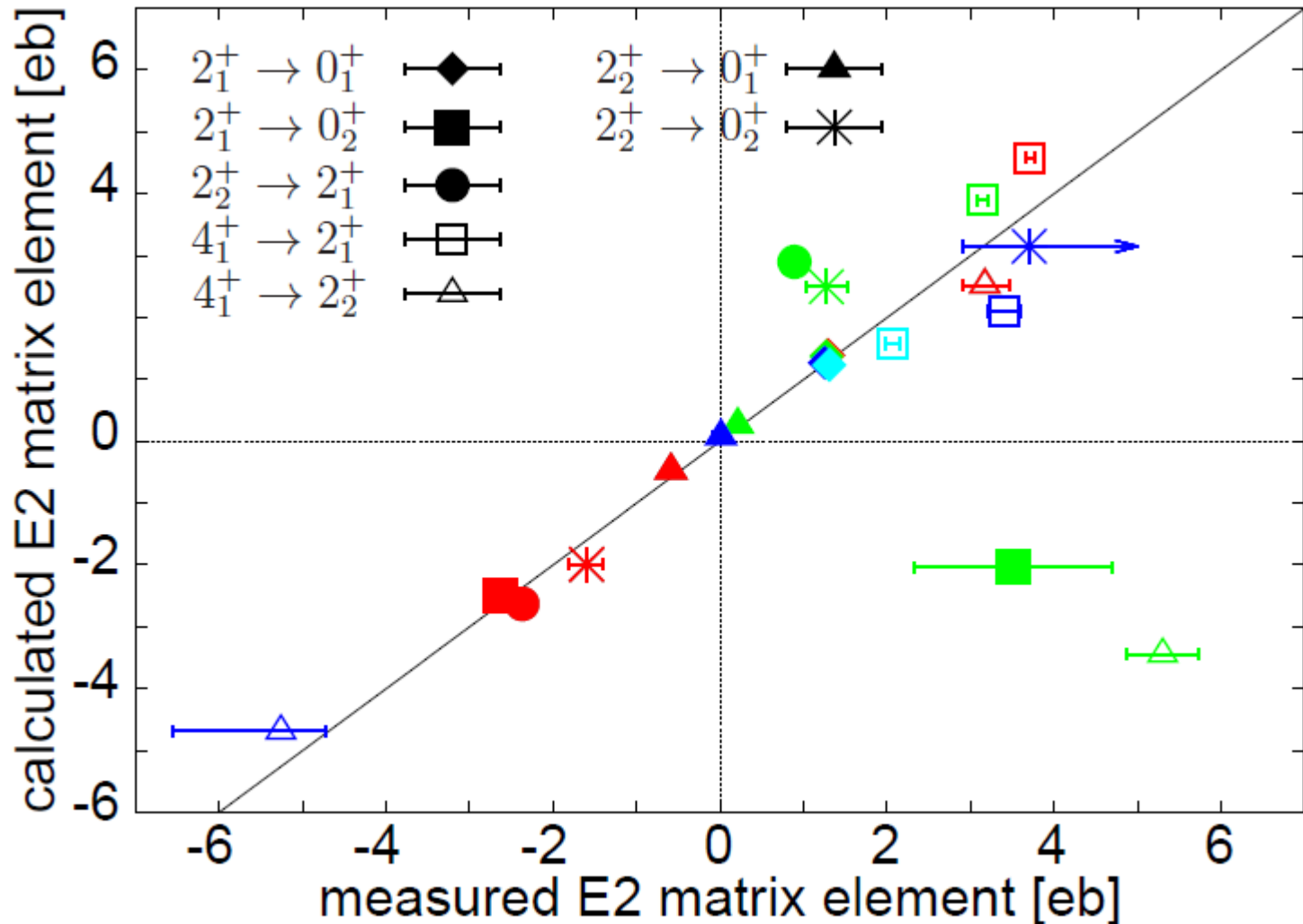
fitted un-mixed ME2's:

$A = \langle 0^+_S || E2 || 2^+_S \rangle = 1.2 \text{ eb}$

$B = \langle 0^+_D || E2 || 2^+_D \rangle = 3.3 \text{ eb}$

$C = \langle 2^+_S || E2 || 2^+_S \rangle = 1.8 \text{ eb}$

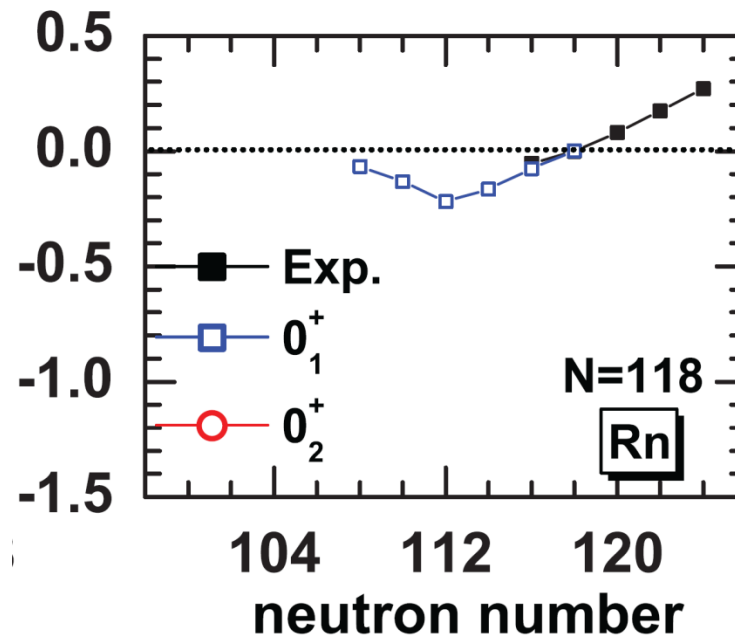
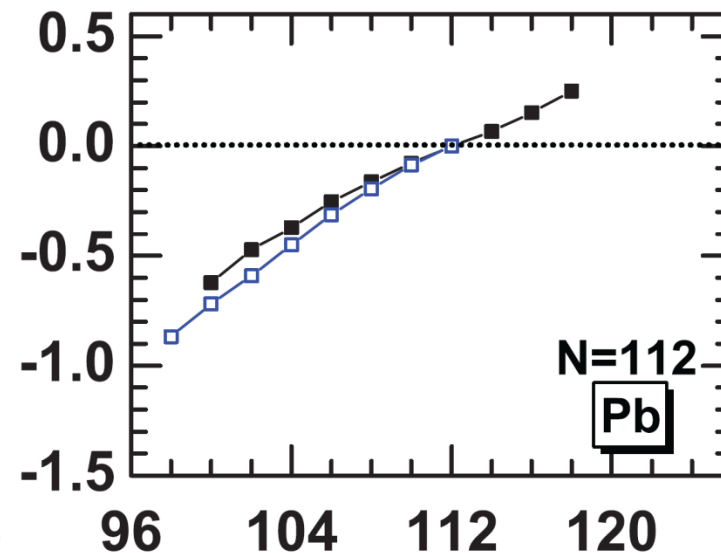
$D = \langle 2^+_D || E2 || 2^+_D \rangle = -4.0 \text{ eb}$



Comparison with beyond mean-field calculations

$$\delta \langle r_{ch}^2 \rangle \text{ (fm}^2\text{)}$$

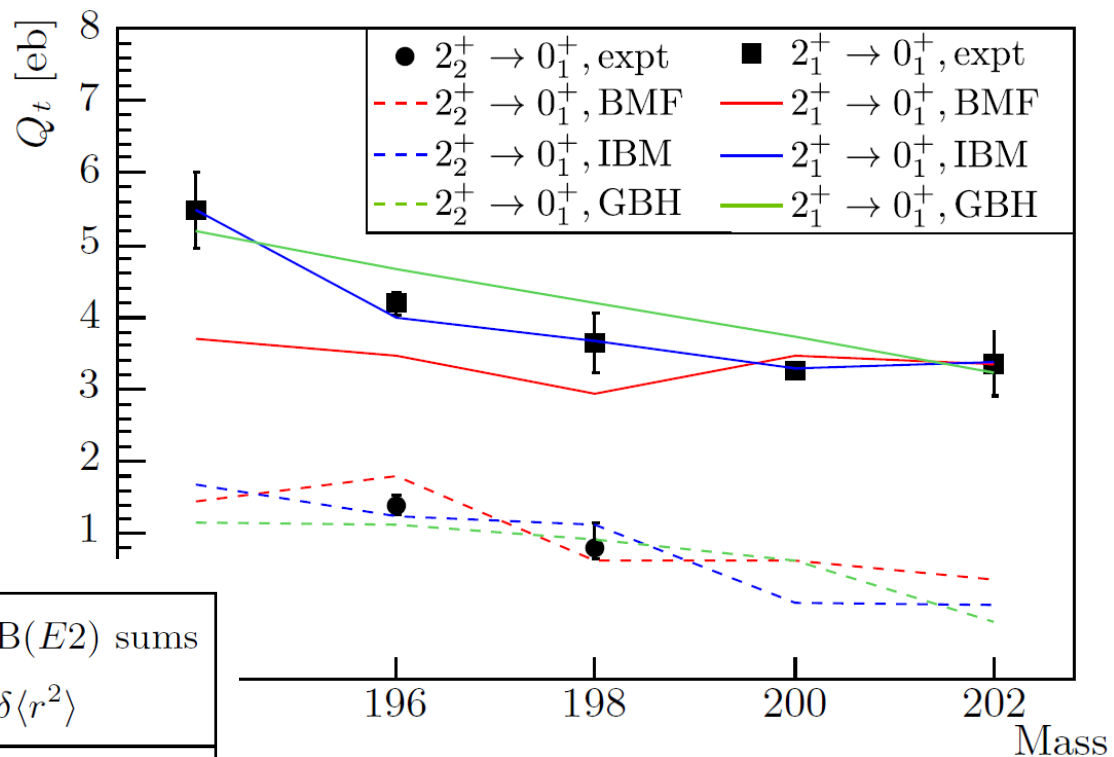
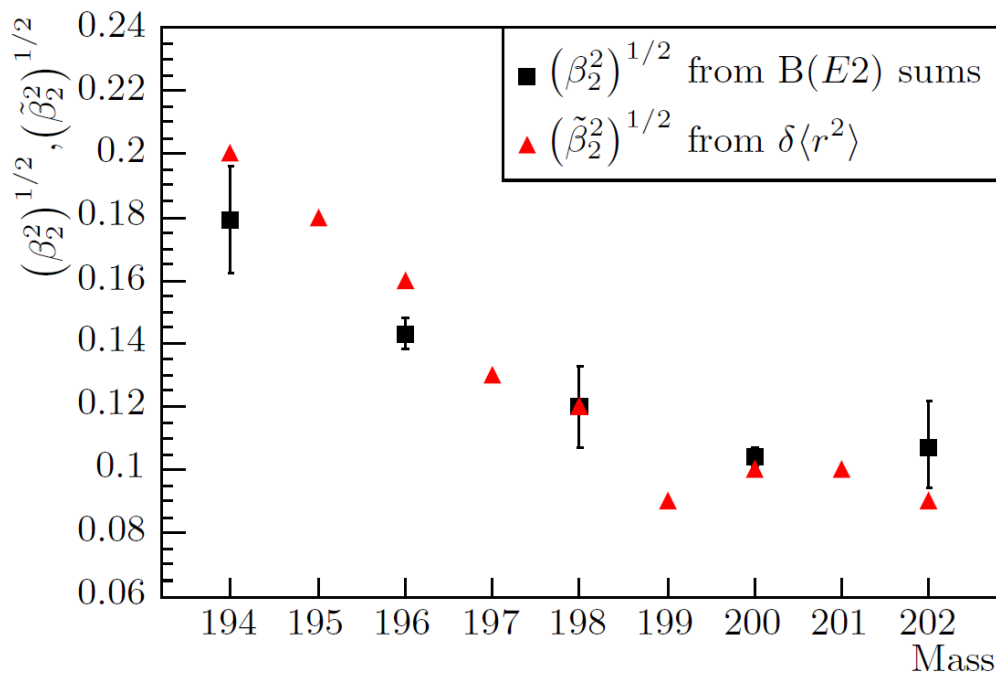
$$\delta \langle r_{ch}^2 \rangle \text{ (fm}^2\text{)}$$



Polonium isotopes

$$\langle r^2 \rangle_A \approx \langle r^2 \rangle_A^{sph} \left(1 + \frac{5}{4\pi} \tilde{\beta}_2^2 \right)$$

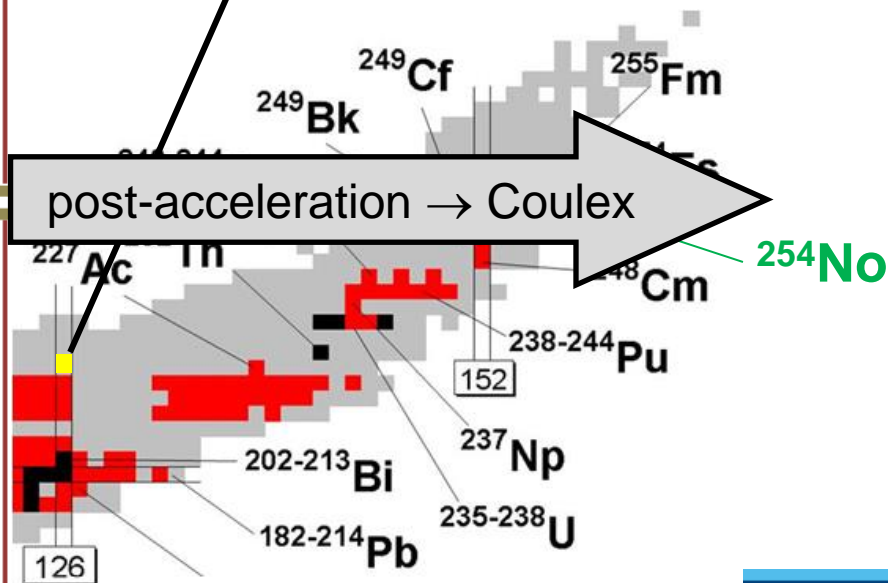
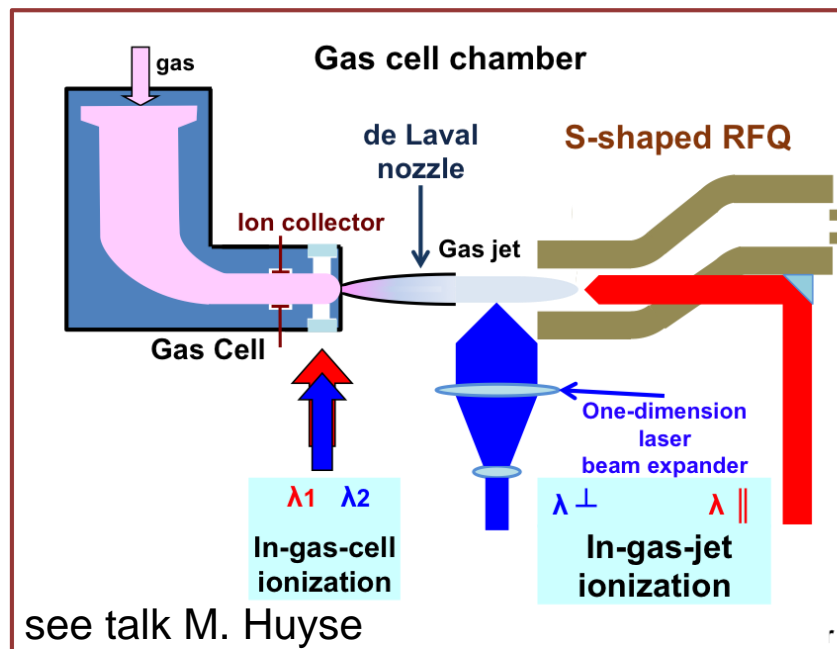
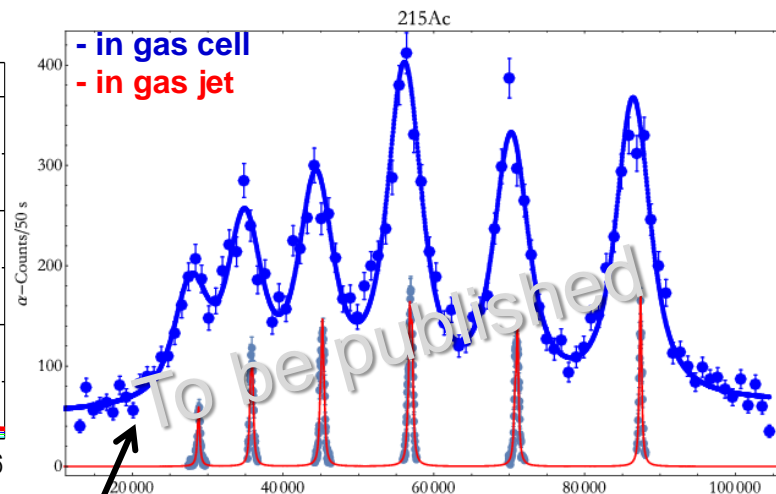
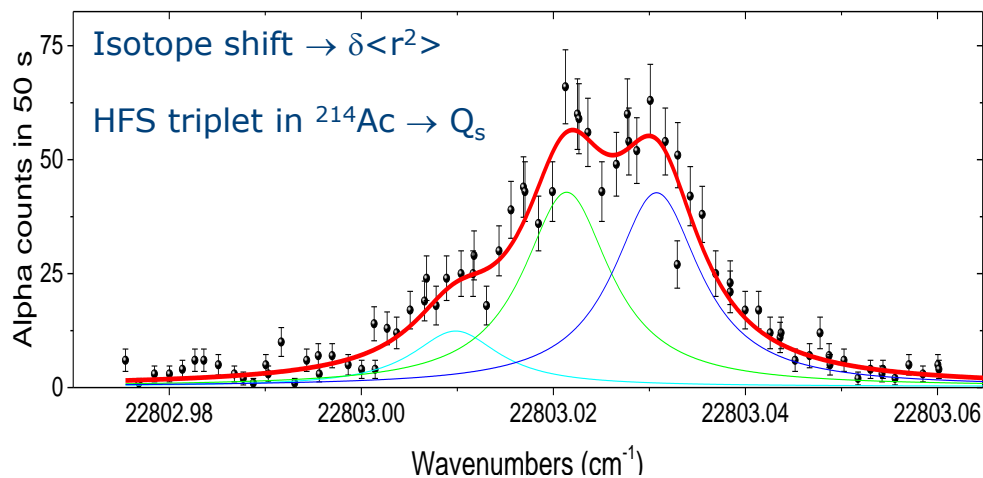
$$\sum_i |\langle 0_1^+ || E2 || 2_i^+ \rangle|^2 = \left(\frac{3}{4\pi} ZeR_0^2 \right)^2 \beta_2^2$$



$$Q_t(I_i \rightarrow I_f) = \frac{\langle I_f || E2 || I_i \rangle}{\langle I_f 020 | I_i 0 \rangle} \cdot \sqrt{\frac{16\pi}{5(2I_f + 1)}}$$

In Gas Laser Ionization Spectroscopy (IGLIS)

^{215}Ac $T_{1/2} = 0.17\text{ s}$ $J_{\pi} = (9/2^{-})$



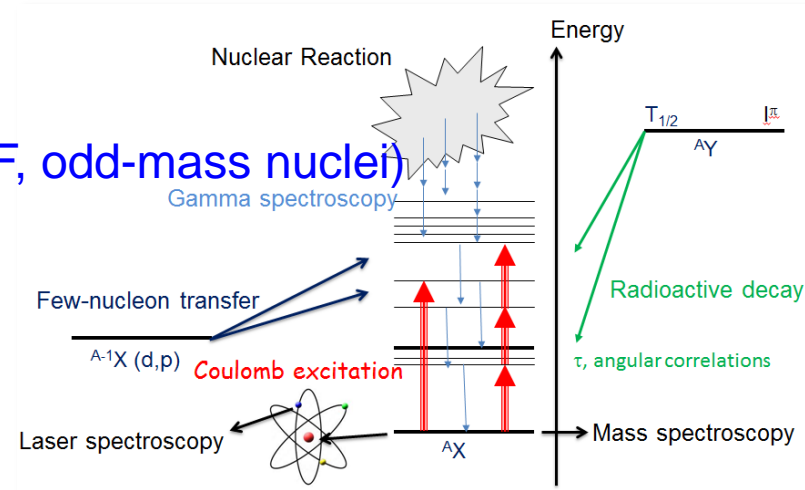
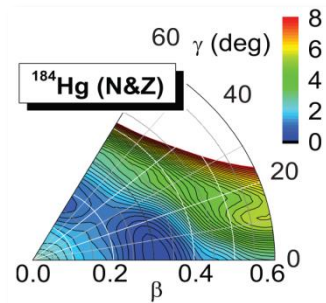
Conclusion

- Shape coexistence is firmly established (took >40 y) but the **underlying microscopic mechanisms** needs further experimental and theoretical input



- Need for **more precise** experimental data (e.g E2 matrix elements), **nucleon transfer reactions** and need to **combine** all information available (→ consistent picture)

- Further **theoretical** developments (triaxial DOF, odd-mass nuclei)



- New **opportunities** with the new facilities (see talks M. Huyse and Y. Blumenfeld)
 - actinide region: laser spectroscopy and possibly post-accelerated beams

Reflections



La plage des savants, Cargèse (1991)