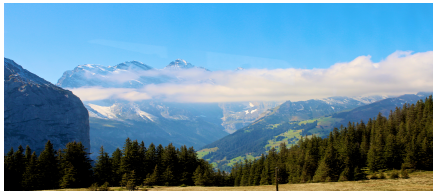


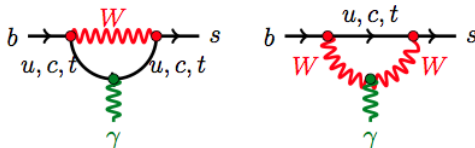
Photon polarisation in $b \rightarrow s\gamma$ transition at LHCb

Zhirui XU (EPFL)



Photon polarisation in $b \rightarrow s\gamma$ transition

- Transitions driven by FCNC represent pure quantum effects within the SM



- Loop-driven B decays are more sensitive to the presence of New Physics beyond SM.
- The SM photon in $b \rightarrow s\gamma$ is predominantly left-handed

$$\bar{s}\Gamma_{\mu}^{b \rightarrow s\gamma} b = \frac{e}{(4\pi)^2} \frac{g^2}{2M_W^2} V_{ts}^* V_{tb} F_2 \bar{s} i\sigma_{\mu\nu} q^{\nu} \left(m_b \frac{1 + \gamma_5}{2} + m_s \frac{1 - \gamma_5}{2} \right) b$$

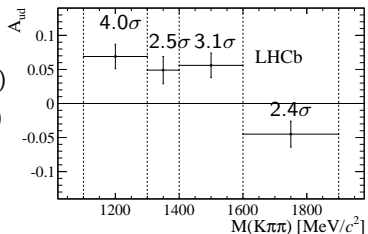
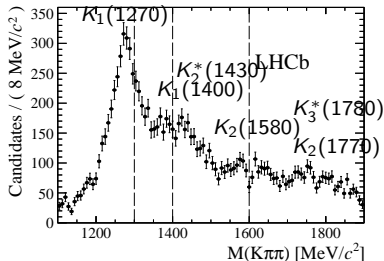
$b_R \rightarrow s_L \gamma_L$

$b_L \rightarrow s_R \gamma_R$

- The right-handed contribution can be significantly enlarged due to new physics.

Photon polarisation in radiative B decays

- Measuring the photon polarisation
 - The time-dependent CP-asymmetry in $B_{(s)} \rightarrow f^{\text{CP}} \gamma$: $B_s^0 \rightarrow \phi \gamma$, $B^0 \rightarrow K_S^0 \pi^0 \gamma$
 - Angular correlations among the three-body decay products of the excited kaons in $B \rightarrow K_{\text{res}}(P_1 P_2 P_3) \gamma$: $B \rightarrow K_1(K \pi \pi) \gamma$, $B \rightarrow \phi K \gamma$
 - Transverse asymmetry in $B^0 \rightarrow K^*(892)^0 l^+ l^-$
 - Direct measurement of the photon polarisation in baryons decays: $\Lambda_b \rightarrow \Lambda^{(*)} \gamma$, $\Xi_b \rightarrow \Xi^{(*)} \gamma$
- Photon polarisation in $b \rightarrow s \gamma$ transition first observed in $B \rightarrow K \pi \pi \gamma$ [PRL 112, 161801 (2014)].



Time dependent CP asymmetry in $B_s^0 \rightarrow \phi\gamma$

The time-dependent decay rate for initial B_s^0 and \bar{B}_s^0 decaying into $\phi\gamma$:

$$\Gamma_{B_s^0 \rightarrow \phi\gamma} \propto |A|^2 e^{-\Gamma_s t} \left(\cosh \frac{\Delta\Gamma_s t}{2} - A^\Delta \sinh \frac{\Delta\Gamma_s t}{2} + C \cos \Delta m_s t - S \sin \Delta m_s t \right)$$

$$\Gamma_{\bar{B}_s^0 \rightarrow \phi\gamma} \propto |A|^2 e^{-\Gamma_s t} \left(\cosh \frac{\Delta\Gamma_s t}{2} - A^\Delta \sinh \frac{\Delta\Gamma_s t}{2} - C \cos \Delta m_s t + S \sin \Delta m_s t \right)$$

respectively, where

$$C = \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2} \approx 0 \text{ (SM)}$$

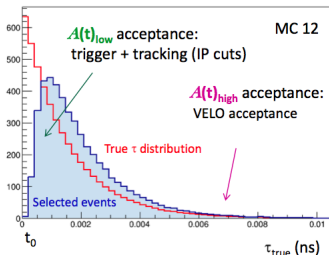
$$S \approx \sin 2\psi \sin \varphi_s \text{ (}\varphi_s \text{ is the } B_s \text{ mixing phase)}$$

$$A^\Delta \approx \sin 2\psi \cos \varphi_s$$

ψ is given by the fraction of “wrongly”-polarized photons and defined as

$$\tan \psi = \left| \frac{A(\bar{B}_s^0 \rightarrow \phi\gamma_R)}{A(\bar{B}_s^0 \rightarrow \phi\gamma_L)} \right|$$

- Physics: $\Gamma_{B_s^0 \rightarrow \phi \gamma} + \Gamma_{\bar{B}_s^0 \rightarrow \phi \gamma} \propto e^{-\Gamma_{B_s^0} t} \left(\cosh \frac{\Delta\Gamma_s t}{2} - A^\Delta \sinh \frac{\Delta\Gamma_s t}{2} \right)$
- Acceptance: trigger, selection and reconstruction requirements



- Key in the photon polarization measurement;
- Need to be precisely determined or controlled.

- Resolution: dominated by the photon momentum, from Monte Carlo
- Background: the B mass distribution

Tracking:

$$\Delta p/p \sim 0.4\% \text{ at } 5\text{GeV}$$

$$\sigma_{\text{IP}} \sim 20\mu\text{m} \text{ for high-}p_T \text{ tracks and } \sigma_{\tau} \sim 45\text{fs}$$

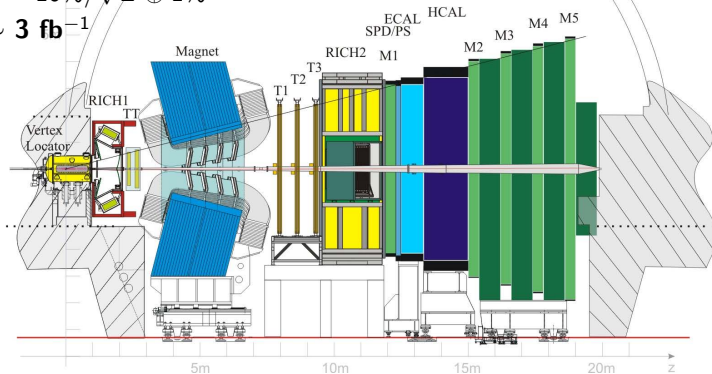
Particle identification:

$$\pi/K \text{ separation over } 2\text{-}100\text{GeV} (\epsilon_K \sim 90\% \text{ for } \sim 5\% \pi \rightarrow K \text{ mid-id})$$

Calorimeter system:

$$\sigma_E/E \sim 10\%/\sqrt{E} \oplus 1\%$$

Run I: $\sim 3 \text{ fb}^{-1}$



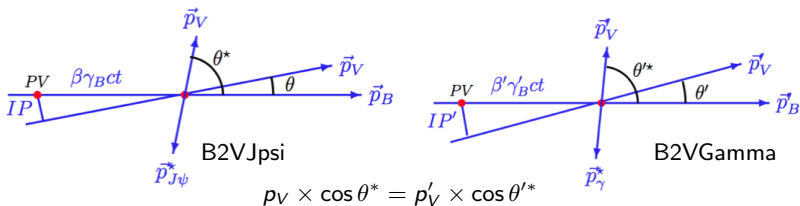
- The acceptance function $A(t)$ can be parametrized as:

$$A(t) = \frac{[a(t - t_0)]^n}{1 + [a(t - t_0)]^n} \cdot e^{-\delta\Gamma t}$$

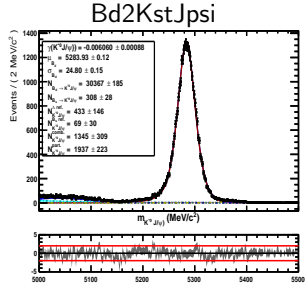
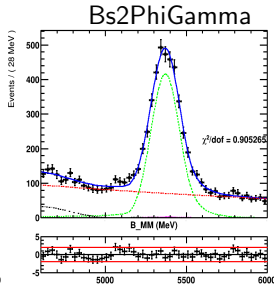
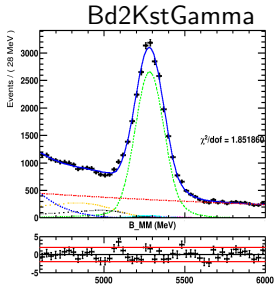
$A(t)_{\text{low}}$
 $A(t)_{\text{high}}$

- Similar radiative decay: $B^0 \rightarrow K^{*0}\gamma$
 - Similar topology, same trigger for the photon, similar resolution, large statistics
 - Large background
 - Avoid selection cuts to introduce difference in acceptance
 - RL (Ratio Loose, 30000) and 4500 $B_s^0 \rightarrow \phi\gamma$
 - RT (Ratio Tight, 20000) and 3500 $B_s^0 \rightarrow \phi\gamma$
 - and DF (Direct Fit, 22000) and 3500 $B_s^0 \rightarrow \phi\gamma$

- Background free channel: $B^0 \rightarrow K^{*0} J/\psi$
 - Kinematic difference: mass difference between photon and J/ψ and the helicity difference between two channels.
 - $A(t)_{\text{high}}$: constrained from Monte Carlo in the final fit
 - $A(t)_{\text{low}}$: two geometrical variables related to vertex displacement — A 2D kinematical reweighting: helicity angle $\cos \theta_H$ and $p_V \cdot \cos \theta$ and then scale the related variables by a constant scaling factor



- Events selected: DF, 30000 (3500 $B_s^0 \rightarrow \phi \gamma$)



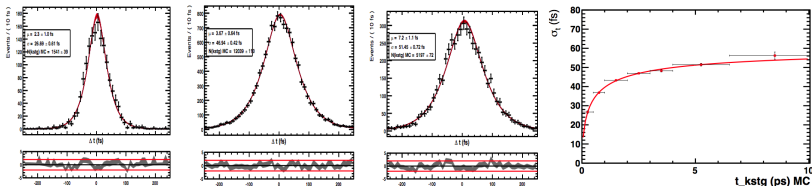
- $B^0 \rightarrow K^{*0} \gamma$ and $B_s^0 \rightarrow \phi \gamma$
 - Use sWeights from mass fits:
the correlation between mass and proper time \Rightarrow systematics
 - Direct background subtraction: crosscheck
- $B^0 \rightarrow K^* J/\psi$: background free in the signal region

- The proper time bias is reduced to about 5 fs and constant over the different calorimeter regions after applying the photon calibration.
- The Apollonios parametrization of the proper time resolution:

$$R(t) = e^{-b\sqrt{1+\frac{(t-\mu)^2}{\sigma^2}}}$$

- And the evolution of the proper time resolution σ as a function of the reconstructed proper time is fitted with

$$\sigma(t) = \sigma_0 \times \frac{t^n}{1 + \alpha t^n}$$



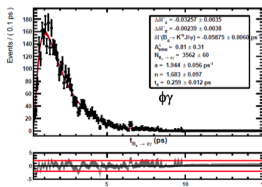
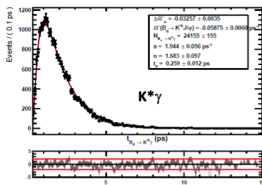
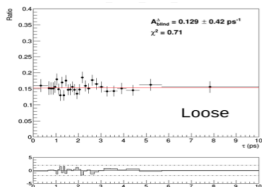
- Two approaches:

- Ratio fit: $B_s^0 \rightarrow \phi\gamma / B^0 \rightarrow K^{*0}\gamma$

$$\frac{d\Gamma_{B_s^0 \rightarrow \phi\gamma}/dt}{d\Gamma_{B^0 \rightarrow K^{*0}\gamma}/dt} \propto \frac{R_{B_s^0}(t, t') \otimes N_{B_s^0} \cdot \epsilon_{B_s^0}(t') \cdot e^{-\Gamma_{B_s^0} t'} (\cosh \frac{\Delta\Gamma_s t'}{2} - A^\Delta \sinh \frac{\Delta\Gamma_s t'}{2})}{R_{B^0}(t, t') \otimes N_{B^0} \cdot \epsilon_{B^0}(t') \cdot e^{-\Gamma_{B^0} t'}}$$

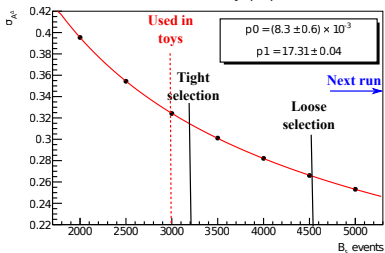
- Direct fit: unbinned simultaneous ML fit to

- $B_s^0 \rightarrow \phi\gamma$ and $B^0 \rightarrow K^{*0}\gamma$
- $B_s^0 \rightarrow \phi\gamma$ and $B^0 \rightarrow K^{*0}J/\psi$
- $B_s^0 \rightarrow \phi\gamma$ and $B^0 \rightarrow K^{*0}\gamma$ and $B^0 \rightarrow K^{*0}J/\psi$



Measurement of A^Δ : blind analysis

Statistical sensitivity (proper-time)



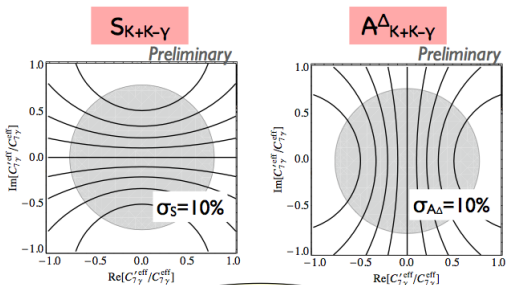
- The blinding method: add random number to result, i.e. A^Δ value
- Statistical uncertainty is worsened by using sPlot for background subtraction
- Systematics uncertainty is in progress: fit model, mass-time correlation, scaling and reweighting, calorimeter calibration, etc.

1	Introduction	1
1.1	Key measurements	3
1.2	The LHCb trigger	4
2	The tree-level determination of γ	8
3	Charmless charged two-body B decays	58
4	Measurement of mixing-induced CP violation in $B_s^0 \rightarrow J/\psi\phi$	150
5	Analysis of the decay $B_s^0 \rightarrow \mu^+\mu^-$	229
6	Analysis of the decay $B^0 \rightarrow K^{*0}\mu^+\mu^-$	275
7	Analysis of $B_s^0 \rightarrow \phi\gamma$ and other radiative B decays	313

Roadmap for selected key measurements of LHCb

Next $B_s^0 \rightarrow \phi\gamma$ study

- Combining S and A^Δ of $B_s^0 \rightarrow \phi\gamma$ will lead to a very strong constraint on the C_7 and C_7' plane



LHCb

from Emi Kou

- Instead of a combined sample, the measurement of S requires tagged sample.

The End

