

Magnetic measurement challenges for compact & low-consumption magnets

M Buzio on behalf of the Magnetic Measurement Section

(CERN, Technology Department, Magnet, Cryostats and Superconductors Group)

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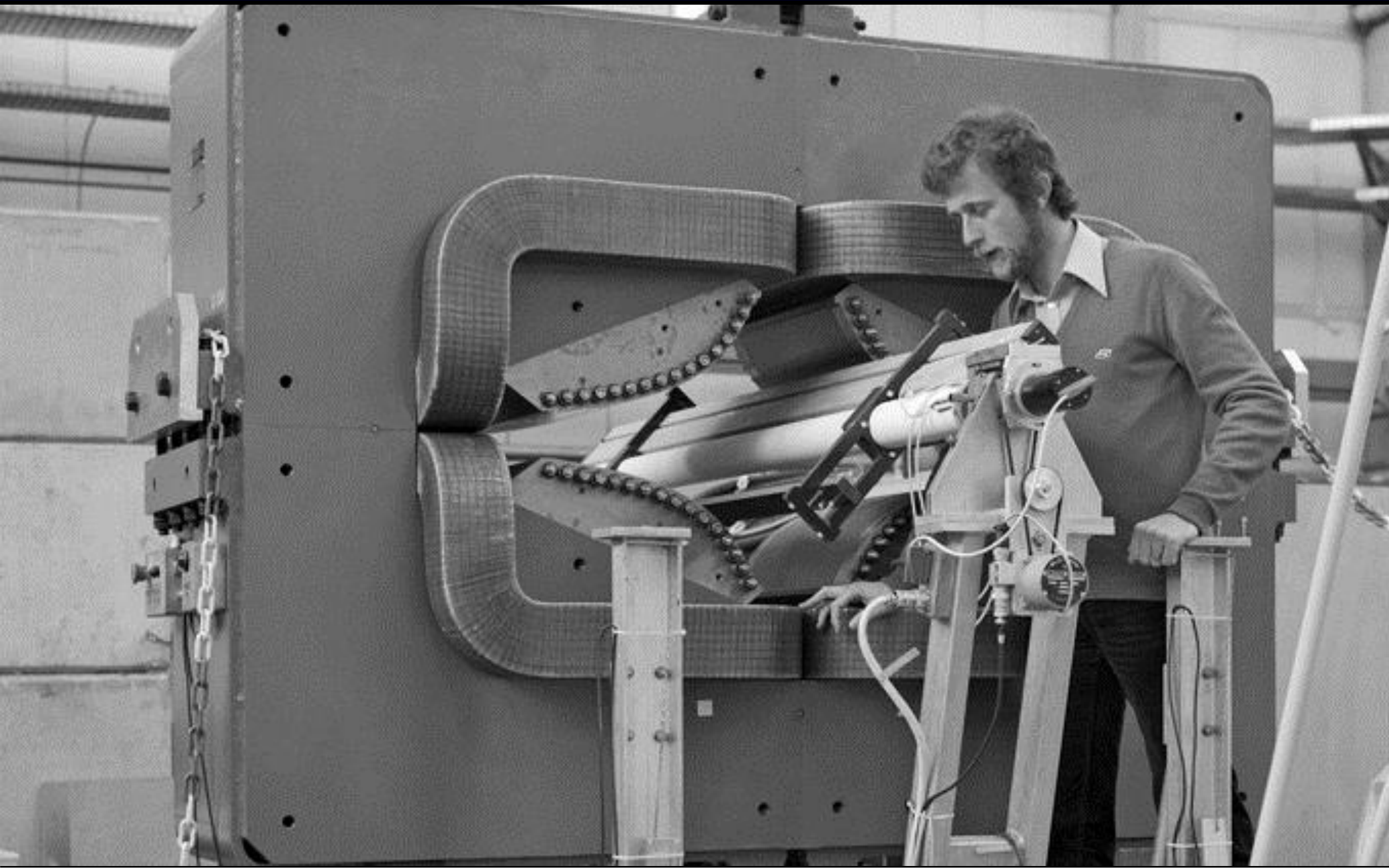
1 – Introduction

Review of major measurement techniques:
Hall probes, fixed/rotating coils, stretched wires

2 – Challenges

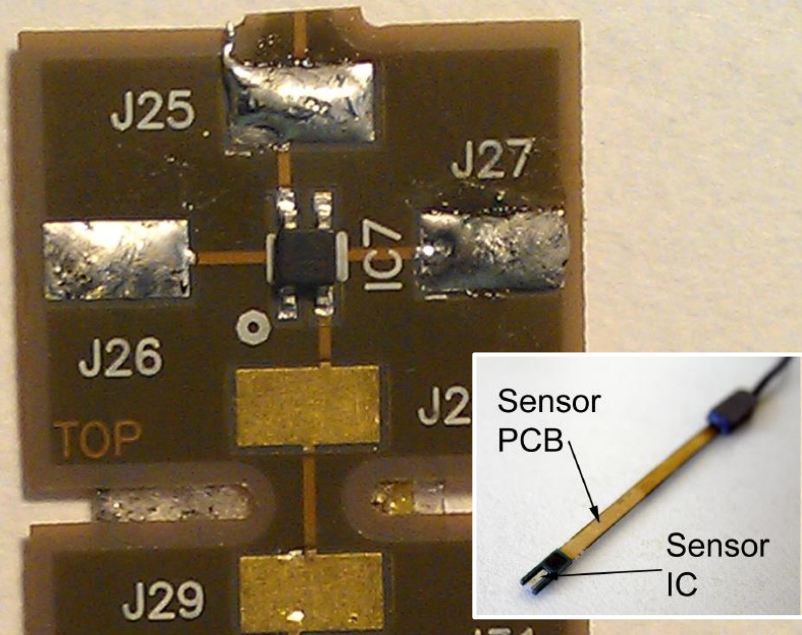
Main issues arising small apertures & low RMS and their solution

3 – Summary & References

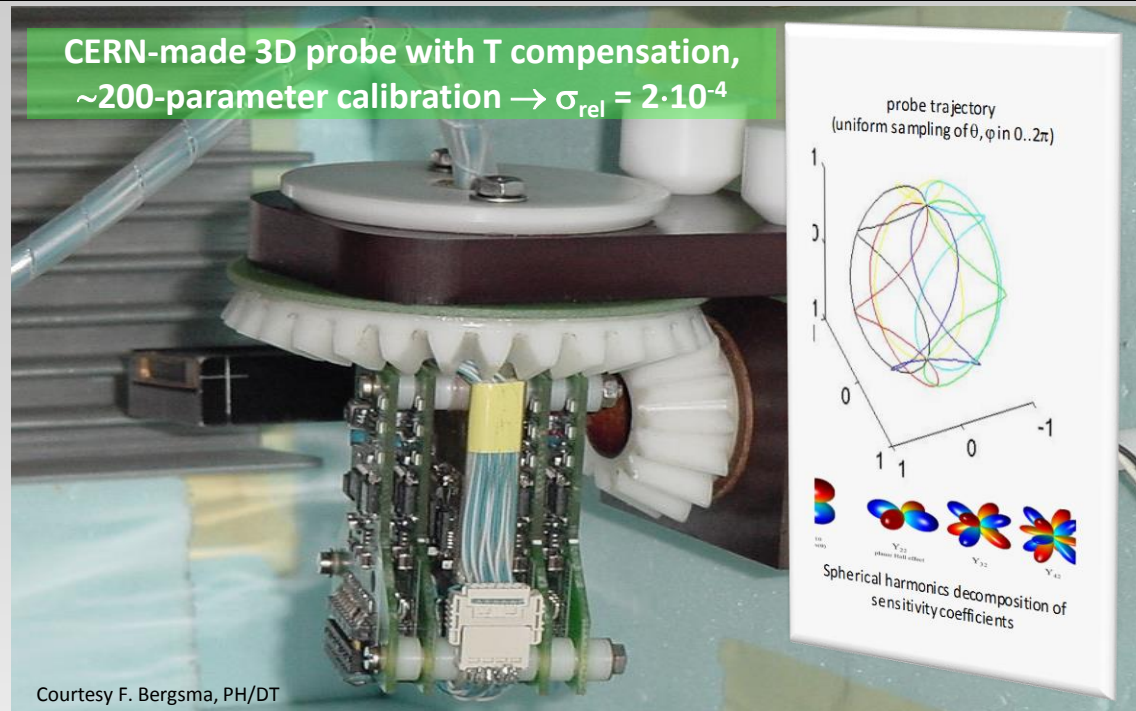


Measurement methods

Commercial Hall probe sensors



CERN-made 3D probe with T compensation,
 ~ 200 -parameter calibration $\rightarrow \sigma_{rel} = 2 \cdot 10^{-4}$

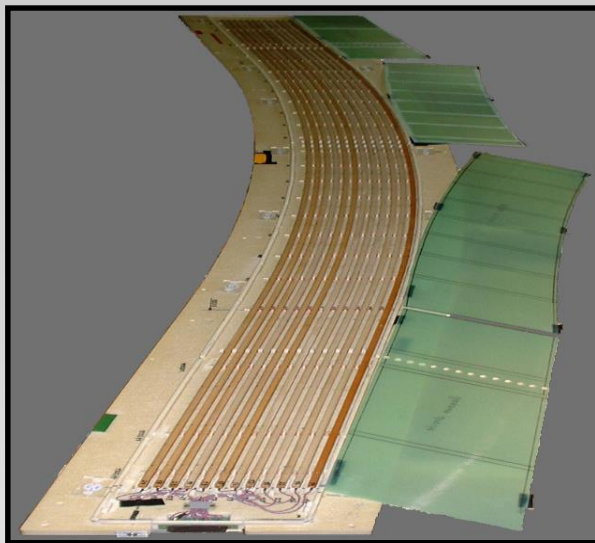


Advantages

- easy to use, readily available on the market (typical accuracy 1%)
- $\sim 1 \text{ mm}^2$ sensor area \rightarrow best suited to detailed maps (e.g. fringe fields)

Drawbacks

- uncertainty better than a few $\sim 10^{-3}$ require complex and frequent calibration + temperature control or compensation
- small probe sensitive to local effects \rightarrow precise integrals require *many* points
- precise positioning in translation and angle requires expensive mechanics
- errors in multipolar/fast ramping fields



Advantages

- natural choice for time-varying/integral fields (S/N improves with increasing size, B , dB/dt)
- DC rotating coil: one turn (0.1~1 s) → full characterization of the field integral within the spanned volume: **field strength, harmonics, direction and axis**
- typical uncertainty: absolute 10^{-4} (straight or PCB coils), harmonics 10^{-5} , resolution 10^{-6}

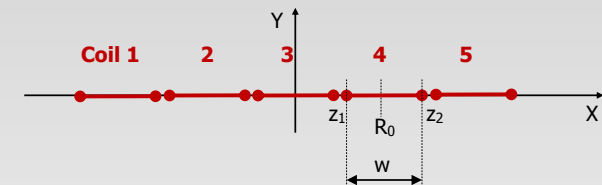
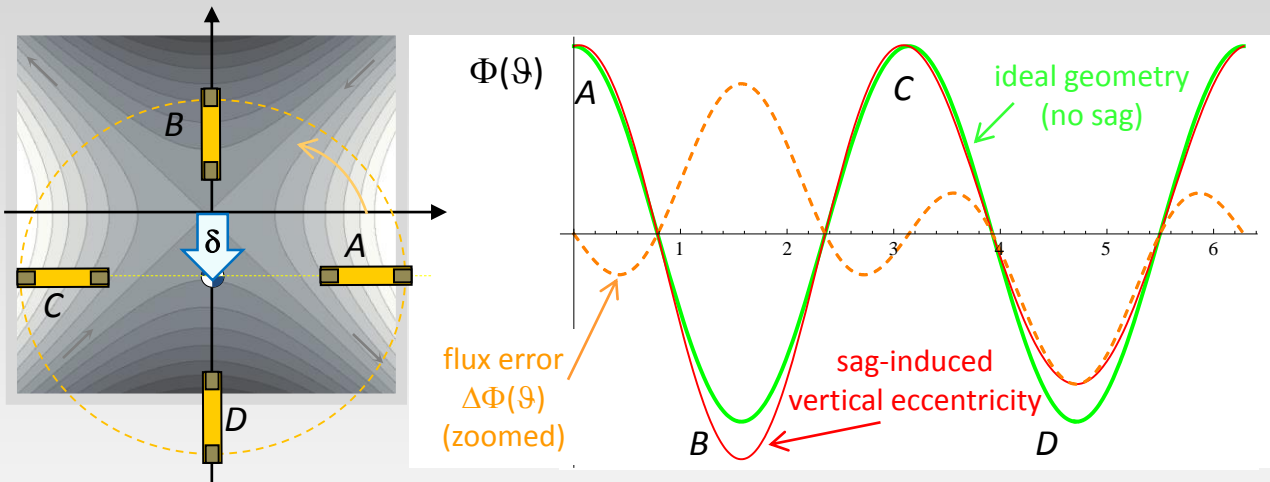
Drawbacks

- **not commercially available**, specialized in-house winding (or PCB design) required
- special techniques required for **strongly curved, or large aspect-ratio** gap magnets
- **expensive mechanics** (non-magnetic, non-conducting shaft, motor) and **top-quality electronics** (digital integrators, programmable amplifiers, angular encoders) necessary to get good results



Coil bucking

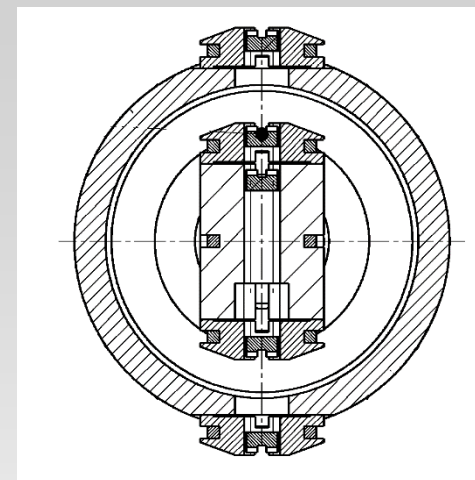
- The accuracy of higher harmonics measured by individual coils may be affected by geometry errors
- Solution = **coil bucking** (or compensation): suitable linear combinations of coil signals cancel out the sensitivity to the main (and lower) harmonics → **robustness to mechanical imperfections**
- Example: in a perfect quadrupole, average gravity-induced sag δ on a radial coil → flux error including mainly B_1 and B_3 components. A four-coil series/anti-series combination cancels out B_2 sensitivity → error-free harmonic measurement



Sensitivity coefficient	Coil 1	Coil 2	Coil 3	Coil 4	Coil 5 (spare)	Bucked coil: Linear combination 1-2-3+4
κ_1	A	A	A	A	A	0
κ_2	-2Aw	-Aw	0	Aw	2Aw	0
κ_3	$\frac{49}{12}Aw^2$	$\frac{13}{12}Aw^2$	$\frac{1}{12}Aw^2$	$\frac{13}{12}Aw^2$	$\frac{49}{12}Aw^2$	$4Aw^2$

- Arbitrary *static* coil imperfections: no major concern (effective geometry can be calibrated)
- Position- or time-dependent **transversal imperfections** → **errors \propto harmonic n =main order**
- Position- or time-dependent **torsional imperfections** → **errors \propto harmonic n =main order -1**
- Coil design objective: $\kappa_{\text{main}} = \kappa_{\text{main}-1} = 0$, maximize $|\kappa_n|$ with $n > \text{main order}$
- Additional benefit: **common mode rejection, improved S/N** (requires separate amplification)

5× tangential coils



- large number of coils made to pick well-matched equivalent surfaces
- coil parallelism measured inside a large dipole
- rotation radius measured inside a large quadrupole

retro-reflector

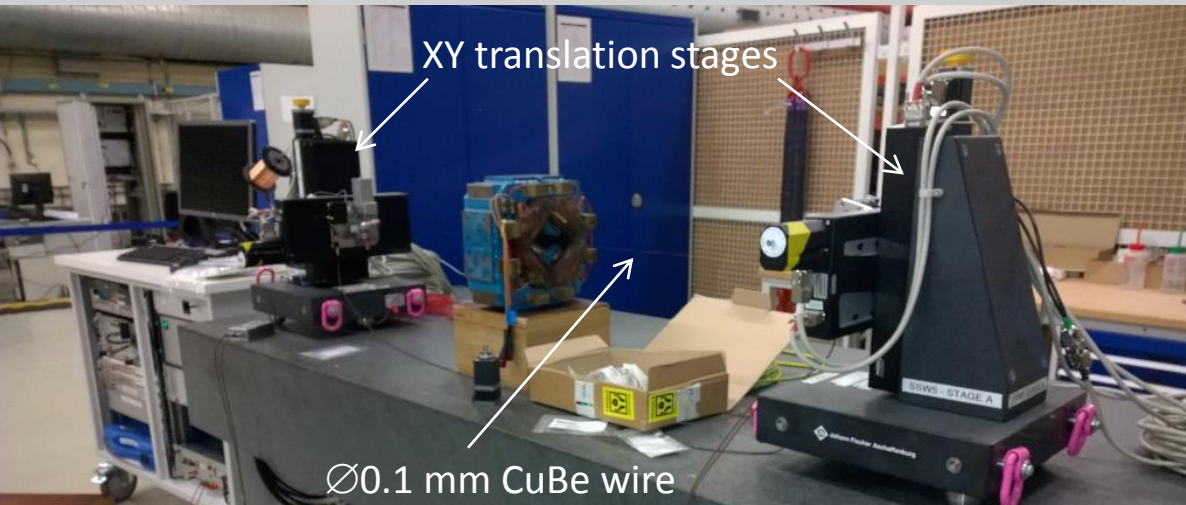
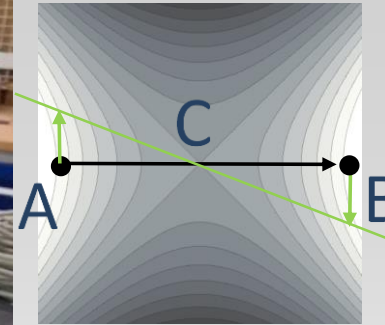
Ø150 mm

G10 support tube

Single Stretched Wire

3x multi-mode systems at CERN (based on FNAL's units used as a reference for LHC cryomagnets)

classic measurement
in a quadrupole:
if $\int_A^B B_y dx = 0$,
then C=magnetic center



Advantages

- unique flexibility: the same sensor adapts to **any size, shape and length** of magnet gap (limited by the range of the translation stages)
- unique capability to measure **longitudinal center + pitch and yaw axis** angles in lenses and solenoids (counter-directional wire movements)
- unique **sub-micron sensitivity for axis localization** in vibrating mode at resonance
- **metrological reference** for integrated field strength, axis and direction in high-field magnets
- Very promising ongoing R&D:
 - integrated harmonics in vibrating mode
 - longitudinal field profiles from measurements in vibrating mode at multiple frequencies

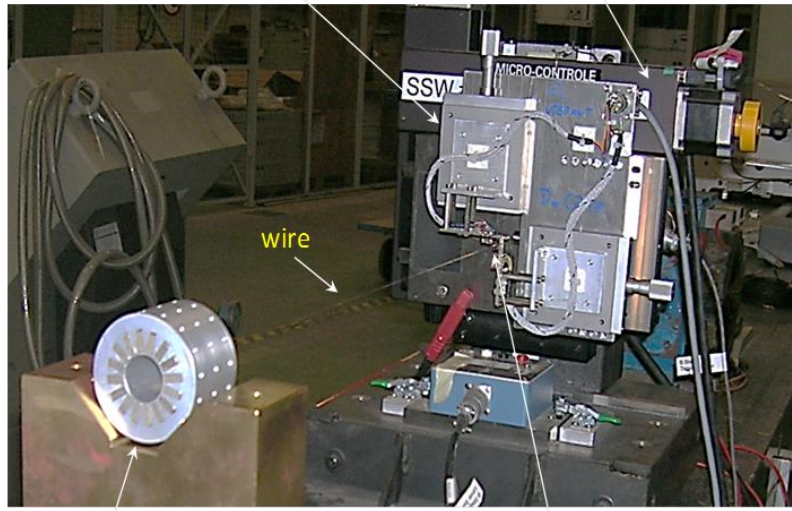
Drawbacks

- Equivalent to 1 turn-coil only → **low sensitivity** of field integrals in short/weak magnets



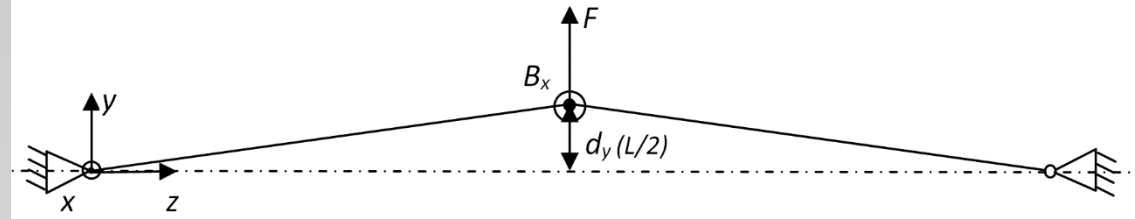
Recent development: oscillating stretched wire

XY micrometric stages XY long-range translation stages



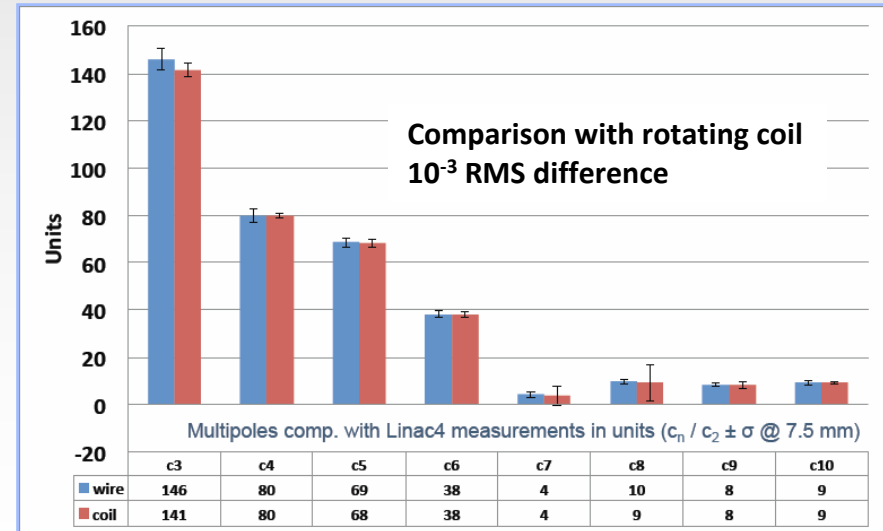
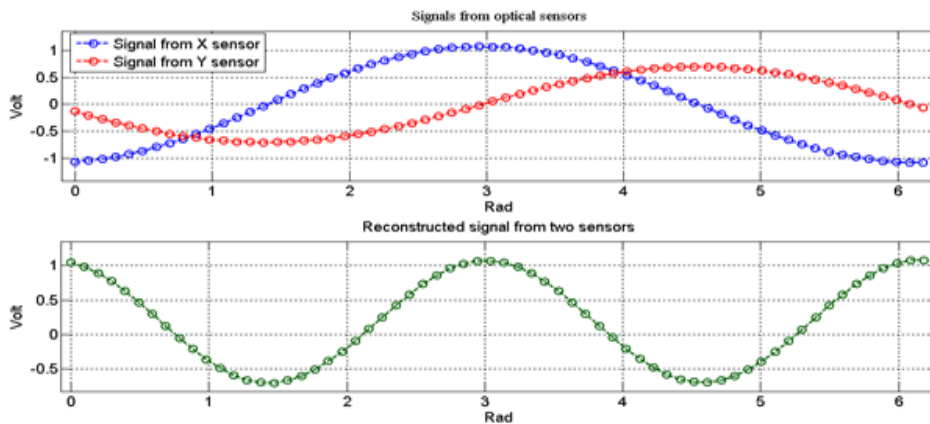
Linac4 R1 PMQ

XY optical wire position detectors



- Stretched-wire system with AC current passed through the wire \rightarrow Lorentz force \rightarrow oscillation \propto BdL
- zero amplitude = wire on magnetic axis
- **ideal method for very small aperture** magnets (CLIC)
- integrated harmonics by stepwise scan around a circle (quasi-static regime \rightarrow insensitive to frequency fluctuations)

XY vibration amplitude = $f(\vartheta)$



Challenges

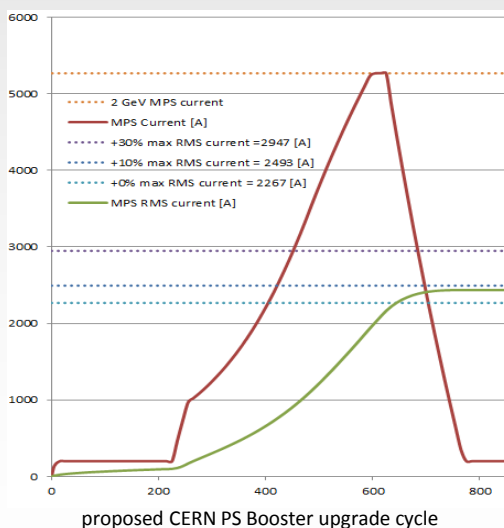
Excitation source (Cu, PM or SC), current density and power consumption: not important *per se* (assuming of course that temperature effects are controlled !)

Small gaps



- Hall probes (very) difficult
- Critical mechanical tolerances
- Reduced n. of turns, weak signals

Low RMS



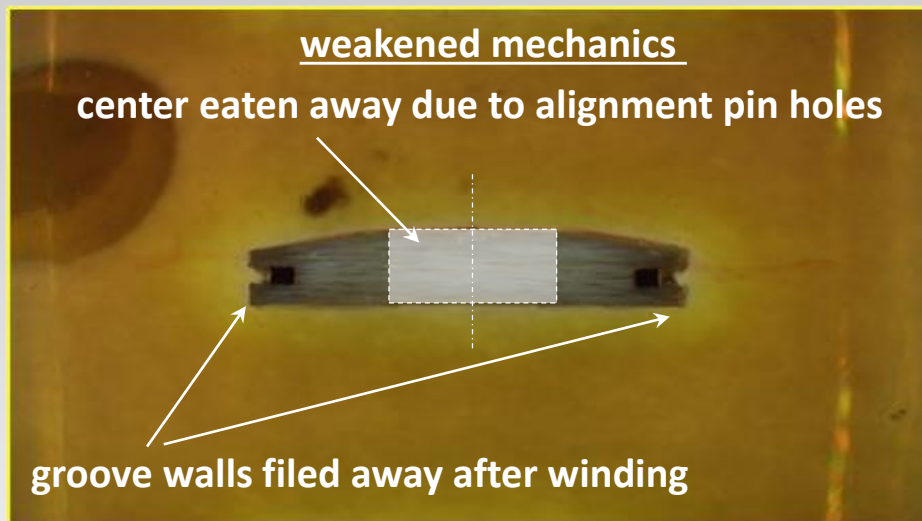
fast ramps

low flat-bottom

- Hall probes hardly possible at all (bandwidth, dynamic range)
- Good coil signals, but RC complicated (remanent lost)
- Extrapolation from low field measurements
- current/magnet cycle reproducibility issues

Small gaps

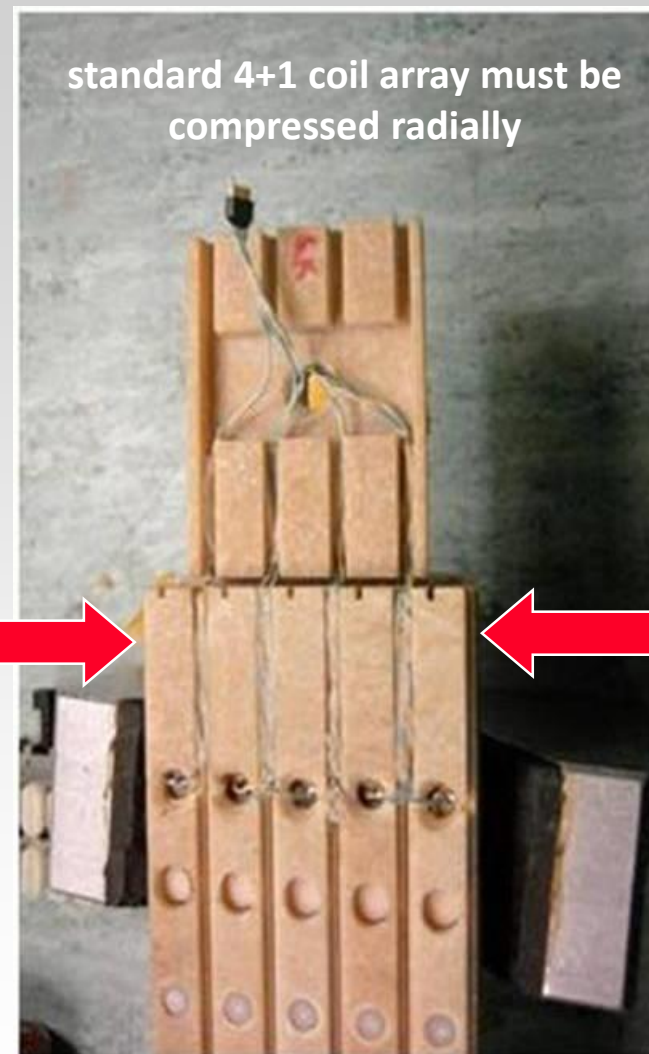
Ø38 mm coil array approaches the limit of traditional fabrication techniques



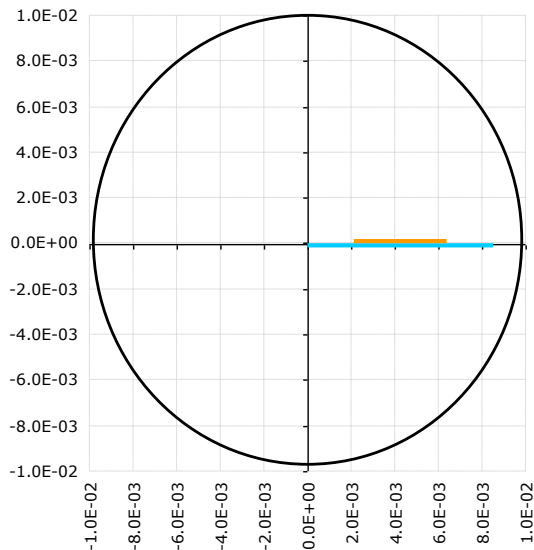
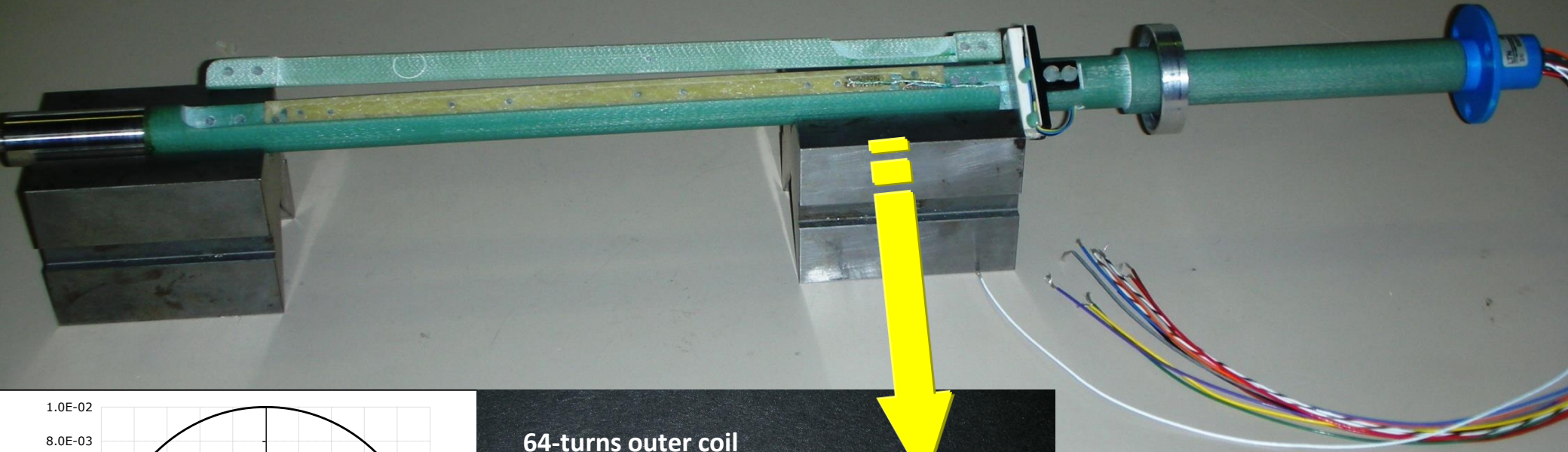
up to 4/5 waste (also due to fragile 20-wire cable) !



additional size constraint: standard connection PCB



Head currently used to measure prototype and pre-series Linac 4 quads



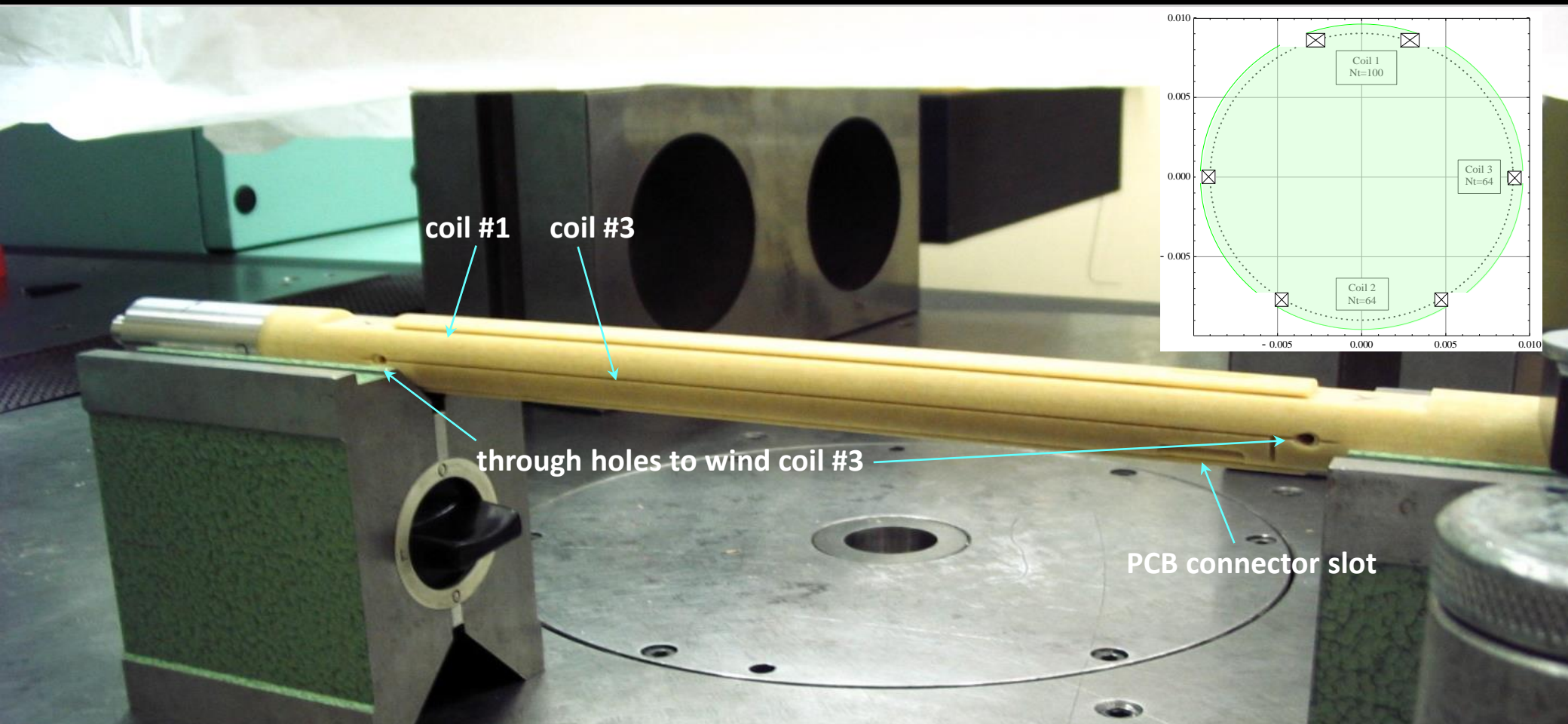
64-turns outer coil
(absolute measurement)

32-turns inner coil
(compensated measurement)

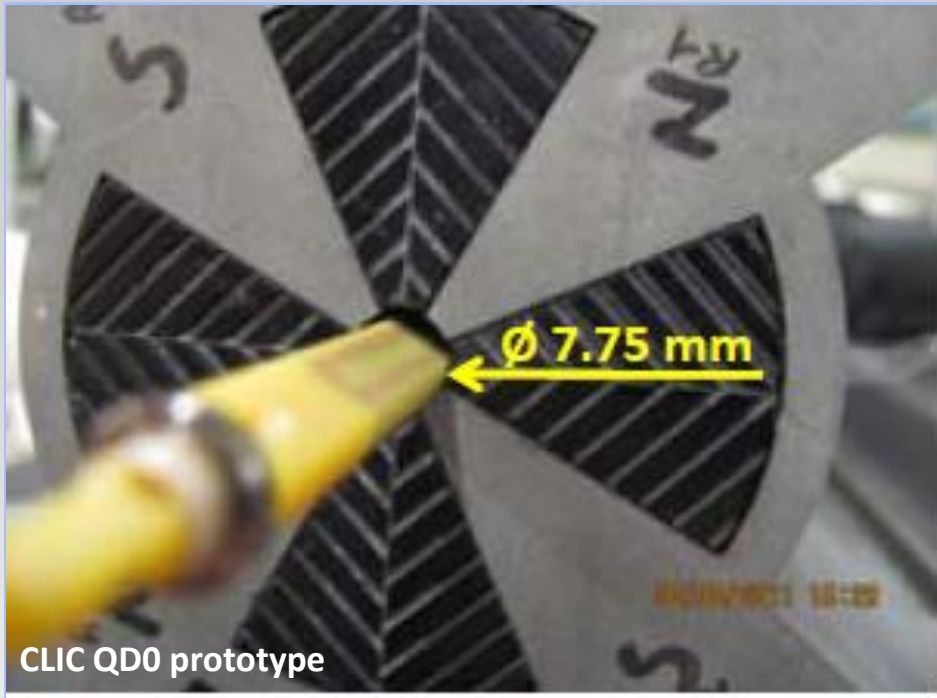
- $\varnothing 19 \times 200$ mm head based on Linac2 design (flat multi-wire cable higher winding density)
- 2 nested coils with $B_1 + B_2$ bucking (same A_{TOT} , R_0)
- longitudinal variations up to **6% width, 20 mrad twist**



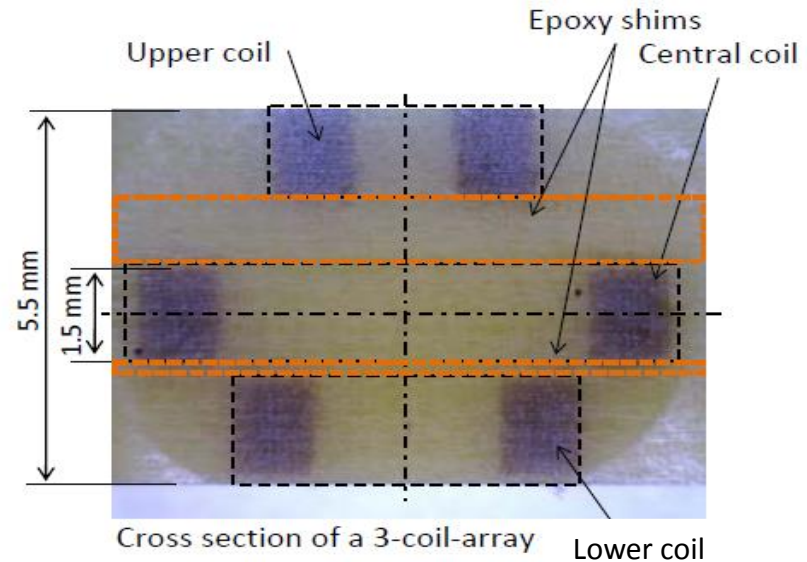
Linac 4 test bench – Coil prototype #2



- Skillful manual winding operation required
- Effective area: longitudinal variations **0.2%**
- Effective radius: longitudinal variations **1%** (outer coils), 0.1 mm (central coils)
- Customary average area/radius calibration does not work very well: a more laborious *in-situ* calibration inside the target magnet is essential to get accurate results

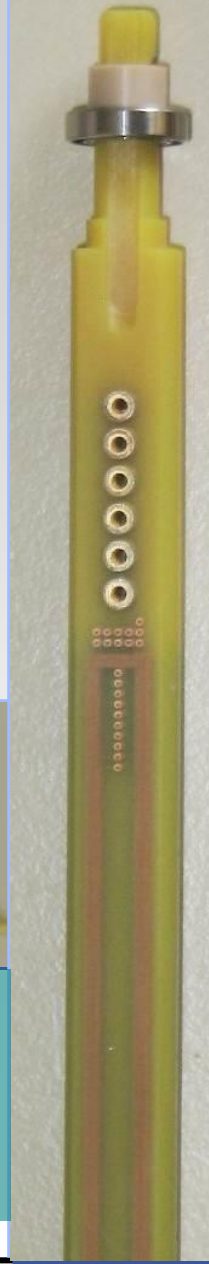


3x 64-layer PCB stack coils



200-turn coils (1/6 density w.r.t. 30 μ m wire)

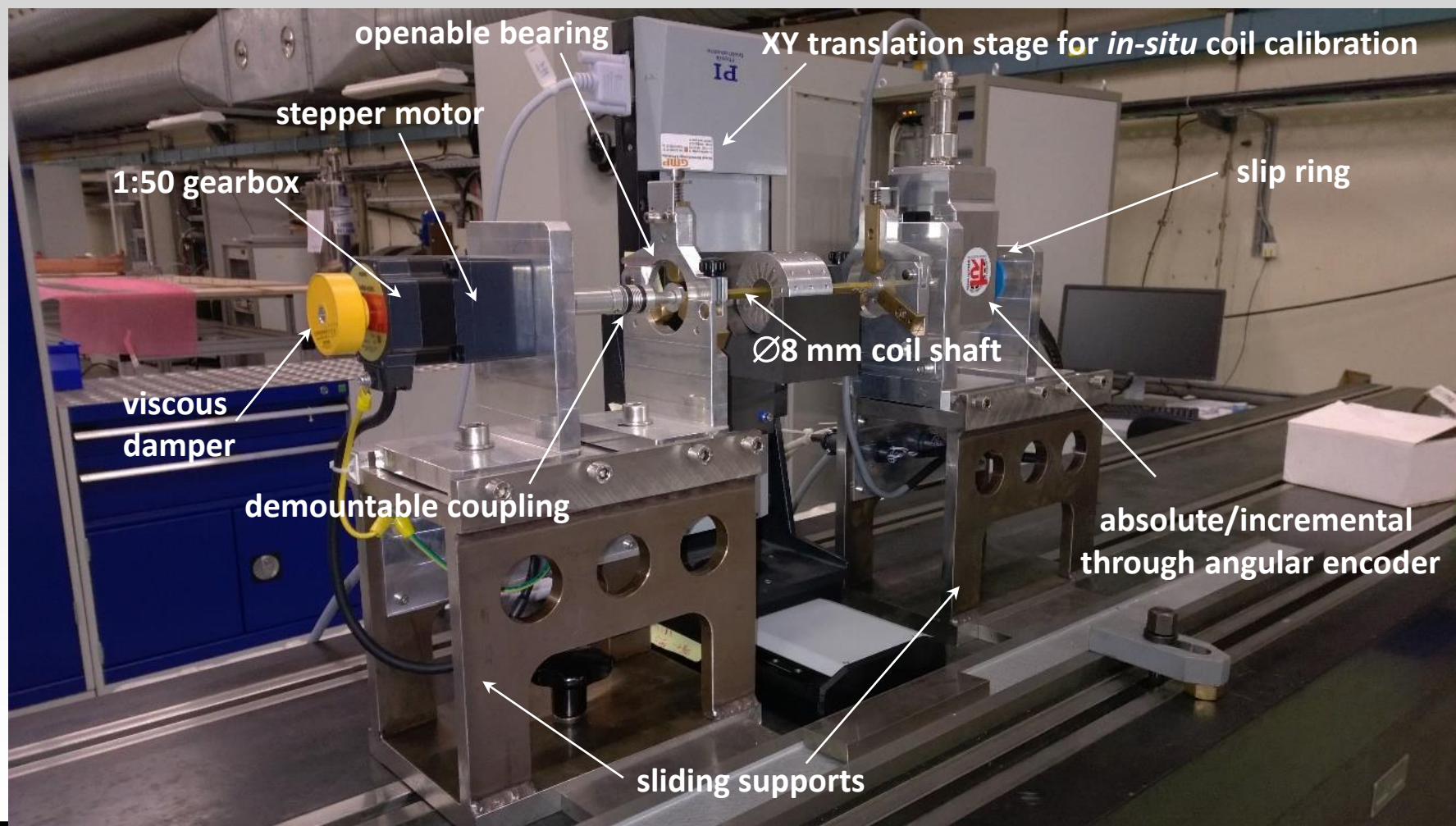
3-coil dipole- and quadrupole- bucked array can be chained to measure long magnets at once



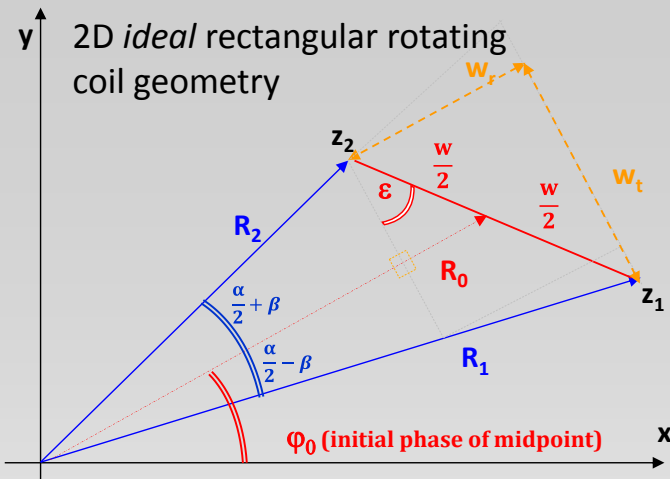
- Industrial PCB production; difficult machining to install ball bearings required
- Effective area: reproducibility $3 \cdot 10^{-4}$, longitudinal variations 0.5% (sag, twist ?)
- Effective radius: variations up to 3% (outer coils), 0.8 mm (central coil)
- Analog compensation cannot work

Linac4/CLIC harmonic coil test bench

- Developed for small-aperture permanent-magnet and fast-pulsed quads
- $\text{Ø}8/19$ mm, 150 to 400 mm long quadrupole-bucked coils
- Harmonic measurements in DC (continuously rotating coil) or fast-pulsed (stepwise rotating) mode.
- *In-situ calibration technique* to improve accuracy despite geometrical coil imperfections



Random errors in rotating coil measurements



$$\Phi(\vartheta) = \Phi(0) - \int_{t(0)}^{t(\vartheta)} V_{coil} dt = \Re \left(\sum_{n=1}^{\infty} \underbrace{\frac{N_T L}{n} (z_2^n - z_1^n)}_{K_n} \frac{C_n}{r_{ref}^{n-1}} e^{in\vartheta} \right)$$

integration constant:
lost with fixed coil measurements,
irrelevant (unphysical) for rotating coils

integration bounds set by precise angular encoder
→ rotation speed fluctuations have negligible effects

measured flux depends on both the field and coil geometry

Fourier component of the flux

absolute:

$$C_n = \frac{2}{N} r_{ref}^{n-1} \frac{\Psi_{n+1}}{A_c r_c^{n-1}} \text{ [T @ } r_{ref} \text{]}$$

Coil area

Coil rotation radius

$$\left(\frac{\sigma_{Cn}}{C_n} \right)^2 = \left(\frac{\sigma_{Ac}}{A_c} \right)^2 + (n-1)^2 \left(\frac{\sigma_{rc}}{r_c} \right)^2 + \left(\frac{\sigma_{\Psi}}{\Psi} \right)^2$$

$\sim 10^{-4}$

$10^{-3} \sim 10^{-4}$

$10^{-3} \sim 10^{-4}$

dominant term for higher harmonics

normalized:
(from absolute signal)

$$c_n = 10^4 \frac{C_n}{C_{main}} \text{ [units @ } r_{ref} \text{]}$$

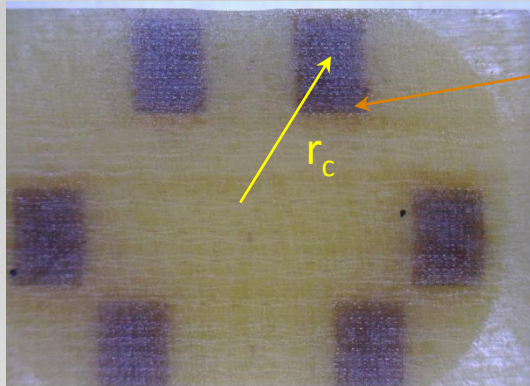
no impact of A_c

$$\left(\frac{\sigma_{cn}}{c_n} \right)^2 \approx (n - main)^2 \left(\frac{\sigma_{rc}}{r_c} \right)^2 + \left(\frac{\sigma_{\Psi}}{\Psi_{n+1}} \right)^2$$

$10^{-3} \sim 10^{-4}$

$10^{-2} \sim 10^{-3}$

dominant term



number of turns available for the coils i.e.

$$A_c \propto \emptyset^2$$

mechanical manufacturing tolerances,
static and dynamic deformations, vibrations,
alignment, temperature drifts $\rightarrow \sigma_{rc}$ constant

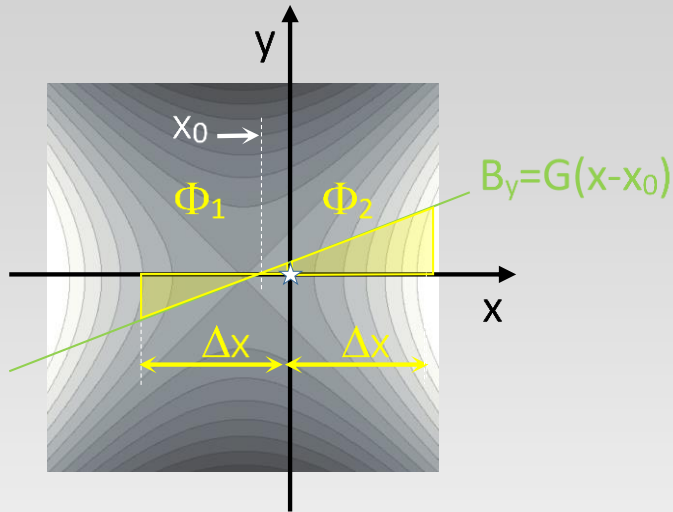
$$\left(\frac{\sigma_{Cn}}{C_n}\right)^2 = \left(\frac{\sigma_{Ac}}{A_c}\right)^2 + (n-1)^2 \left(\frac{\sigma_{rc}}{r_c}\right)^2 + \left(\frac{\sigma_{\Psi}}{\Psi}\right)^2$$

$$r_c \approx \frac{\emptyset}{3} \text{ (typically)}$$

Integrated flux change \sim constant
(peak pole field, rotation/translation speed,
magnet length being equal)

**absolute harmonic coefficient uncertainty scales with:
harmonic order, $\emptyset^{-1} \sim \emptyset^{-2}$**

Wire moved in two steps of width $\Delta x \propto \emptyset$ inside a quadrupole of unknown gradient G
 Coordinate frame offset from magnetic axis by unknown amount x_0
 Flux Φ integrated over the wire length L_w



$$\begin{cases} \Phi_1 = L_w \int_{-\Delta x}^0 B_y dx & G = \frac{\Phi_1 - \Phi_2}{\Delta x^2} \\ \Phi_2 = L_w \int_0^{\Delta x} B_y dx & x_0 = \frac{\Delta x}{2} \frac{\Phi_1 + \Phi_2}{\Phi_1 - \Phi_2} \end{cases}$$

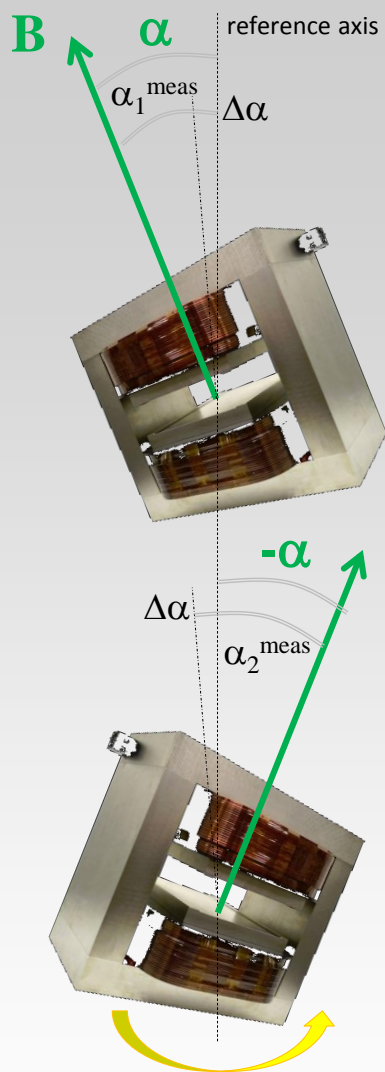
$$\begin{cases} \left(\frac{\sigma_G}{G}\right)^2 = \frac{1}{2} \left(\frac{\sigma_\Phi}{\Phi}\right)^2 + 2 \left(\frac{\sigma_{\Delta x}}{\Delta x}\right)^2 \\ \left(\frac{\sigma_{x_0}}{x_0}\right)^2 = \left(\frac{\sigma_\Phi}{\Phi}\right)^2 + \left(\frac{\sigma_{\Delta x}}{\Delta x}\right)^2 \end{cases}$$

uncertainty of gradient and magnetic axis scales with \emptyset^{-1}

Caveat: increasing too much $\Delta x \rightarrow$
 get too close to the poles, harmonic
 errors perturb the result:

$$\frac{\frac{1}{2}(\int G_x dl + \int G_y dl)}{\int G dl} = 1 + \frac{2}{3} \left(\frac{\Delta x}{r_{ref}}\right)^4 \frac{b_6}{10^4} + \frac{2}{5} \left(\frac{\Delta x}{r_{ref}}\right)^8 \frac{b_{10}}{10^4} + \dots$$





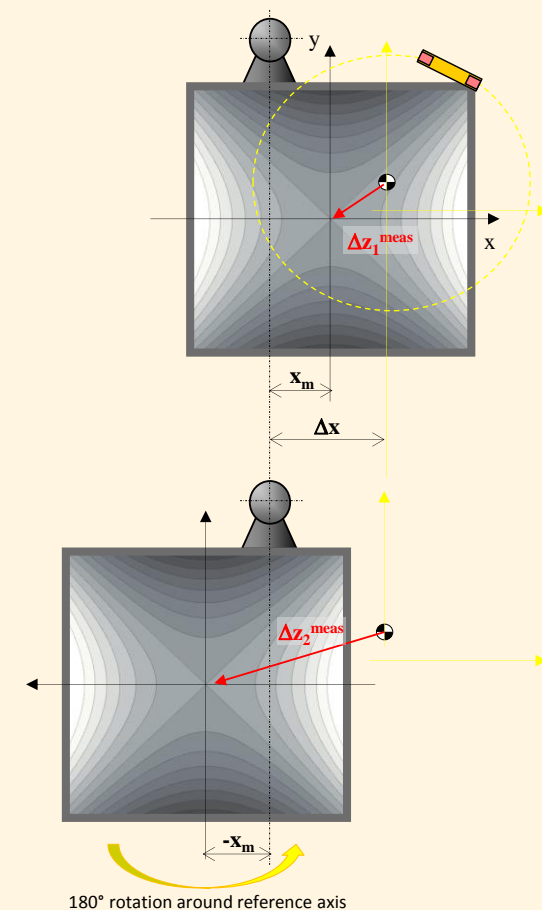
example of field direction calibration:
180° rotation around reference axis
(gravity or mechanical support)

- Simple cases: transverse and angular offsets
- Repeat the measurement by flipping the magnet around a vertical axis
- Reversing y is seldom possible, vertical offset is much harder to determine

$$\begin{cases} \alpha_1^{meas} = +\alpha - \Delta\alpha \\ \alpha_2^{meas} = -\alpha - \Delta\alpha \end{cases} \Rightarrow$$

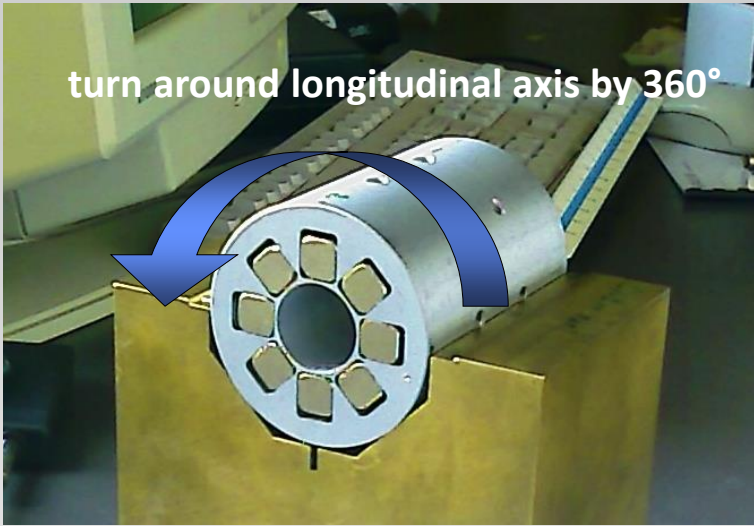
$$\begin{cases} \alpha = \frac{\alpha_1^{meas} - \alpha_2^{meas}}{2} \\ \Delta\alpha = -\frac{\alpha_1^{meas} + \alpha_2^{meas}}{2} \end{cases}$$

Systematic axis offset correction

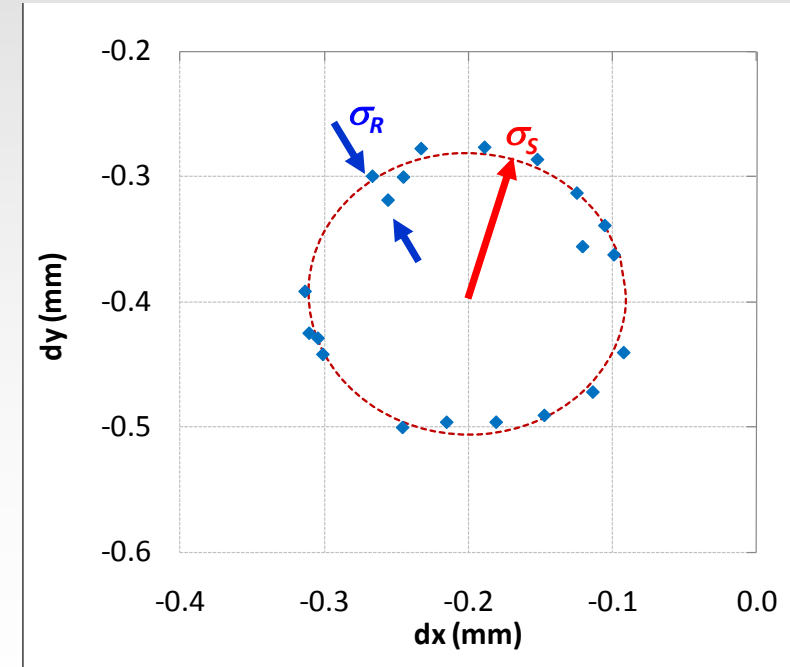
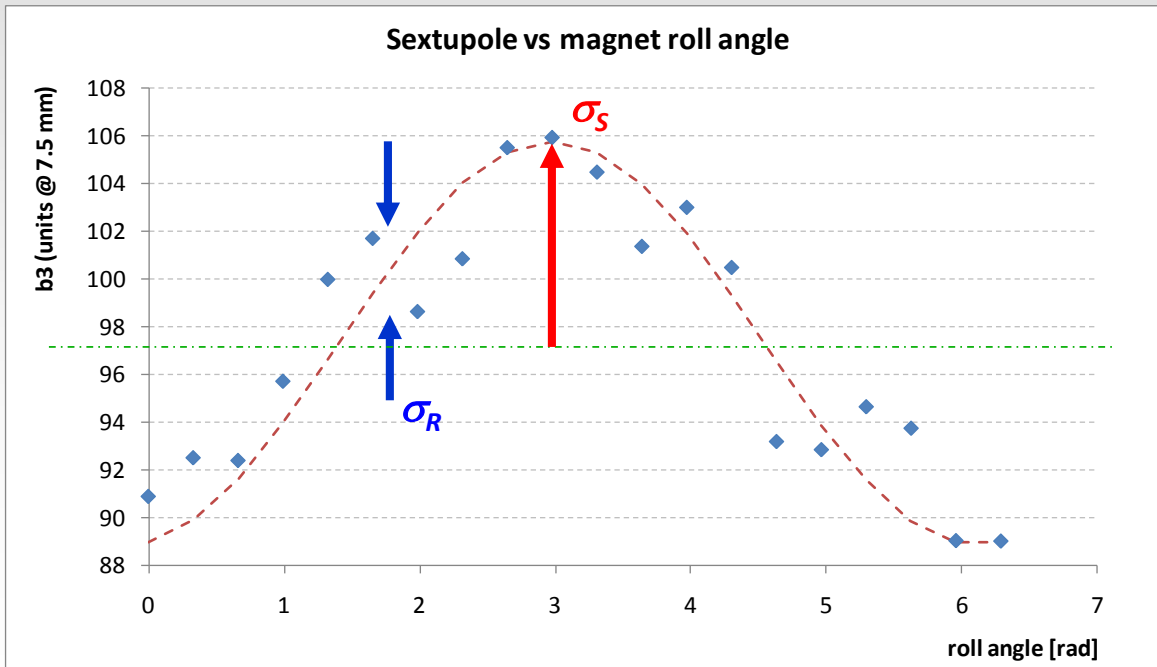


$$\begin{cases} \Delta x_1^{meas} = \Delta x - x_m \\ \Delta x_2^{meas} = \Delta x + x_m \end{cases} \Rightarrow \begin{cases} x_m = \frac{\Delta x_2^{meas} - \Delta x_1^{meas}}{2} \\ \Delta x = \frac{\Delta x_2^{meas} + \Delta x_1^{meas}}{2} \end{cases}$$

Estimation of systematic errors

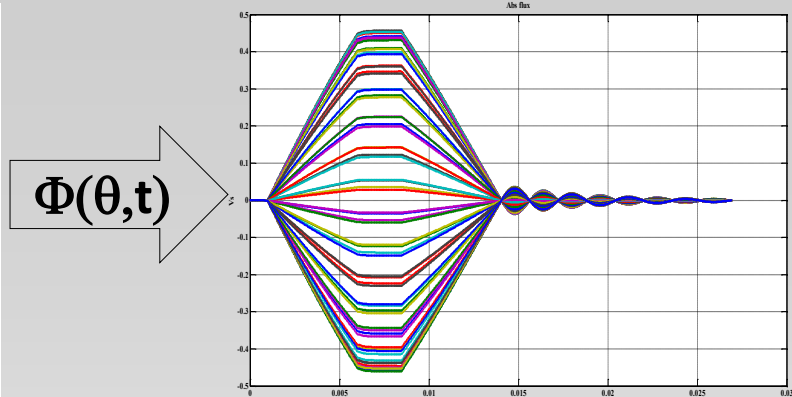
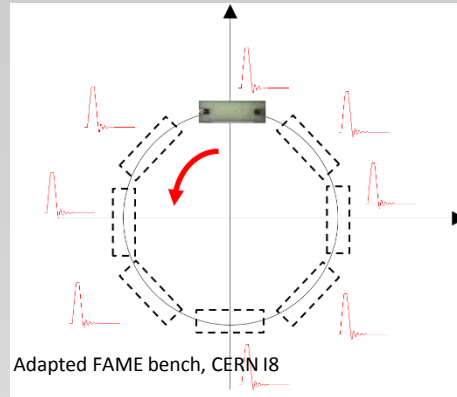


- random uncertainty: electrical and mechanical noise
- systematic uncertainty: mechanical coil imperfections = $f(\theta)$ (weight-induced sag, ball-bearing eccentricity, torsional vibrations ...)
- “true” value – may be averaged from any two measurements 180° apart

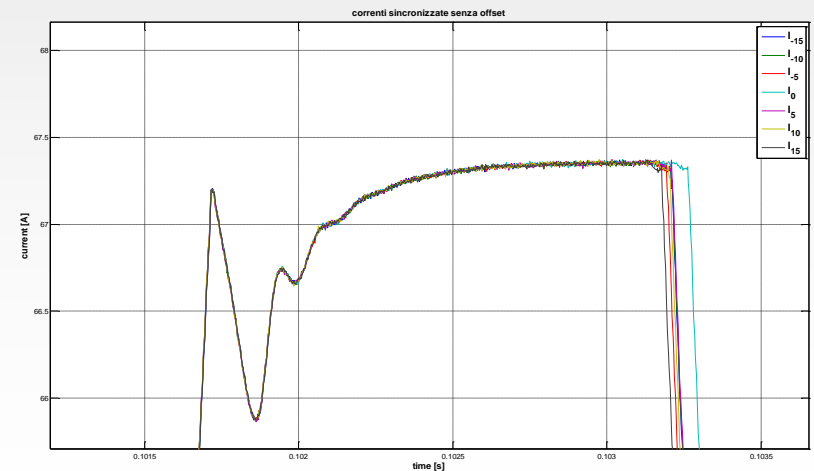
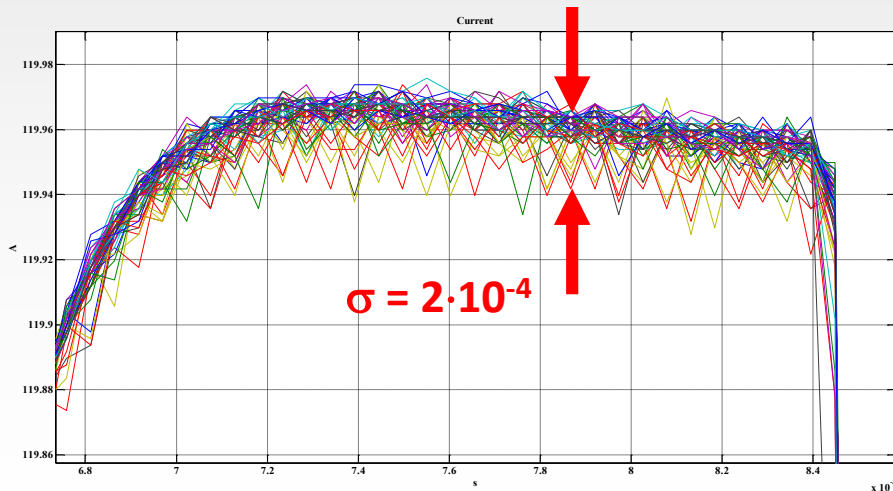


Low RMS

Stepwise pulsed-mode harmonic coils

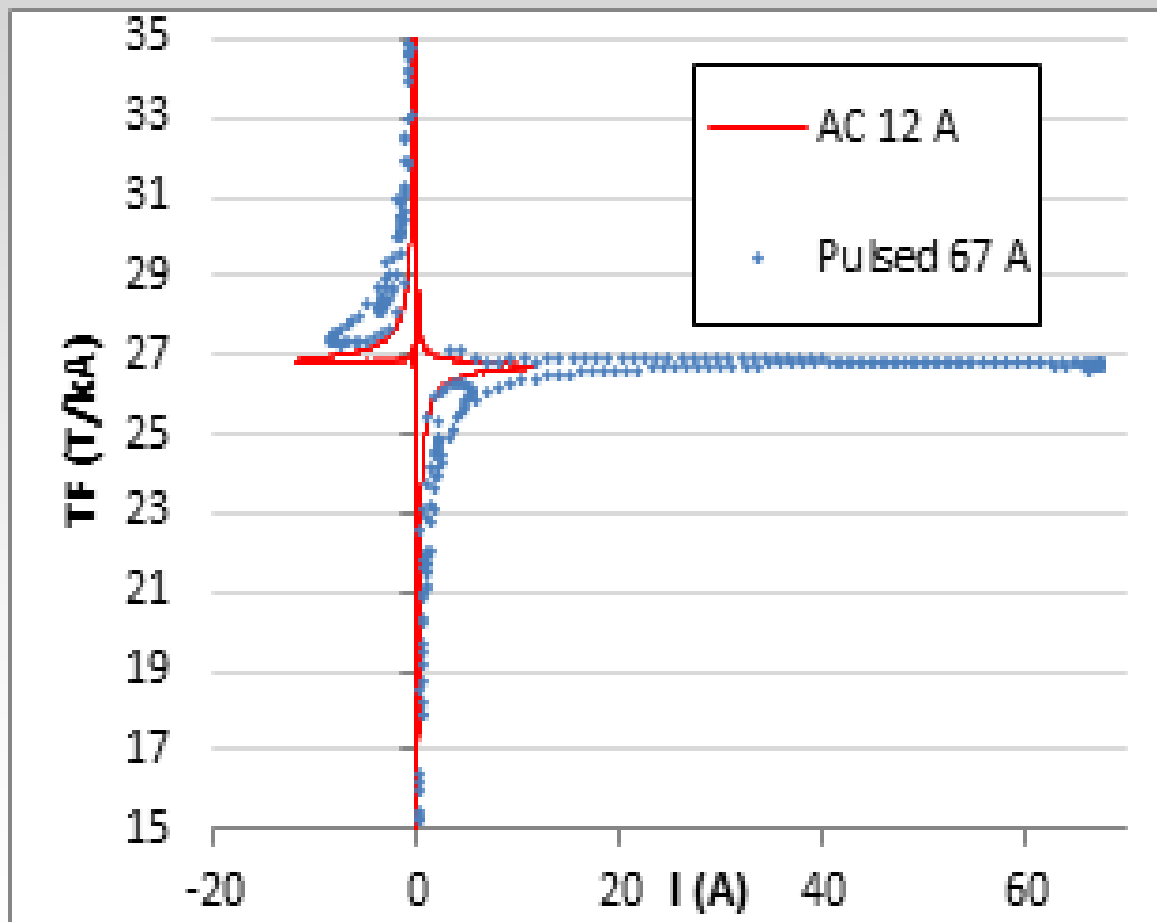


- New technique being developed to measure dynamically in the nominal powering conditions
- Requires: precise angular positioning, repeatability of power cycles
- Alternative solutions being evaluated:
 - scaling/time shifting $I(t)$ and $\Phi(t)$ to recover errors by post-processing
 - differential mode measurement (additional fixed coil as a reference for scaling)
 - simultaneous rotation/current pulsing (reasonable for pulse lengths 0.1~1 s)



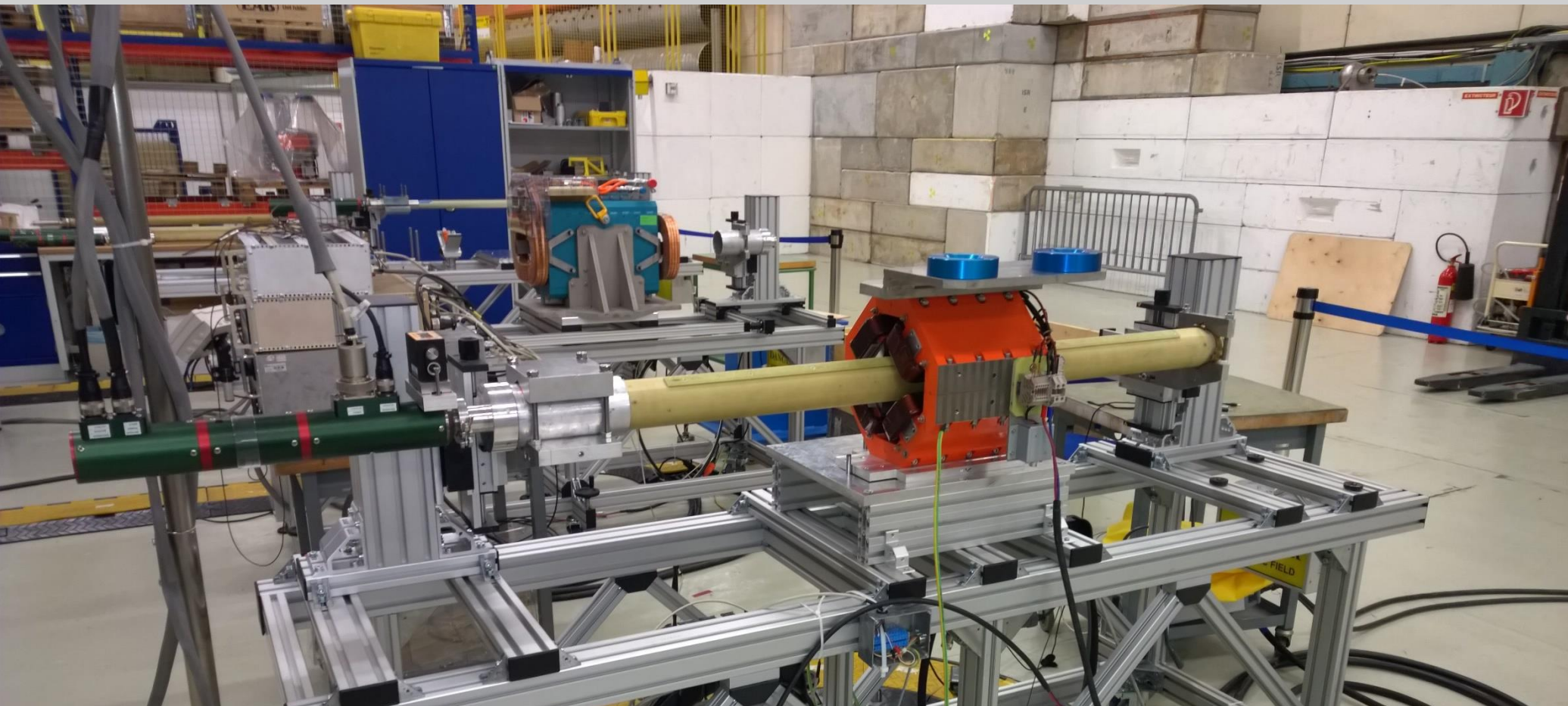
Extrapolation of low-field measurements

- Simplest solution: use standard DC rotating coils or DC/AC stretched wire → extrapolate to high fields (flat-top conditions must be equivalent to DC when the beam passes i.e. eddy current must decayed)
- Transfer function is perturbed by eddy current losses + remanent field effects
- So far used only for Linac4 EMQs ($\pm 0.5\%$ tolerance)



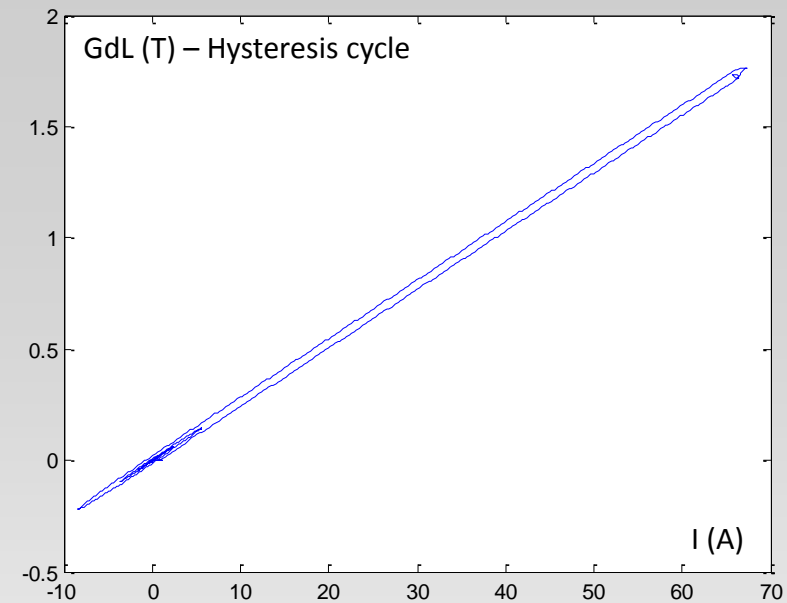
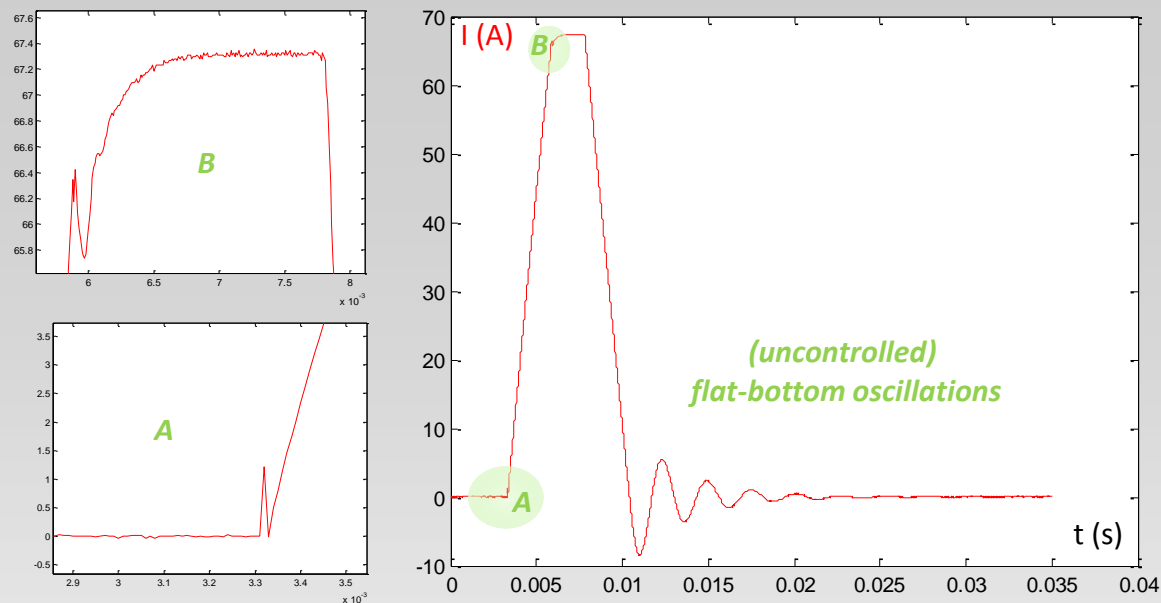
Example:

GdL transfer function of Linac4 EMQ

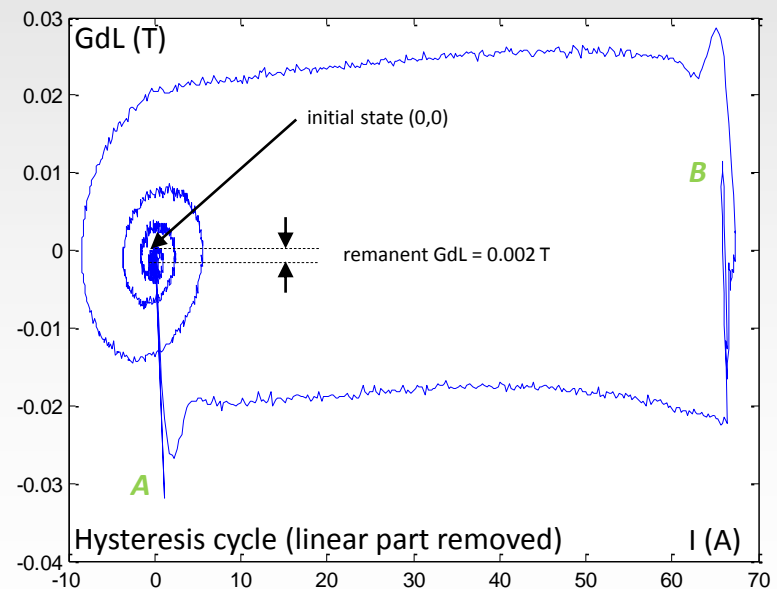


- Example: CERN linac4 TL EMQs: $B_0 \approx 0.05$ mT, $\int G d\ell = 0.14$ T @ 9 A (nominal 120 A), $L_c = 1.2$ m, $L_m = 0.3$ m
 $\rightarrow \Delta z \approx \frac{B_0(L_c - L_m)}{\int G d\ell} \approx 0.3$ mm
- Can be suppressed by flipping around the magnet/inverting the current
- Remanent field in the poles adds up: the effect can also be suppressed taking care about cycling

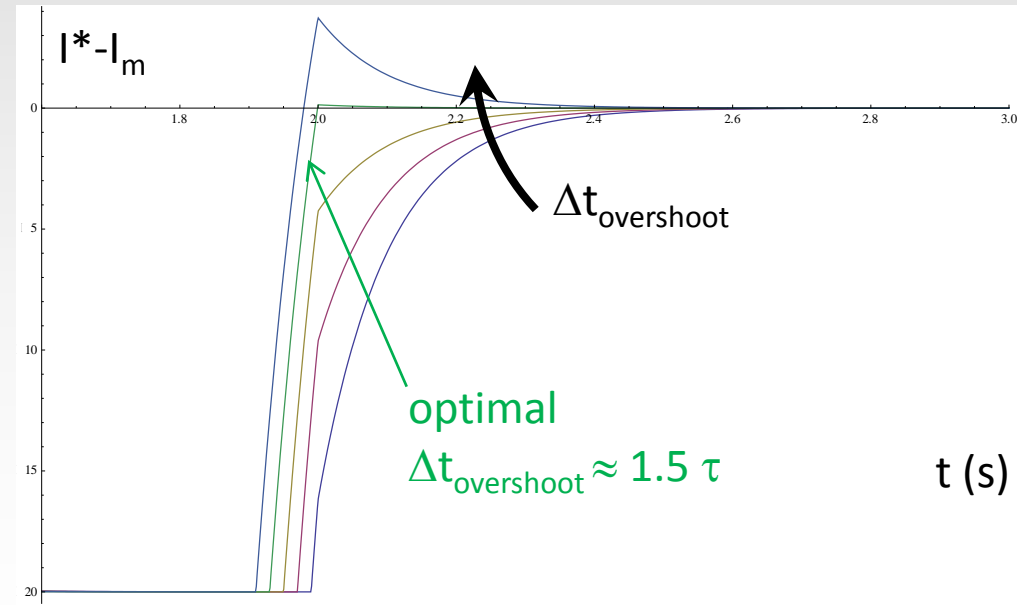
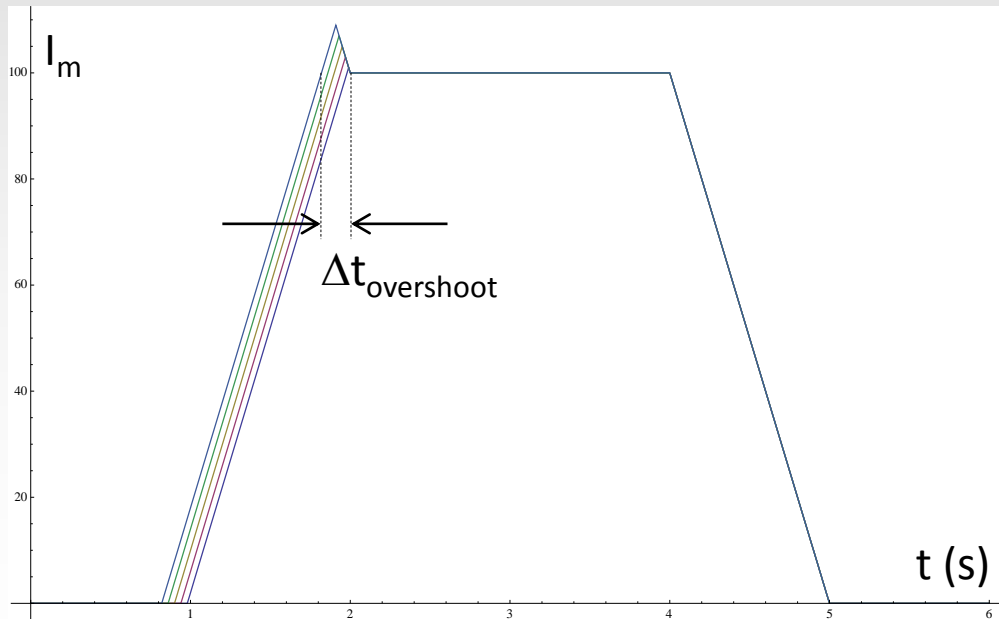
Cycling issues in fast-ramping magnets



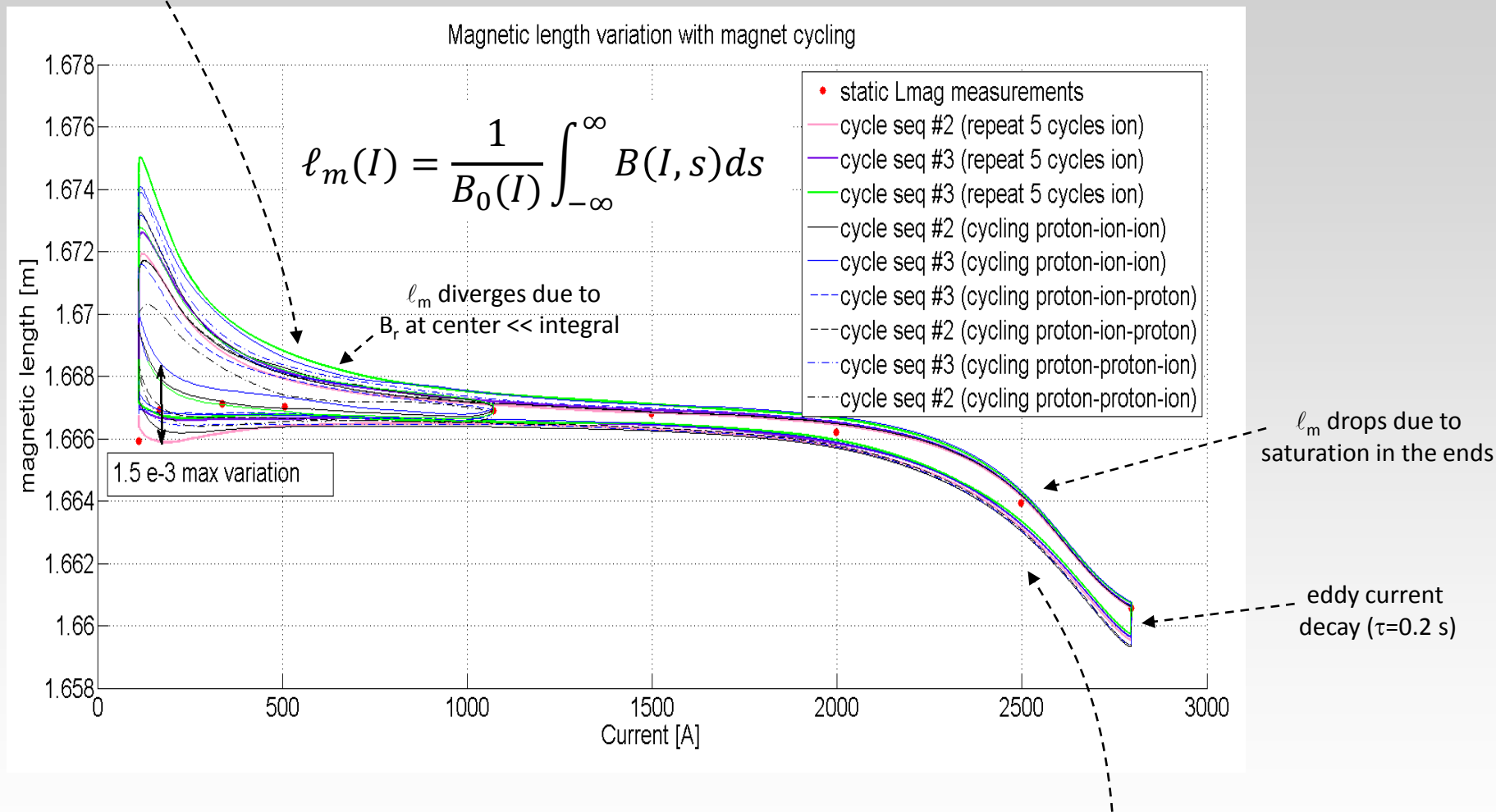
- Example: fast capacitive discharge powering of Linac4 inter-tank EMQs
- current spikes lead minor hysteresis loops \rightarrow field reproducibility degradation
- oscillations at the end of the ramp-down may provide a beneficial free de-gaussing, *if* symmetrical
- the overshoot at the end of the ramp-up may give a more stable flat-top, but makes it less reproducible



- Eddy currents can be partially, totally or over-canceled by a linear excitation current overshoot at the end of ramp-up
- Example: stable flat-top reached at the time cost of $\sim 1.5\tau$ (to be compared with exponential decay time $\sim 3\tau$)
- Caveats:
 - power converter needs high dV/dt ;
 - the maximum working point is increased considerably, at the risk of saturation



large fluctuations due to history-dependent residual field
reproducibility degrades at low field



saturation tends to erase previous magnetic history
→ better reproducibility at high field

Courtesy G. Golluccio

Summary

- Measurement systems for **small gaps** ($\varnothing < 40$ mm) currently being developed at CERN
- Main limitations:
 - mechanical tolerances for rotating coil systems
 - mechanical tolerances and various non-linear effects for stretched-wire systems
- reasonable performance at $\varnothing 20$ mm, still work to do at $\varnothing 8$ mm
- many sources of systematic errors can be corrected by flipping the magnet around or upside/down

- Techniques for **fast cycled magnets** are being developed too (Linac EMQs)
- Open issues: precise extrapolation of main component, dynamic behavior of field harmonics/magnetic center
- The impact of new current waveforms on magnetic performance should be evaluated early in the design cycle:
 - extending dynamic range at the bottom degrades reproducibility
 - uncontrolled transients degrade the reproducibility
 - prototype magnets should ideally be tested along with their prototype power supplies !

- [1] M. Buzio, S. Sanfilippo *et al.*, SMALL-DIAMETER ROTATING COILS FOR FIELD QUALITY MEASUREMENTS IN QUADRUPOLE MAGNETS, XX IMEKO TC-4 International Symposium, September 15–17, 2014, Benevento, Italy
- [2] S. Kasaei *et al.*, MAGNETIC CHARACTERIZATION OF FAST-PULSED ELECTROMAGNETIC QUADRUPOLES FOR LINAC4 AT CERN, Proceedings of Linac14, CERN, September 2014