

# Compact CW Recirculating Electron Accelerator – eFFAG (3–9 MeV)

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*Acknowledgements: Passport Systems,  
Particle Accelerator Corp*

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# Introduction – Outline

Primary Application presented: Cargo scanning

Compact, mobile, economical system

Purpose

Physics of cargo interrogation

Radiography (transmission through cargo of photon and neutron beams)

Neutron induced reactions

Photon induced reactions

The role that accelerators play

Photon generation

- Hadron induced reactions
- Bremsstrahlung
- Monochromatic photon sources

Important characteristics of electron accelerators

Novel approaches to accelerators

# Identified Threats easily transported in containers

## Explosives

Examples of extensive damage

- Oklahoma, Lebanon, Lockerbie, Halifax (1917)

Important elements: N, O, Cl, Na, S, K, P (and fulminates)

## Toxic Substances

Mustard gas ( $C_4H_8Cl_2S$ ), Sarin ( $C_4H_{10}FO_2P$ ), Phosgene ( $CCl_2O$ ), etc.

## Dirty Bombs

$^{137}Cs$ ,  $^{60}Co$ , etc.

Shielding materials Pb, W, Fe, etc.

## Special Nuclear materials (SNM)

$^{235}U$ ,  $^{239}Pu$ ,  $^{237}Np$

## Weapons of Mass Destruction

$^{235}U$ ,  $^{239}Pu$ ,  $^{237}Np$ , Explosives, Tamper materials, etc.

# Cargo interrogation

examination must be performed rapidly

non-intrusive

Therefore must involve penetrating radiation; i.e. neutrons and/or photons

requires an electron or neutron particle accelerator

Only photon-producing electron accelerators are considered here

One alternative: cosmic muons have been used to identify

SNM such as  $^{235}\text{U}$  and  $^{239}\text{Pu}$

heavy metal shields such as lead or tungsten (through multiple scattering differences)

the technique does not identify specific isotopes

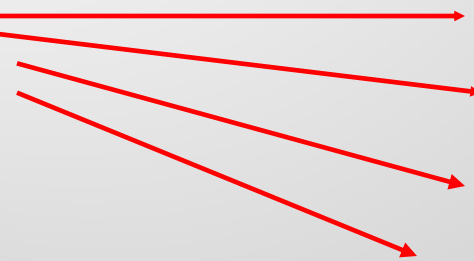
Does not provide comprehensive identification of explosives and toxic substances

# Threat-based materials can be detected using photons via nuclear reactions

**BEAM**

**MEASURED PARTICLE**

**Photons**



***Photons:***

**Transmission**

**Nuclear Resonance Fluorescence**

**Effective Z measurement**

***Neutrons:***

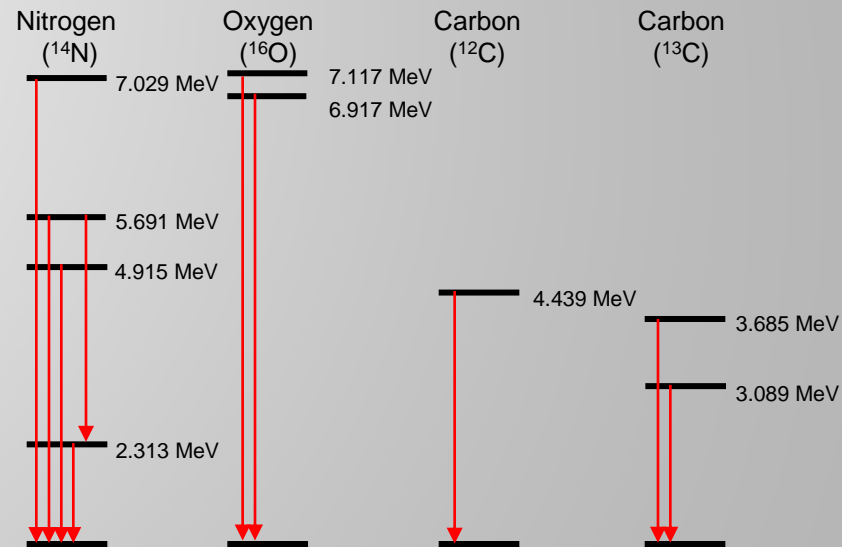
**Photon Induced Fission  
(Prompt Neutrons)**

# Identification of Nuclear States

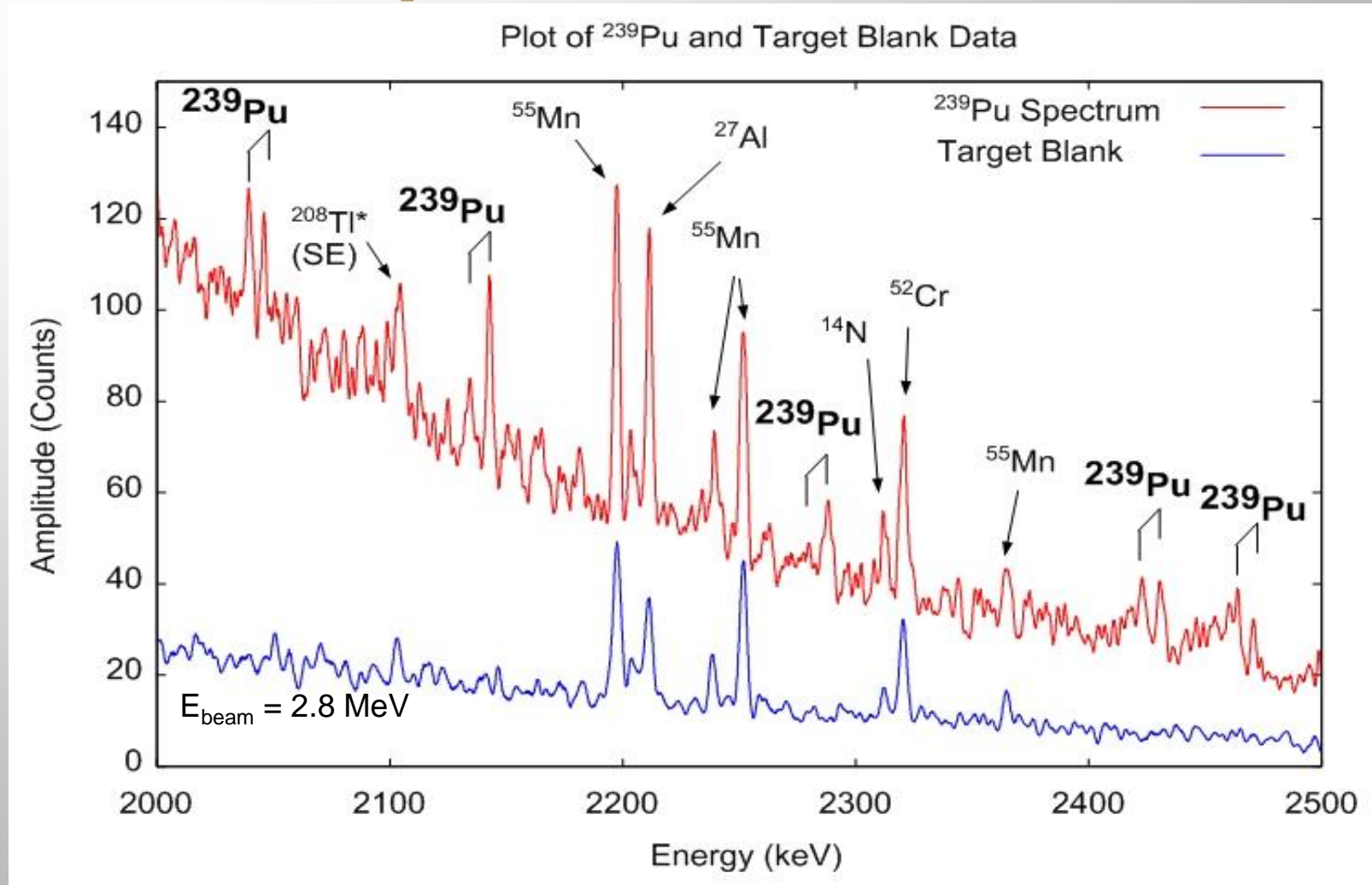
## Nuclear Resonance Fluorescence

*Broad spectrum bremsstrahlung* photons produced by a 5-9 MeV electron accelerator are collimated and scan the container. *Different nuclear species* are resonantly excited by specific photon energies and then re-emit photons when the “excited” state decays. The photons are Doppler-shifted and therefore not re-absorbed.

*Using photon spectroscopy* the contents can be identified and a 3-D image reconstructed with proper detector arrays



# NRF Spectrum from $^{239}\text{Pu}$



Measurements performed with LLNL

# Exploiting NRF and non-resonant backgrounds

## Photons emitted by NRF

1-8 MeV (hence a high energy electron accelerator/photons are required)

Nearly isotropic

- Detect back-emitted photons
- Reduction in background (forward Compton scattering)

## Effective Z measurement

Non-resonant background from multiple processes

- Multiple processes again require high energy incident photons

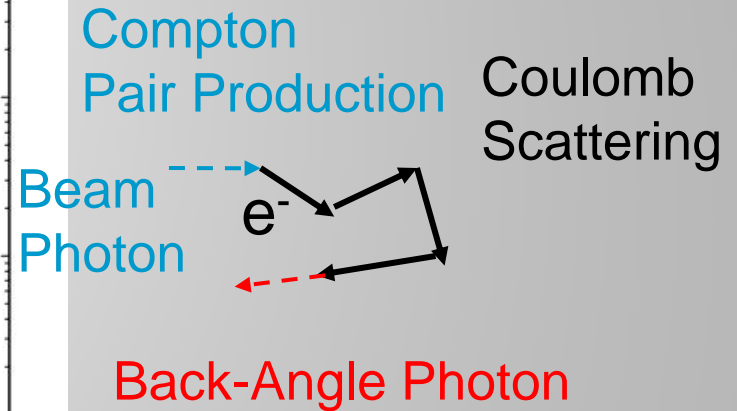
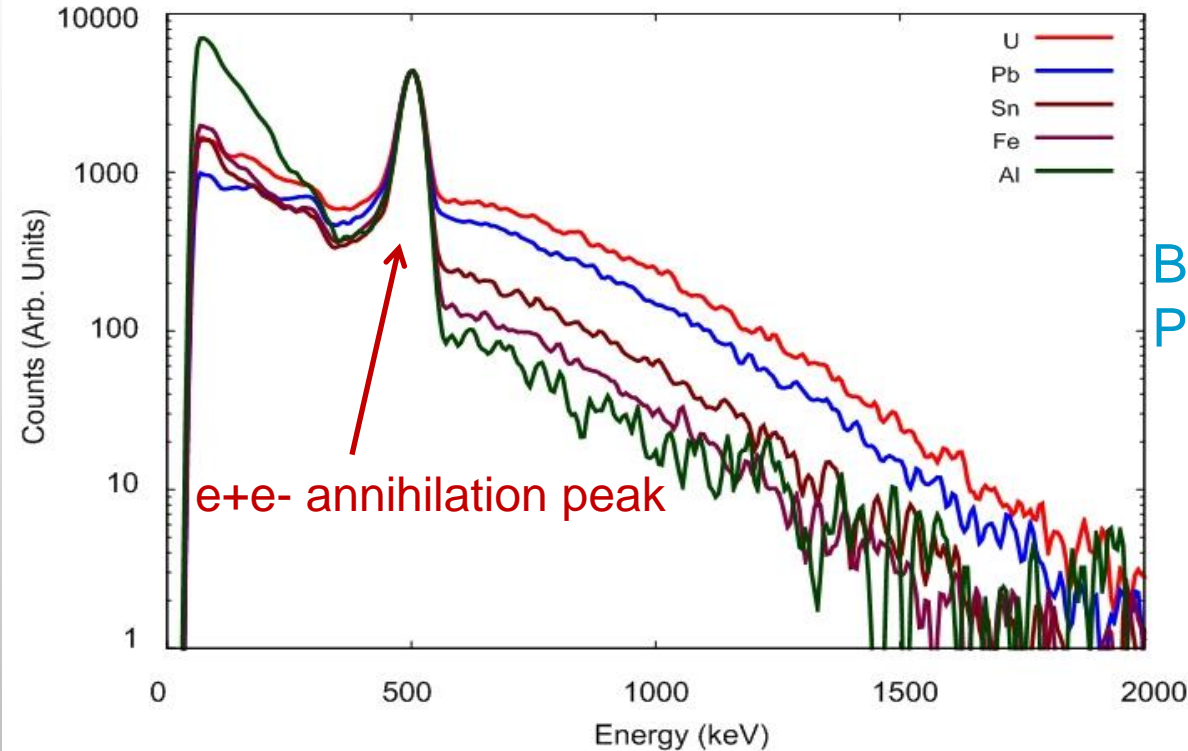
Simple count of back-angle photons

Count proportional to  $Z^\alpha$  to when normalized to annihilation peak



# Effective Z Determination

Plot of Different Z materials. Normalized 511



Fast detection of high Z contents

Single Compton Scatter with  $E_{\text{beam}} = 5 \text{ MeV}$

- $E_{\text{photon}} (120^\circ) = 320 \text{ keV}$

- $E_{\text{photon}} (180^\circ) = 240 \text{ keV}$

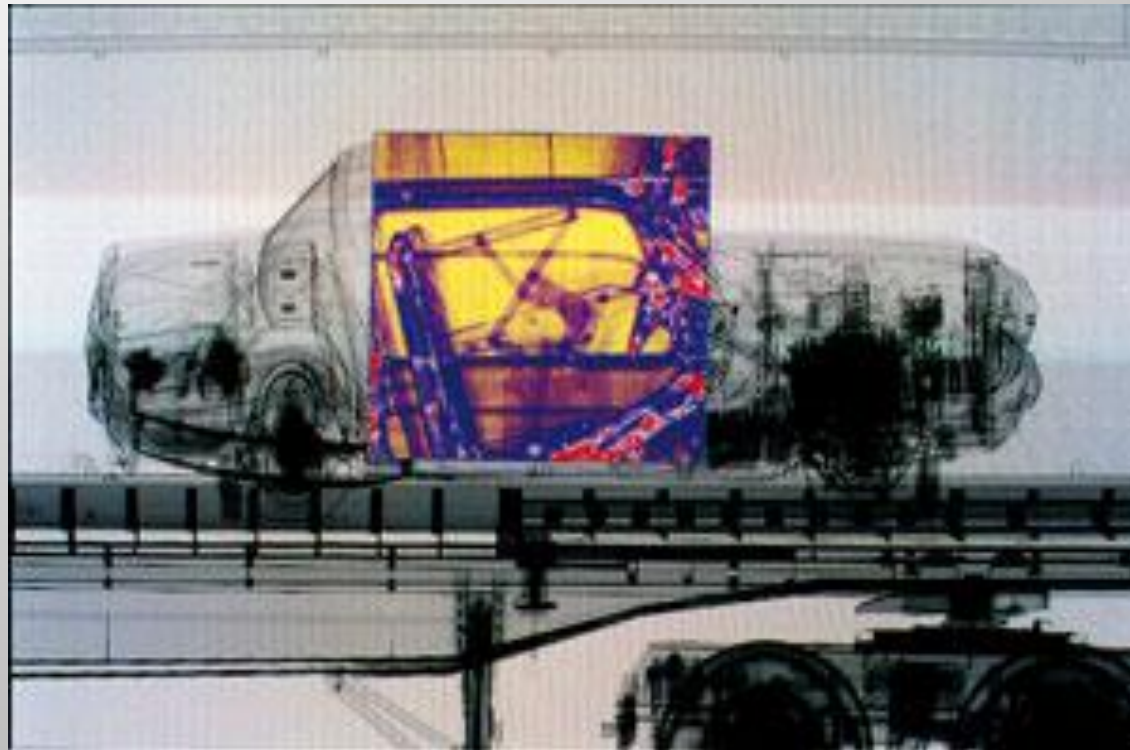
Pair Production  $\rightarrow$   $e^+e^-$  annihilation ( $\sim Z^2$ )

- $E_{\text{photon}} = 511 \text{ keV}$

# Transmission Radiograph with NRF

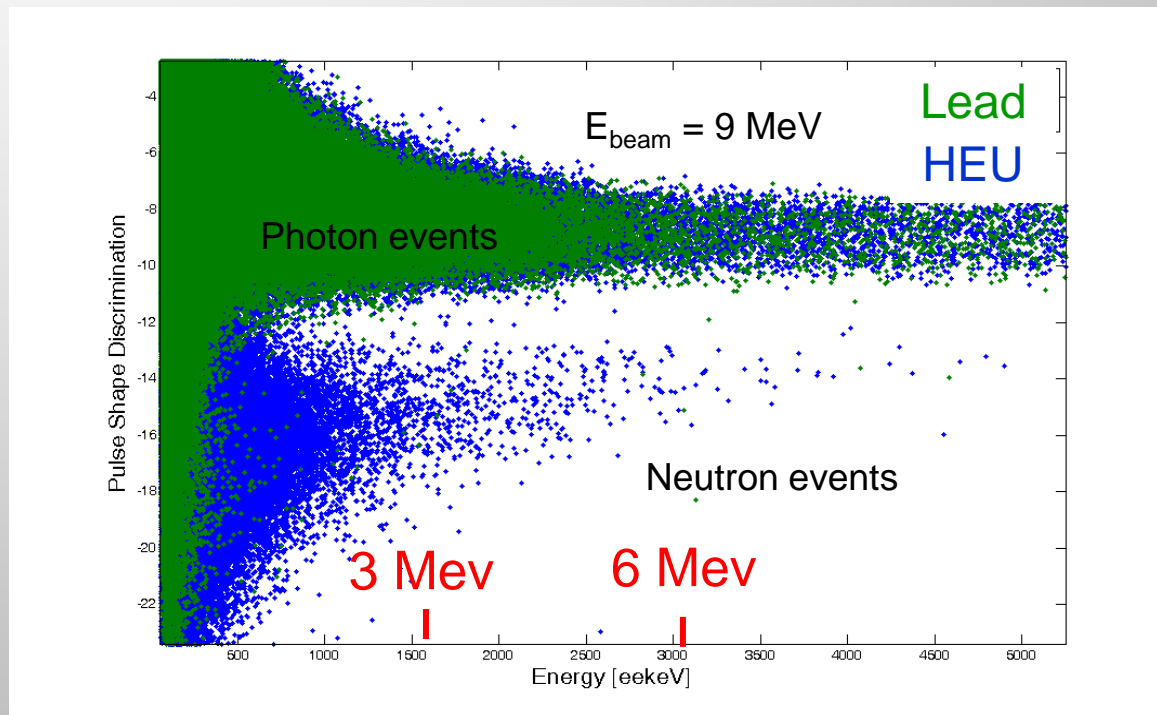
Nuclear resonance lines are now “missing” from spectrum and can be identified.

Photon spectroscopy combined with transmission provides 2D isotopic composition to the radiograph.



# Prompt Neutrons from Photofission

- Prompt, high energy neutrons provide unique signal for fissile material and actinides
- Photon events are distinguished from neutrons by pulse-shape discrimination
- Neutron energy distribution independent of incident photon energy



# Summarizing criteria

Energy of photons:  $\sim 3 - 9$  MeV

3 MeV: NRF signatures for SNM is generally below this energy

6 MeV: Good yield for effective Z and good penetration for radiography

9 MeV: NRF from oxygen; well above photofission thresholds for actinides, increased penetration; excellent yield effective Z determination

Intensity: Adjustable from zero to several mA

Must dynamically adjust to different cargo loadings to minimize dose

Continuous time distribution (high duty cycle) desirable

Single photon/neutron counting

Signal to Noise

# List of References

*Muon Radiography; Detecting Nuclear Contraband*, Brian Fishbine, Los Alamos Research Quaterly, Spring 2003

*Gamma-ray and Neutron Radiography as Part of a Pulsed Fast Neutron Analysis Inspection System-* J. Rynes, J. Bendahan, T. Gozani, R. Loveman, J. Stevenson and C. Bell; [Nuclear Instruments and Methods in Physics Research Section A: Accelerators, Spectrometers, Detectors and Associated Equipment, Volume 422, Issues 1-3](#), 11 February 1999, Pages 895-899: And references therein.

*Imaging and Radiography with Nuclear Resonance Fluorescence and Effective-Z (EZ-3D™) Determination; SNM Detection Using Prompt Neutrons from Photon Induced Fission;* William Bertozzi, Richard Hasty, Alexei Klimenko, Stephen E. Korbly, Robert J. Ledoux and William Park; **CP1099**, Application of Accelerators in Research and Industry: 20<sup>th</sup> International Conference, edited by F.D.McDaniel and B.L. Doyle; AIP Conference Proceedings **1099**

W. Bertozzi, J. A. Caggiano, W. K. Hensley, M. S. Johnson, S. E. Korbly, R. J. Ledoux, D. P. McNabb, E. B. Norman, W. H. Park, and G. A. Warren, Phys. Rev. C **78** 041601(C) (2008)

C. P. Sargent, W. Bertozzi, P. T. Demos, J. L. Matthews, and W. Turchinets, *Prompt Neutrons from Thorium Photofission*, Phys. Rev. **137**, B89 - B101 (1965)



# Higher energy, higher intensity accelerators are required

Low duty cycle RF LINACs traditionally used for transmission radiography (measurement of energy deposition for photons  $\geq$  MeV)

duty cycle is typically limited to a value of approximately  $10^{-3}$  (microsecond beam pulses at  $\sim 300$  Hz rates)

Innovative technologies that provide material discrimination capabilities rely on single event counting and spectral analysis require much higher duty cycles

Higher duty cycle DC beams can be obtained from electrostatic machines, but require a large footprint limiting their mobility and broad application.

High initial and operating costs for commercial high duty cycle machines (e.g. IBA Rhodotron)

# Applied accelerator types

1. *Synchrotrons*
  - Pulsed magnetic fields
  - Swept-frequency accelerating systems
  - Separated components
  - Long component-free straight sections
  - Size scales with momentum/charge to mass ratio.
  - Current limited -  $\mu\text{A}$
2. *Cyclotrons*
  - Fixed magnetic fields
  - DC beams, fixed-frequency accelerating system
  - Relativistic energies: swept-frequency accelerating systems
    - Synchrocyclotron.
  - Monolithic pole pieces - no straight sections
  - Very large at high energy
  - High currents (mA)
3. *Linacs*
  - Longest footprint
  - Costly - no recirculating beam or "reuse" of components
  - Ultra-high currents - 100 mA
4. *FFAGs*
  - Fixed magnetic fields
  - DC beams well into the relativistic regime
  - Compact
  - Separated components
  - Synchrotron-like dynamics allow component-free straight sections
  - High current - tens of mA

## Synchrotron

*Low Current ( $< \text{mA}$ )*

*High Energy (TeV)*

*Pulsed Beam*

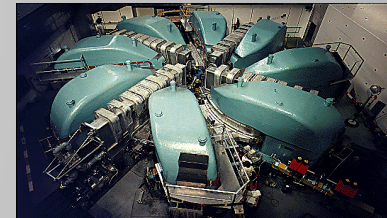


## Cyclotron

*High Current ( $< \text{A}$ )*

*Low Energy (600MeV)*

*CW beam*



## Linac

*High Current, High Energy*

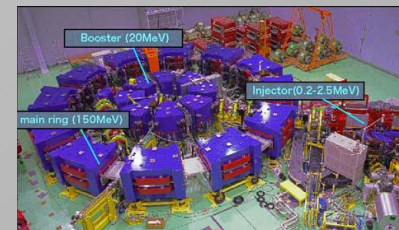
*Pulsed or continuous beam: Large, expensive*



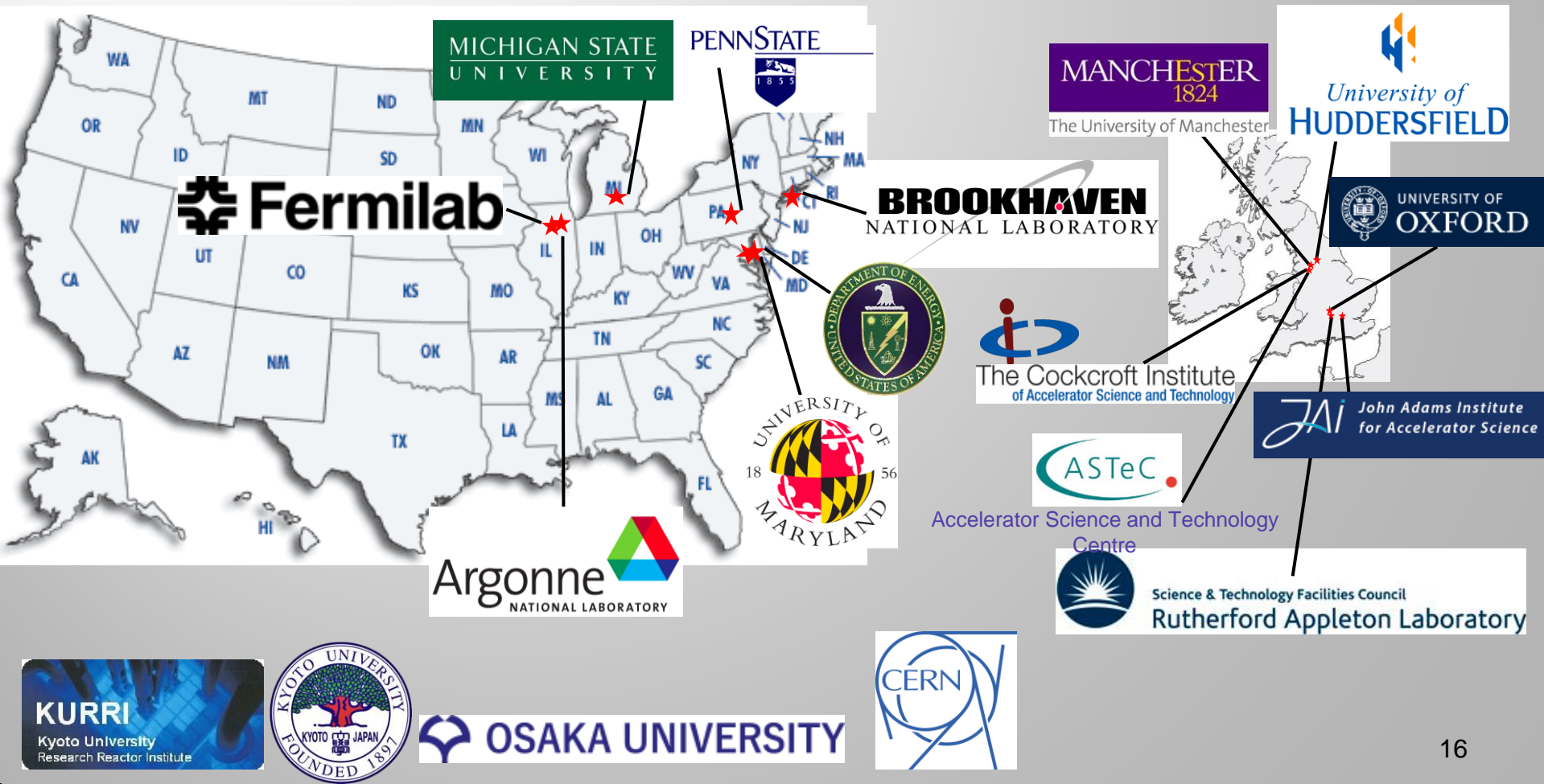
## FFAG

*High Current ( $< \text{A}$ ), High Energy (few GeV)*

*CW beam, compact*



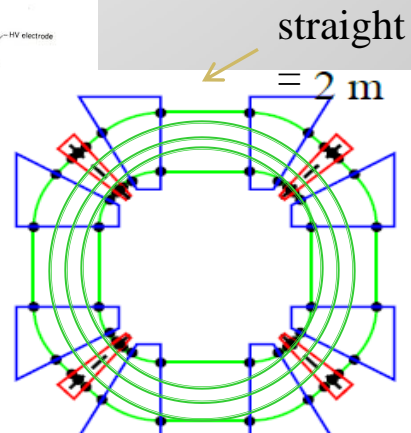
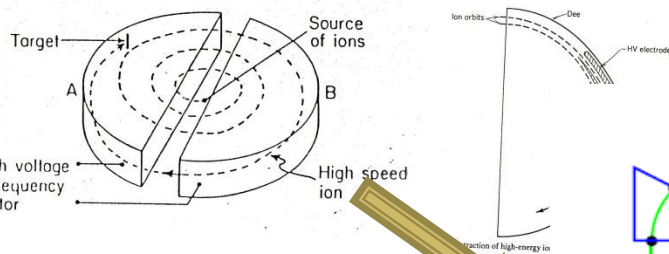
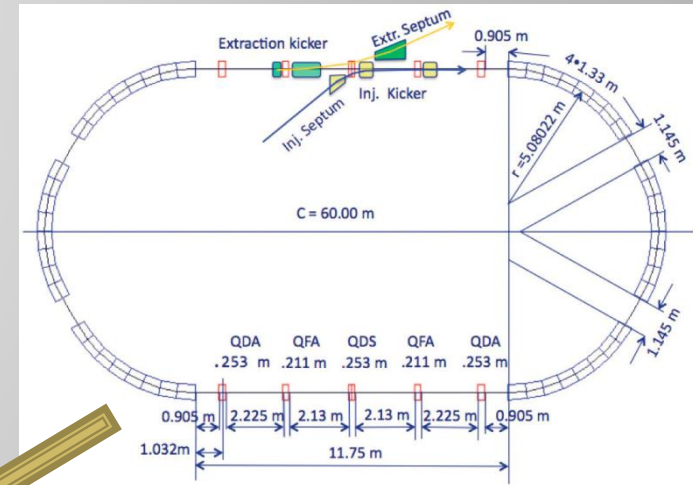
# The international FFAG collaboration





# Next generation high duty cycle Accelerators

- Apply a “synchrotron” strong-focusing field profile to each “cyclotron” orbit
- Strong-focusing allows
  - Strong RF acceleration
  - Low -loss profile of the synchrotron
  - CW beam to high energies in compact structure
  - Avoidance of unstable beam regions
    - constant machine tune



**CONSTANT MACHINE TUNES AT NEAR AND FAR RELATIVISTIC ENERGIES**

**UNLIKE A CYCLOTRON**

# Focusing in Conventional Accelerators

- The FFAG combines all forms of transverse beam (envelope) confinement in an arbitrary CF magnet:
  - For the horizontal, the three terms are

$$1/f_F = \underbrace{k_F l}_{\text{Synchrotron and linac}} + \underbrace{\frac{\mathcal{G}}{\rho_F} + \frac{\eta}{\rho_F}}_{\text{cyclotron}}$$

with  $\mathcal{G}$  is the sector bend angle,  $\eta$  the edge angle (edge angle is assume small so tangent is approximated), length,  $l$ , is the F half - magnet length and  $k_F$  is the "local" gradient for an arbitrary order field.

- For the vertical only the quadrupole gradient,  $k_D l$ , and the edge term are available
- The different focusing terms can be varied independently to optimize machine parameters such as footprint, aperture, and tune in a FFAG

# Quick Guide to FFAGs

## Simplest Dynamical Definition:

FFAG is ~ a cyclotron with a gradient; beam confinement is via:

- Strong alternating-gradient (AG) focusing, both planes: *radial sector FFAG*
  - normal/reversed gradients alternate (like a synchrotron)
- Gradient focusing in horizontal, edge focusing in vertical: *spiral sector FFAG*
  - vertical envelope control is through edge focusing (like a cyclotron)
  - the normal gradient increases edge focusing with radius / momentum (unlike a cyclotron)

A cyclotron is a dipole-only FFAG with no reverse gradients

## Types of FFAGs:

### Scaling:

- B field follows a scaling law as a function of radius -  $r^k$  (k a constant;) present-day scaling FFAGs: Y. Mori, Kyoto University Research Reactor Institute

### Nonscaling:

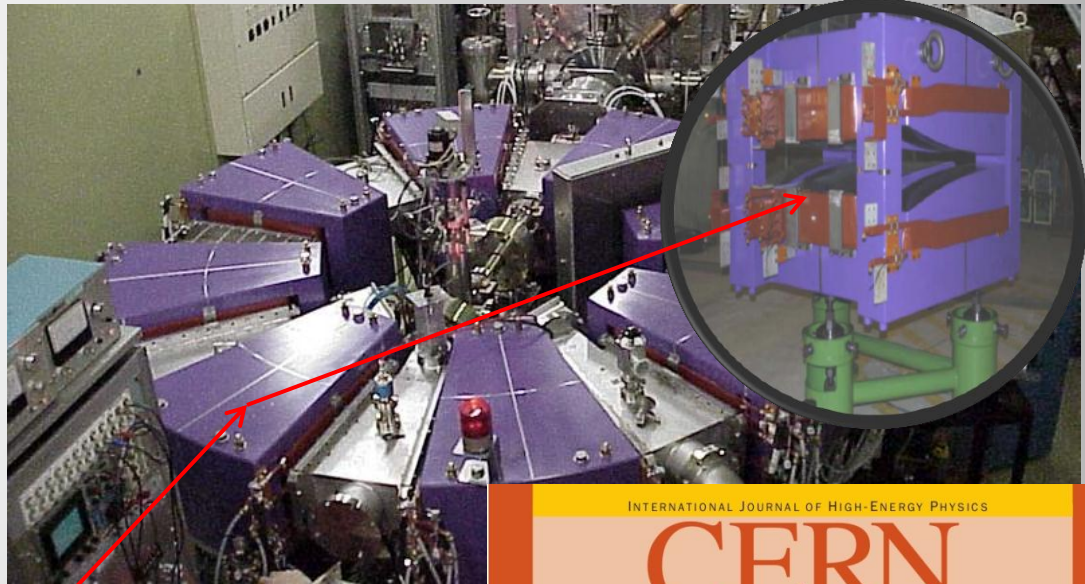
- Linear (quadrupole) gradient; beam parameters generally vary with energy (EMMA FFAG, Daresbury Laboratory, first nonscaling FFAG)
- Nonlinear-gradient; beam parameters such as machine tune can be fixed (as in a synchrotron)

# FFAGs and their Variations

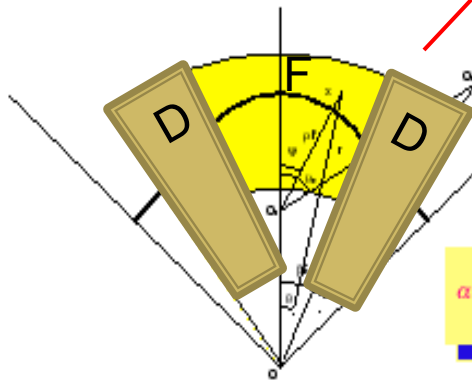
**Scaling FFAGs** (spiral or radial-sector) are characterized by geometrically similar orbits of increasing radius, imposing a constant tune (field and derivative gradient scale identically with  $r$ ). Magnetic field follows the law  $B \propto r^k$ , with  $r$  as the radius, and  $k$  as the constant field index.

$$B = B_0 \left( \frac{r}{r_0} \right)^k = B_0 \left( 1 + \frac{k}{r_0} x + \frac{k(k-1)}{2! r_0^2} x^2 + \dots \right)$$

**Field expansion:**  $k$  determines multipole order;  
 Comments: the lower the  $k$  value, the more slowly field increases with  $r$  and the larger the horizontal aperture, but the more linear the field composition and dynamics.



**Spiral Sector:** example: more compact; positive bend field only. Vertical focusing controlled by edge crossing angle.



**Radial Sector:** example: This is a triplet DFD cell; there are also FDF, FODO and doublets. In a radial sector the D is the negative of the F field profile, but shorter.

$$\alpha = \frac{1}{k+1} : \text{momentum compaction factor}$$

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# CERN COURIER

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**ANR**

**RACCAM**

## A route to rapid acceleration

**CERN**  
LHC gets onto the starting blocks p5

**LHC FOCUS**  
Nobel expectations at Lindau meeting p29

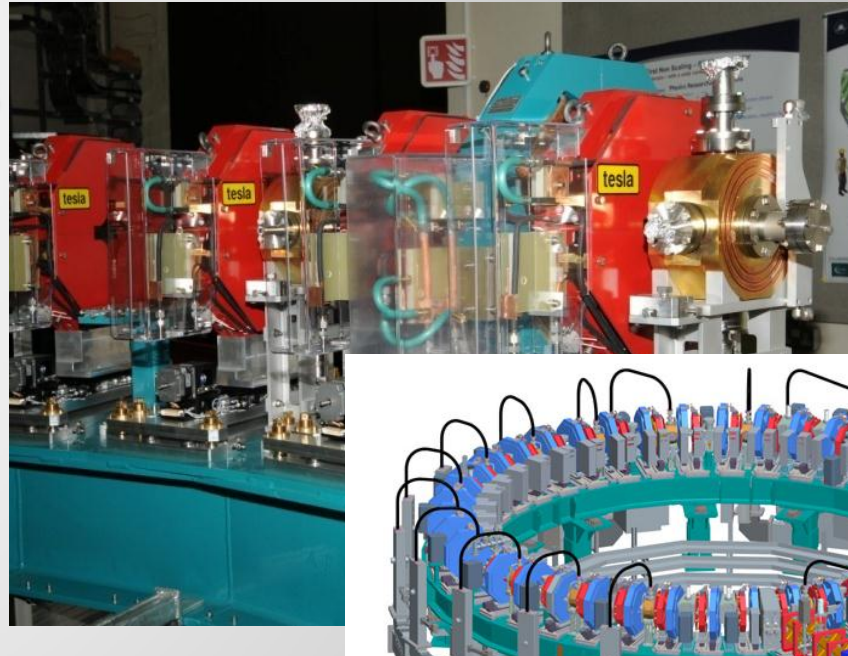
**ENERGY**  
Chris Llewellyn-Smith looks to the future p33



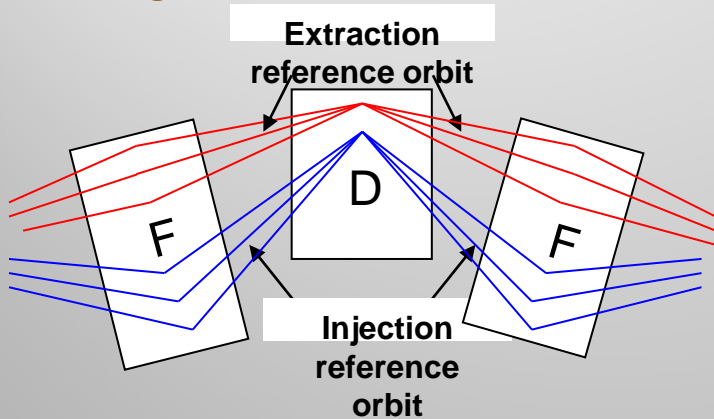
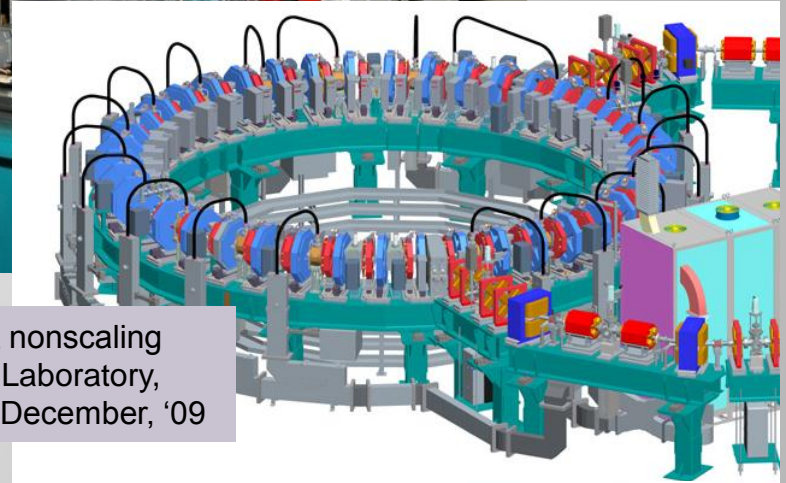
# Linear nonscaling FFAGs for rapid acceleration

## Linear-field, nonscaling FFAG.

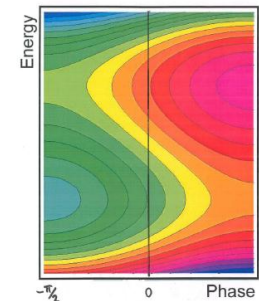
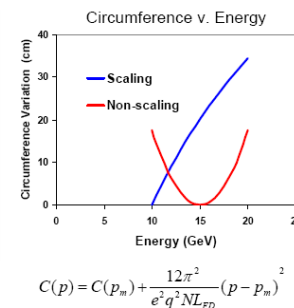
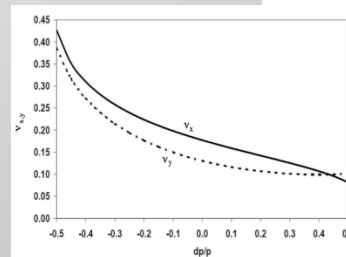
Ultra-compact magnet aperture, proposed and developed for High Energy Physics (Neutrino Factories and Muon Colliders), relaxes optical parameters and aims only for stable acceleration. In general they are not suitable for an accelerator with a modest acceleration system and accelerate only over a factor of 2-3 range in momentum.



EMMA – world’s first nonscaling FFAG, @Daresbury Laboratory, commissioning, late December, ‘09



Cartoon of orbit compaction: nonsimilar orbits, nonconstant tune, resonance crossing



**Characteristics– tune sweep/unit cell, parabolic pathlength on momentum (small radial apertures); serpentine (rapid) acceleration – beam “phase-slips”, crossing the peak 3 times, accelerating between rf buckets**

# Tune-stable nonscaling FFAGs for slower acceleration

## Tune-stable, nonscaling FFAGs

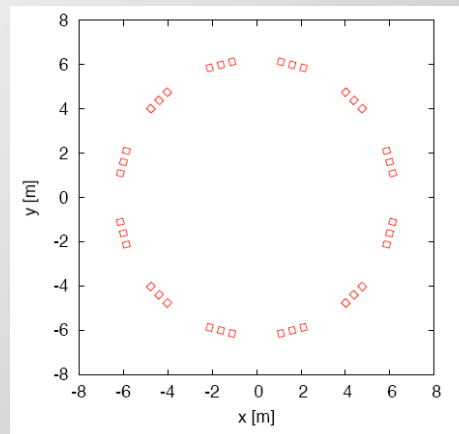
• Tune is strongest indicator of stable particle motion – allowing particles execute periodic motion eventually returning to the same transverse position relative to a reference orbit. Constraining the tune can be sufficient to design a stable machine.

• Release of other linear optical parameter allows flexibility and optimization both in cost and complexity of the accelerator design; i.e. simpler magnets, strong vertical focusing, for example

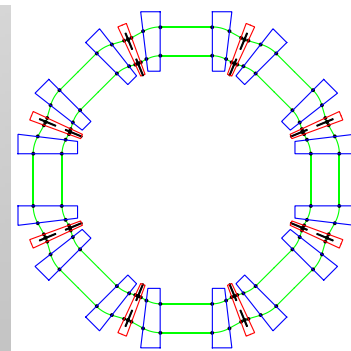
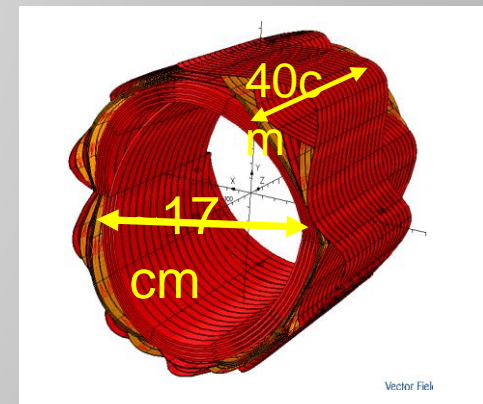
• ***Tune Stable Nonscaling FFAGs have either linear or nonlinear field profiles and/or edge contours***

## Two lattice approaches

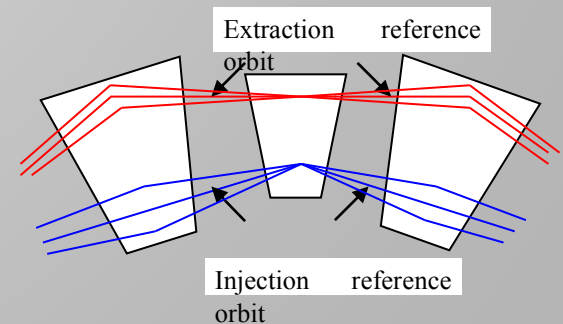
- 1) Machida version - which uses a scaling law truncated at decapole, rectangular magnets, (not discussed here, see PAMELA project) and
- 2) Johnstone version – The most general form of a radial sector : allowing independent, unconstrained field and edge profiles between two combined-function magnets.



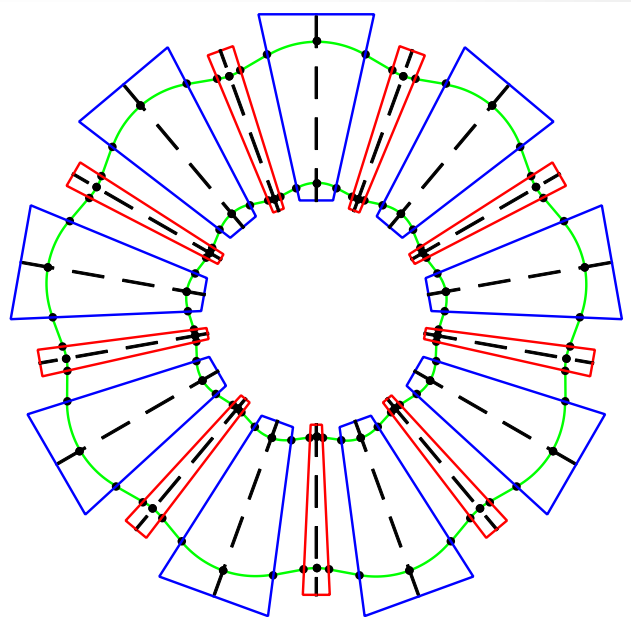
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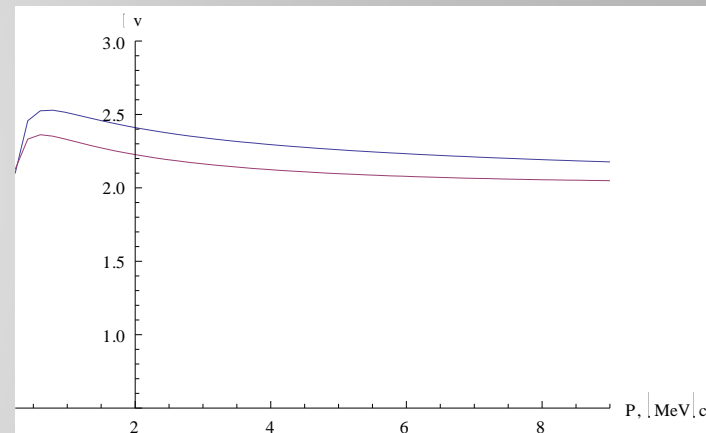
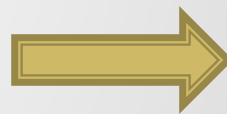
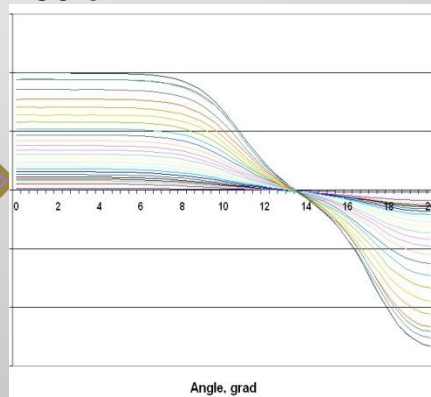
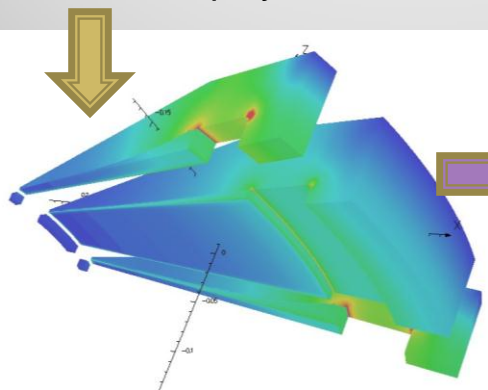
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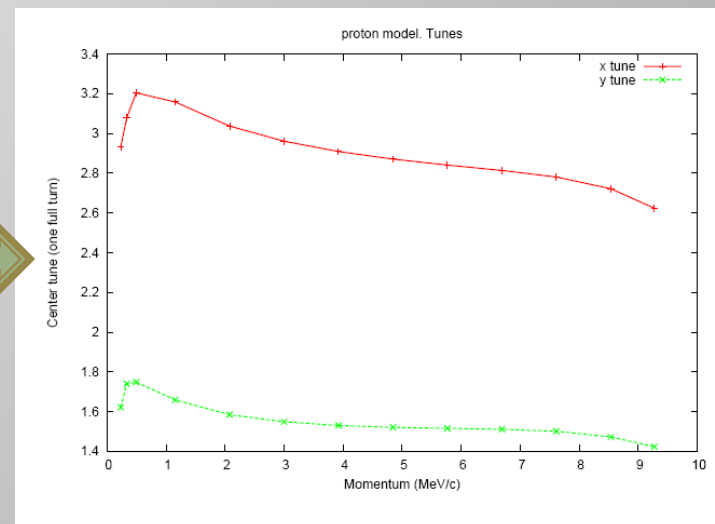
# A 50 keV to 9 MeV Compact eFFAG



Full 9-MeV ring with injection and extraction orbits displayed. Outer radius <math>< 50\text{ cm}</math>



**COMPARISON** Tune Results:  
hard edge vs. extended fringe field

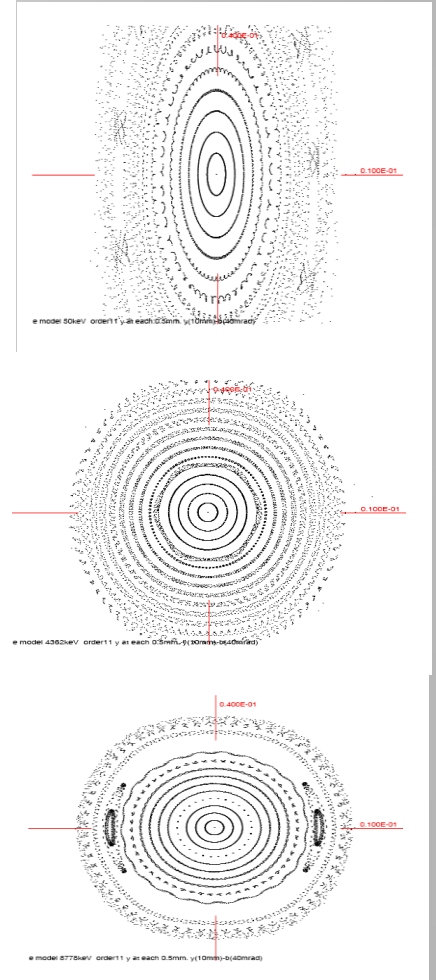
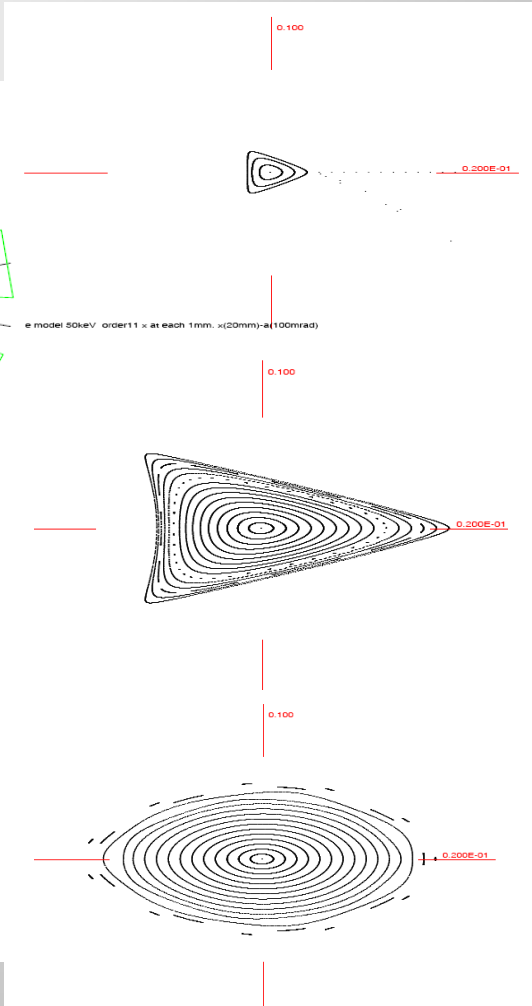
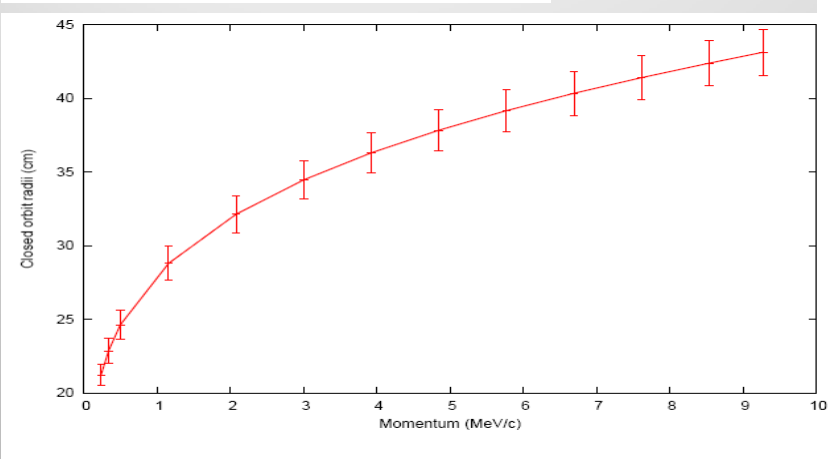
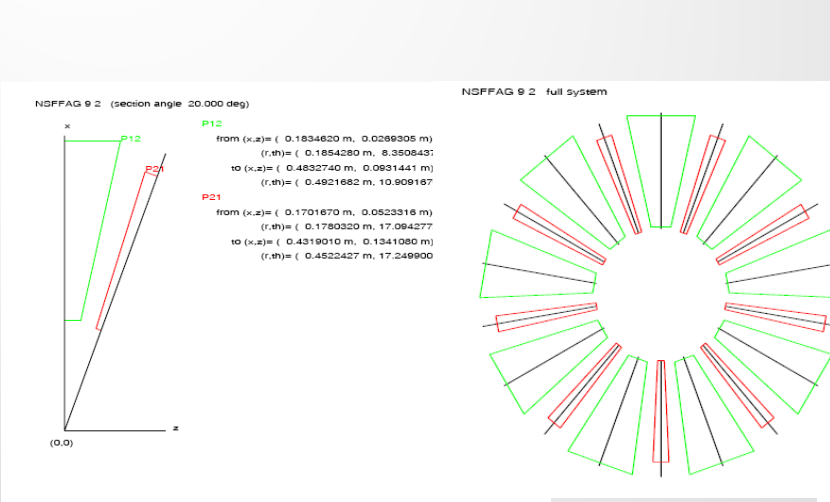


**TOSCA** Magnet Design and Field profiles.

**COSY INFINITY** Tune Results

Originally designed for induction acceleration ( $\sim$ kHz cycle time)

# Simulation: COSY Results: Electron Ring

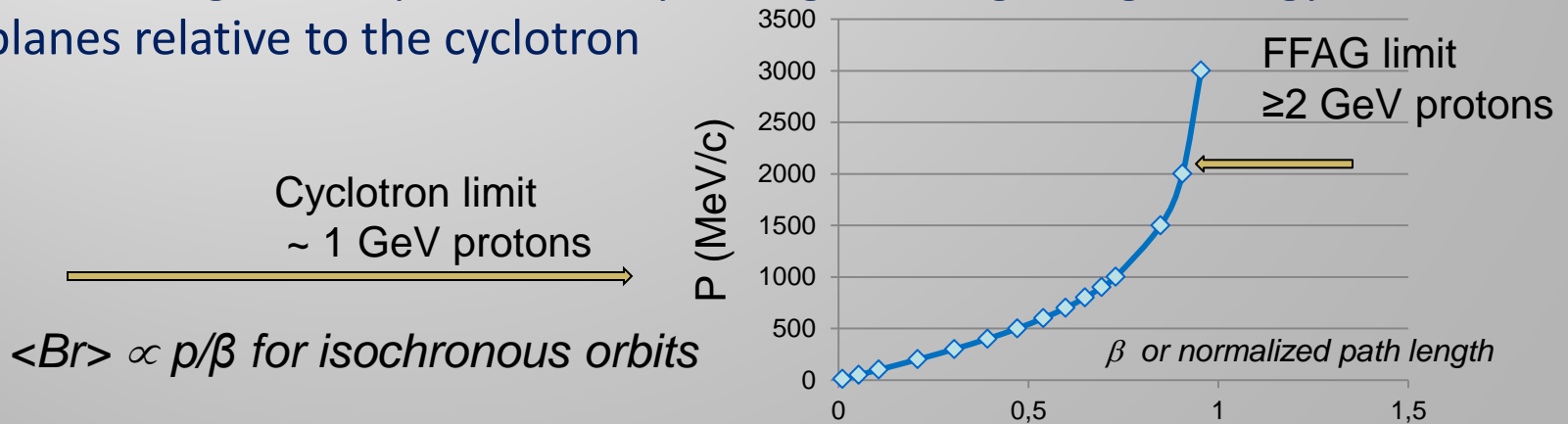


COSY results and tracking at 3 energies: 50 keV, 4.4 MeV, and 9 MeV machine. Except for the tune change required at horizontal injection, the DA is very large.



# Relativistic CW (DC-beam) FFAGs

- ✓ NS FFAG can maintain isochronous orbits at relativistic energies
  - ✓ Pathlength of isochronous orbits are proportional to velocity
  - ✓ Orbits as a function of momentum follow, therefore the B field must scale with velocity
  - ✓ At relativistic energies, momentum is an increasingly nonlinear function of velocity; therefore B field transitions from a linear slope to nonlinear, non-relativistic to relativistic as an approximate function of radius.
  - ✓ **THIS HAS BEEN ACHIEVED IN RECENT NONLINEAR NS FFAG DESIGNS**
  - ✓ Nonlinear field expansion + edge angle can constrain the tune
  - ✓ Nonlinear gradient provides very strong focusing at high energy in **both** planes relative to the cyclotron



# The significance of CW Accelerators

A CW accelerator implies:

Fixed magnetic fields

- 50 Hz is the  $\sim$  practical technical limit for pulsed magnet systems
  - Stored power and expense of pulsed supplies can be commercially prohibitive

The simplicity of fixed-frequency rf

- the rotational frequency of orbits is a constant at all energies

## Consequences of non-isochronous orbits

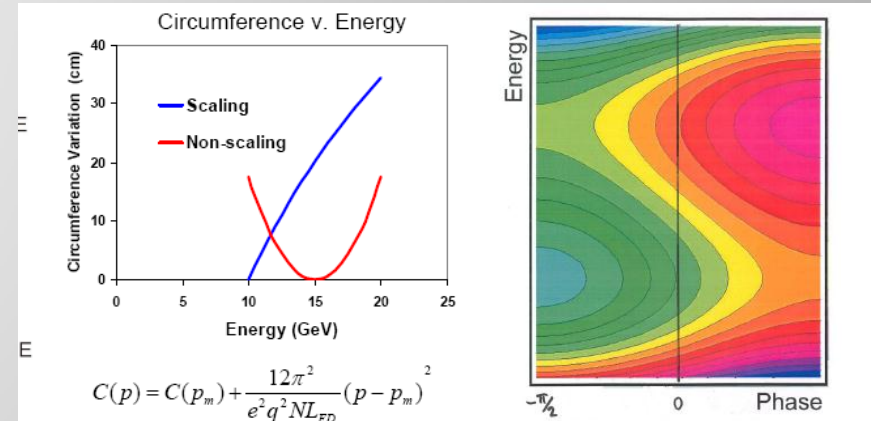
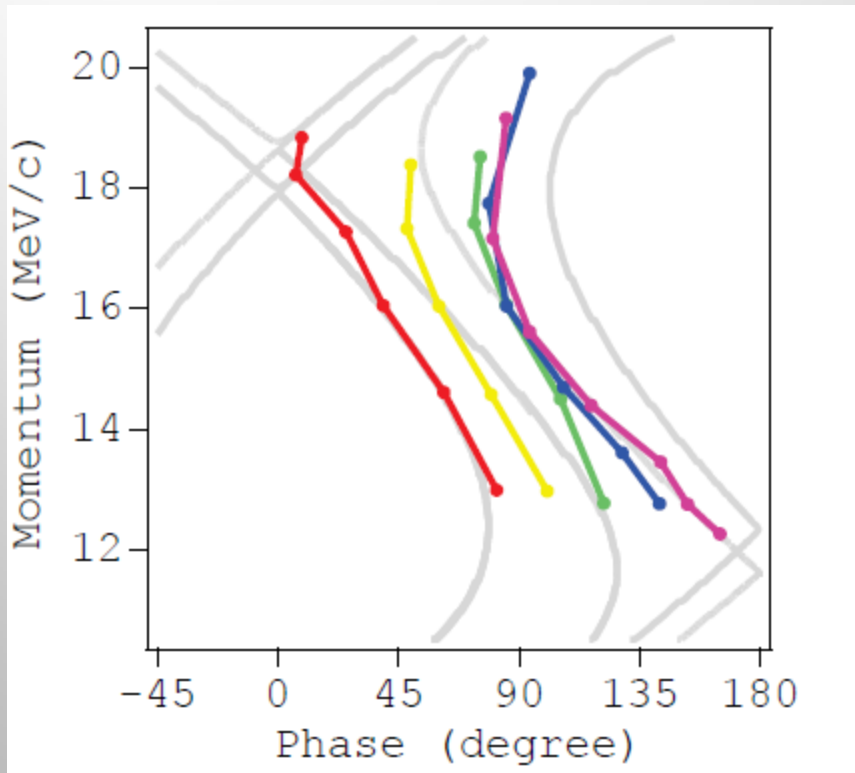
Beam is pulsed at the rf sweep rate, not continuous

Swept-frequency rf (rf timing is changed to match the revolution time of the beam – the synchrotron and synchro-cyclotron)

- 50-100 Hz sweep rate for rf frequencies  $\geq$  tens of MHz
- KHz sweep rate for broad-band rf ( $\sim$ MHz)
  - slow acceleration, high power consumption

# Serpentine channel Acceleration in a nonscaling FFAG

## Acceleration in serpentine channel



$$\delta T \sim \frac{2 V_{min}}{h \Delta E} \sim \frac{1}{turns}$$

example

The machine which takes 100 turns to the top energy needs 1% level of isochronous.

Longitudinal phase space trajectories of beams with five different initial phases in EMMA. All of these cases clearly demonstrate acceleration within the serpentine channel.

# Serpentine CW acceleration in 50 keV – 9 MeV eFFAG

To be isochronous, constant revolution time

$$\langle R_{ext} \rangle / \langle R_{inj} \rangle = \beta_{ext} / \beta_{inj}$$

For this design

$$\langle R_{ext} \rangle \sim 43 \text{ cm}, \beta_{ext} = 0.4127$$

$$\langle R_{inj} \rangle \sim 19 \text{ cm}, \beta_{inj} = 0.9985$$

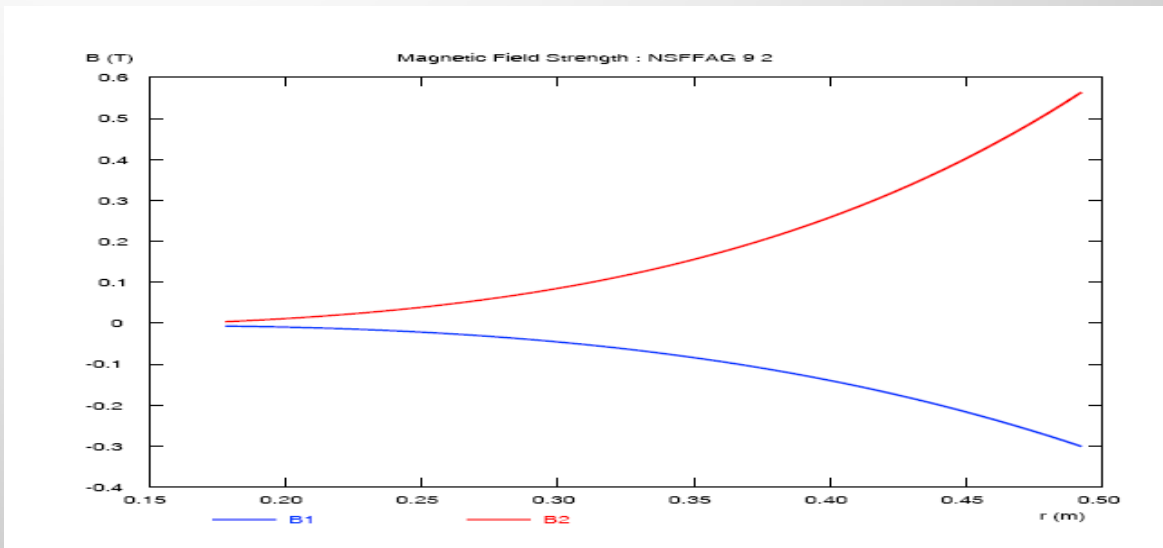
$$\beta_{ext} / \beta_{inj} = 0.4133$$

$$\langle R_{ext} \rangle / \langle R_{inj} \rangle = 0.4418$$

*Isochronous to  $\sim \pm 1-2\%$*

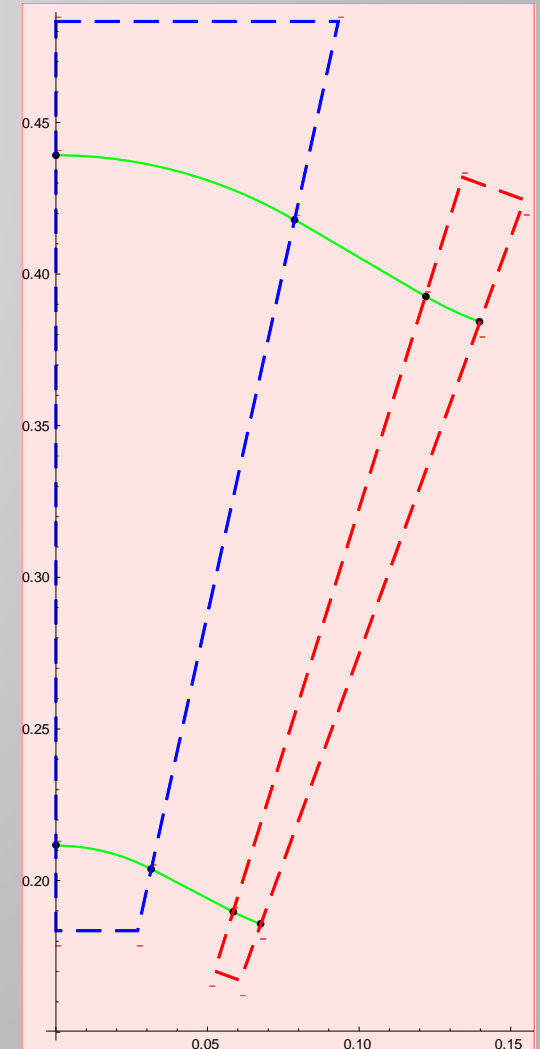
Conclusion: serpentine channel can be used for 50-100 turn acceleration as designed

# Magnetic components and fields for a half cell



Magnetic field profile F and D

Half cell injection and extraction orbits



# RF acceleration system and estimated output power

A single 100-200 keV cavity gives 45-90 acceleration turns

Single Dee Cavity as in cyclotron

Dummy Dee on other side

Can use the serpentine channel for acceleration

Duty cycle:  $1\text{ ns}/10\text{ ns} \sim 10\%$

Factor of 100 increase

Space charge limited to  $\sim 10^9$  electrons/RF micro bunch

For fundamental, 100 MHz cavity (10 ns bunch spacing)

1.6 mA,  $\sim 140$  kW CW beam

# Innovation: permanent magnets

The level of specified magnetic fields ( $<0.4$  T) permits the use permanent magnets. Permanent magnets have the following advantages:

- Do not need power supplies;
- Do need a cooling system and water;
- Low magnet system cost and simple industrialization;
- No hysteresis effects, remnant fields, degaussing, etc.;
- More field concentration under magnet poles;
- Magnets without coupling to AC fields.

However the following issues require control or resolution

- Magnets calibration;
- Thermal effects control;
- Adjust or correct the field.

For this magnet system, permanent magnets based on: ceramic ferrites, SmCo<sub>5</sub> or NdFeB permanent magnets could be used.



# SmCo chosen for permanent magnet material

Although NdFeB magnets B fields up to 1.35 T.

have lower temperature stability than SmCo

must be rust protected by thin Ni coating.

SmCo with B fields just over 1T.

most stable to the temperature change and radiation.

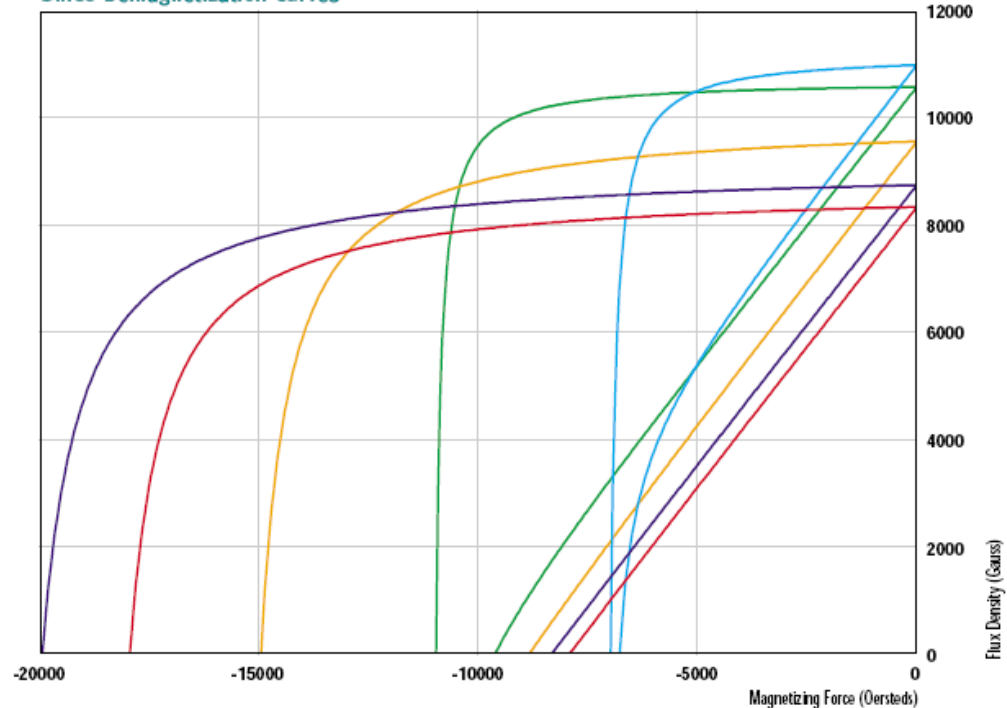
Used in linear accelerators.

high energy concentration in a small volume

MMPA Material Properties for SmCo

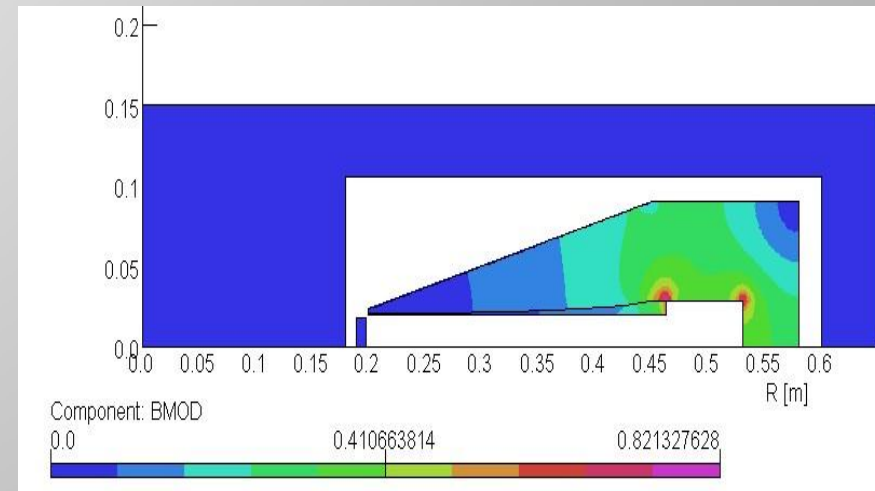
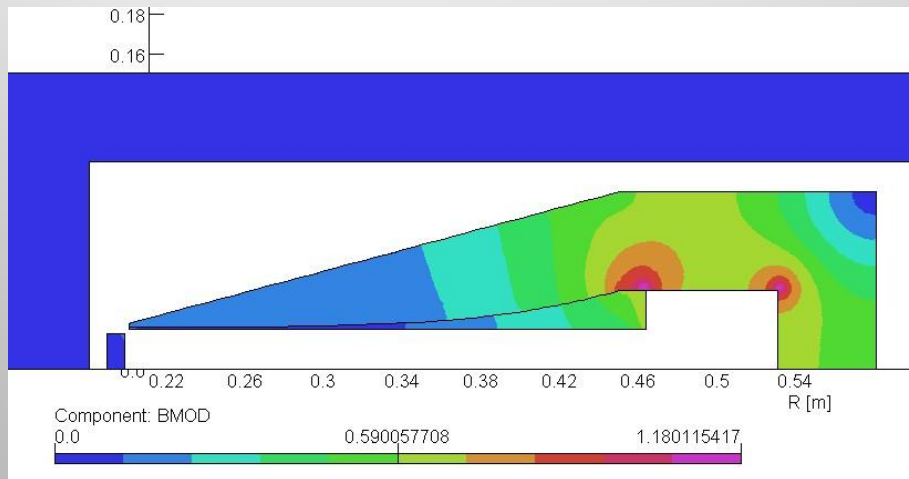
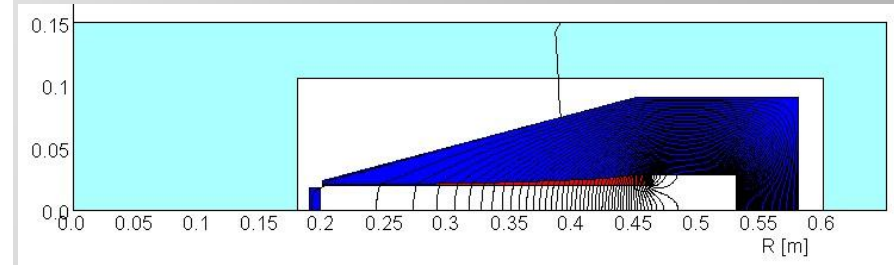
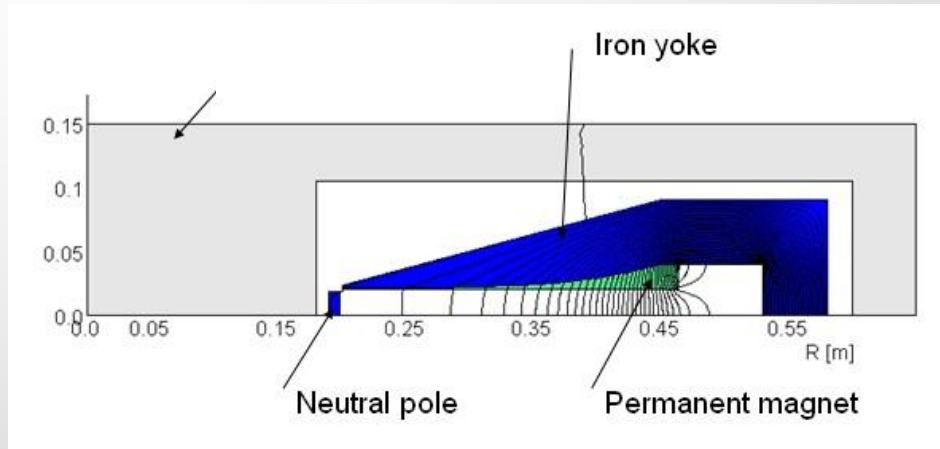
Material	$(BH)_{max}$ (MGOe)	$B_r$ (Gauss)	$H_c$ (Oersteds)	$H_{c2}$ (Oersteds)
SmCo 16/18 (1-5)	16	8300	7500	18000
SmCo 18/20 (1-5)	18	8700	8000	20000
SmCo 22/15 (1-5)	22	9500	9000	15000
SmCo 26/11 (2-17)	26	10500	9000	11000
SmCo 28/7 (2-17)	28	10900	6500	7000

SmCo Demagnetization Curves





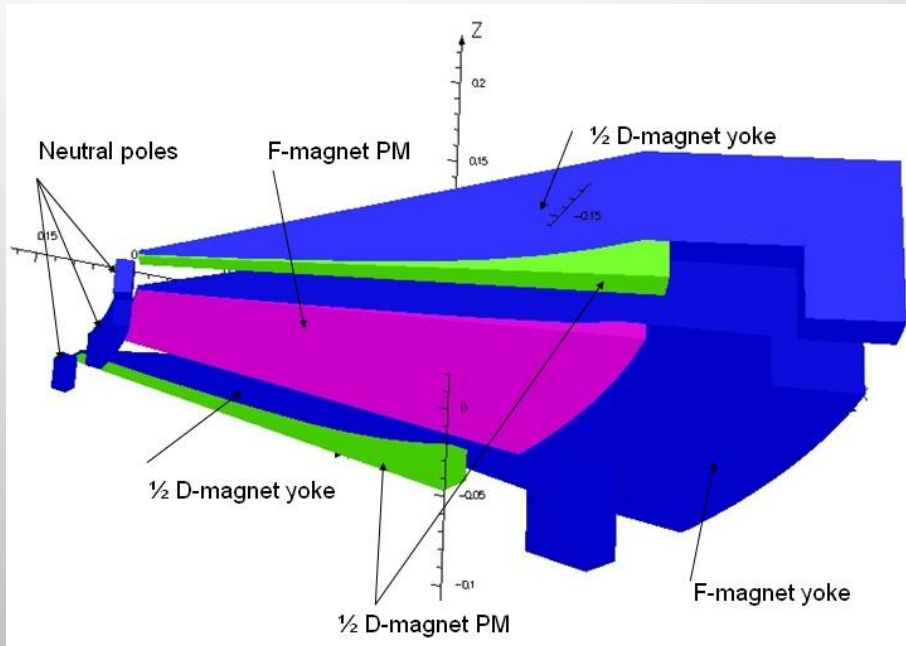
# F and D Magnet design and fluxes



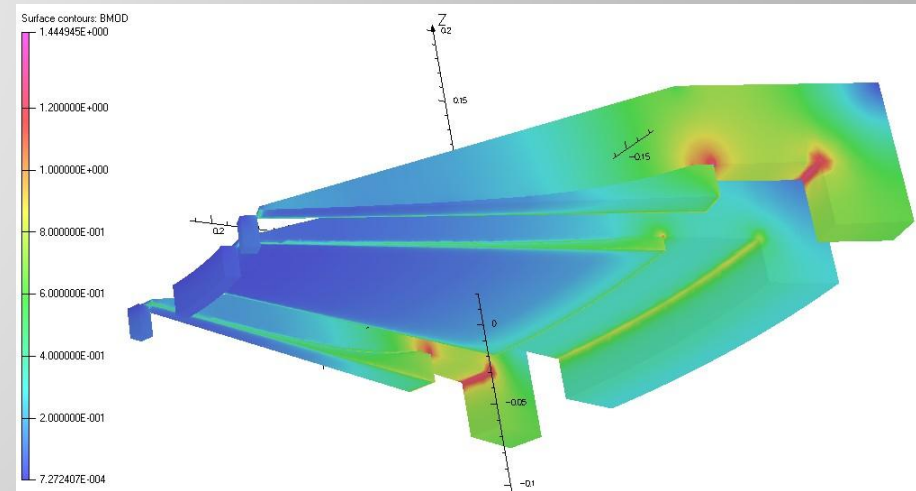
D-magnet geometry and flux lines (top) and flux density in poles (bottom)

F-magnet geometry and flux lines (top) and flux density in poles (bottom)

# Overview of magnet system



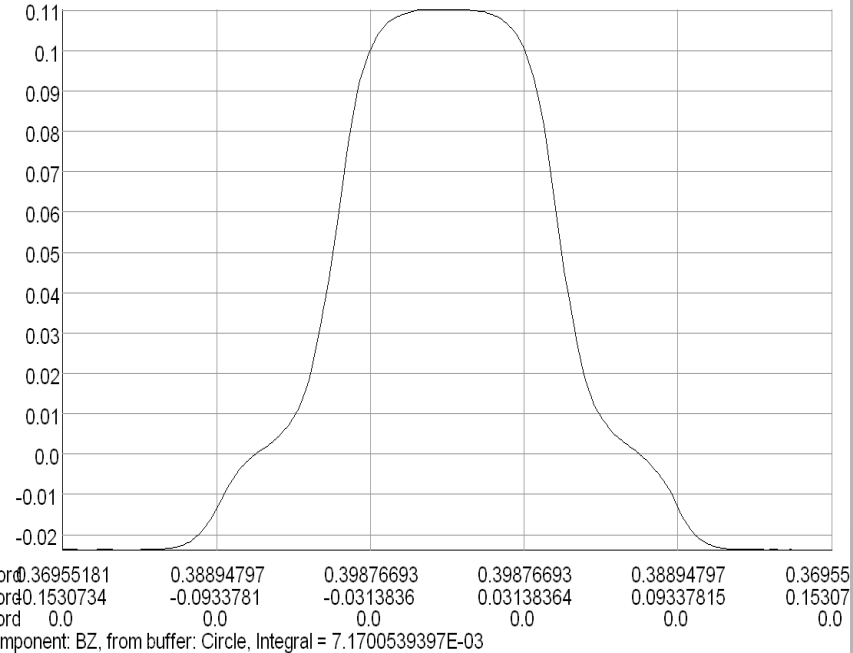
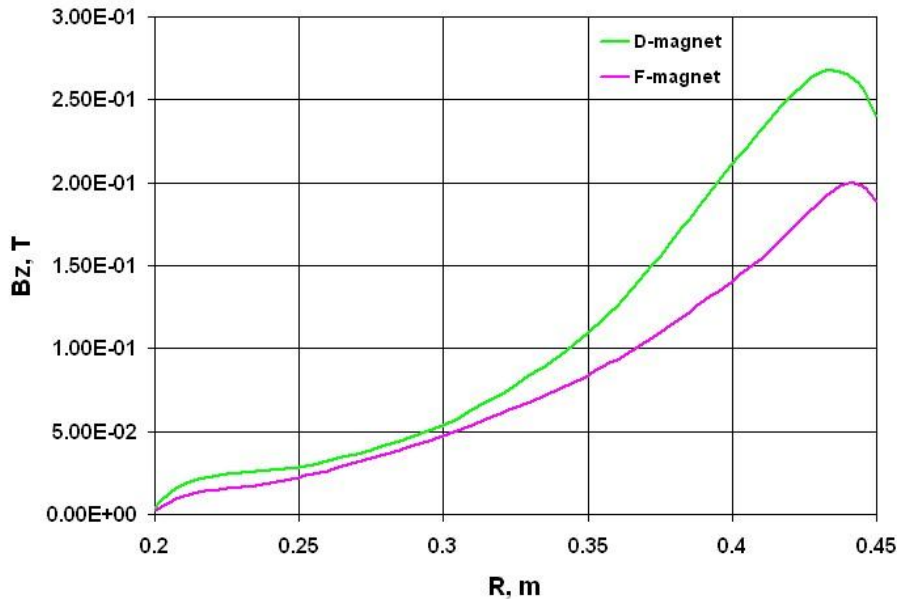
One cell of 3D Model geometry view (upper part from middle plane).



Flux density distribution in the magnet yoke and SmCo5 PM material.

# Magnetic field profiles

SmCo5, gap 40 mm



$B_z$  field distribution in radial direction for the D and F magnet vertical center planes (SmCo5, 40 mm gap)

$B_z$  field in the middle plane on arc  $\pm 22.5^\circ$  with radius 0.4 m (extraction).

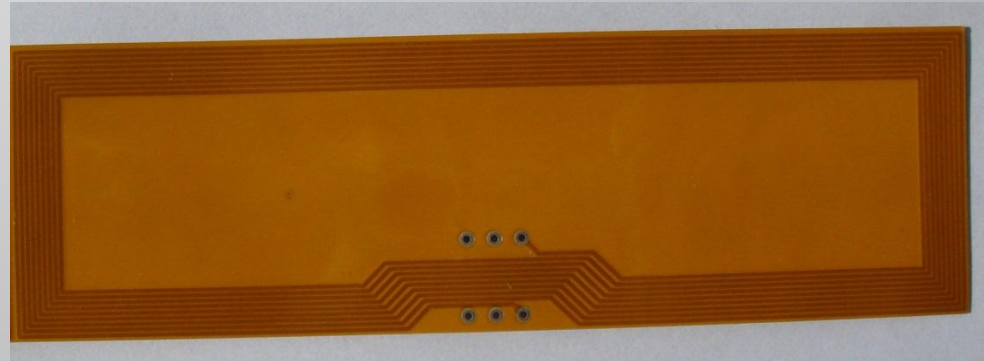
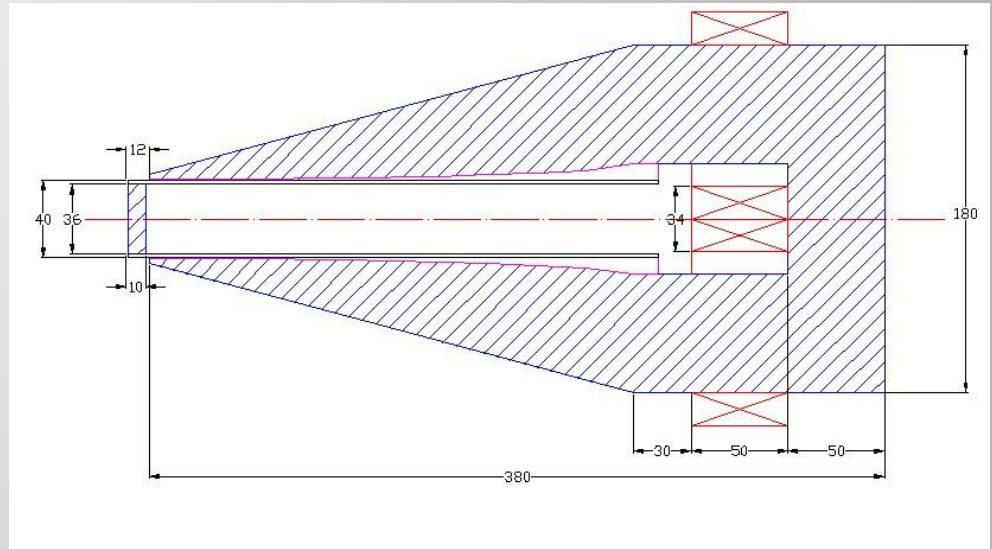
# Correction coils

The main correction coils mounted on the return yokes and installed through the magnet air gap.

They correct the field peak value.

There may be installed also pole coils for fine field tuning at injection radius.

For this, a printed board technique could be used as shown



# Strong market case for compact CW effags

Compact, 1m diameter, transportable

High current - 1-2 mA projected

No power supplies

Inexpensive components