

Nuclear pairing from realistic forces

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Research interests

- Dispersive treatment of πK scattering with G. Colangelo (Bern U.), work in progress.
- Nuclear pairing problem with P. Finelli (Bologna U.) and J. Holt (Washington U.),
Published on PRC 90 044003 (2014)

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Summary

- 1 Nuclear pairing...
- 2 ..from realistic forces
- 3 Results
- 4 Conclusions

Nuclear pairing: introduction

Consider a gas (infinite system) of particles with spin $\frac{1}{2}$ (neutrons):
at $T=0$, density ρ (or Fermi momentum k_F),
the ground state is the Fermi sea $|0\rangle$.

But...

(Cooper theorem)

if we switch on an attractive interaction, $|0\rangle$ is an unstable state.

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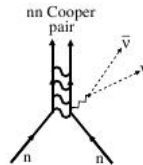
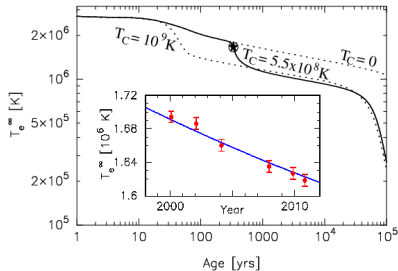
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Neutron pairing and neutron stars

Fast cooling of SN 1987A explained by nuclear pairing:



BCS theory: an overview

$$H = \sum_{s=\pm\frac{1}{2}\mathbf{k}} \varepsilon_{\mathbf{k}} c_{s\mathbf{k}}^{\dagger} c_{s\mathbf{k}} + \sum_{\mathbf{k}} V_{\mathbf{k}\mathbf{k}'} c_{\uparrow\mathbf{k}}^{\dagger} c_{\downarrow-\mathbf{k}}^{\dagger} c_{\downarrow-\mathbf{k}'} c_{\uparrow\mathbf{k}'}$$

In the BCS approximation:

$$\begin{pmatrix} b_{s\mathbf{k}} \\ b_{-s\mathbf{k}}^{\dagger} \end{pmatrix} = \begin{pmatrix} u_{s\mathbf{k}} & v_{s\mathbf{k}} \\ -v_{s\mathbf{k}}^* & u_{s\mathbf{k}}^* \end{pmatrix} \begin{pmatrix} c_{s\mathbf{k}} \\ c_{-s\mathbf{k}}^{\dagger} \end{pmatrix}$$

$$H_{BCS} = \sum_{s=\pm\frac{1}{2}\mathbf{k}} E_{\mathbf{k}} b_{s\mathbf{k}}^{\dagger} b_{s\mathbf{k}}$$

BCS theory: an overview

To solve the problem, we must satisfy:

$$E_{\mathbf{k}} = \sqrt{\varepsilon_{\mathbf{k}}^2 + \Delta_{\mathbf{k}}^2} \quad \Delta_{\mathbf{k}} = \sum_{\mathbf{k}'} \frac{V_{\mathbf{k}\mathbf{k}'}}{2E_{\mathbf{k}'}} \Delta_{\mathbf{k}'}$$

Nonlinear integral equation for the gap $\Delta_{\mathbf{k}}$

\implies numerical techniques.

How do we choose the interaction $V_{\mathbf{k}\mathbf{k}'}$?

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Realistic nuclear forces

Recipe for a modern nuclear force (ChPT):

- find out the most appropriate d.o.f. (N, π);
- write down the **most general** Lagrangian \mathcal{L} allowed by (eventually broken) QCD symmetries;
- expand \mathcal{L} in powers of momenta or masses ($Q/\Lambda_{QCD} \ll 1$);
- introduce a cutoff to remove UV divergences;
- fit the free parameters to reproduce the low energy experimental data;
- check the cutoff-dependence.

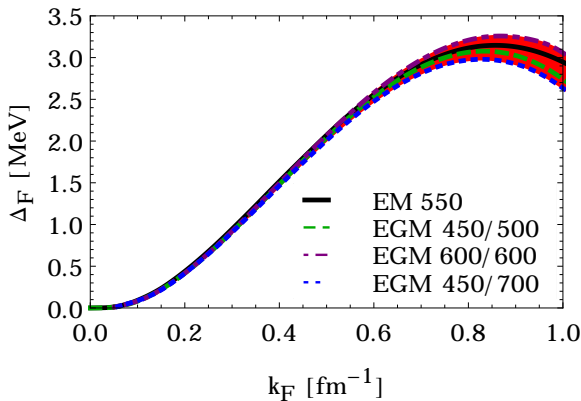
Results for Δ_{k_F} as a function of k_F in neutron matter.

The interaction is decomposed in partial waves
In every density region just one partial wave included.

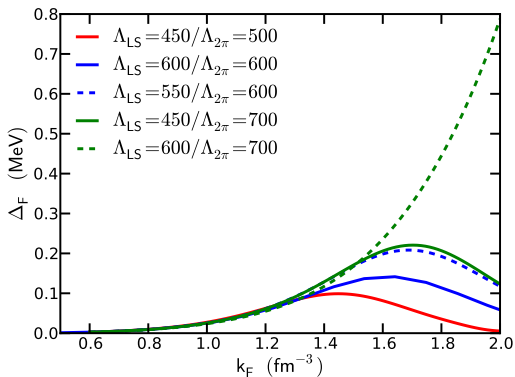
$$V_{\mathbf{k}\mathbf{k}'} = \sum_{l'l'mm'} V_{l'l'}(|\mathbf{k} - \mathbf{k}'|) Y_{lm}^*(\hat{\mathbf{k}}) Y_{l'm}(\hat{\mathbf{k}'})$$

We have used two different interactions (different regularization schemes and NR reduction).

Low density BCS gap (S wave)



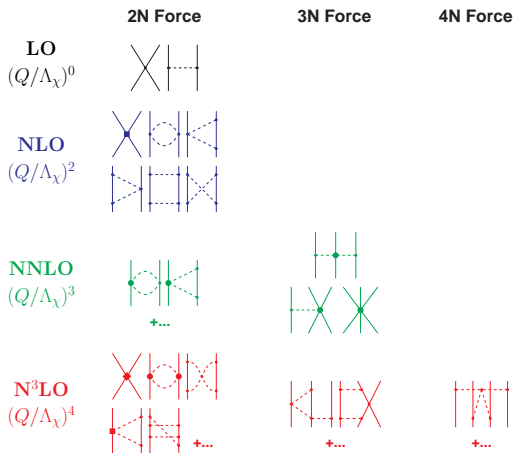
Higher density BCS gap (P-F waves)



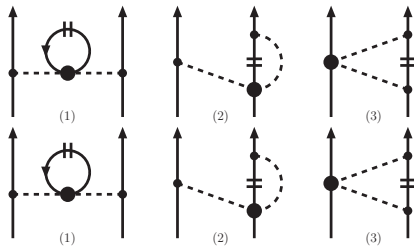
Conclusions

- Nuclear pairing is fundamental to understand supernovae physics.
- Thanks to *ChPT* we are now able to obtain informations about nuclear physics from fundamental (QCD) principles + experiment.
- At low densities, we obtained a good cutoff independence.
- Further corrections are needed to better understand the nuclear pairing.
- The finite temperature extension will be soon available.

Structure of the nuclear interaction

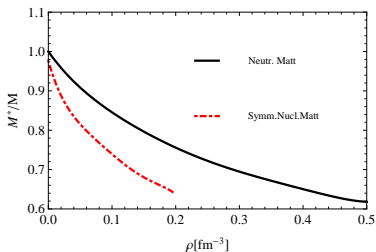


NNN interaction included via den. dep. NN int.:



Holt *et al.* PRC 81 024002 (2010)

Self-energy corrections included via effective mass (density functional):



Holt, J.W. et al. Eur.Phys.J. A47 (2011) 128