

Implications of the broad charm resonance measurement in $B \rightarrow K\ell\ell$

Lyon and RZ 1406.0566v1 (v2 to appear)

CP³ Origins
Cosmology & Particle Physics



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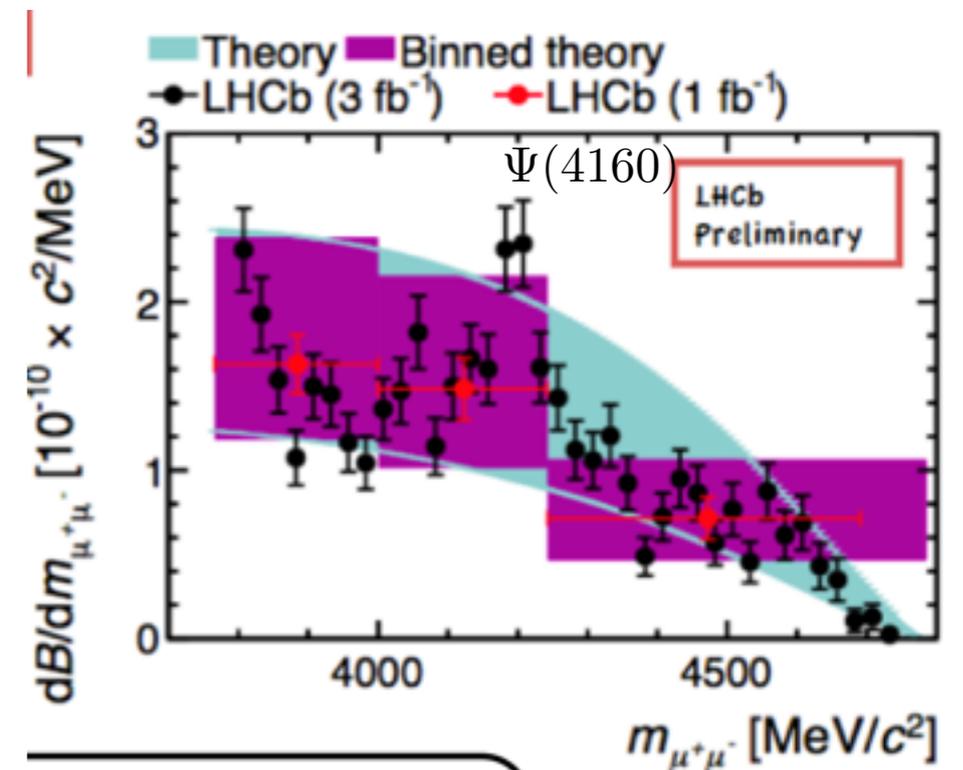
16 October 2014 — Implication of LHCb measurements (CERN)

Overview and main messages...

1) *assessment*

- A. **factorisation** (failure)
- B. **fits** probing **non-fac.** contributions

question: is it SM(QCD) or not?



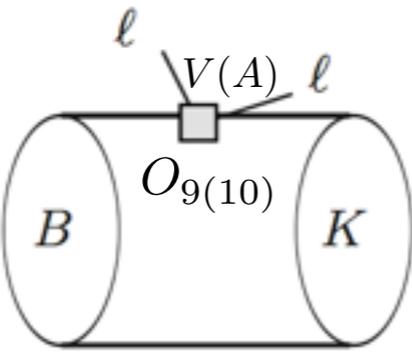
2) *implications and analysis*

- A. **right-handed currents** versus non-universal **polarisation**
- B. how surprising charm resonances
 - i) in view **$B \rightarrow J/\Psi K$** & ii) quark-hadron duality
- C. **phase of J/Ψ unknown** (experiment)
impacts low q^2 e.g. $\text{Br}(B \rightarrow K\ell\ell)$ and P_5'

Basics of $B \rightarrow K^{(*)} \ell \ell$ for this talk

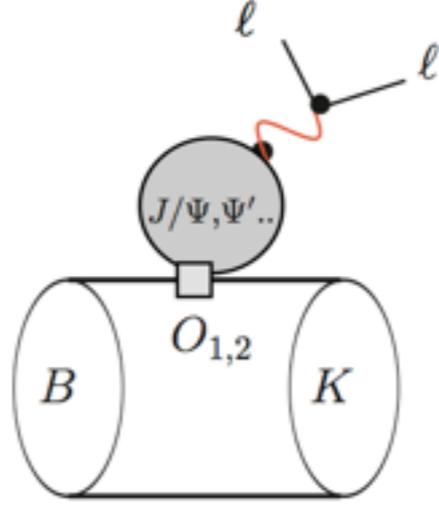
- main actors of this talk (same quantum numbers!)

short distance



electroweak penguin (also $O_{7..}$)

long distance



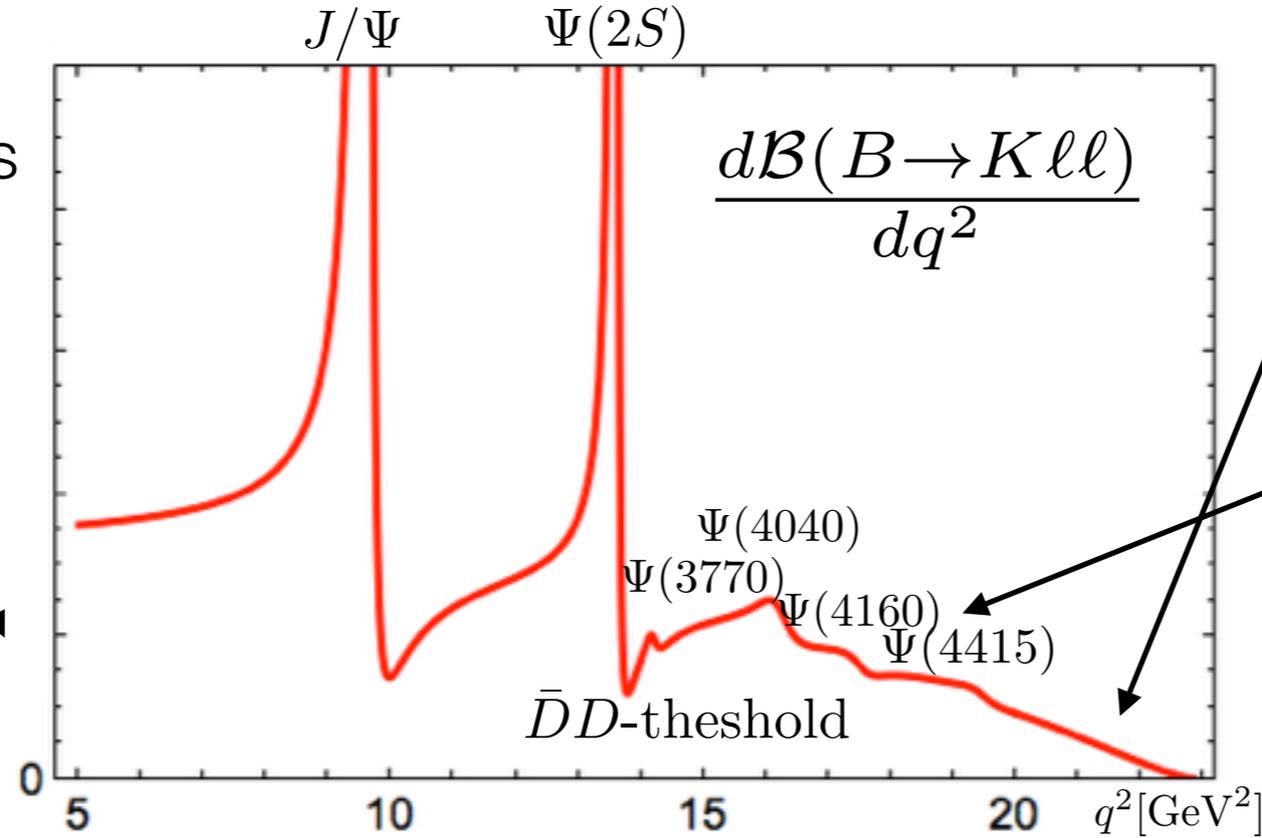
4-quark operators (also $O_{3..6}$)

K **fast**:

- **light-cone** methods
LCSR, QCDF/SCET

K **slow**:

- high- q^2 “**OPE**”
- **endpoint relations**

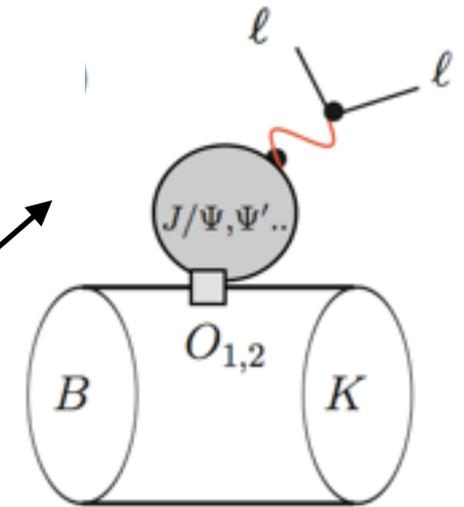
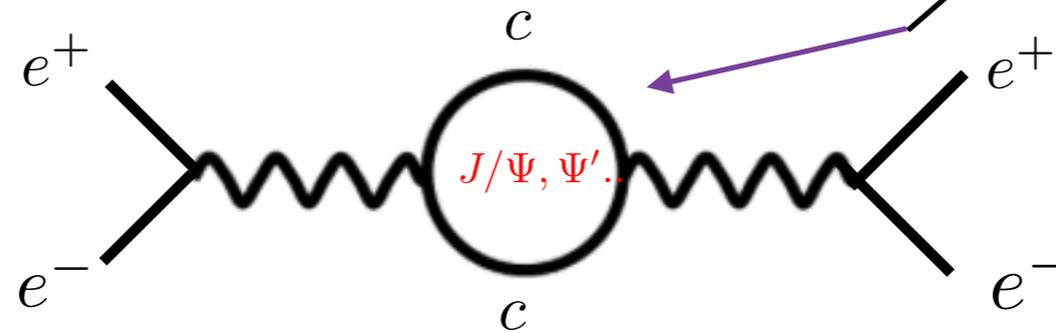


diagnostic shape
for charm

$O_{7,9}^2$ -dominates narrow resonances O_9^2 -dominates
 O_2 - $O_{7,9}$ -interference $(O_2)^2$ -effect O_2 - O_9 -interference

A. does factorisation describe $B \rightarrow Kll$ data?

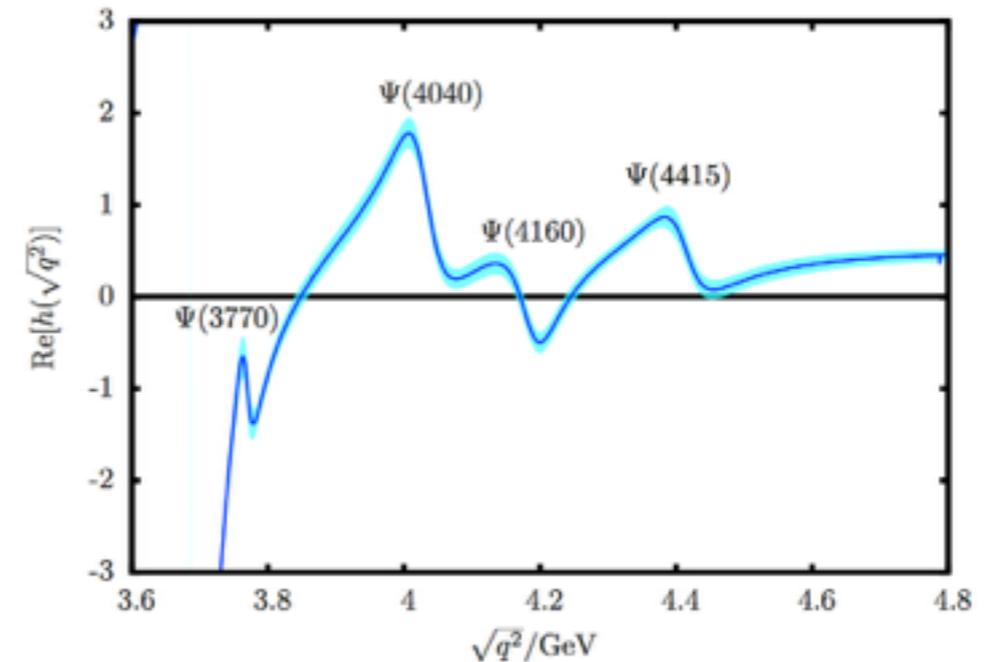
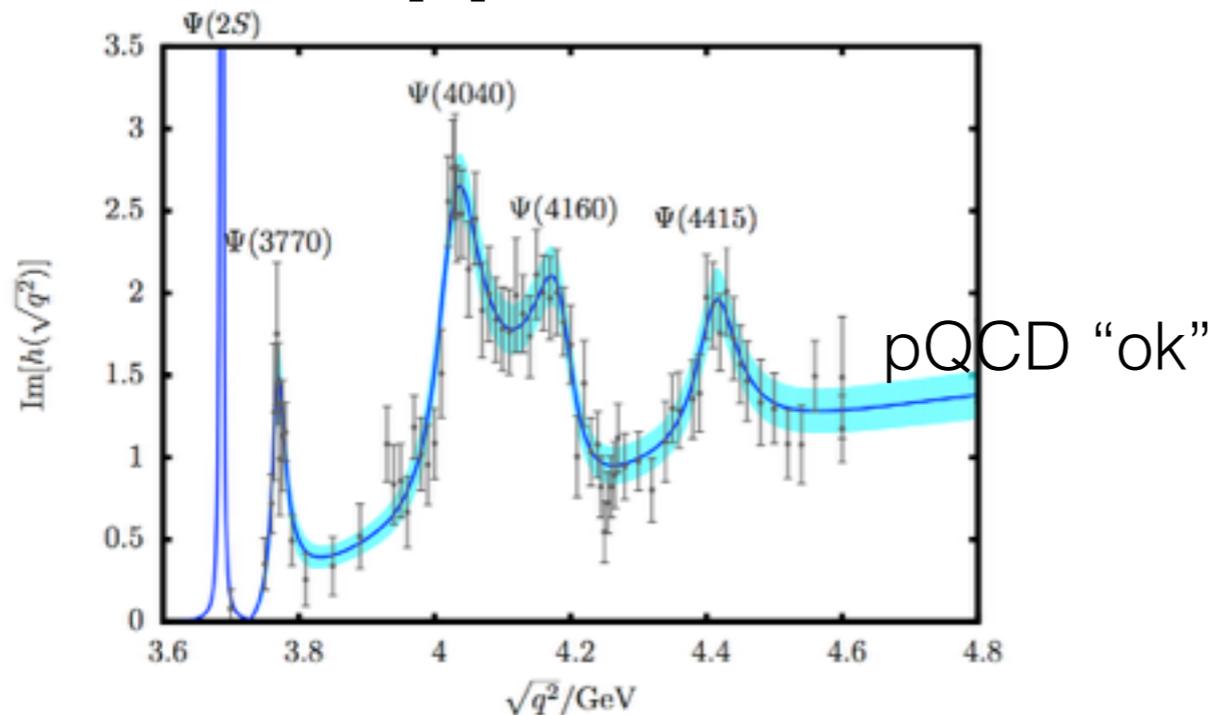
- vac. pol. $h(q^2)$ from $e^+e^- \rightarrow \text{hadrons}$ as for $(g-2)$



Disc $\sim \text{Im}[h]$; BESII-data'PLB08

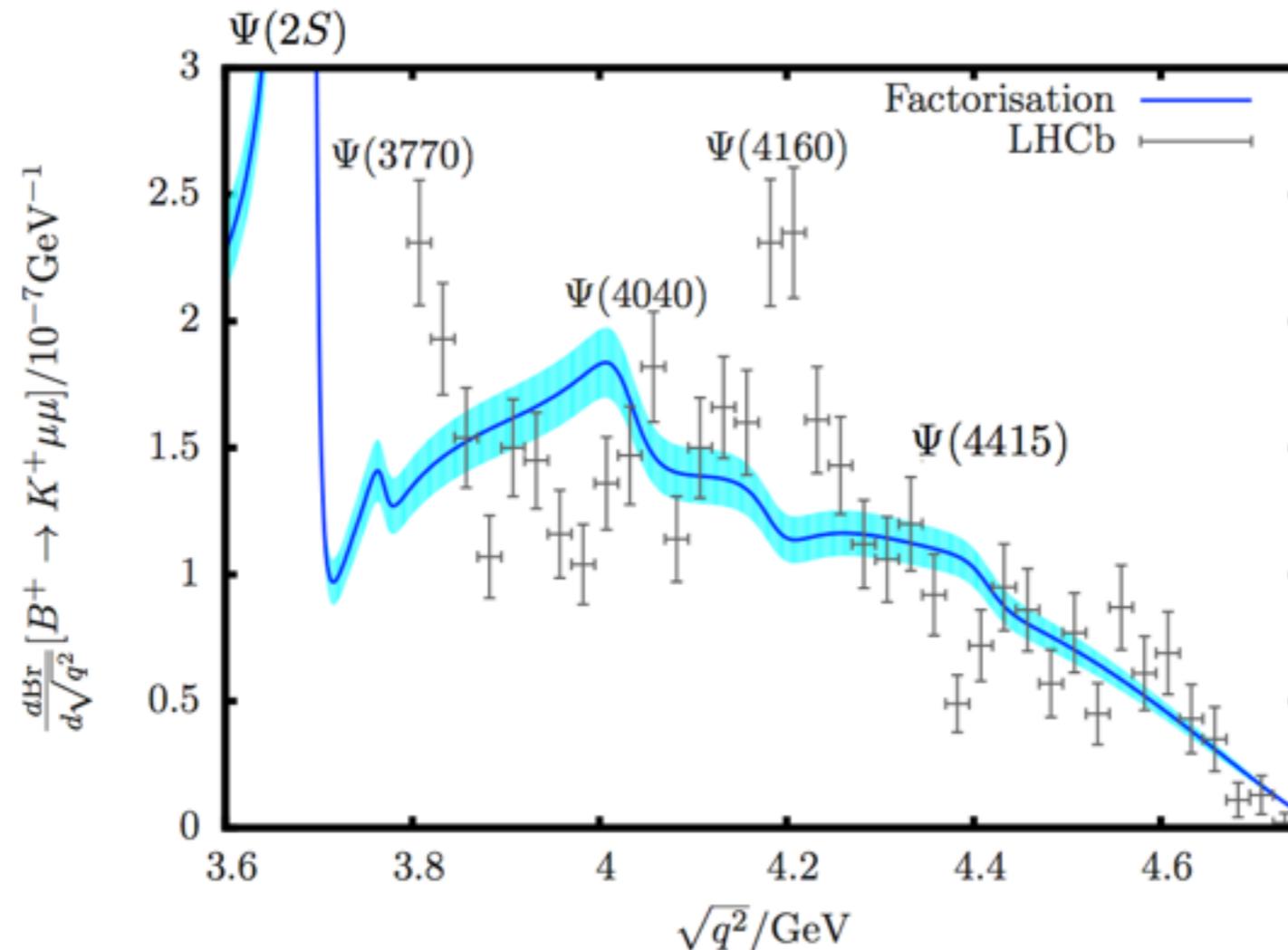


$\text{Re}[h]$ **dispersion relation**



our $\chi^2/\text{dof} = 1.015$

Factorisation (BESII-data) applied to $B \rightarrow K\ell\ell$ at high q^2



**clear failure of
factorisation**

clarifying status of factorisation of importance since:

- factorisation used estimate of “duality violations”
- perturbative factorisation used in most high- q^2 OPE predictions

B. probing non-factorisable effects

- think resonances described Breit-Wigner

N.B. 1) location of **pole**
& 2) **residue** are **physical!**

$$\mathcal{A}(B \rightarrow K \ell \ell) |_{q^2 \simeq m_\Psi^2} = \frac{\mathcal{A}(B \rightarrow \Psi K) \mathcal{A}^*(\Psi \rightarrow \ell \ell)}{q^2 - m_\Psi^2 + i m_\Psi \Gamma_\Psi} + \dots$$

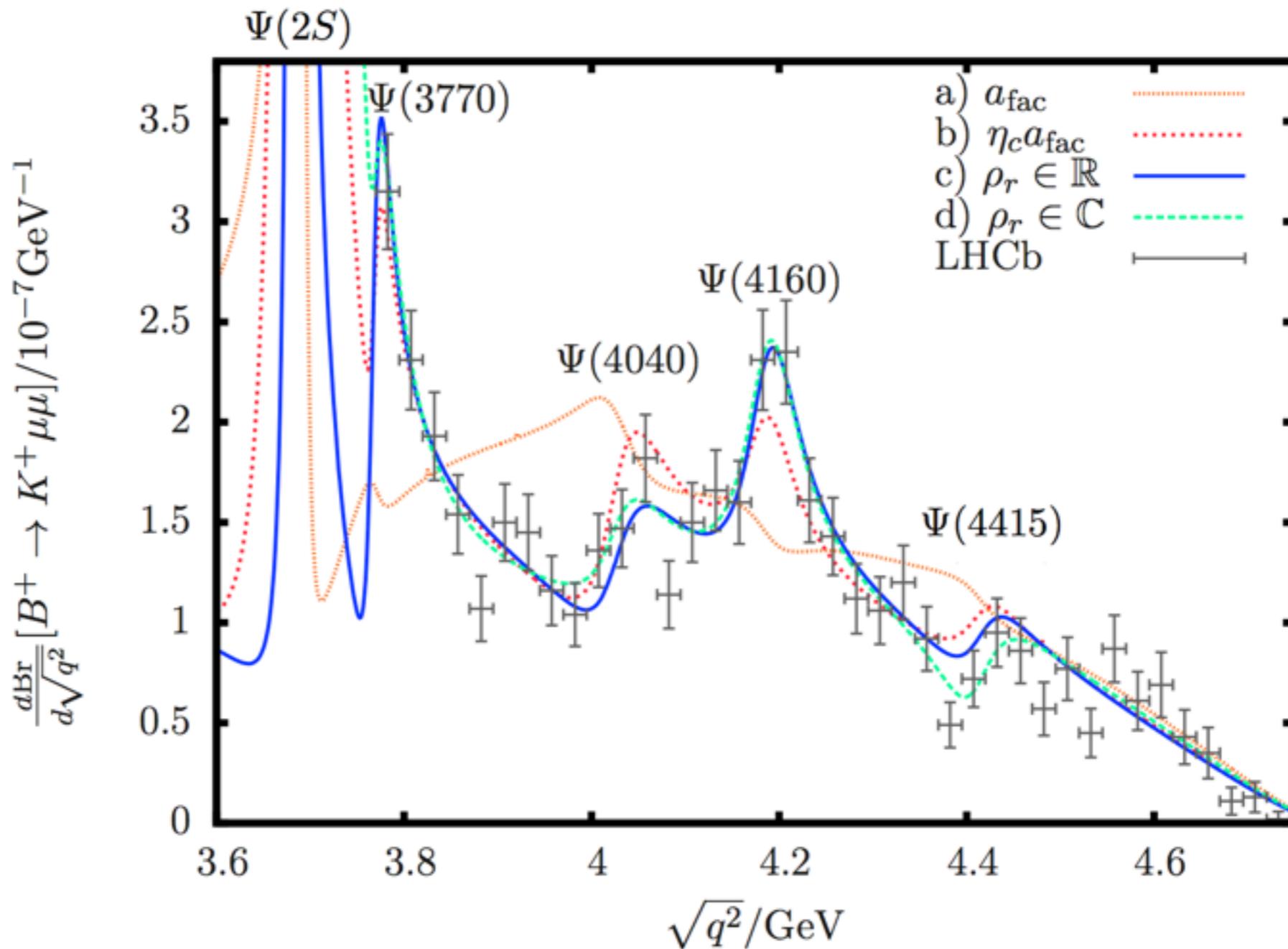
- idea: correct for Ψ -production (residue physical)

$$\begin{aligned} \mathcal{A}(B \rightarrow \Psi K) |_{\text{fac}} &\sim f_+^{B \rightarrow K}(q^2) \mathcal{A}(\Psi \rightarrow \ell \ell) \\ &\rightarrow f_+^{B \rightarrow K}(q^2) \underbrace{\eta_\Psi}_{1+\text{non-fac}} \mathcal{A}(\Psi \rightarrow \ell \ell) \sim \mathcal{A}(B \rightarrow \Psi K) \end{aligned}$$

- fits η_Ψ : b) global (scaled)fac; c) real-variable; d) complex-variable

only option d) sensible a priori

results ...



Fit	η_B	η_c	$\rho_{\Psi(2S)}$	$\rho_{\Psi(3770)}$	$\rho_{\Psi(4040)}$	$\rho_{\Psi(4160)}$	$\rho_{\Psi(4415)}$	$\chi^2/\text{d.o.f.}$	d.o.f.	pts	p-value
■ a)	0.98	$\equiv 1$	$\equiv 1$	$\equiv 1$	$\equiv 1$	$\equiv 1$	$\equiv 1$	3.59	99	117	$\simeq 10^{-30}$
■ b)	1.08	-2.55	$\equiv 1$	$\equiv 1$	$\equiv 1$	$\equiv 1$	$\equiv 1$	1.334	98	117	1.5%
■ c)	0.81	$\equiv 1$	-1.3	-7.2	-1.9	-4.6	-3.0	1.169	94	117	12%
■ d)	1.06	$\equiv 1$	$3.8-5.1i$ $6.4e^{-i53.3^\circ}$	$-0.1-2.3i$ $2.0e^{-i92^\circ}$	$-0.5-1.2i$ $1.3e^{-i111^\circ}$	$-3.0-3.1i$ $4.3e^{-i135^\circ}$	$-4.5+2.3i$ $5.1e^{i153^\circ}$	1.124	89	117	20%

2) Implications and contemplation

real question is whether this is SM(QCD) or not

right-handed currents (RHC) vs non-universal polarisation in $B \rightarrow K^* \Pi$

- issue imminent from structure of **helicity amplitudes**

$$H_0^V \sim (C_9 - C_9') \hat{H}_0^V(q^2) + \dots, \quad H_{\parallel}^V \sim (C_9 - C_9') \hat{H}_{\parallel}^V(q^2) + \dots, \quad H_{\perp}^V \sim \sqrt{\lambda_{K^*}} (C_9 + C_9') \hat{H}_{\perp}^V(q^2) + \dots,$$

RHC $C_9' \neq 0$ intertwined polarisation effects $0, \parallel, \perp$

- question: do fac and non-fac depend differently on polarisation?

$$\frac{|H_0^V|}{|H_{\parallel}^V|} \stackrel{?}{\sim} \frac{|f_0^V|}{|f_{\parallel}^0|} \quad \text{for some } q^2, \quad f \text{ form factor}$$

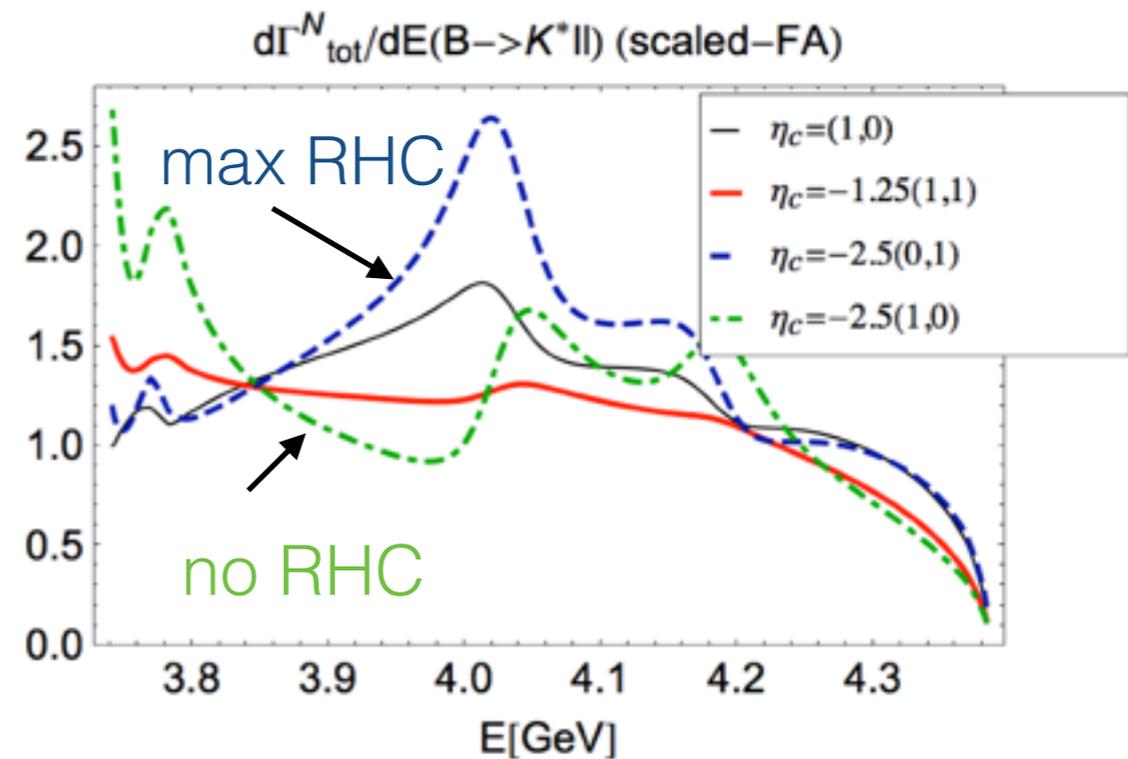
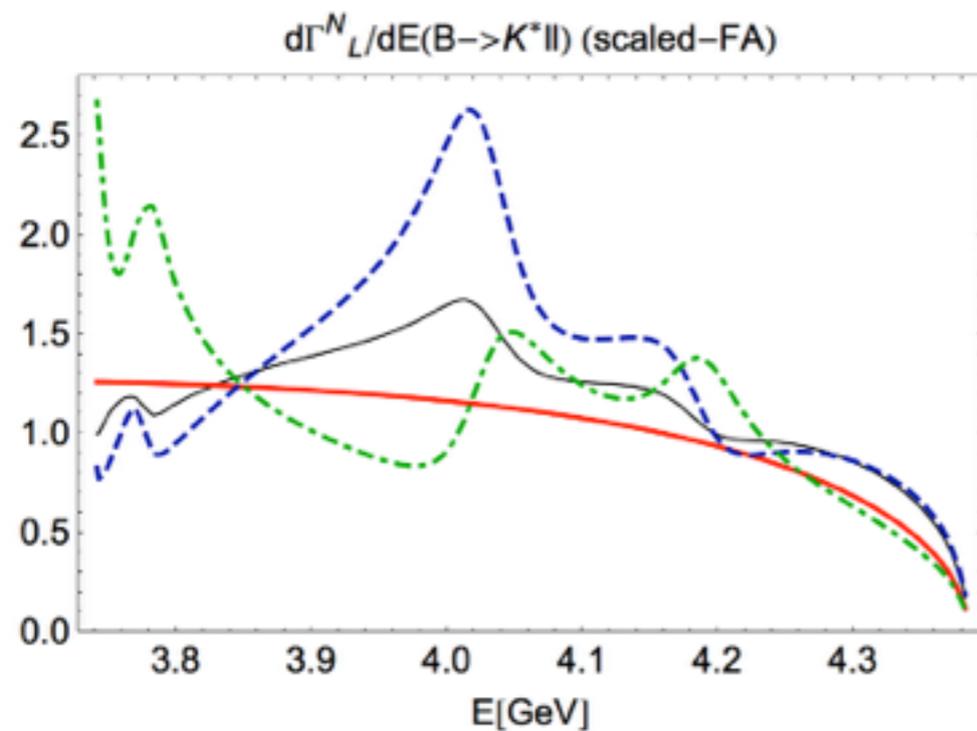
S-state: J/Ψ ok, $\Psi(2S)$ okish,

P-state: χ_{c1} broken

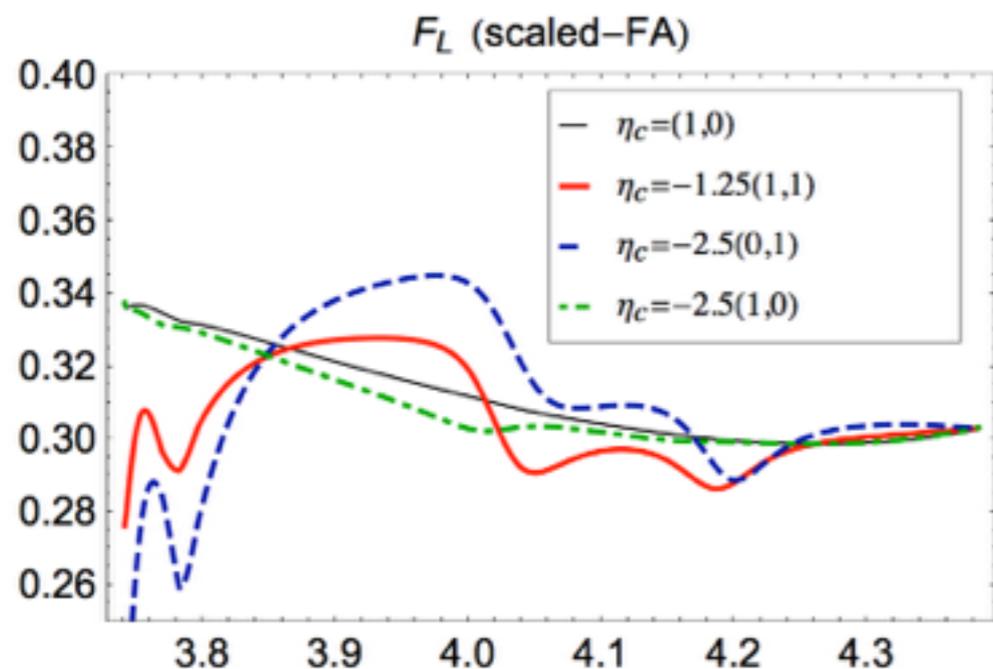
D-state: $\Psi(3370), \Psi(4160)$? — experimentally accessible

what is the pattern?

- if **polarisation universal** then $\text{Br}_X(B \rightarrow K^* \ell \ell)$ $X=L, \text{tot}$ good observable to test for right-handed currents*



- if **polarisation universal** and **no RHC** then resonance effect minimal in class of observables **Hiller and RZ'13**



$\frac{1}{3}$

e.g. **black** and **green** curve nearly **identical** even though green curve has 2.5 as much resonances!
N.B. endpoint all curves asymptotes $\frac{1}{3}$

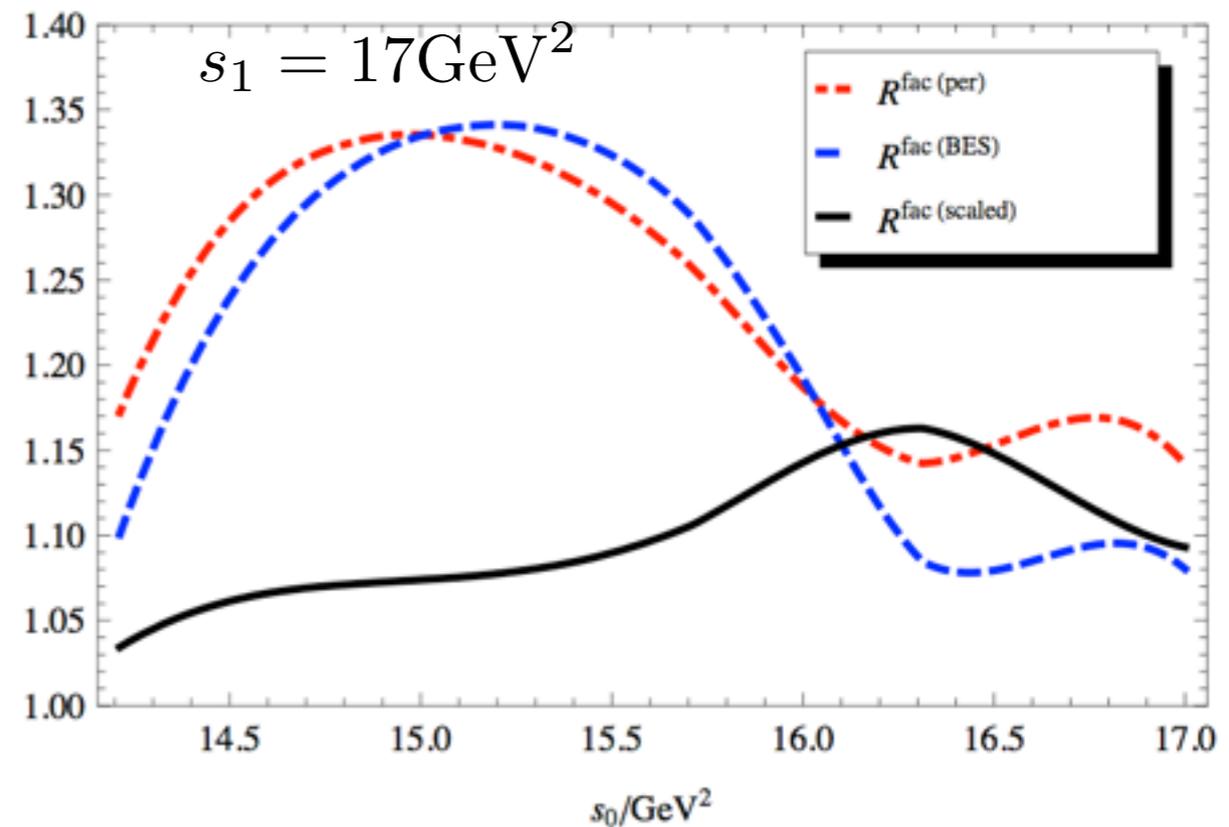
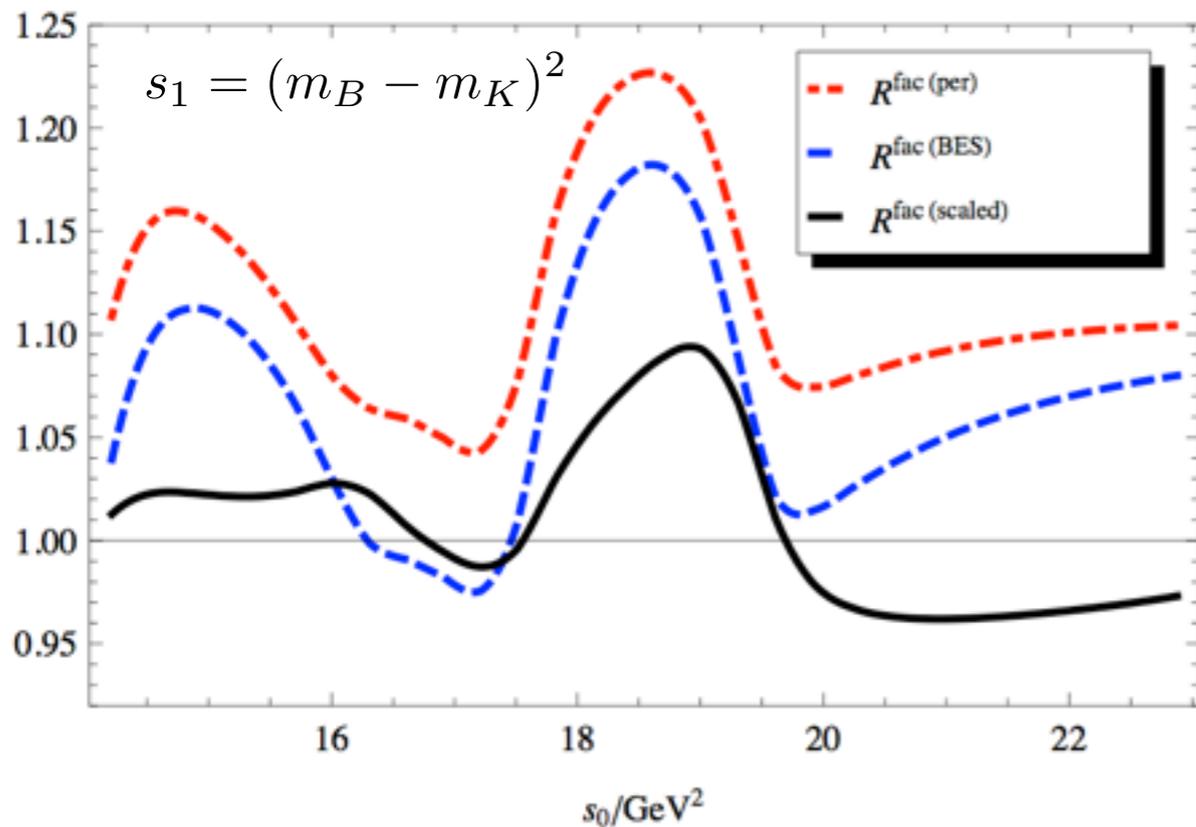
* assumes effect same magnitude in $B \rightarrow K^* \ell \ell$ (could be bit smaller or larger in reality)

Br(B → Kll) high q²: a priori and a posteriori

- ratio of Br(B → Kll) using
 - factorisation perturbative (no resonances)
 - factorisation (BES-data)
 vs data as function lower bin bdry s₀

$$\frac{\text{Br}(B^+ \rightarrow K^+ \ell\ell)_{[s_0, s_1]}^{i), ii)}}{\text{Br}(B^+ \rightarrow K^+ \ell\ell)_{[s_0, s_1]}^{fit-d)}}$$

basically as good as data (by construction) →

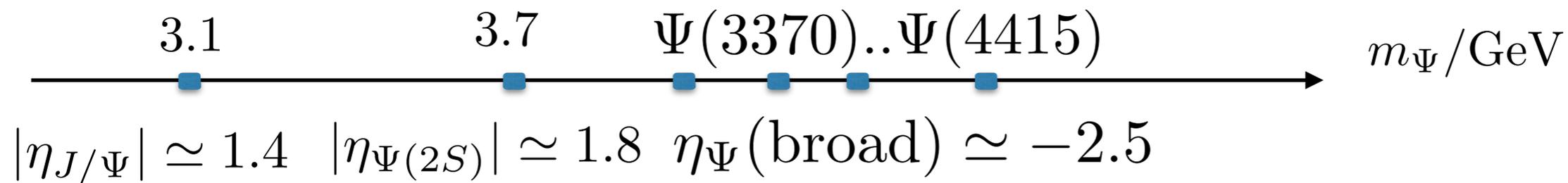


if assume that effect is **SM** ⇒

one has to give some thought to this issues when making predictions

Factorisation as a function of m_Ψ

- for a long time there has been information on $B \rightarrow J/\Psi K$ and $B \rightarrow \Psi(2S)K$ quantify failure of factorisation by scale factor $\eta_\Psi = 1 + \text{non-fac}$



- whereas corr. to J/Ψ , $\Psi(2S)$ could be 40%, 80% “only”, the **new** information is that on **average** there is **350%** correction for **broad $\Psi(3770) - \Psi(4415)$**
- N.B magnitude 2.5 not a big surprise but that they
 - all have “same sign”** & ii) **sign negative****challenges quark-hadron duality*** (nominal core 50%)

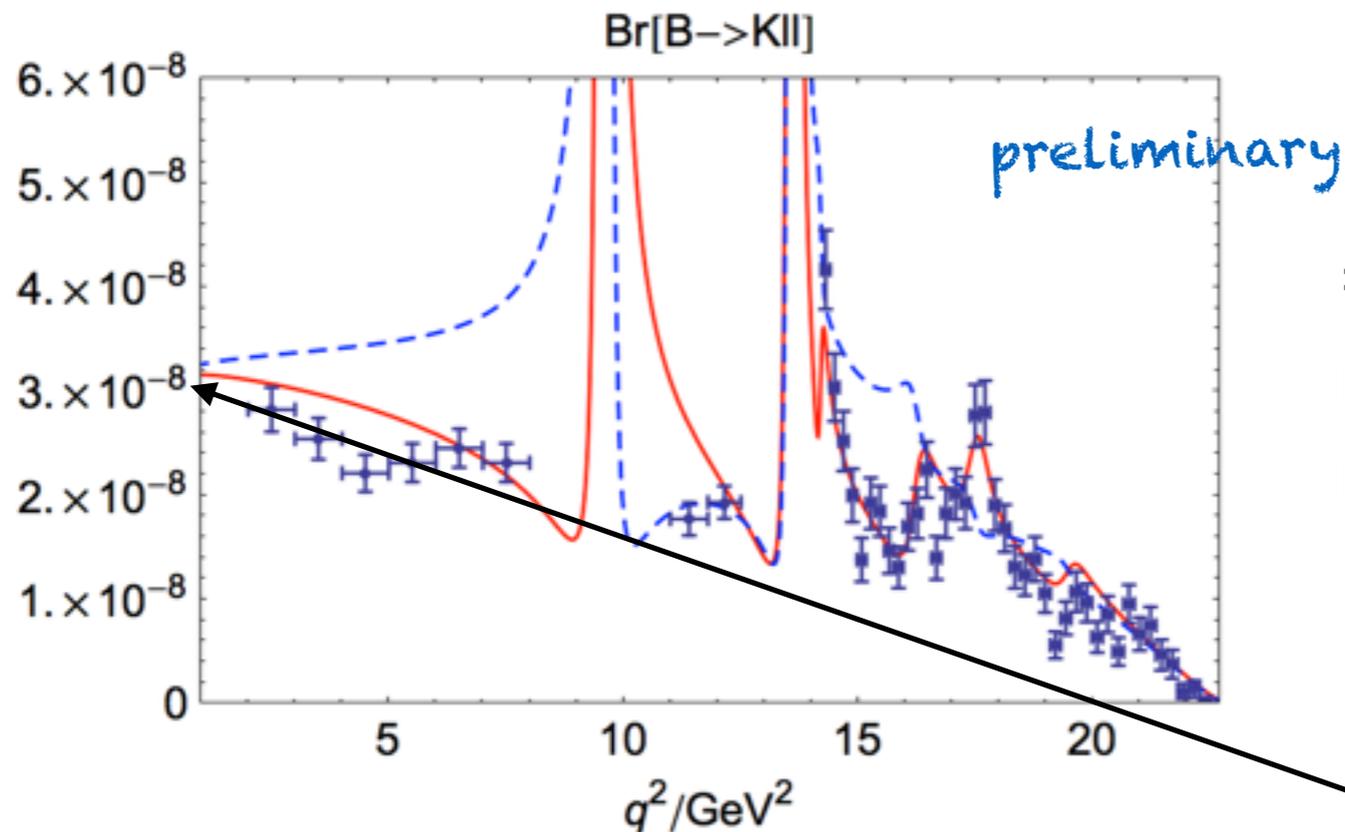
is it really SM(QCD)?
or new bccs-operators?
- N.B. phase of J/Ψ is unknown $\Psi(2S)$ marginal knowledge from fit phase influences low q^2 observables in bin close to J/Ψ

* more in paper and in backup slides (no time for discussion here)

the unknown J/Ψ phase

$$\eta_{J/\Psi} = |\eta_{J/\Psi}| e^{i\delta_{J/\Psi}} \simeq 1.4 e^{i\delta_{J/\Psi}}$$

- to match/fit slope of pQCD charm $\delta_{J/\psi} \simeq 0$ e.g. Khodjamirian et al'10 and others
- let's change phase to $\delta_{J/\psi} \simeq \pi$ and compare with $\text{Br}(B \rightarrow Kll)$



\Rightarrow empirically $\delta_{J/\psi} \simeq \pi$
 not absurd (even slightly favoured)
 not as conclusive as high q^2

- $\delta_{J/\psi} \simeq \pi$ matched charm amplitude to SM at $q^2 = 0$
 well but then slope of charm amplitude has wrong sign!
 \Rightarrow more precise data binning

relation to P_5'

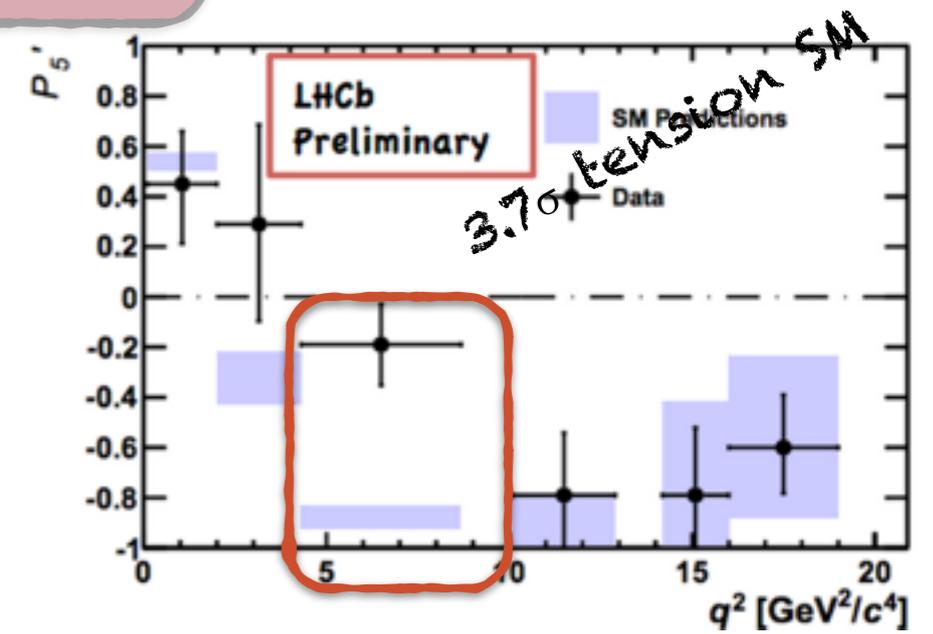
preliminary sketch

- change to $\delta_{J/\psi} \approx \pi$ will affect other observables and the decisive q^2 -spectrum is the thing to look out for
- let's define "average shift" of C_9 as function of the bin

$$\Delta C_9|_{s_0, s_1}^K = a_{\text{fac}} \int_{s_0}^{s_1} ds (h_c^{\delta_{J/\psi}=\pi}(s) - h_c^{pQCD}(s))$$

$[s_0, s_1]/\text{GeV}^2$	$[4.3, 8.68]$	$[2, 4.3]$	$[1, 6]$
$\Delta C_9 _{s_0, s_1}^K$	-2	-0.5	-0.6

- $[4.3, 8.68]: (\Delta C_9, \Delta C_9') = (-1.6, -1.4) \Rightarrow \Delta P_5' \approx -0.4$
 recall that $\Delta C_9'$ not significant
 hence $\Delta C_9 \approx -2$ can close to LHCb-value*



* effect larger in $B \rightarrow K^* l l$ as breaking of factorisation stronger $B \rightarrow J/\psi K^*$ than in $B \rightarrow J/\psi K$

conclusions and summary

- **factorisation** approximation **fails** spectacularly (resonances same sign) \Rightarrow **pressure SM(QCD)** in view of quark-hadron duality (backup slides)
- can there be **new physics in bsc operators?** (contrived)
what are the constraints? $b \rightarrow s\gamma$, $\Delta\Gamma_s$, ..
- **If** polarisation universality is established (reasonable level) \Rightarrow test for **right-handed currents** in bccs-operators in $\text{Br}_{L,\text{tot}}(B \rightarrow K^* \ell \ell)$
- How distinguish short-distance from long-distance* new physics?
If non form factor q^2 -dependence \Rightarrow long-distance new physics*
- change in $\delta_{J/\psi} \approx \pi$ (empirically unknown) fits **shape** and magnitude of $\text{Br}(B \rightarrow K \ell \ell)$ low q^2 and also looks promising for P_5'
- whereas charm troubles short-distance new physics searches
 - i) of course there is room for short-distance new physics in C_9^{eff}
 - ii) progress in form factor correlations (backup) should help in searches
 - iii) charm resonances are lepton-universal \Rightarrow no relation to R_K

Thanks for your attention

* long distance new physics: $\delta H_{NP} = \delta C O$ which contributes to long-distance physics

backup slides

comment on form factor correlations

Use of equation of motion for form factors

- Consider QCD e.o.m./Ward-identity (*study correction Isgur-Wise relations*)

Grinstein Pirjol'04

$$i\partial^\nu (\bar{s} i\sigma_{\mu\nu}(\gamma_5)b) = -(m_s \pm m_b)\bar{s}\gamma_\mu(\gamma_5)b + i\partial_\mu(\bar{s}(\gamma_5)b) - 2\bar{s}i\overleftarrow{D}_\mu(\gamma_5)b$$

- Evaluate on $\langle K^* | \dots | B \rangle$ get 4 independent equations e.g.

$$T_1(q^2) - \frac{(m_b + m_s)}{m_B + m_{K^*}} V(q^2) + \mathcal{D}_1(q^2) = 0$$

- 1) any determination of form factors must satisfy e.o.m.
2) Correlation function lattice/LCSR are compatible e.o.m. up to irrelevant contact terms

Hambrock, Hiller, Schacht, Zwicky '13
Bharucha, Straub, Zwicky'14 (to appear)

$$T_1(q^2) - \frac{(m_b + m_s)}{m_B + m_{K^*}} V(q^2) + \mathcal{D}_1(q^2) = 0$$

- 1) denote $F(q^2)_{s_0^F, M_F^2}$, s_0^F threshold, M_F^2 Borel parameter
then compatible with eom $s_0^{T_1} = s_0^V = s_0^{\mathcal{D}_1}$ and $M_{T_1}^2 = M_V^2 = M_{\mathcal{D}_1}^2$
- 2) observe $T_1, V \gg \mathcal{D}_1$ (5% maximal) over q^2 -range $[0, 15] \text{ GeV}^2$ *

- even associate 40% uncertainty to \mathcal{D}_1 then ratio

$$r_{\perp} = \frac{(m_b + m_s)}{m_B + m_{K^*}} \frac{V(q^2)}{T_1(q^2)} \quad \text{determined up to 2\%}$$

Crucial for $B \rightarrow K^* \parallel$ pheno as determines zero of helicity amplitude

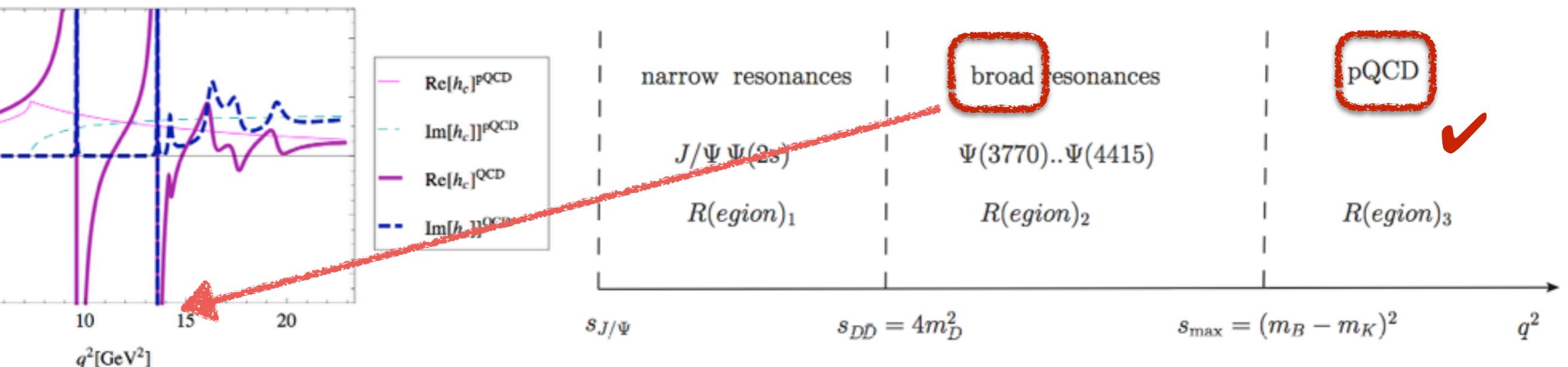
* means that s_0 and M^2 of T_1 and V highly correlated

Quark-hadron duality

- colloquially: “when smeared (integrated) over large enough interval quark and hadrons lead to quantitatively similar results”
- formal level only approach (known to me) via dispersion relations (DR)
non-fac. correction obey same (verified) DR as factorisable part

$$H^{V,X}(s) = H^{V,X}(s_0) + \frac{(s - s_0)}{2\pi i} \int_{s_{J/\Psi}}^{\infty} \frac{dt}{t - s_0} \frac{\text{Disc}[H^{V,X}](t)}{t - s - i0}, \quad X \in \{\text{fac}, \text{cor}\}.$$

factorisable charmloop $h_c(q^2)$

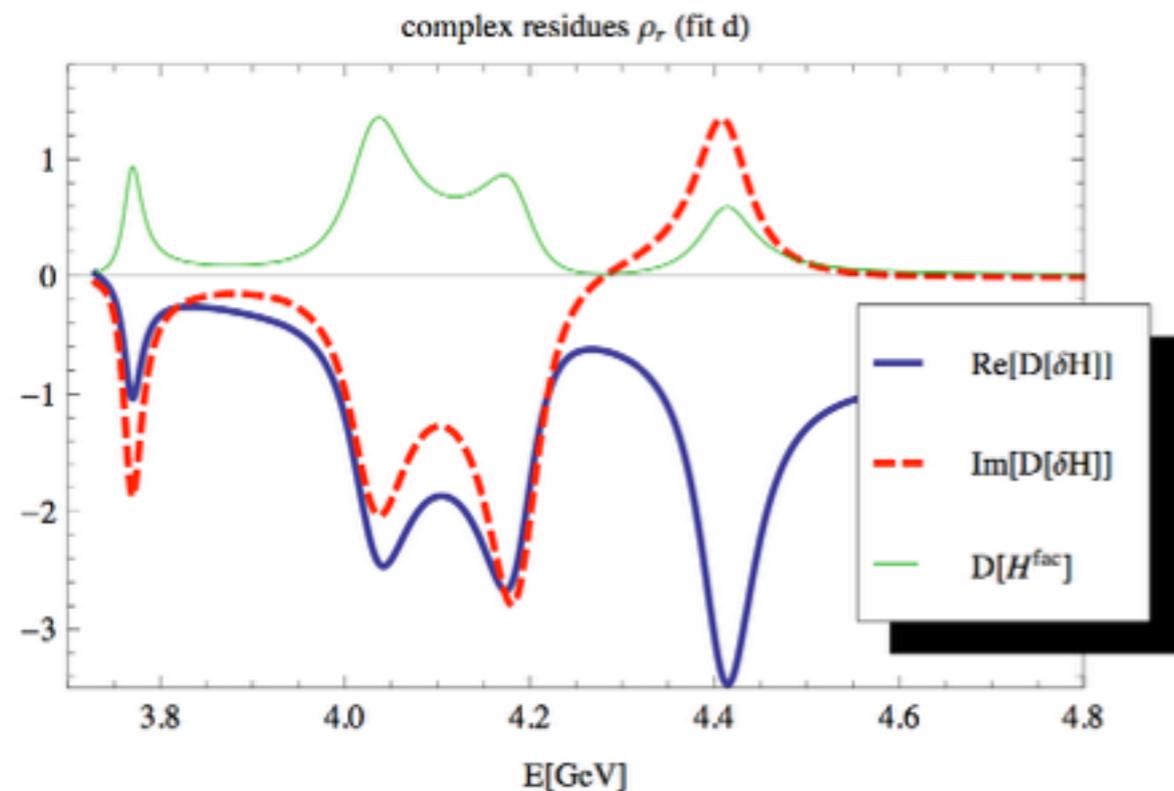


second conclusions of phase 2:

- factorisable part duality over open charm is ok less 10% error
- could expect same pattern for non-factorisable part

possible pitfall — (non)-positivity

- factorisable part is positive definite (related cross section)
non-factorisable part isn't (that's what we tested for in fit d)



- fit d) effect of cancellations 20% instead of 350% its 280%
 \Rightarrow that's not it!

third conclusions of phase 2:

- in our analysis we have not come across any signs that QCD can explain this effect. Yet charm-physics has a reputation of being notoriously difficult.
- extending the duality interval over J/Ψ and $\Psi(2S)$ resonances is desirable in view of the results how to improve