

Theory uncertainty of γ from tree decays

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Partially based on
[Brod, Zupan JHEP 1401 (2014) 051]

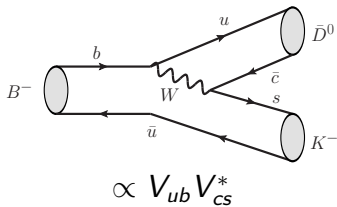
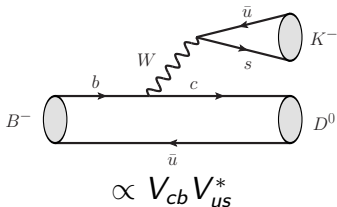
The main point

- $\gamma \equiv \arg(-V_{ud}V_{ub}^*/V_{cd}V_{cb}^*)$
- Determination from $B \rightarrow DK$ theoretically extremely clean
 - Fit hadronic parameters from data
 - Irreducible theory uncertainty is tiny

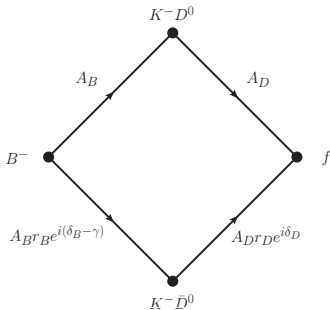
Outline

- γ from tree decays
- Theory uncertainty
- Electroweak corrections

γ from tree decays – general idea



- $b \rightarrow c\bar{u}s, b \rightarrow u\bar{c}s$
- no penguin contribution
- interference from common D^0, \bar{D}^0 final states



Several choices for final state f in D decay

- CP eigenstates (e.g. $D \rightarrow K^+K^-, \pi^+\pi^-; K_S\pi^0$)

[Gronau, London, Wyler 1990/91]

- Flavor states (e.g. $D \rightarrow \pi^-K^+, \pi^+K^-$) [Atwood, Dunietz, Soni 1997]

- Many-body final states (e.g. $D \rightarrow K_S K^+K^-, K_S\pi^+\pi^-$)

[Giri, Grossman, Soffer, Zupan 2003; Poluektov 2004]

- Many variants:

- Use $D^* \rightarrow D\pi^0, D\gamma$ [Bondar, Gershon 2004]

- Many-body B final states

[Aleksan, Petersen, Soffer 2002, Gershon 2008; Gershon, Poluektov 2009]

- Neutral B_d, B_s [Aleksan, Dunietz, Kayser 1992; Kayser, London 2000; Atwood, Soni 2003;

Fleischer 2003; Gronau et al. 2004]

- . . .

- Have $\sim n_D n_B$ measurements, $\sim n_D + n_B$ unknowns

- Test for systematics or NP by comparing different methods

Why measure γ ?

- Important SM input parameter
- Theoretically very clean
 - Hadronic parameters can be fitted from data
- Unlikely that there is NP at current precision
[Brod, Lenz; work in progress]
- In the (far) future can be used to search for NP at high scales
- \Rightarrow what is the theory uncertainty?

Theory uncertainty

- Several sources can induce a shift in γ
- Dalitz plot errors
 - \Rightarrow More statistics
- $D - \bar{D}$, $K - \bar{K}$, $B_{(s)} - \bar{B}_{(s)}$ mixing
[E.g. Bondar et al., 1004.2350; Rama 1307.4384; Gronau, Grossman, Surujon, Zupan 0702011]
- Direct CP violation in D decays
[E.g. Martone, Zupan 1212.0165; Wang 2012; Bhattacharya et al. 2013; Bondar et al., 2013]
- Higher-order electroweak effects

$D - \bar{D}$ mixing

- In SM, $D - \bar{D}$ mixing is CP conserving
 - \Rightarrow the effect is small
- $D - \bar{D}$ mixing dilutes interference term [Grossman, Soffer, Zupan 2005]
- $\Delta\gamma \propto \epsilon_f \sim \mathcal{O}(x_D^2/r_f^2, y_D^2/r_f^2) < 1^\circ$
 - Practically shift is still linear [Rama 2014]
- In model-indep. Dalitz analysis no change needed
- When using information from correlated $D\bar{D}$ pairs from charm factories, shift is linear [Bondar, Poluektov, Vorobiev 2010]
 - Numerically still small, $\Delta\gamma \leq 2.9^\circ$
- $D - \bar{D}$ mixing can be included exactly if x_D, y_D precisely measured!

Direct CP violation in SCS D decays

$$\Delta\mathcal{A}_{CP} = (-0.253 \pm 0.104)\%$$

[HFAG May 2014]

- Modify amplitudes D decay amplitudes.

$$A_f \equiv A(D^0 \rightarrow f) = A_f^T [1 + r_f e^{i(\delta_f + \phi_f)}],$$

$$\bar{A}_f \equiv A(\bar{D}^0 \rightarrow f) = A_f^T [1 + r_f e^{i(\delta_f - \phi_f)}].$$

$$a_f^{\text{dir}} \approx -2r_f \sin \delta_f \sin \phi_f, \quad r_f = \mathcal{O}(10^{-3})$$

Including direct CP violation

- B decay amplitude gets modified:

$$A(B^\pm \rightarrow f_D K^\pm) = A_B A_f^T [1 + r_B^\pm e^{i(\delta'_B \pm \gamma \pm \delta\gamma)}]$$

- Good news: can still solve the system
 - Unknowns: $2n_{\text{SCS}} + 3n_B + 1$
 - Observables: $2n_B(n_{\text{CA}} + n_{\text{SCS}})$
- Shift symmetry $\gamma \rightarrow \gamma + \phi$, $\alpha_f \rightarrow \alpha_f - \phi$
 - $\alpha_f \equiv \arg(A_f/\bar{A}_f) = -a_f^{\text{dir}} \cot \delta_f$
- Cannot extract γ from $B \rightarrow DK$ alone without assumptions

Including direct CP violation

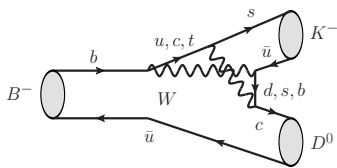
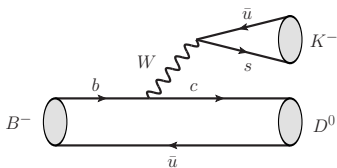
- What is the size of the effect?

$$\delta\gamma = \mathcal{O}(r_f/r_B), \quad \delta'_B - \delta_B = \mathcal{O}(r_f/r_B), \quad r_B^\pm - r_B = \mathcal{O}(r_f)$$

- $r_B(DK) = \mathcal{O}(10\%)$, $\delta\gamma = \mathcal{O}(\text{few } \%)$
- $r_B(D\pi) = \mathcal{O}(0.5\%)$, $\delta\gamma = \mathcal{O}(1)$

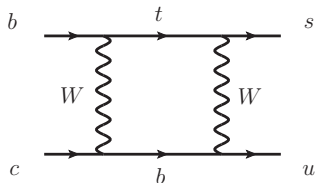
Irreducible theory uncertainty

- Many effects (mixing, CP violation) can be incorporated by modifying equations and extracted from data
- QED radiative corrections \rightarrow CP conserving
- Electroweak corrections
 - No effect from Z exchange
 - No effect from vertex corrections
 - **Box diagrams can change CKM structure**

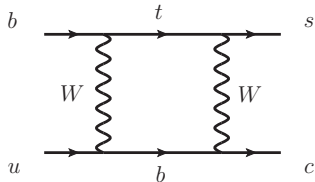


Box diagrams

- $b \rightarrow u\bar{c}s$:
 - tree level $\sim V_{ub} V_{cs}^*$
 - box diagram $\sim (V_{tb} V_{ts}^*)(V_{ub} V_{cb}^*)$
- same weak phase, no shift in γ



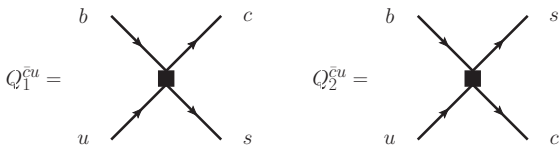
- $b \rightarrow c\bar{u}s$:
 - tree level $\sim V_{cb} V_{us}^*$
 - box diagram $\sim (V_{tb} V_{ts}^*)(V_{cb} V_{ub}^*)$
- different weak phase, induces $\delta\gamma$



Effective Hamiltonians at scale m_b

$$\mathcal{H}_{\bar{c}u}^{(0)} = \frac{G_F}{\sqrt{2}} V_{cb} V_{us}^* [C_1(\mu_b) Q_1^{\bar{c}u} + C_2(\mu_b) Q_2^{\bar{c}u}],$$

$$\mathcal{H}_{\bar{u}c}^{(0)} = \frac{G_F}{\sqrt{2}} V_{ub} V_{cs}^* [C_1(\mu_b) Q_1^{\bar{u}c} + C_2(\mu_b) Q_2^{\bar{u}c}].$$



- Electroweak corrections induce a shift:

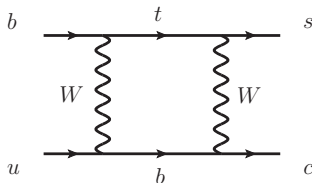
$$\mathcal{H}_{\bar{c}u}^{(1)} = \frac{G_F}{\sqrt{2}} V_{cb} V_{us}^* [(C_1 + \Delta C_1)(\mu_b) Q_1^{\bar{c}u} + (C_2 + \Delta C_2)(\mu_b) Q_2^{\bar{c}u}]$$

Direct Matching

- Integrate out W , top, bottom simultaneously

$$\Delta C_2 = - \left| \frac{\alpha}{4\pi \sin^2 \theta_w} \frac{V_{tb} V_{ts} V_{ub}}{V_{us}} \right| F(x_t, y_b) e^{i\gamma}$$

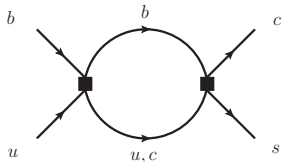
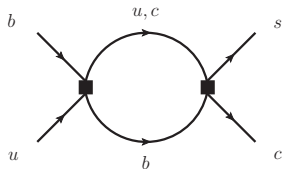
$$\Delta C_1 = 0$$



- $F(x_t, y_b)$ function of $x_t = m_t^2/M_W^2$, $y_b = m_b^2/M_W^2$
- $F(x_t, y_b) \rightarrow 0$ for $x_t \rightarrow 0$, $y_b \rightarrow 0$
- $F(x_t, y_b) \simeq 2y_b \log y_b$
- $|\Delta C_2| = (5.9 \pm 0.3) \times 10^{-8} \quad \propto \lambda^4 \times \alpha/4\pi \times F$

Summation of the log

- Logarithmic term dominates $|\Delta C_2|$
- Summation of the leading-log (QCD) effects
- Matching in two steps:
 - 1 At scale m_t integrate out W , top
 - 2 At scale m_b match to four-flavor effective theory



Solving the RGE

- Resulting RGE can be solved with standard methods (cf. ϵ_K)

$$\mu \frac{d}{d\mu} \tilde{C}_k = \sum_l \tilde{C}_l \gamma_{lk} + \sum_{ij} C_i C_j \hat{\gamma}_{ij,k},$$

- Matching to effective four-flavor theory yields

$$\Delta C_k(\mu_b) = 2m_b^2 \frac{\sqrt{2}G_F}{16\pi^2} \left| \frac{V_{tb} V_{ts} V_{ub}}{V_{us}} \right| e^{i\gamma} \tilde{C}_k^{(0)}(\mu_b).$$

- Expansion about $\mu = M_W$ reproduces the leading log
- Numerically

- $|\Delta C_1| = (4.5 \pm 0.2) \times 10^{-9}$

- $|\Delta C_2| = (4.3 \pm 0.2) \times 10^{-8}$

The shift $\delta\gamma$

- Need the ratio of matrix elements

$$r_A = \frac{\langle K^- D^0 | Q_2^{\bar{c}u} | B^- \rangle}{\langle K^- D^0 | Q_1^{\bar{c}u} | B^- \rangle} = \frac{f_D F_0^{B \rightarrow K}(0)}{f_K F_0^{B \rightarrow D}(0)} = 0.4$$

- Setting $\gamma = 68^\circ$, this yields

$$\delta\gamma = \frac{\text{Im}\Delta C_1}{C_1 + C_2 r_A} + \frac{\text{Im}\Delta C_2}{C_1/r_A + C_2} = 2.0 \times 10^{-8}$$

- Nonlocal contributions expected to lead to uncertainty of a factor of a few

$$\delta\gamma/\gamma \lesssim \mathcal{O}(10^{-7})$$

Summary

- γ from $B \rightarrow DK$ is theoretically extremely clean,

$$\delta\gamma/\gamma \lesssim \mathcal{O}(10^{-7})$$

- Important SM input parameter
- Search for new physics
- Questions:
 - Size of e/w corrections for $B \rightarrow D\pi$
 - Is there really no NP in tree decays?
 - ...

Backup

Summation of the log

- Using GIM and keeping only relevant terms, can write

$$\mathcal{H}_{\text{eff}}^{f=5} = \frac{G_F}{\sqrt{2}} \sum_{\substack{u_{1,2}=u,c \\ d_{1,2}=s,d,b}} V_{u_1 d_2} V_{u_2 d_1}^* \sum_{i,j=1}^2 C_i(\mu) Z_{ij} Q_j^{(u_1 d_2; d_1 u_2)}$$

$$+ 2G_F^2 V_{cb} V_{us}^* \cdot \left| \frac{V_{tb} V_{ts} V_{ub}}{V_{us}} \right| e^{i\gamma} \left[\sum_{i,j,k=1}^2 C_i C_j \hat{Z}_{ij,k} \tilde{Q}_k + \sum_{l,k=1}^2 \tilde{C}_l \tilde{Z}_{lk} \tilde{Q}_k \right]$$

- Effective dim. 8 operators

$$\tilde{Q}_1 = \frac{m_b^2}{\mu^{2\epsilon} g_s^2} (\bar{s}u)_{V-A} (\bar{c}b)_{V-A}, \quad \tilde{Q}_2 = \frac{m_b^2}{\mu^{2\epsilon} g_s^2} (\bar{s}b)_{V-A} (\bar{c}u)_{V-A}.$$

- Double insertions of dim. 6 operators generate \tilde{C}